# Notes on determining equations

Johannes Borgqvist

October 11, 2022

#### Contents

6	The fourth determining equation	6
	$5.3  l=2  \dots  \dots  \dots  \dots  \dots  \dots$	
	$5.2  l = 1  \dots  \dots  \dots  \dots  \dots  \dots  \dots  \dots  \dots $	
	5.1 $l=0$	
5	The third determining equation	3
4	The second determining equation	2
3	The first determining equation	2
2	Four determining equations	2
1	Introduction	1

#### 1 Introduction

We analyse the following second order travelling wave ODE

$$\frac{\mathrm{d}}{\mathrm{d}z} \left( \frac{1}{u(z)^{\ell}} \frac{\mathrm{d}u}{\mathrm{d}z} \right) + c \frac{\mathrm{d}u}{\mathrm{d}z} + f(u(z)) = 0.$$

with  $f(u) = c_0 + c_1 u + c_2 u^2 + c_3 u^3$  and  $\ell$  is an arbitrary index. This is a generalisation of the model in [1] that describes *cell migration in an epithelial tissue*, and this particular model is retrieved by the choices  $\ell = 2$  and f(u) = u(1 - u).

By denoting derivatives by du/dz = u' this ODE can be re-written as follows:

$$u'' - \frac{\ell(u')^2}{u} + cu'u^{\ell} + u^{\ell}f(u) = 0.$$
 (1)

Now, we are interested in an infinitesimal generator of the Lie group

$$X = \xi(z, u)\partial_z + \eta(z, u)\partial_u$$

which has a second prolongation given by

$$X^{(2)} = \xi(z, u)\partial_z + \eta(z, u)\partial_u + \eta^{(1)}(z, u, u')\partial_{u'} + \eta^{(2)}(z, u, u', u'')\partial_{u''}.$$

Here, the two prolonged infinitesimals  $\eta^{(1)}$  and  $\eta^{(2)}$  are given by [2]

$$\eta^{(1)}(z, u, u') = \eta_z + (\eta_u - \xi_z)u' - \xi_u(u')^2, \qquad (2)$$

$$\eta^{(2)}(z, u, u', u'') = \eta_{zz} + (2\eta_{zu} - \xi_{zz})u' + (\eta_{uu} - 2\xi_{zu})(u')^2 - \xi_{uu}(u')^3 + \{\eta_u - 2\xi_z - 3\xi_u u'\}u''.$$
(3)

The linearised symmetry condition for our ODE of interest is given by

$$\eta^{(2)} + \left(cu^{\ell} - \frac{2\ell u'}{u}\right)\eta^{(1)} + \left(\frac{\ell(u')^{2}}{u^{2}} + c\ell u^{\ell-1}u' + u^{\ell-1}\left(\ell f(u) + u\frac{\mathrm{d}f}{\mathrm{d}u}\right)\right)\eta = 0$$
whenever  $u'' - \frac{\ell(u')^{2}}{u} + cu'u^{\ell} + u^{\ell}f(u) = 0.$  (4)

Now, we will try to solve the linearised symmetry condition.

## 2 Four determining equations

By calculating the prolonged infinitesimals  $\eta^{(1)}$  and  $\eta^{(2)}$ , plugging these into this linearised symmetry condition, and then organising the resulting equation in terms of powers of u' results in the following four so called *determining equations* 

$$(u')^3: \quad \xi_{uu} + \frac{\ell}{u} \xi_u = 0, \tag{5}$$

$$(u')^{2}: \quad 2c\xi_{u}u^{\ell} + \eta_{uu} - 2\xi_{zu} - \frac{(\eta_{u} - 2\xi_{z})\ell}{u} + \frac{\ell\eta}{u^{2}}$$
 = 0, (6)

$$u': cu^{\ell}\xi_z + 3u^{\ell}\xi_u f(u) + 2\eta_{zu} - \xi_{zz} - \frac{2\ell}{u}\eta_z + c\ell u^{\ell-1}\eta = 0,$$
 (7)

1: 
$$u^{\ell-1} \left( \ell f(u) + u \frac{\mathrm{d}f}{\mathrm{d}u} \right) \eta + c u^{\ell} \eta_z + \eta_{zz} + u^{\ell} f(u) (2\xi_z - \eta_u) = 0.$$
 (8)

Now, we will treat these four equations systematically, and solve them one by one. We will use our friend SymPy to do this.

## 3 The first determining equation

Ok, so we are interested in the following PDE

$$\xi_{uu} + \frac{\ell}{u}\xi_u = 0.$$

This one we solve by hand which gives us the following equation

$$\xi(z,u) = A(z) \left(\frac{1}{1-\ell}\right) u^{1-\ell} - B(z) \tag{9}$$

where  $A, B \in \mathcal{C}^{\infty}(\mathbb{R})$ .

## 4 The second determining equation

Ok, so we are interested in solving the following PDE

$$2c\xi_u u^{\ell} + \eta_{uu} - 2\xi_{zu} - \frac{(\eta_u - 2\xi_z)\ell}{u} + \frac{\ell\eta}{u^2} = 0$$

for the unknown tangent  $\eta(z, u)$ . We plugged this PDE into Wolphram Alpha, and out came the following suggested solution

$$\eta(z,u) = 2cA(z) \left(\frac{1}{\ell-2}\right) u^2 + A'(z) \left(\frac{2\ell-1}{(\ell-1)^2}\right) u^{2-\ell} + 2B'(z) \left(\frac{\ell}{\ell-1}\right) u \ln(u) + C(z) u^{\ell} + D(z) u$$
(10)

where  $C, D \in \mathcal{C}^{\infty}(\mathbb{R})$  are two new arbitrary functions.

## 5 The third determining equation

Ok, thus far we have landed at the following infinitesimals:

$$\xi(z,u) = A(z) \left(\frac{1}{1-\ell}\right) u^{1-\ell} - B(z), \tag{11}$$

$$\eta(z,u) = 2cA(z) \left(\frac{1}{\ell-2}\right) u^2 + A'(z) \left(\frac{2\ell-1}{(\ell-1)^2}\right) u^{2-\ell} + 2B'(z) \left(\frac{\ell}{\ell-1}\right) u \ln(u) + C(z) u^{\ell} + D(z) u.$$
(12)

Next, we want to use following PDE

$$cu^{\ell}\xi_{z} + 3u^{\ell}\xi_{u}f(u) + 2\eta_{zu} - \xi_{zz} - \frac{2\ell}{u}\eta_{z} + c\ell u^{\ell-1}\eta = 0$$

in order to get equations that we can solve for the four unknown functions  $A, B, C, D \in \mathcal{C}^{\infty}(\mathbb{R})$ . So what we can do here is to plug in our unknown functions, and the we see that the above equation entails finding the roots of a polynomial in u. Hence, the equation decomposes into a set of sub equations which we can solve individually.

Here, we are going to assume a general cubic reaction term

$$f(u) = c_0 + c_1 u + c_2 u^2 + c_3 u^3. (13)$$

where  $c_0, c_1, c_2, c_3 \in \mathbb{R}$  are four arbitrary constants.

By plugging in the tangents, into the third linearised symmetry condition, we get the following

equation:

$$2A(z)c^{2}l^{4}u^{4}u^{l} - 6A(z)c^{2}l^{3}u^{4}u^{l} \\ + 6A(z)c^{2}l^{2}u^{4}u^{l} - 2A(z)c^{2}lu^{4}u^{l} \\ - 4A(z)cl^{4}u^{4} + 20A(z)cl^{3}u^{4} - 36A(z)cl^{2}u^{4} \\ + 28A(z)clu^{4} - 8A(z)cu^{4} + 3A(z)c_{0}l^{4}u^{3} \\ - 15A(z)c_{0}l^{3}u^{3} + 27A(z)c_{0}l^{2}u^{3} - 21A(z)c_{0}lu^{3} + 6A(z)c_{0}u^{3} \\ + 3A(z)c_{1}l^{4}u^{4} - 15A(z)c_{1}l^{3}u^{4} + 27A(z)c_{1}l^{2}u^{4} \\ - 21A(z)c_{1}lu^{4} + 6A(z)c_{1}u^{4} + 3A(z)c_{2}l^{4}u^{5} \\ - 15A(z)c_{2}l^{3}u^{5} + 27A(z)c_{2}l^{2}u^{5} - 21A(z)c_{2}lu^{5} + 6A(z)c_{2}u^{5} + 3A(z)c_{3}l^{4}u^{6} \\ - 15A(z)c_{3}l^{3}u^{6} + 27A(z)c_{3}l^{2}u^{6} - 21A(z)c_{3}lu^{6} + 6A(z)c_{3}u^{6} \\ - 7A''(z)l^{3}u^{4}u^{-1} + 24A''(z)l^{2}u^{4}u^{-1} - 23A''(z)lu^{4}u^{-1} + 6A''(z)l^{4}u^{3} \\ + A'(z)cl^{3}u^{4} - A'(z)cl^{2}u^{4} - 3A'(z)clu^{4} + 2A'(z)cu^{4} - 4B''(z)l^{5}u^{3}\ln(u) \\ + 20B''(z)l^{4}u^{3}\ln(u) + 3B''(z)l^{4}u^{3} \\ - 36B''(z)l^{3}u^{3}\ln(u) - 11B''(z)l^{3}u^{3} + 28B''(z)l^{2}u^{3}\ln(u) \\ + 11B''(z)l^{2}u^{3} - 8B''(z)lu^{3}\ln(u) - B''(z)lu^{3} - 2B''(z)u^{3} \\ + 2B'(z)cl^{5}u^{3}u^{l}\ln(u) \\ - 8B'(z)cl^{4}u^{3}u^{l}\ln(u) + 5B'(z)cl^{4}u^{3}u^{l} \\ - 4B'(z)cl^{2}u^{3}u^{l}\ln(u) + 9B'(z)cl^{2}u^{3}u^{l} \\ - 7B'(z)clu^{3}u^{l} + 2B'(z)cu^{3}u^{l} + C(z)cl^{5}u^{2}u^{2} - 5C(z)cl^{4}u^{2}u^{2} \\ + 9C(z)cl^{3}u^{2}u^{2} - 7C(z)cl^{2}u^{2}u^{2} + 2C(z)clu^{2}u^{2} + D(z)cl^{5}u^{3}u^{l} \\ - 7D(z)cl^{2}u^{3}u^{l} + 2D(z)clu^{3}u^{l} - 2D'(z)l^{5}u^{3} \\ + 12D'(z)l^{4}u^{3} - 28D'(z)l^{3}u^{3} \\ + 32D'(z)l^{2}u^{3} - 18D'(z)lu^{3} + 4D'(z)u^{3} = 0.$$

We see that in Equation (14), we have a lot of terms involving the function  $\ln(u)$ . So, we can actually further split this equation into two smaller equations based on the terms involving  $\ln(u)$  and the terms not involving  $\ln(u)$ . Here, is the equation where the terms involve  $\ln(u)$ :

$$-4B''(z)l^{5}u^{3} + 20B''(z)l^{4}u^{3} - 36B''(z)l^{3}u^{3}$$

$$+28B''(z)l^{2}u^{3} - 8B''(z)lu^{3} + 2B'(z)cl^{5}u^{3}u^{l}$$

$$-8B'(z)cl^{4}u^{3}u^{l} + 10B'(z)cl^{3}u^{3}u^{l} - 4B'(z)cl^{2}u^{3}u^{l} = 0.$$

$$(15)$$

Here, is the equation where the terms do not involve ln(u):

$$2A(z)c^{2}l^{4}u^{4}u^{l} - 6A(z)c^{2}l^{3}u^{4}u^{l} + 6A(z)c^{2}l^{2}u^{4}u^{l}$$

$$-2A(z)c^{2}lu^{4}u^{l} - 4A(z)cl^{4}u^{4} + 20A(z)cl^{3}u^{4}$$

$$-36A(z)cl^{2}u^{4} + 28A(z)clu^{4} - 8A(z)cu^{4}$$

$$+3A(z)c_{0}l^{4}u^{3} - 15A(z)c_{0}l^{3}u^{3} + 27A(z)c_{0}l^{2}u^{3}$$

$$-21A(z)c_{0}lu^{3} + 6A(z)c_{0}u^{3} + 3A(z)c_{1}l^{4}u^{4}$$

$$-15A(z)c_{1}l^{3}u^{4} + 27A(z)c_{1}l^{2}u^{4} - 21A(z)c_{1}lu^{4}$$

$$+6A(z)c_{1}u^{4} + 3A(z)c_{2}l^{4}u^{5} - 15A(z)c_{2}l^{3}u^{5}$$

$$+27A(z)c_{2}l^{2}u^{5} - 21A(z)c_{2}lu^{5} + 6A(z)c_{2}u^{5} +$$

$$3A(z)c_{3}l^{4}u^{6} - 15A(z)c_{3}l^{3}u^{6} + 27A(z)c_{3}l^{2}u^{6}$$

$$-21A(z)c_{3}lu^{6} + 6A(z)c_{3}u^{6} - 7A''(z)l^{3}u^{4}u^{-l}$$

$$+24A''(z)l^{2}u^{4}u^{-l} - 23A''(z)lu^{4}u^{-l} + 6A''(z)u^{4}u^{-l}$$

$$+2A'(z)cl^{3}u^{4} - A'(z)cl^{2}u^{4} - 3A'(z)clu^{4}$$

$$+2A'(z)cu^{4} + 3B''(z)l^{4}u^{3} - 11B''(z)l^{3}u^{3}$$

$$+11B''(z)l^{2}u^{3} - B''(z)lu^{3} - 2B''(z)u^{3}u^{l}$$

$$-7B'(z)clu^{3}u^{l} + 2B'(z)cu^{3}u^{l} + 9B'(z)cl^{2}u^{3}u^{l}$$

$$-7B'(z)clu^{3}u^{l} + 2B'(z)cl^{3}u^{2}l^{l} - 7C(z)cl^{2}u^{2}u^{2l}$$

$$+2C(z)clu^{2}u^{2l} + D(z)cl^{3}u^{3}u^{l} - 5D(z)cl^{4}u^{3}u^{l}$$

$$+9D(z)cl^{3}u^{3}u^{l} - 7D(z)cl^{2}u^{3}u^{l} + 2D(z)clu^{3}u^{l}$$

$$-2D'(z)l^{5}u^{3} + 12D'(z)l^{4}u^{3} - 28D'(z)l^{3}u^{3}$$

$$+32D'(z)l^{2}u^{3} - 18D'(z)lu^{3} + 4D'(z)u^{3} = 0.$$

Next, we will plug in the values l = 0, 1, 2 and then further split these equations in terms of the monomials.

#### **5.1** l = 0

$$u^{3}: 6A(z)c_{0} - 2B''(z) + 2B'(z)c + 4D'(z) = 0,$$

$$u^{4}: -8A(z)c + 6A(z)c_{1} + 6A''(z) + 2A'(z)c = 0,$$

$$u^{5}: 6A(z)c_{2} = 0,$$

$$u^{6}: 6A(z)c_{3} = 0.$$

From these determining equations, we see that

$$A(z) = 0 \forall z \in \mathbb{R}$$

and that

$$D(z) = \frac{1}{2} (B'(z) + B(z)) + K \quad K \in \mathbb{R}.$$

#### 5.2 l=1

In the case when l = 1, all sub equations were set to 0, so we gain no extra information from the third determining equation.

#### 5.3 l = 2

In the case when l = 2, all sub equations were set to 0, so we gain no extra information from the third determining equation.

## 6 The fourth determining equation

In summary, we have two different infinitesimal depending on whether l = 0 or if l > 0. In the former case, the infinitesimals are given by

$$\xi(z, u) = -B(z),\tag{17}$$

$$\eta(z,u) = -2B'(z)u\ln(u) + C(z) + \left(\frac{1}{2}(B'(z) + B(z)) + K\right)u, \quad K \in \mathbb{R}.$$
 (18)

In the latter case, when l > 0 the infinitesimals are still given by

$$\xi(z,u) = A(z) \left(\frac{1}{1-\ell}\right) u^{1-\ell} - B(z), \tag{19}$$

$$\eta(z,u) = 2cA(z) \left(\frac{1}{\ell-2}\right) u^2 + A'(z) \left(\frac{2\ell-1}{(\ell-1)^2}\right) u^{2-\ell} + 2B'(z) \left(\frac{\ell}{\ell-1}\right) u \ln(u) + C(z)u^{\ell} + D(z)u.$$
(20)