

# Math Methods Assignment #5

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1. (a) From the definition of orthogonality given by Arfken (page 105), to show that they are orthogonal we must show that  $g_{ij} = \frac{\partial \mathbf{r}}{\partial q_i} \cdot \frac{\partial \mathbf{r}}{\partial q_j} = 0$  for all  $i \neq j$ :

$$\frac{\partial \mathbf{r}}{\partial q_\eta} \cdot \frac{\partial \mathbf{r}}{\partial q_\phi} = \left( \sinh \eta \cos \hat{\phi} + \cosh \eta \sin \hat{\phi} \right) \cdot \left( -\cosh \eta \sin \hat{\phi} + \sinh \eta \cos \hat{\phi} \right) = 0$$

$$\frac{\partial \mathbf{r}}{\partial q_\eta} \cdot \frac{\partial \mathbf{r}}{\partial q_z} = \left( \sinh \eta \cos \hat{\phi} + \cosh \eta \sin \hat{\phi} \right) \cdot \left( \hat{k} \right) = 0$$

$$\frac{\partial \mathbf{r}}{\partial q_\phi} \cdot \frac{\partial \mathbf{r}}{\partial q_z} = \left( -\cosh \eta \sin \hat{\phi} + \sinh \eta \cos \hat{\phi} \right) \cdot \left( \hat{k} \right) = 0$$

- (b) The metric tensor is defined as  $ds^2 = g_{ij}dq_i dq_j$ . Since we know the coordinate system is orthogonal we only need to calculate the diagonal terms:

$$ds^2 = \begin{bmatrix} d\eta & d\phi & dz \end{bmatrix} \begin{bmatrix} \sinh^2 \eta \cos^2 \phi + \cosh^2 \eta \sin^2 \phi & 0 & 0 \\ 0 & \cosh^2 \eta \sin^2 \phi + \sinh^2 \eta \cos^2 \phi & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} d\eta \\ d\phi \\ dz \end{bmatrix}$$

- (c) Defining  $g_{ii} = h_i^2$  we can use the relation from Arfken:  $\nabla f = \hat{q}_i \frac{1}{h_i} \frac{\partial f}{\partial q_i}$ :

$$\nabla f = \frac{\hat{\eta}}{\sinh \eta \cos \phi + \cosh \eta \sin \phi} \frac{\partial f}{\partial \eta} + \frac{\hat{\phi}}{\cosh \eta \sin \phi + \sinh \eta \cos \phi} \frac{\partial f}{\partial \phi} + \hat{z} \frac{\partial f}{\partial z}$$

- (d) Using equation (2.21) from Arfken:

$$\nabla \cdot \mathcal{E} = \frac{\frac{\partial}{\partial \eta} \mathcal{E}_\eta + \frac{\partial}{\partial \phi} \mathcal{E}_\phi + \frac{\partial}{\partial z} \mathcal{E}_z (\cosh \eta \sin \phi + \sinh \eta \cos \phi)}{(\cosh \eta \sin \phi + \sinh \eta \cos \phi)}$$

- (e) Using equation (2.22) from Arfken:

$$\nabla^2 \cdot \mathcal{E} = \frac{\frac{\partial^2 f}{\partial \eta^2} + \frac{\partial^2 f}{\partial \phi^2} + \frac{\partial}{\partial z} [(\cosh \eta \sin \phi + \sinh \eta \cos \phi)^2 \frac{\partial f}{\partial z}]}{(\cosh \eta \sin \phi + \sinh \eta \cos \phi)^2}$$

2.

$$\nabla^2 \phi = \frac{2\epsilon^4 - 6r^2 \epsilon^2}{r(r^2 + \epsilon^2)^3}$$

(a) Looking at the limits in steps:

$$\lim_{\epsilon \rightarrow 0} \nabla^2 \phi(r_0) = \frac{0}{r_0 (r_0^2)^3}$$

$$\lim_{r_0 \rightarrow 0} 0 \rightarrow 0$$

(b) Again, looking at the limits in steps:

$$\lim_{\epsilon \rightarrow 0} \nabla^2 \phi(r_0) = \frac{2\epsilon^4}{0} \rightarrow \infty$$

$$\lim_{r_0 \rightarrow 0} \infty \rightarrow \infty$$

(c) At all points other than zero we know from part (a) that when  $\epsilon \rightarrow 0$   $\phi(r) \rightarrow \frac{1}{r}$  the limit of the laplacian approaches 0. But for the point at zero, we can treat  $\epsilon$  as a number infinitely close to the origin, since we have already evaluated  $\phi(r)$  at  $r = r_0 \rightarrow \infty$  which as we approach the origin ( $\epsilon \rightarrow 0$ ). Since  $\nabla^2 \frac{1}{r}$  is zero except for the origin, it is a function that has the same behaviour as the delta function, and thus  $\int f(r) \nabla^2 \frac{1}{r} d^3r = f(0)$ .

3. (a) Starting from the definition of polar coordinates:

$$\begin{aligned} x &= \rho \cos \varphi \\ y &= \rho \sin \varphi \end{aligned}$$

$$\begin{aligned} \nabla f &= \frac{\hat{\varphi}}{\frac{\partial \mathbf{r}}{\partial \varphi} \cdot \frac{\partial \mathbf{r}}{\partial \varphi}} \frac{\partial f}{\partial \varphi} + \frac{\hat{\rho}}{\frac{\partial \mathbf{r}}{\partial \rho} \cdot \frac{\partial \mathbf{r}}{\partial \rho}} \frac{\partial f}{\partial \rho} \\ \nabla f &= \frac{\hat{\varphi}}{\rho (\cos^2 \varphi + \sin^2 \varphi)} \frac{\partial f}{\partial \varphi} + \frac{\hat{\rho}}{\cos^2 \varphi + \sin^2 \varphi} \frac{\partial f}{\partial \rho} \\ \nabla f &= \frac{\hat{\varphi}}{\rho} \frac{\partial f}{\partial \varphi} + \hat{\rho} \frac{\partial f}{\partial \rho} \end{aligned}$$

(b)

$$\begin{aligned} \nabla \times \vec{g} &= \frac{\partial}{\partial \rho} \frac{\hat{\varphi}}{\rho} \frac{\partial f}{\partial \varphi} + \frac{\partial}{\partial \varphi} \hat{\rho} \frac{\partial f}{\partial \rho} \\ \nabla \times \vec{g} &= \left( -\frac{1}{\rho^2} \frac{\partial f}{\partial \varphi} + \frac{1}{\rho} \frac{\partial}{\partial \rho} \frac{\partial f}{\partial \varphi} \right) \hat{\varphi} + \hat{\rho} \frac{\partial}{\partial \varphi} \frac{\partial f}{\partial \rho} \end{aligned}$$

(c) Since the gradient of scalar field is conservative, the line integral will be zero, since it's a loop and will start and end in the same points.<sup>1</sup>

(d) This agrees with the question statement.

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<sup>1</sup><https://math.stackexchange.com/questions/1435044/line-integral-of-conservative-field-in-polar-coordinates>