General Formulas

$$\begin{array}{l} e^{i\theta} = \cos\theta + i\sin\theta \\ -\frac{\hbar^2}{2m} \frac{\delta\Psi(x,t)}{\delta x^2} + V(x)\Psi(x,t) = i\hbar \frac{\delta\Psi(x,t)}{\delta t} \\ \sin\alpha + \sin\beta = 2\cos\frac{\alpha-\beta}{2}\sin\frac{\alpha+\beta}{2} \end{array}$$

Light

 $\begin{array}{lll} E_0 = & & n = IndexOfRefraction \\ k = & Wavenumber & a = SingleSlitWidth \\ \omega = & Wavelength & h = Planck'sConstant \\ T = & Period & K = KineticEnergy \\ \omega = & OrdinaryFrequency & W = WorkFunction \\ \phi = & PhaseShift & p = momentum \\ c = & SpeedOfLightVacuum \\ \end{array}$

$$\begin{split} \mathcal{E} &= \mathcal{E}_0 \cos(kx - \omega t) \ _{(1.1)} \\ k &= \frac{2\pi}{\lambda} \ _{(1.2)} \\ \omega &= \frac{2\pi}{T} = 2\pi\nu \ _{(1.3)} \\ \nu &= 1/T \\ \mathcal{E} &= \mathcal{E}_0 \cos(kx - \omega t + \phi) \ _{(1.6)} \\ e^{i\theta} &= \cos\theta + i \sin\theta \ _{(1.7)} \\ \omega &= kc \ _{(1.11)} \\ \omega\nu &= c \ _{(1.12)} \\ \frac{\delta^2 \mathcal{E}}{\delta x^2} - \frac{n^2}{n} \frac{\delta^2 \mathcal{E}}{\delta t^2} = 0 \ _{(1.13)} \\ \lambda\nu &= \frac{c}{n} \ _{(1.14)} \\ a \sin\theta &= n\lambda \ _{(minima)} \\ E &= h\nu \ _{(1.18)} \\ K &= h\nu - W \ _{(1.19)} \\ h\nu_0 &= hc/\lambda_0 = W \ _{(1.20)} \\ p &= \frac{h}{\lambda} \ _{(1.21)} \\ \lambda' - \lambda &= \frac{h}{mc} (1 - \cos\theta) \ _{(1.28)}$$
 Compton

The First Principle of Quantum Mechanics The probability of an event = z^*z (1.32)

The Second Principle of Quantum Mechanics To determine the probability amplitude for a process that can be viewed as taking place in a series of steps we multiply the probability amplitudes for each of these steps.

$$z = z_a z_b \cdots (1.38)$$

The Third Principle of Quantum Mechanics If there are multiple ways that an event can occur we add the amplitudes for each of these ways.

$$z = z_1 + z_2 + \cdots (1.47)$$

$$\phi = kx$$

$$z = x + iy = r \cos \phi + ir \sin \phi = re^{i\phi}$$

$$z^* = x - iy = r \cos \phi - ir \sin \phi = re^{-i\phi}$$

*In this section we assume a free particle, V(x)=0

Wave Mechanics

$$\begin{split} j = & Probability Current \\ & \langle x \rangle = Expectation Value \\ \end{split} \\ \lambda &= \frac{h}{p} \ (2.1) \ \text{de Broglie wavelength} \\ d & \sin \theta = n \lambda \ (2.3) \ (maxima) \\ x_{n+1} - x_n &= \frac{L \lambda}{d} \ (2.4) \\ 2d & \sin \theta = n \lambda \ (2.5) \ \text{Bragg relation} \\ - \frac{\hbar^2}{2m} \frac{\delta \Psi(x,t)}{\delta x^2} + V(x) \Psi(x,t) &= i \hbar \frac{\delta \Psi(x,t)}{\delta t} \ (2.6) \\ - \frac{\hbar^2}{2m} \frac{\delta \Psi(x,t)}{\delta x^2} &= i \hbar \frac{\delta \Psi(x,t)}{\delta t} \ (2.7) \\ \frac{\delta^2 \mathcal{E}}{\delta x^2} &= \frac{n^2}{c} \frac{\delta^2 \mathcal{E}}{\delta t^2} \ (2.8) \\ E &= h \nu - \frac{h}{2\pi} 2\pi \nu = \hbar \omega \ (2.9) \\ p &= \frac{h}{\lambda} &= \frac{h}{2\pi} \frac{2\pi}{\lambda} = \hbar k \ (2.10) \end{split}$$

 $E = pc_{(2.12)}$ $\hbar\omega = \frac{\hbar^2 \dot{k}^2}{2m} \ (2.15)$ $p = \frac{h}{\lambda} = \hbar k_{(2.16)}$ $E = h\nu = \hbar \omega_{(2.17)}$ $E=rac{p^2}{2m}_{(2.18)} \ |\Psi(x,t)|^2 dx$ =the probability of finding the $particle\ between\ x\ and\ x+dx\ at\ the\ time\ t$ if a measurement of the particle's position is carried out $|\Psi(x,t)|^2$ probability density $\int_{-\infty}^{\infty} |\Psi(x,t)|^2 dx = 1 \ \ _{(2.19)}$ $\frac{\delta|\Psi|^2}{\delta t} = \frac{\Psi^* \Psi}{\delta t} = \Psi^* \frac{\delta \Psi}{\delta t} + \Psi \frac{\delta \Psi^*}{\delta t} (2.20)$ $j_x(x,t) = \frac{\hbar}{2mi} (\Psi^* \frac{\delta \Psi}{\delta t} + \Psi \frac{\delta \Psi^*}{\delta t}) (2.24)$ $\frac{d}{dt} \int_{-\infty}^{\infty} |\Psi(x,t)|^2 dx = -j_x(x,t)|_{-\infty}^{\infty} = 0$ $\Psi(x,t) = \int_{-\infty}^{\infty} A(k)e^{i(kx-\omega t)}dk \ _{(2.29)}$ $\Delta x \Delta k \ge \frac{1}{2}$ (2.30) $\Delta x \Delta p_x \geq \frac{\hbar}{2} \ (2.31) \ \text{Heisenberg}$ $v_{ph} = \frac{\omega}{k} = \frac{2\pi\nu}{(2\pi/\lambda)} = \lambda\nu \ (2.33)$ The phase velocity is the speed at which a point on the wave, such as a crest, moves. $v_{ph} = \frac{\omega}{k} = \frac{\hbar\omega}{\hbar k} = \frac{E}{p} = \frac{mv^2/2}{mv} \frac{v}{2}$ (2.34) $v_g = \frac{d\omega}{dk} \ _{(2.36)}$ The group velocity is the speed of a localized packet of waves that has been generated by superposing many waves together $\Psi(x,t) = \int_{-\infty}^{\infty} A(k) e^{i(kx - \omega t)} dk$ (2.37) $\omega \cong \omega_0 + v_g(\widetilde{k} - k_0)$ (2.39) Dispersion relation is the relationship between ω and $k \langle x \rangle = \int_{-\infty}^{\infty} x |\Psi(x,t)|^2 dx$ The average values $\langle x \rangle$ are referred to as the expectation values $\langle x^2 \rangle = \int_{-\infty}^{\infty} x^2 |\Psi(x,t)|^2 dx$ (2.55) $(\Delta x)^2 = \langle x^2 \rangle - \langle x \rangle^2 \ _{(2.56)}$ Δx , the standard deviation, is also called the uncertainty $(\Delta p_x)^2 = \langle p_x^2 \rangle - \langle p_x \rangle^2$ (2.57) $\frac{d\langle p_x \rangle}{dt} = \left\langle -\frac{\delta V}{\delta x} \right\rangle \,_{(2.64)}$

 $\hbar\omega = \hbar k c_{(2.11)}$

The Time-Independent Schrödinger Equation

*In this section we assume V(x) is independent of t $\delta_{nm} = KroneckerDelta$ $\psi_a = Eigenfunction$ T = TransmissionCoef.

$$\begin{split} &\Psi(x,t) = \psi(x)f(t)_{(3.2)} \\ &\frac{\delta^2 \Psi(x,t)}{\delta x^2} = f(t) \frac{d^2 \psi(x)}{dx^2}_{(3.3)} \\ &\frac{\delta \Psi(x,t)}{\delta t} = \psi(x) \frac{df(t)}{dt}_{(3.4)} \\ &\frac{df(t)}{dt} = \frac{-iE}{\hbar} f(t)_{(3.8)} \\ &-\frac{\hbar^2}{2m} \frac{\delta \psi(x)}{\delta x^2} + V(x) \psi(x) = E \psi(x)_{(3.9)} \\ &f(t) = f(0) e^{-iEt/\hbar}_{(3.10)}_{(3.11)} \\ &E = \hbar \omega_{(3.12)} \\ &\Psi(x,t) = \psi(x) e^{-iEt/\hbar}_{(3.13)}_{(3.14)} \end{split}$$

$$V(x) = \begin{cases} 0, & 0 < x < L. \\ \infty, & \text{elsewhere.} \end{cases}$$

$$-\frac{\hbar^2}{2m} \frac{\delta \psi}{\delta x^2} = E \psi_{(3.16)} \ 0 < x < L$$

$$k^2 = \frac{2mE}{\hbar^2} \ (3.17)$$

$$\psi(x) = A \sin kx + B \cos kx_{(3.21)} \ 0 < x < L$$

$$k_n = \frac{n\pi}{L} \ (3.26)$$

$$E_n = \frac{\hbar k_n^2}{2m} = \frac{n\hbar^2 \pi^2}{2mL^2} \ (3.27)$$

$$\psi(x) = A_n \sin \frac{n\pi}{L} \ (3.28) \ 0 < x < L$$

$$\psi_n(x) = \begin{cases} \sqrt{\frac{2}{L}} \sin \frac{n\pi x}{L}, & 0 < x < L. \\ 0, & \text{elsewhere.} \end{cases}$$

$$\Psi(x) = c_1 \psi_1(x) + c_2 \psi_2(x)$$
(3.38)
$$c_1(t) = \frac{1}{\sqrt{2}} e^{-iE_1 t/\hbar}$$
(3.39)
$$\Psi = \sum_{n=1}^{\infty} c_n \psi_n(x)$$
(3.40)

$$\delta_{nm} = \begin{cases} 1, & m = n. \\ 0, & m \neq n. \end{cases}$$

 $\int_{-\infty}^{\infty} \psi_x^*(x) \psi_n(x) dx = \delta_{nm \ (3.49)}$ $|c_n|^2 = P_n \ _{(3.59)}$

The above is the probability of obtaining E_n if a measurement of the energy of a particle with wave function Ψ is carried out

$$\begin{split} \langle E \rangle &= \sum_{n=1}^{\infty} |c_n|^2 E_{n \ (3.61)} \\ A_{op} \psi_a &= a \psi_a \ (3.63) \\ x_{op} &= x \ (3.64) \\ p_{xop} &= \frac{\hbar}{i} \frac{\delta}{\delta x} \ (3.65) \\ E_{op} &= \frac{(p_{xop})^2}{2m} + V(x_{op}) \ (3.71) \\ H &\equiv E_{op} &= -\frac{\hbar^2}{2m} \frac{\delta^2}{\delta x^2} + V(x) \ (3.72) \\ \langle E \rangle &= \int_{-\infty}^{\infty} \Psi^* H \Psi dx \ (3.81) \end{split}$$

One-Dimensional Potentials

$$V(x) = \begin{cases} 0, & |x| < a/2. \\ V_0, & |x| > a/2. \end{cases}$$

$$k = \frac{\sqrt{2mE}}{\hbar}$$

$$\psi(x) = Ae^{ikx} + Be^{-ikx} |x| < a/2$$

$$\kappa = \frac{\sqrt{2m(V_0 - E)}}{\hbar} > 0$$

$$\psi(x) = Ce^{\kappa x} + De^{-\kappa x} |x| > a/2$$

$$\psi(x) = \begin{cases} Ce^{\kappa x}, & x \le -a/2. \\ 2A\cos kx, & -a/2 \le x \le a/2. \\ Ce^{-\kappa x}, & x \ge a/2. \end{cases}$$

$$V(x) = \begin{cases} 0, & x < 0. \\ V_0, & x > 0. \end{cases}$$

$$\begin{split} k &= \frac{\sqrt{2mE}}{\hbar} \\ k_0 &= \sqrt{k^2 - \frac{2mV_0}{\hbar^2}} \\ \psi(x) &= \begin{cases} Ae^{ikx} + Be^{-ikx}, & x < 0. \\ Ce^{ik_0x}, & x > 0. \end{cases} \\ j_x &= \begin{cases} \frac{\hbar k}{m} (|A|^2 - |B|^2), & x < 0. \\ \frac{\hbar k_0}{m} |C|^2, & x > 0. \end{cases} \end{split}$$