

UniDyn--Study-02.nb

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Abstract: This demonstration notebook loads the **UniDyn** package, loads the **Inv** function, and runs the **Inv** function unit tests. Below we write a new function for carrying out unitary rotations, **Evolver2[]**, that works on a case-by-case basis. It is approximately 25x faster than **Evolver[]**, and handles more cases.

Set the path to the package

Check the Mathematica version number .

```
In[96]:= $VersionNumber
```

```
Out[96]:= 12.3
```

Tell *Mathematica* the path to the directory containing the packages.

EDIT THE FOLLOWING PATH STRINGS:

```
In[97]:= $UniDynPath =
```

```
"/Users/jam99/Dropbox/MarohnGroup__Software_Library/UniDyn/  
  unidyn";
```

```
In[98]:= $UniDynStudyPath =
```

```
"/Users/jam99/Dropbox/MarohnGroup__Software_Library/UniDyn/  
  studies";
```

YOU SHOULD NOT NEED TO EDIT ANYTHING FROM HERE ONWARDS.

Load and test the packages

Append the package path to the system path. Before trying to load the package, ask *Mathematica* to find it. This is a test that we directed *Mathematica* to the correct directory. The output of this command should be the full system path to the UniDyn.m file.

```
In[99]:= $Path = AppendTo[$Path, $UniDynPath];
FindFile["UniDyn`"]
```

```
Out[100]= /Users/jam99/Dropbox/MarohnGroup__Software_Library/UniDyn/unidyn/UniDyn.m
```

Now that we are confident that the path is set correctly, load the package. Setting the global `$VerboseLoad` variable to `True` will print out the help strings for key commands in the package.

```
In[101]:= $VerboseLoad = True;
Needs["UniDyn`"]
```

```
In[103]:= $Path = AppendTo[$Path, $UniDynStudyPath];
FindFile["Inv`"]
FindFile["SpinBoson`"]
```

```
Out[104]= /Users/jam99/Dropbox/MarohnGroup__Software_Library/UniDyn/studies/Inv.m
```

```
Out[105]= /Users/jam99/Dropbox/MarohnGroup__Software_Library/UniDyn/studies/SpinBoson.m
```

```
In[106]:= $VerboseLoad = True;
Needs["Inv`"]
Needs["SpinBoson`"]
```

Extend the directory to the `/studies` directory and run all the tests found there.

```
In[109]:= SetDirectory[$UniDynStudyPath];
```

```
In[110]:= fn = FileNames["*-tests.m"];
test$report = TestReport /@ fn;
TableForm[Table[test$report[[k]], {k, 1, Length[test$report]}]]
```

```
... SpinSingle$CreateOperators: Spin operators already exist.
```

```
... SpinSingle$CreateOperators: Adding spin commutations relations.
```

```
... SpinSingle$CreateOperators: Angular momentum  $L = 1/2$ . Adding operator simplification rules.
```

```
... OscSingle$CreateOperators: Oscillator operators already exist.
```

```
... OscSingle$CreateOperators: Adding oscillator commutations relations.
```

```
Out[112]//TableForm=
```

TestReportObject[ 	Title: Test Report: Evolve2-tests.m Success rate: 100% Tests run: 22]
TestReportObject[ 	Title: Test Report: Inv-tests.m Success rate: 100% Tests run: 22]
TestReportObject[ 	Title: Test Report: SpinBoson-tests.m Success rate: 100% Tests run: 2]

Create quantum optics operators

```
In[113]:= Clear[Ix, Iy, Iz, I$P, I$m, aR, aL];
SpinBoson$CreateOperators[Ix, Iy, Iz, I$P, I$m, aR, aL];
```

... SpinSingle\$CreateOperators: Spin operators already exist.

... SpinSingle\$CreateOperators: Adding spin commutations relations.

... SpinSingle\$CreateOperators: Angular momentum $L = 1/2$. Adding operator simplification rules.

... OscSingle\$CreateOperators: Oscillator operators already exist.

... OscSingle\$CreateOperators: Adding oscillator commutations relations.

... SpinBoson\$CreateOperators: Creating operators.

... SpinBoson\$CreateOperators: Adding Ip and Im commutations relations.

... SpinBoson\$CreateOperators: Adding Ip and Im simplification rules.

... SpinBoson\$CreateOperators: Adding aL and aR normal ordering rule.

```
In[115]:= Clear[Δ, ω1, g];
CreateScalar[{Δ, ω1, g}];
$Assumptions = {Δ ∈ Reals, ω1 ∈ PositiveReals, g ∈ PositiveReals};
```

Are these assumptions enough? We can see the following examples that Sqrt[] involving the real parameters will not simplify correctly if there is more than one real parameter in the Sqrt[]. To simplify the Sqrt[] with multiple real parameters inside, it is enough to assume that one of the parameters is positive.

```
In[118]:= PowerExpand[Sqrt[-Δ2]]
```

Out[118]= $i \Delta$

```
In[119]:= PowerExpand[Sqrt[-Δ2 - ω12]]
```

Out[119]= $\sqrt{-\Delta^2 - \omega_1^2}$

```
In[120]:= PowerExpand[Sqrt[-Δ2 - ω12], Assumptions → $Assumptions]
```

Out[120]= $i \sqrt{\Delta^2 + \omega_1^2}$

```
In[121]:= Simplify[Sqrt[-Δ2 - ω12]]
```

Out[121]= $i \sqrt{\Delta^2 + \omega_1^2}$

A rotation involving spin raising and lowering operators that succeeds

We can now evolve the raising and lowering operators .

```
In[122]:= {Evolver[λ Iz, t, I$ p], Evolver[λ Iz, t, I$ m]}
```

```
Out[122]:= {e-i t λ I$ p, ei t λ I$ m}
```

Two rotation involving mixed spin and harmonic oscillator operators that fail

Here is a unitary rotation that Evolver[] cannot resolve.

```
In[123]:= terms = Evolver[λ Mult[Iz, aL], t, I$ p, quiet → False]
```

$$\rho \text{ matrix} = \begin{pmatrix} I\$ p & & & \\ -i \lambda \text{Mult}[I\$ p, aL] & & & \\ -\lambda^2 \text{Mult}[I\$ p, aL, aL] & & & \\ i \lambda^3 \text{Mult}[I\$ p, aL, aL, aL] & & & \\ \lambda^4 \text{Mult}[I\$ p, aL, aL, aL, aL] & & & \end{pmatrix}$$

... Evolver: Unrecognized evolution

```
Out[123]:= {I$ p, -i λ Mult[I$ p, aL], -λ2 Mult[I$ p, aL, aL],  
i λ3 Mult[I$ p, aL, aL, aL], λ4 Mult[I$ p, aL, aL, aL, aL]}
```

The ratio of the 1st to the 0th term is independent of I\$ p.

```
In[124]:= Mult[Inv[terms[[1]]], terms[[2]]]
```

```
Out[124]:= -i aL λ
```

The ratio of a higher term to the next lower term is likewise independent of I\$ p.

```
In[125]:= {Mult[Inv[terms[[2]]], terms[[3]]], Mult[Inv[terms[[3]]], terms[[4]]]}
```

```
Out[125]:= {-i aL λ, -i aL λ}
```

Here is another unitary rotation that Evolver[] cannot resolve.

```
In[126]:= terms = Evolver[2 I g (Mult[Iz, aR] - Mult[Iz, aL]), t, I$ p] // Simplify
```

... Evolver: Unrecognized evolution

```
Out[126]:= {I$ p, -2 g (Mult[I$ p, aL] - Mult[I$ p, aR]),  
-4 g2 (I$ p - Mult[I$ p, aL, aL] + 2 Mult[I$ p, aR, aL] - Mult[I$ p, aR, aR]),  
8 g3 (3 Mult[I$ p, aL] - 3 Mult[I$ p, aR] - Mult[I$ p, aL, aL, aL] +  
3 Mult[I$ p, aR, aL, aL] - 3 Mult[I$ p, aR, aR, aL] + Mult[I$ p, aR, aR, aR]),  
16 g4 (3 I$ p - 6 Mult[I$ p, aL, aL] + 12 Mult[I$ p, aR, aL] - 6 Mult[I$ p, aR, aR] +  
Mult[I$ p, aL, aL, aL, aL] - 4 Mult[I$ p, aR, aL, aL, aL] + 6 Mult[I$ p, aR, aR, aL, aL] -  
4 Mult[I$ p, aR, aR, aR, aL] + Mult[I$ p, aR, aR, aR, aR])}
```

Notice that again the ratio of the 1st to the 0th term is independent of I\$ p.

```
In[127]:= {Mult[Inv[terms[[1]]], terms[[2]]]}
```

```
Out[127]:= {-2 (aL - aR) g}
```

Create a new Evolver algorithm

Pass 1

Strip down the existing algorithm, and print out some useful intermediate results.

```
In[128]:= VisualComplexity := (Count[ToBoxes[#], Except[" " | "(" | ")"], _String], Infinity] &
Options[Evolver2] = {quiet → True};

In[130]:= Evolver2[H$sym_, t$sym_, rho$sym$0_, opts : OptionsPattern[]] :=

Module[{},

Clear[rho$sym];
rho$sym[0] = rho$sym$0;

Do[
rho$sym[k + 1] = (-I Comm[H$sym, rho$sym[k]] /. Mult → SortedMult)
// FullSimplify[#, ComplexityFunction → VisualComplexity] &,
{k, 0, 4}
];

If[OptionValue[quiet] == False,
Print[" $\rho_n =$  ", rho$sym[#] & /@ {0, 1, 2, 3, 4} // MatrixForm];
Print["divisions = ", {Mult[Inv[rho$sym[0]], rho$sym[1]],
Mult[Inv[rho$sym[0]], rho$sym[2]],
Mult[Inv[rho$sym[1]], rho$sym[3]]}];
Print["commutators = ", {rho$sym[1],
Comm[rho$sym[0], rho$sym[1]],
Comm[rho$sym[0], rho$sym[2]],
Comm[rho$sym[1], rho$sym[3]]}]]

]
```

The 0th case to consider is when the Hamiltonian commutes with the initial condition. If this is the case we should just stop.

```
In[131]:= terms = Evolver2[Δ I$ p, t, I$ p, quiet → False]
```

$$\rho_n = \begin{pmatrix} I\$ p \\ 0 \\ 0 \\ 0 \\ 0 \end{pmatrix}$$

Power: Infinite expression $\frac{1}{0}$ encountered.

Infinity: Indeterminate expression 0 ComplexInfinity encountered.

divisions = {0, 0, Indeterminate}

commutators = {0, 0, 0, 0}

The 1st case is when the first commutator commutes with the initial condition.

```
In[132]:= terms = Evolver2[g Mult[Iz, aL], t, I$ p, quiet → False]
```

$$\rho_n = \begin{pmatrix} I\$ p \\ -i g \text{Mult}[I\$ p, aL] \\ -g^2 \text{Mult}[I\$ p, aL, aL] \\ i g^3 \text{Mult}[I\$ p, aL, aL, aL] \\ g^4 \text{Mult}[I\$ p, aL, aL, aL, aL] \end{pmatrix}$$

divisions = {-i aL g, -g² Mult[aL, aL], -g² Mult[aL, aL]}

commutators = {-i g Mult[I\$ p, aL], 0, 0, 0}

The 2nd case is when the second commutator commutes with the initial condition.

```
In[133]:= Evolver2[Δ Iz, t, Ix, quiet → False]
```

$$\rho_n = \begin{pmatrix} Ix \\ Iy \Delta \\ -Ix \Delta^2 \\ -Iy \Delta^3 \\ Ix \Delta^4 \end{pmatrix}$$

divisions = {Δ Mult[Inv[Ix], Iy], -Δ², -Δ²}

commutators = {Iy Δ, i Iz Δ, 0, 0}

The 3rd case is when the third commutator commutes with the first commutator.

```
In[134]:= Evolver2[Δ Iz + ω1 Ix, t, Iz, quiet → False]
```

$$\rho_n = \begin{pmatrix} Iz \\ -Iy \omega 1 \\ \omega 1 (Ix \Delta - Iz \omega 1) \\ Iy \omega 1 (\Delta^2 + \omega 1^2) \\ \omega 1 (-Ix \Delta + Iz \omega 1) (\Delta^2 + \omega 1^2) \end{pmatrix}$$

divisions = {-ω1 Mult[Inv[Iz], Iy], ω1 (-ω1 + Δ Mult[Inv[Iz], Ix]), -Δ² - ω1²}

commutators = {-Iy ω1, i Ix ω1, i Iy Δ ω1, 0}

Pass 2

Now implement cases, writing down the solution to the differential equations by hand.

```
In[135]:= Evolver2::unsolvable = "Unrecognized evolution";

Evolver2[H$sym_, t$sym_, rho$sym$0_, opts : OptionsPattern[]] :=

Module[{rho$sym, divisions, commutators,  $\rho$ , w},

rho$sym[0] = rho$sym$0;

Do[
  rho$sym[k + 1] = (-I Comm[H$sym, rho$sym[k]] /. Mult  $\rightarrow$  SortedMult)
    // FullSimplify[#, ComplexityFunction  $\rightarrow$  VisualComplexity] &,
  {k, 0, 4}
];

If[OptionValue[quiet] == False, Print[" $\rho^{(n)}$  = ", rho$sym[#] & /@ {0, 1, 2, 3}]];

If[rho$sym[1] == 0,
  If[OptionValue[quiet] == False, Print["Case 0"]];
  Return[rho$sym[0]]];

If[Comm[rho$sym[1], H$sym] == 0,
  If[OptionValue[quiet] == False, Print["Case 1L"]];
  Return[rho$sym[0] + rho$sym[1] t$sym]];

commutators = {rho$sym[1],
  Comm[rho$sym[0], rho$sym[1]],
  Comm[rho$sym[0], rho$sym[2]],
  Comm[rho$sym[1], rho$sym[3]]};

If[OptionValue[quiet] == False, Print["commutators = ", commutators]];

divisions = {Mult[Inv[rho$sym[0]], rho$sym[1]],
  Mult[Inv[rho$sym[0]], rho$sym[2]],
  Mult[Inv[rho$sym[1]], rho$sym[3]]};

If[OptionValue[quiet] == False, Print["divisions = ", divisions]];

If[commutators[[2]] == 0,
  If[OptionValue[quiet] == False, Print["Case 1E"]];
```

```

    ρ = Mult[Exp[divisions[[1]] t$sym], rho$sym[0]];
    Return[ρ];

    If[commutators[[3]] == 0,
      If[OptionValue[quiet] == False, Print["Case 2E"]];
      w = PowerExpand[Sqrt[-divisions[[2]]]];
      If[OptionValue[quiet] == False, Print["w = ", w]];
      ρ = Mult[Cos[w t$sym], rho$sym[0]] + Mult[Sin[w t$sym], commutators[[1]] / w];
      Return[ρ];

    If[commutators[[4]] == 0,
      If[OptionValue[quiet] == False, Print["Case 3E"]];
      w = PowerExpand[Sqrt[-divisions[[3]]]];
      If[OptionValue[quiet] == False, Print["w = ", w]];
      ρ = rho$sym[0] + Mult[Sin[w t$sym], rho$sym[1] / w] +
        Mult[1 - Cos[w t$sym], rho$sym[2] / w^2];
      Return[ρ];

    Message[Evolver2::unsolvable];
    Return[{rho$sym[#] & /@ {0, 1, 2, 3}, commutators, divisions}]
  ]

```

The 0a-th case is when the Hamiltonian commutes with the density operator.

```
In[137]:= Evolver2[λ I$ρ, t, I$ρ, quiet → False]
```

$$\rho^{(n)} = \{I\$ρ, 0, 0, 0\}$$

Case 0

```
Out[137]:= I$ρ
```

The 0b-th case is when the Hamiltonian commutes with the first commutator.

```
In[138]:= Evolver2[I g  $\frac{1}{\sqrt{2}}$  (aR - aL), 1,  $\frac{1}{\sqrt{2}}$  (aL + aR), quiet → False] // Simplify
```

$$\rho^{(n)} = \left\{ \frac{aL + aR}{\sqrt{2}}, -g, 0, 0 \right\}$$

Case 1L

```
Out[138]=  $\frac{aL + aR}{\sqrt{2}} - g$ 
```

The 1st case is when the first commutator commutes with the initial condition.

```
In[139]:= Evolver2[g Mult[Iz, aL], t, I$ρ, quiet → False]
```


$$\rho^{(n)} = \{I\$, -i g \text{Mult}[I\$, aL], -g^2 \text{Mult}[I\$, aL, aL], i g^3 \text{Mult}[I\$, aL, aL, aL]\}$$

$$\text{commutators} = \{-i g \text{Mult}[I\$, aL], 0, 0, 0\}$$

$$\text{divisions} = \{-i aL g, -g^2 \text{Mult}[aL, aL], -g^2 \text{Mult}[aL, aL]\}$$

Case 1E

$$\text{Out[139]} = \text{Mult}[e^{-i aL g t}, I\$]$$

Compare to

$$\text{In[140]} := \text{Evolver}[g \text{Mult}[Iz, aL], t, I\$]$$

... Evolver: Unrecognized evolution

$$\text{Out[140]} = \{I\$, -i g \text{Mult}[I\$, aL], -g^2 \text{Mult}[I\$, aL, aL], i g^3 \text{Mult}[I\$, aL, aL, aL], g^4 \text{Mult}[I\$, aL, aL, aL, aL]\}$$

The 2nd case is when the second commutator commutes with the initial condition.

$$\text{In[141]} := \text{Evolver2}[\Delta Iz, t, Ix, \text{quiet} \rightarrow \text{False}]$$

$$\rho^{(n)} = \{Ix, Iy \Delta, -Ix \Delta^2, -Iy \Delta^3\}$$

$$\text{commutators} = \{Iy \Delta, i Iz \Delta, 0, 0\}$$

$$\text{divisions} = \{\Delta \text{Mult}[\text{Inv}[Ix], Iy], -\Delta^2, -\Delta^2\}$$

Case 2E

$$w = \Delta$$

$$\text{Out[141]} = Ix \cos[t \Delta] + Iy \sin[t \Delta]$$

Compare to

$$\text{In[142]} := \text{Evolver}[\Delta Iz, t, Ix]$$

$$\text{Out[142]} = Ix \cos[t \Delta] + Iy \sin[t \Delta]$$

The 3rd case is when the third commutator commutes with the first commutator.

$$\text{In[143]} := \text{Evolver2}[\Delta Iz + \omega 1 Ix, t, Iz, \text{quiet} \rightarrow \text{False}] // \text{Collect}[\#, \{Ix, Iy, Iz\}, \text{Simplify}] \&$$

$$\rho^{(n)} = \{Iz, -Iy \omega 1, \omega 1 (Ix \Delta - Iz \omega 1), Iy \omega 1 (\Delta^2 + \omega 1^2)\}$$

$$\text{commutators} = \{-Iy \omega 1, i Ix \omega 1, i Iy \Delta \omega 1, 0\}$$

$$\text{divisions} = \{-\omega 1 \text{Mult}[\text{Inv}[Iz], Iy], \omega 1 (-\omega 1 + \Delta \text{Mult}[\text{Inv}[Iz], Ix]), -\Delta^2 - \omega 1^2\}$$

Case 3E

$$w = \sqrt{\Delta^2 + \omega 1^2}$$

$$\text{Out[143]} = \frac{Iz (\Delta^2 + \omega 1^2 \cos[t \sqrt{\Delta^2 + \omega 1^2}])}{\Delta^2 + \omega 1^2} + \frac{2 Ix \Delta \omega 1 \sin[\frac{1}{2} t \sqrt{\Delta^2 + \omega 1^2}]^2}{\Delta^2 + \omega 1^2} - \frac{Iy \omega 1 \sin[t \sqrt{\Delta^2 + \omega 1^2}]}{\sqrt{\Delta^2 + \omega 1^2}}$$

Compare to

```
In[144]:= Evolver[Δ Iz + ω1 Ix, t, Iz] // Collect[#, {Ix, Iy, Iz}, Simplify] &
```

$$\text{Out[144]} = \frac{Iz \left(\Delta^2 + \omega 1^2 \cos \left[t \sqrt{\Delta^2 + \omega 1^2} \right] \right)}{\Delta^2 + \omega 1^2} + \frac{2 Ix \Delta \omega 1 \sin \left[\frac{1}{2} t \sqrt{\Delta^2 + \omega 1^2} \right]^2}{\Delta^2 + \omega 1^2} - \frac{Iy \omega 1 \sin \left[t \sqrt{\Delta^2 + \omega 1^2} \right]}{\sqrt{\Delta^2 + \omega 1^2}}$$

Speed test

```
In[145]:= timing1 = Timing[
  Evolver[Δ Iz + ω1 Ix, t, Iz]
]
```

```
Out[145]= {0.026488,
```

$$\frac{\Delta (Iz \Delta + Ix \omega 1) + \omega 1 (-Ix \Delta + Iz \omega 1) \cos \left[t \sqrt{\Delta^2 + \omega 1^2} \right] - Iy \omega 1 \sqrt{\Delta^2 + \omega 1^2} \sin \left[t \sqrt{\Delta^2 + \omega 1^2} \right]}{\Delta^2 + \omega 1^2}}$$

```
In[146]:= timing2 = Timing[
  Evolver2[Δ Iz + ω1 Ix, t, Iz]
]
```

```
Out[146]= {0.00099, Iz + \frac{\omega 1 (Ix \Delta - Iz \omega 1)}{\Delta^2 + \omega 1^2} - \frac{\omega 1 (Ix \Delta - Iz \omega 1) \cos \left[ t \sqrt{\Delta^2 + \omega 1^2} \right]}{\Delta^2 + \omega 1^2} - \frac{Iy \omega 1 \sin \left[ t \sqrt{\Delta^2 + \omega 1^2} \right]}{\sqrt{\Delta^2 + \omega 1^2}}}
```

The ratio of the execution times is

```
In[147]:= Round[timing1[[1]] / timing2[[1]], 1]
```

```
Out[147]= 27
```

For this example of **Case 3E** evolution, **Evolver2[]** is approximately **28x faster** than **Evolver[]**.

Electron transfer test cases

This first case my old code handled .

```
In[148]:= { Evolver[-2 I g Mult[Iz, aR - aL], 1, aL],
  Evolver2[-2 I g Mult[Iz, aR - aL], 1, aL]}
```

```
Out[148]= {aL + 2 g Iz, aL + 2 g Iz}
```

These next two more complicated cases, **Evolver[]** failed to handle.

Evolver2[] handles these two cases of non-trivial spin-boson evolution!

```
In[149]:= { Evolver[I g (Mult[Iz, aR] - Mult[Iz, aL]), t, I$P],
  Evolver2[I g (Mult[Iz, aR] - Mult[Iz, aL]), t, I$P]}
```

... **Evolver**: Unrecognized evolution

```
Out[149]:= {{I$P, -g (Mult[I$P, aL] - Mult[I$P, aR]),
  -g^2 (I$P - Mult[I$P, aL, aL] + 2 Mult[I$P, aR, aL] - Mult[I$P, aR, aR]),
  g^3 (3 Mult[I$P, aL] - 3 Mult[I$P, aR] - Mult[I$P, aL, aL, aL] +
    3 (Mult[I$P, aR, aL, aL] - Mult[I$P, aR, aR, aL]) + Mult[I$P, aR, aR, aR]),
  g^4 (3 I$P - 6 Mult[I$P, aL, aL] + 12 Mult[I$P, aR, aL] - 6 Mult[I$P, aR, aR] + Mult[
    I$P, aL, aL, aL, aL] - 4 Mult[I$P, aR, aL, aL, aL] + 6 Mult[I$P, aR, aR, aL, aL] -
    4 Mult[I$P, aR, aR, aR, aL] + Mult[I$P, aR, aR, aR, aR])}, Mult[E^(-(aL-aR) g t), I$P]}
```

```
In[150]:= { Evolver[-2 I g Mult[Iz, aR - aL], t, I$P],
  Evolver2[-2 I g Mult[Iz, aR - aL], t, I$P]}
```

... **Evolver**: Unrecognized evolution

```
Out[150]:= {{I$P, 2 g (Mult[I$P, aL] - Mult[I$P, aR]),
  -4 g^2 (I$P - Mult[I$P, aL, aL] + 2 Mult[I$P, aR, aL] - Mult[I$P, aR, aR]),
  -8 g^3 (3 Mult[I$P, aL] - 3 Mult[I$P, aR] - Mult[I$P, aL, aL, aL] +
    3 (Mult[I$P, aR, aL, aL] - Mult[I$P, aR, aR, aL]) + Mult[I$P, aR, aR, aR]),
  16 g^4 (3 I$P - 6 Mult[I$P, aL, aL] + 12 Mult[I$P, aR, aL] - 6 Mult[I$P, aR, aR] + Mult[
    I$P, aL, aL, aL, aL] - 4 Mult[I$P, aR, aL, aL, aL] + 6 Mult[I$P, aR, aR, aL, aL] -
    4 Mult[I$P, aR, aR, aR, aL] + Mult[I$P, aR, aR, aR, aR])}, Mult[E^(2 (aL-aR) g t), I$P]}
```

Create two spin operators (both $I = 1/2$)

```
In[151]:= Clear[Ix, Iy, Iz, Sx, Sy, Sz];
CreateOperator[{{Ix, Iy, Iz}, {Sx, Sy, Sz}}];
SpinSingle$CreateOperators[Ix, Iy, Iz, 1/2];
SpinSingle$CreateOperators[Sx, Sy, Sz, 1/2];
```

... **SpinSingle\$CreateOperators**: Spin operators already exist.

... **SpinSingle\$CreateOperators**: Adding spin commutations relations.

... **SpinSingle\$CreateOperators**: Angular momentum $L = 1/2$. Adding operator simplification rules.

... **SpinSingle\$CreateOperators**: Spin operators already exist.

... **SpinSingle\$CreateOperators**: Adding spin commutations relations.

... **SpinSingle\$CreateOperators**: Angular momentum $L = 1/2$. Adding operator simplification rules.

```
In[155]:= Clear[Δ, J];
CreateScalar[{Δ, J}];
$Assumptions = {Δ ∈ Reals, J ∈ Reals};
```

This rotation shows up when calculating evolution of an I_x under a heteronuclear J coupling.

For the $I = 1/2$ case, both **Evolver[]** and **Evolver2[]** know what to do.

```
In[158]:= { Evolver[J Mult[Sz, Iz], t, Ix],
           Evolver2[J Mult[Sz, Iz], t, Ix]}
```

```
Out[158]= { Ix Cos[ $\frac{J t}{2}$ ] + 2 Mult[Iy, Sz] Sin[ $\frac{J t}{2}$ ], Ix Cos[ $\frac{J t}{2}$ ] + 2 Mult[Iy, Sz] Sin[ $\frac{J t}{2}$ ] }
```

Create two spin operators (one $I = 1/2$)

```
In[159]:= Clear[Ix, Iy, Iz, Sx, Sy, Sz];
           CreateOperator[{{Ix, Iy, Iz}, {Sx, Sy, Sz}}];
           SpinSingle$CreateOperators[Ix, Iy, Iz, 1/2];
           SpinSingle$CreateOperators[Sx, Sy, Sz];
```

... SpinSingle\$CreateOperators: Spin operators already exist.

... SpinSingle\$CreateOperators: Adding spin commutations relations.

... SpinSingle\$CreateOperators: Angular momentum $L = 1/2$. Adding operator simplification rules.

... SpinSingle\$CreateOperators: Spin operators already exist.

... SpinSingle\$CreateOperators: Adding spin commutations relations.

... SpinSingle\$CreateOperators: No angular momentum L defined.

```
In[163]:= Clear[Δ, J];
           CreateScalar[{Δ, J}];
           $Assumptions = {Δ ∈ Reals, J ∈ Reals};
```

For the general- I case, **Evolver[]** cannot handle the following rotation.

```
In[166]:= Evolver[J Mult[Sz, Iz], t, Ix]
```

... Evolver: Unrecognized evolution

```
Out[166]= { Ix, J Mult[Iy, Sz], -J^2 Mult[Ix, Sz, Sz],
           -J^3 Mult[Iy, Sz, Sz, Sz], J^4 Mult[Ix, Sz, Sz, Sz, Sz] }
```

Evolver2[] handles the general I -spin heteronuclear J -coupling evolution nicely.

```
In[167]:= Evolver2[J Mult[Sz, Iz], t, Ix, quiet → False]
```

$$\rho^{(n)} = \{I_x, J \text{ Mult}[I_y, Sz], -J^2 \text{ Mult}[I_x, Sz, Sz], -J^3 \text{ Mult}[I_y, Sz, Sz, Sz]\}$$

$$\text{commutators} = \{J \text{ Mult}[I_y, Sz], i J \text{ Mult}[I_z, Sz], 0, 0\}$$

$$\text{divisions} = \{J \text{ Mult}[\text{Inv}[I_x], I_y, Sz], -J^2 \text{ Mult}[Sz, Sz], -J^2 \text{ Mult}[Sz, Sz]\}$$

Case 2E

$$w = J \sqrt{\text{Mult}[Sz, Sz]}$$

$$\text{Out[167]} = \text{Mult}\left[\text{Cos}\left[J t \sqrt{\text{Mult}[Sz, Sz]}\right], I_x\right] + \text{Mult}\left[\text{Sin}\left[J t \sqrt{\text{Mult}[Sz, Sz]}\right], \frac{\text{Mult}[I_y, Sz]}{\sqrt{\text{Mult}[Sz, Sz]}}\right]$$

We could now think about using the Lowdin projection operator theorem to expand functions like

$$\text{Cos}\left[J t \sqrt{\text{Mult}[Sz, Sz]}\right] \text{ and } \text{Sin}\left[J t \sqrt{\text{Mult}[Sz, Sz]}\right].$$

Neato!

Create two spin operators (both general I) [unrecognized evolution]

```
In[168]:= Clear[Ix, Iy, Iz, Sx, Sy, Sz];
CreateOperator[{{Ix, Iy, Iz}, {Sx, Sy, Sz}}];
SpinSingle$CreateOperators[Ix, Iy, Iz];
SpinSingle$CreateOperators[Sx, Sy, Sz];
```

... SpinSingle\$CreateOperators: Spin operators already exist.

... SpinSingle\$CreateOperators: Adding spin commutations relations.

... SpinSingle\$CreateOperators: No angular momentum L defined.

... SpinSingle\$CreateOperators: Spin operators already exist.

... SpinSingle\$CreateOperators: Adding spin commutations relations.

... SpinSingle\$CreateOperators: No angular momentum L defined.

```
In[172]:= Clear[Δ, J];
CreateScalar[Δ, J];
$Assumptions = {Δ ∈ Reals, J ∈ Reals};
```

For the general-I case, **Evolver2[]** cannot handle the following rotation.

We can see that the commutators quickly get out of hand.

```
In[175]:= Evolver2[J Mult[Sz, Iz, Iy], t, Ix]
```

... Evolver2: Unrecognized evolution

$$\text{Out[175]} = \left\{ \left\{ I_x, J \left(\text{Mult}[I_y, Iy, Sz] - \text{Mult}[Iz, Iz, Sz] \right), \right. \right. \\ \left. -J^2 \left(\text{Mult}[Ix, Iy, Iy, Sz, Sz] + \text{Mult}[Iy, Ix, Iy, Sz, Sz] + \right. \right. \\ \left. \left. \text{Mult}[Iz, Ix, Iz, Sz, Sz] + \text{Mult}[Iz, Iz, Ix, Sz, Sz] \right), \right. \\ \left. J^3 \left(2 \text{Mult}[Ix, Ix, Iy, Iy, Sz, Sz, Sz] + \text{Mult}[Ix, Iy, Ix, Iy, Sz, Sz, Sz] + \right. \right. \\ \left. \left. \text{Mult}[Ix, Iy, Iz, Iz, Sz, Sz, Sz] + \text{Mult}[Iy, Ix, Iz, Iz, Sz, Sz, Sz] \right) \right\}$$


```

Sz, Sz, Sz, Iz, Sz] + 2 i Mult[Iz, Iy, Iy, Ix, Iy, Sz, Sz, Sz, Sz] -
2 (-i Mult[Ix, Iy, Iy, Iy, Sz, Sz, Sz, Iz, Sz] - i Mult[Iy, Ix, Iy, Iy, Sz,
Sz, Sz, Iz, Sz] - i Mult[Iy, Iy, Ix, Iy, Sz, Sz, Sz, Iz, Sz] - i Mult[Iy, Iy,
Iy, Ix, Sz, Sz, Sz, Iz, Sz] - i Mult[Iz, Ix, Iy, Iy, Iy, Sz, Sz, Sz, Sz] -
i Mult[Iz, Iy, Ix, Iy, Iy, Sz, Sz, Sz, Sz] - i Mult[Iz, Iy, Iy, Ix,
Iy, Sz, Sz, Sz, Sz] - i Mult[Iz, Iy, Iy, Iy, Ix, Sz, Sz, Sz, Sz]) -
i Mult[Iz, Iy, Iz, Ix, Sz, Sz, Sz, Iz, Sz] - i Mult[Iz, Iz, Ix, Iy,
Iz, Sz, Sz, Sz, Sz] - i Mult[Iz, Iz, Ix, Iz, Iy, Sz, Sz, Sz, Sz] -
i Mult[Iz, Iz, Iy, Ix, Iz, Sz, Sz, Sz, Sz] -
i Mult[Iz, Iz, Iy, Iz, Ix, Sz, Sz, Sz, Sz] -
2 (i Mult[Iz, Iz, Ix, Iy, Sz, Sz, Sz, Iz, Sz] +
i Mult[Iz, Iz, Iy, Ix, Sz, Sz, Sz, Iz, Sz] + i Mult[Iz, Iz, Iz, Ix, Iy,
Sz, Sz, Sz, Sz] + i Mult[Iz, Iz, Iz, Iy, Ix, Sz, Sz, Sz, Sz]))},
{J (Mult[Inv[Ix], Iy, Iy, Sz] - Mult[Inv[Ix], Iz, Iz, Sz]),
-j^2
(Mult[Iy, Iy, Sz, Sz] +
Mult[Inv[Ix], Iy, Ix, Iy, Sz, Sz] +
Mult[Inv[Ix], Iz, Ix, Iz, Sz, Sz] +
Mult[Inv[Ix], Iz, Iz, Ix, Sz, Sz]),
j^2 (2 Mult[Inv[Mult[Iy, Iy, Sz] - Mult[Iz, Iz, Sz]], Ix, Ix, Iy, Iy, Sz, Sz, Sz] +
Mult[Inv[Mult[Iy, Iy, Sz] - Mult[Iz, Iz, Sz]],
Ix, Iy, Ix, Iy, Sz, Sz, Sz] +
Mult[Inv[Mult[Iy, Iy, Sz] - Mult[Iz, Iz, Sz]],
Iy, Ix, Ix, Iy, Sz, Sz, Sz] -
2 Mult[Inv[Mult[Iy, Iy, Sz] - Mult[Iz, Iz, Sz]], Iy, Iy, Iy, Iy, Sz, Sz, Sz] -
Mult[Inv[Mult[Iy, Iy, Sz] - Mult[Iz, Iz, Sz]],
Iz, Ix, Ix, Iz, Sz, Sz, Sz] -
Mult[Inv[Mult[Iy, Iy, Sz] - Mult[Iz, Iz, Sz]],
Iz, Ix, Iz, Ix, Sz, Sz, Sz] -
2 (Mult[Inv[Mult[Iy, Iy, Sz] - Mult[Iz, Iz, Sz]], Iz, Iz, Ix, Ix, Sz, Sz, Sz] -
Mult[Inv[Mult[Iy, Iy, Sz] - Mult[Iz, Iz, Sz]], Iz, Iz, Iz, Iz, Sz, Sz, Sz]))}

```

Harmonic oscillator test case

```
In[176]:= Clear[aL$sym, aR$sym, ω, Q$sym, P$sym, H$sym, Q, P];
```

```
CreateScalar[ω, delta$x$sym, delta$p$sym];
```

```
OscSingle$CreateOperators[aL$sym, aR$sym];
```

```
... OscSingle$CreateOperators: Creating oscillator operators.
```

```
... OscSingle$CreateOperators: Adding oscillator commutations relations.
```

```
In[179]:= H$sym = ω (Mult[aL$sym, aR$sym] + Mult[aR$sym, aL$sym]) / 2;
```

```
In[180]:= CreateOperator[{Q, P}];
{Q$sym, P$sym} = {(aR$sym + aL$sym) / Sqrt[2], I (aR$sym - aL$sym) / Sqrt[2]};
QP$rules = {aR$sym → (Q - I P) / Sqrt[2], aL$sym → (Q + I P) / Sqrt[2]};
```

Check that we recover the expected Hamiltonian written in terms of (P,Q).

```
In[183]:= H$sym /. QP$rules // Simplify
```

```
Out[183]=  $\frac{1}{2} \omega (\text{Mult}[P, P] + \text{Mult}[Q, Q])$ 
```

Here is a test of **Evolver[]** from Evolve-tests.m. It gives the expected result.

```
In[184]:= ExpToTrig[Expand[Evolver[H$sym, t$sym, Q$sym]]] /. QP$rules // Simplify
```

```
Out[184]= Q Cos[t$sym ω] - P Sin[t$sym ω]
```

Here is the same test applied to **Evolver2[]**. It also gives the expected result.

```
In[185]:= ExpToTrig[Expand[Evolver2[H$sym, t$sym, Q$sym]]] /. QP$rules // Simplify
```

```
Out[185]= Q Cos[t$sym ω] - P Sin[t$sym ω]
```

Check that Evolver2[] evolves the raising and lower operators correctly. It does.

```
In[186]:= {Evolver2[H$sym, t$sym, aR$sym], Evolver2[H$sym, t$sym, aL$sym]}
```

```
Out[186]= {aR$sym e-i t$sym ω, aL$sym ei t$sym ω}
```

In the following cell we use Evolve[] expand Q\$sym into a sum of two operators, evolve these operators separately under the Hamiltonian using Evolver2[], add the results together, and convert raising and lower operators into (Q,P) operators. This procedure also gives the correct result.

```
In[187]:= Evolve[H$sym, t$sym, Q$sym] /. Evolve → Evolver2 /. QP$rules // FullSimplify
```

```
Out[187]= Q Cos[t$sym ω] - P Sin[t$sym ω]
```

Let's dig a little deeper into how Evolver2[] gets the right answer. We see that the 0th and 2nd commutators are proportional. The algorithm correctly identifies this as a 2E case. However, when it tries to “divide” the operators, it cannot simplify the resulting expressions.

```
In[188]:= Evolver2[H$sym, t$sym, Q$sym, quiet → False]
```


$$\rho^{(n)} = \left\{ \frac{aL\$sym + aR\$sym}{\sqrt{2}}, \frac{i (aL\$sym - aR\$sym) \omega}{\sqrt{2}}, -\frac{(aL\$sym + aR\$sym) \omega^2}{\sqrt{2}}, -\frac{i (aL\$sym - aR\$sym) \omega^3}{\sqrt{2}} \right\}$$

$$\text{commutators} = \left\{ \frac{i (aL\$sym - aR\$sym) \omega}{\sqrt{2}}, -i \omega, 0, 0 \right\}$$

$$\begin{aligned} \text{divisions} = & \{ i \omega (\text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aL\$sym] - \text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aR\$sym]), \\ & -\omega^2 (\text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aL\$sym] + \text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aR\$sym]), \\ & -\omega^2 (\text{Mult}[\text{Inv}[aL\$sym - aR\$sym], aL\$sym] - \text{Mult}[\text{Inv}[aL\$sym - aR\$sym], aR\$sym]) \} \end{aligned}$$

Case 2E

$$w = \omega \sqrt{\text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aL\$sym] + \text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aR\$sym]}$$

$$\begin{aligned} \text{Out[188]} = & \frac{1}{\sqrt{2}} (\text{Mult}[\text{Cos}[t\$sym \omega \\ & \sqrt{\text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aL\$sym] + \text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aR\$sym]}], \\ & aL\$sym] + \text{Mult}[\text{Cos}[t\$sym \omega \\ & \sqrt{\text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aL\$sym] + \text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aR\$sym]}], \\ & aR\$sym]) + \frac{1}{\sqrt{2}} i \left(\text{Mult}[\text{Sin}[t\$sym \omega \right. \\ & \sqrt{\text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aL\$sym] + \text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aR\$sym]}], \\ & aL\$sym, \\ & \left. \frac{1}{\sqrt{\text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aL\$sym] + \text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aR\$sym]}} \right] - \text{Mult}[\text{Sin}[t\$sym \omega \\ & \sqrt{\text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aL\$sym] + \text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aR\$sym]}], \\ & aR\$sym, \\ & \left. \frac{1}{\sqrt{\text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aL\$sym] + \text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aR\$sym]}} \right] \end{aligned}$$

In the next step we write the raising and lowering operators in terms of (Q, P).

In[189]:= Evolver2[H\$sym, t\$sym, Q\$sym] /. QP\$rules

$$\begin{aligned} \text{Out[189]} = & \frac{1}{\sqrt{2}} \left(\frac{1}{\sqrt{2}} \left(-i \text{Mult}[\text{Cos}[t\$sym \omega \sqrt{\left(\frac{-i P+Q}{\sqrt{2}} + \frac{i P+Q}{\sqrt{2}} \right)}, P] + \text{Mult}[\text{Inv}\left[\frac{-i P+Q}{\sqrt{2}} + \frac{i P+Q}{\sqrt{2}}\right], Q] \right. \right. \\ & \left. \left. + \frac{i \text{Mult}[\text{Inv}\left[\frac{-i P+Q}{\sqrt{2}} + \frac{i P+Q}{\sqrt{2}}\right], P] + \text{Mult}[\text{Inv}\left[\frac{-i P+Q}{\sqrt{2}} + \frac{i P+Q}{\sqrt{2}}\right], Q]}{\sqrt{2}} \right) \right], P \end{aligned}$$

$$\begin{aligned}
& \frac{i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}} \Bigg), Q, \\
& \frac{1}{\sqrt{\frac{-i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}+\frac{i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}}} \\
& \Bigg] + \frac{1}{\sqrt{2}} \left(i \operatorname{Mult}\left[\operatorname{Sin}\left[t_{\text{sym}} \omega\right.\right.\right. \\
& \left.\left.\left.\sqrt{\left(\frac{-i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}+\right.\right.\right. \\
& \left.\left.\left.\frac{i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}\right)\right], P, \\
& \frac{1}{\sqrt{\frac{-i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}+\frac{i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}}} \\
& \Bigg] + \operatorname{Mult}\left[\operatorname{Sin}\left[t_{\text{sym}} \omega\right.\right. \\
& \left.\left.\sqrt{\left(\frac{-i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}+\right.\right.\right. \\
& \left.\left.\left.\frac{i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}\right)\right], Q, \\
& \frac{1}{\sqrt{\frac{-i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}+\frac{i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}}} \\
& \Bigg] \Bigg)
\end{aligned}$$

Amazingly, this simplifies to the expected result.

```
In[190]:= Evolver2[H$sym, t$sym, Q$sym] /. QP$rules // Simplify  
Out[190]= Q Cos[t$sym  $\omega$ ] - P Sin[t$sym  $\omega$ ]
```