

# UniDyn--Study-02.nb

John A. Marohn  
jam99@cornell.edu  
Cornell University

**Abstract:** This demonstration notebook loads the **UniDyn** package, loads the **Inv** function, and runs the **Inv** function unit tests. Below we write a new function for carrying out unitary rotations, **Evolver2[]**, that works on a case-by-case basis. It is approximately 25x faster than **Evolver[]**, and handles more cases.

---

## Set the path to the package

Check the Mathematica version number.

```
In[96]:= $VersionNumber
Out[96]= 12.3
```

Tell *Mathematica* the path to the directory containing the packages.

EDIT THE FOLLOWING PATH STRINGS:

```
In[97]:= $UniDynPath =
  "/Users/jam99/Dropbox/MarohnGroup__Software_Library/UniDyn/
    unidyn";
In[98]:= $UniDynStudyPath =
  "/Users/jam99/Dropbox/MarohnGroup__Software_Library/UniDyn/
    studies";
```

YOU SHOULD NOT NEED TO EDIT ANYTHING FROM HERE ONWARDS.

---

## Load and test the packages

Append the package path to the system path. Before trying to load the package, ask *Mathematica* to find it. This is a test that we directed *Mathematica* to the correct directory. The output of this command should be the full system path to the UniDyn.m file.

```
In[99]:= $Path = AppendTo[$Path, $UniDynPath];
FindFile["UniDyn`"]

Out[100]= /Users/jam99/Dropbox/MarohnGroup__Software_Library/UniDyn/unidyn/UniDyn.m
```

Now that we are confident that the path is set correctly, load the package. Setting the global \$VerboseLoad variable to True will print out the help strings for key commands in the package.

```
In[101]:= $VerboseLoad = True;
Needs["UniDyn`"]

In[103]:= $Path = AppendTo[$Path, $UniDynStudyPath];
FindFile["Inv`"]
FindFile["SpinBoson`"]

Out[104]= /Users/jam99/Dropbox/MarohnGroup__Software_Library/UniDyn/studies/Inv.m
Out[105]= /Users/jam99/Dropbox/MarohnGroup__Software_Library/UniDyn/studies/SpinBoson.m

In[106]:= $VerboseLoad = True;
Needs["Inv`"]
Needs["SpinBoson`"]
```

Extend the directory to the /studies directory and run all the tests found there.

```
In[109]:= SetDirectory[$UniDynStudyPath];
In[110]:= fn = FileNames["*-tests.m"];
test$report = TestReport /@ fn;
TableForm[Table[test$report [[k]], {k, 1, Length[test$report]}]]
```

... SpinSingle\$CreateOperators: Spin operators already exist.  
... SpinSingle\$CreateOperators: Adding spin commutations relations.  
... SpinSingle\$CreateOperators: Angular momentum L = 1/2. Adding operator simplification rules.  
... OscSingle\$CreateOperators: Oscillator operators already exist.  
... OscSingle\$CreateOperators: Adding oscillator commutations relations.

Out[112]/TableForm=

TestReportObject	[ + ]	<input checked="" type="checkbox"/>	Title: Test Report: Evolve2–tests.m Success rate: 100% Tests run: 22	]
TestReportObject	[ + ]	<input checked="" type="checkbox"/>	Title: Test Report: Inv–tests.m Success rate: 100% Tests run: 22	]
TestReportObject	[ + ]	<input checked="" type="checkbox"/>	Title: Test Report: SpinBoson–tests.m Success rate: 100% Tests run: 2	]

## Create quantum optics operators

```
In[113]:= Clear[Ix, Iy, Iz, I$p, I$m, aR, aL];
SpinBoson$CreateOperators[Ix, Iy, Iz, I$p, I$m, aR, aL];

... SpinSingle$CreateOperators: Spin operators already exist.

... SpinSingle$CreateOperators: Adding spin commutations relations.

... SpinSingle$CreateOperators: Angular momentum L = 1/2. Adding operator simplification rules.

... OscSingle$CreateOperators: Oscillator operators already exist.

... OscSingle$CreateOperators: Adding oscillator commutations relations.

... SpinBoson$CreateOperators: Creating operators.

... SpinBoson$CreateOperators: Adding Ip and Im commutations relations.

... SpinBoson$CreateOperators: Adding Ip and Im simplification rules.

... SpinBoson$CreateOperators: Adding aL and aR normal ordering rule.
```

```
In[115]:= Clear[Δ, ω1, g];
CreateScalar[{Δ, ω1, g}];
$Assumptions = {Δ ∈ Reals, ω1 ∈ PositiveReals, g ∈ PositiveReals};
```

Are these assumptions enough? We can see the following examples that `Sqrt[]` involving the real parameters will not simplify correctly if there is more than one real parameter in the `Sqrt[]`. To simplify the `Sqrt[]` with multiple real parameters inside, it is enough to assume that one of the parameters is positive.

```
In[118]:= PowerExpand[Sqrt[-Δ^2]]
```

```
Out[118]=  $\pm \Delta$ 
```

```
In[119]:= PowerExpand[Sqrt[-Δ^2 - ω1^2]]
```

```
Out[119]=  $\sqrt{-\Delta^2 - \omega_1^2}$ 
```

```
In[120]:= PowerExpand[Sqrt[-Δ^2 - ω1^2], Assumptions → $Assumptions]
```

```
Out[120]=  $\pm \sqrt{\Delta^2 + \omega_1^2}$ 
```

```
In[121]:= Simplify[Sqrt[-Δ^2 - ω1^2]]
```

```
Out[121]=  $\pm \sqrt{\Delta^2 + \omega_1^2}$ 
```

### A rotation involving spin raising and lowering operators that succeeds

We can now evolve the raising and lowering operators .

```
In[122]:= {Evolver[\lambda Iz, t, I$p], Evolver[\lambda Iz, t, I$m]}
Out[122]= {e^{-i t \lambda} I$p, e^{i t \lambda} I$m}
```

## Two rotation involving mixed spin and harmonic oscillator operators that fail

Here is a unitary rotation that Evolver[] cannot resolve.

```
In[123]:= terms = Evolver[\lambda Mult[Iz, aL], t, I$p, quiet → False]
```

$$\rho \text{ matrix} = \begin{pmatrix} I$p \\ -i \lambda \text{Mult}[I$p, aL] \\ -\lambda^2 \text{Mult}[I$p, aL, aL] \\ i \lambda^3 \text{Mult}[I$p, aL, aL, aL] \\ \lambda^4 \text{Mult}[I$p, aL, aL, aL, aL] \end{pmatrix}$$

... Evolver: Unrecognized evolution

```
Out[123]= {I$p, -i \lambda \text{Mult}[I$p, aL], -\lambda^2 \text{Mult}[I$p, aL, aL],
           i \lambda^3 \text{Mult}[I$p, aL, aL, aL], \lambda^4 \text{Mult}[I$p, aL, aL, aL, aL]}
```

The ratio of the 1st to the 0th term is independent of I\$p.

```
In[124]:= Mult[Inv[terms[[1]]], terms[[2]]]
```

```
Out[124]= -i aL \lambda
```

The ratio of a higher term to the next lower term is likewise independent of I\$p.

```
In[125]:= {Mult[Inv[terms[[2]]], terms[[3]]], Mult[Inv[terms[[3]]], terms[[4]]]}
```

```
Out[125]= {-i aL \lambda, -i aL \lambda}
```

Here is another unitary rotation that Evolver[] cannot resolve.

```
In[126]:= terms = Evolver[2 I g (Mult[Iz, aR] - Mult[Iz, aL]), t, I$p] // Simplify
```

... Evolver: Unrecognized evolution

```
Out[126]= {I$p, -2 g (Mult[I$p, aL] - Mult[I$p, aR]),
           -4 g^2 (I$p - Mult[I$p, aL, aL] + 2 Mult[I$p, aR, aL] - Mult[I$p, aR, aR]),
           8 g^3 (3 Mult[I$p, aL] - 3 Mult[I$p, aR] - Mult[I$p, aL, aL, aL] +
                  3 Mult[I$p, aR, aL, aL] - 3 Mult[I$p, aR, aR, aL] + Mult[I$p, aR, aR, aR]),
           16 g^4 (3 I$p - 6 Mult[I$p, aL, aL] + 12 Mult[I$p, aR, aL] - 6 Mult[I$p, aR, aR] +
                  Mult[I$p, aL, aL, aL, aL] - 4 Mult[I$p, aR, aL, aL] + 6 Mult[I$p, aR, aR, aL, aL] -
                  4 Mult[I$p, aR, aR, aR, aL] + Mult[I$p, aR, aR, aR, aR])}
```

Notice that again the ratio of the 1st to the 0th term is independent of I\$p.

```
In[127]:= {Mult[Inv[terms[[1]]], terms[[2]]]}
```

```
Out[127]= {-2 (aL - aR) g}
```

# Create a new Evolver algorithm

## Pass 1

Strip down the existing algorithm, and print out some useful intermediate results.

```
In[128]:= VisualComplexity := (Count[ToBoxes[#], Except[" " | "(" | ")"], _String], Infinity] &
Options[Evolver2] = {quiet → True};

In[130]:= Evolver2[H$sym_, t$sym_, rho$sym$0_, opts : OptionsPattern[]] :=
Module[{},

Clear[rho$sym];
rho$sym[0] = rho$sym$0;

Do[
  rho$sym[k + 1] = (-I Comm[H$sym, rho$sym[k]] /. Mult → SortedMult)
    // FullSimplify[#, ComplexityFunction → VisualComplexity] &,
{k, 0, 4}
];

If[OptionValue[quiet] = False,
Print["ρn= ", rho$sym[#] & /@ {0, 1, 2, 3, 4} // MatrixForm];
Print["divisions = ", {Mult[Inv[rho$sym[0]], rho$sym[1]],
  Mult[Inv[rho$sym[0]], rho$sym[2]],
  Mult[Inv[rho$sym[1]], rho$sym[3]]}];
Print["commutators = ", {rho$sym[1],
  Comm[rho$sym[0], rho$sym[1]],
  Comm[rho$sym[0], rho$sym[2]],
  Comm[rho$sym[1], rho$sym[3]]}]]

]
```

The 0th case to consider is when the Hamiltonian commutes with the initial condition. If this is the case we should just stop.

```
In[131]:= terms = Evolver2[Δ I$p, t, I$p, quiet → False]
```

$$\rho_n = \begin{pmatrix} I$p \\ 0 \\ 0 \\ 0 \\ 0 \end{pmatrix}$$

::: Power: Infinite expression  $\frac{1}{0}$  encountered.

::: Infinity: Indeterminate expression 0 ComplexInfinity encountered.

```
divisions = {0, 0, Indeterminate}
```

```
commutators = {0, 0, 0, 0}
```

The 1st case is when the first commutator commutes with the initial condition.

```
In[132]:= terms = Evolver2[g Mult[Iz, aL], t, I$p, quiet → False]
```

$$\rho_n = \begin{pmatrix} I$p \\ -i g \text{Mult}[I$p, aL] \\ -g^2 \text{Mult}[I$p, aL, aL] \\ i g^3 \text{Mult}[I$p, aL, aL, aL] \\ g^4 \text{Mult}[I$p, aL, aL, aL, aL] \end{pmatrix}$$

```
divisions = {-i g aL, -g^2 Mult[aL, aL], -g^2 Mult[aL, aL]}
```

```
commutators = {-i g Mult[I$p, aL], 0, 0, 0}
```

The 2nd case is when the second commutator commutes with the initial condition.

```
In[133]:= Evolver2[Δ Iz, t, Ix, quiet → False]
```

$$\rho_n = \begin{pmatrix} Ix \\ Iy \Delta \\ -Ix \Delta^2 \\ -Iy \Delta^3 \\ Ix \Delta^4 \end{pmatrix}$$

```
divisions = {Δ Mult[Inv[Ix], Iy], -Δ^2, -Δ^2}
```

```
commutators = {Iy Δ, i Iz Δ, 0, 0}
```

The 3rd case is when the third commutator commutes with the first commutator.

```
In[134]:= Evolver2[Δ Iz + ω1 Ix, t, Iz, quiet → False]
```

$$\rho_n = \begin{pmatrix} Iz \\ -Iy \omega_1 \\ \omega_1 (Ix \Delta - Iz \omega_1) \\ Iy \omega_1 (\Delta^2 + \omega_1^2) \\ \omega_1 (-Ix \Delta + Iz \omega_1) (\Delta^2 + \omega_1^2) \end{pmatrix}$$

```
divisions = {-ω1 Mult[Inv[Iz], Iy], ω1 (-ω1 + Δ Mult[Inv[Iz], Ix]), -Δ^2 - ω1^2}
```

```
commutators = {-Iy ω1, i Ix ω1, i Iy Δ ω1, 0}
```

## Pass 2

Now implement cases, writing down the solution to the differential equations by hand.

```
In[135]:= Evolver2::unsolvable = "Unrecognized evolution";

Evolver2[H$sym_, t$sym_, rho$sym$0_, opts : OptionsPattern[]] :=
Module[{rho$sym, divisions, commutators, ρ, w},
ρho$sym[0] = rho$sym$0;
Do[
  rho$sym[k+1] = (-I Comm[H$sym, rho$sym[k]] /. Mult → SortedMult)
    // FullSimplify[#, ComplexityFunction → VisualComplexity] &,
{k, 0, 4}
];
If[OptionValue[quiet] == False, Print["ρ(n) = ", rho$sym[#] & /@ {0, 1, 2, 3}]];
If[rho$sym[1] == 0,
  If[OptionValue[quiet] == False, Print["Case 0"]];
  Return[rho$sym[0]]];
If[Comm[rho$sym[1], H$sym] == 0,
  If[OptionValue[quiet] == False, Print["Case 1L"]];
  Return[rho$sym[0] + rho$sym[1] t$sym]];
commutators = {rho$sym[1],
  Comm[rho$sym[0], rho$sym[1]],
  Comm[rho$sym[0], rho$sym[2]],
  Comm[rho$sym[1], rho$sym[3]]};
If[OptionValue[quiet] == False, Print["commutators = ", commutators]];

divisions = {Mult[Inv[rho$sym[0]], rho$sym[1]],
  Mult[Inv[rho$sym[0]], rho$sym[2]],
  Mult[Inv[rho$sym[1]], rho$sym[3]]};

If[OptionValue[quiet] == False, Print["divisions = ", divisions]];
If[commutators[[2]] == 0,
  If[OptionValue[quiet] == False, Print["Case 1E"]];
```

```

ρ = Mult[Exp[divisions[[1]] t$sym], rho$sym[0]];
Return[ρ];

If[commutators[[3]] == 0,
  If[OptionValue[quiet] == False, Print["Case 2E"]];
  w = PowerExpand[Sqrt[-divisions[[2]]]];
  If[OptionValue[quiet] == False, Print["w = ", w]];
  ρ = Mult[Cos[w t$sym], rho$sym[0]] + Mult[Sin[w t$sym], commutators[[1]] / w];
  Return[ρ];

If[commutators[[4]] == 0,
  If[OptionValue[quiet] == False, Print["Case 3E"]];
  w = PowerExpand[Sqrt[-divisions[[3]]]];
  If[OptionValue[quiet] == False, Print["w = ", w]];
  ρ = rho$sym[0] + Mult[Sin[w t$sym], rho$sym[1] / w] +
    Mult[1 - Cos[w t$sym], rho$sym[2] / w^2];
  Return[ρ];

Message[Evolver2::unsolvable];
Return[{rho$sym[#] & /@ {0, 1, 2, 3}, commutators, divisions}]
]

```

The 0a-th case is when the Hamiltonian commutes with the density operator.

```
In[137]:= Evolver2[λ I$p, t, I$p, quiet → False]
```

$$\rho^{(n)} = \{I$p, 0, 0, 0\}$$

Case 0

```
Out[137]= I$p
```

The 0b-th case is when the Hamiltonian commutes with the first commutator.

```
In[138]:= Evolver2[I g  $\frac{1}{\sqrt{2}}$  (aR - aL), 1,  $\frac{1}{\sqrt{2}}$  (aL + aR), quiet → False] // Simplify
```

$$\rho^{(n)} = \left\{ \frac{aL + aR}{\sqrt{2}}, -g, 0, 0 \right\}$$

Case 1L

$$\text{Out[138]}= \frac{aL + aR}{\sqrt{2}} - g$$

The 1st case is when the first commutator commutes with the initial condition.

```
In[139]:= Evolver2[g Mult[Iz, aL], t, I$p, quiet → False]
```

$$\rho^{(n)} = \{I\$p, -\text{i} g \text{Mult}[I\$p, aL], -g^2 \text{Mult}[I\$p, aL, aL], \text{i} g^3 \text{Mult}[I\$p, aL, aL, aL]\}$$

$$\text{commutators} = \{-\text{i} g \text{Mult}[I\$p, aL], 0, 0, 0\}$$

$$\text{divisions} = \{-\text{i} aL g, -g^2 \text{Mult}[aL, aL], -g^2 \text{Mult}[aL, aL]\}$$

Case 1E

$$\text{Out}[139]= \text{Mult}[e^{-\text{i} aL g t}, I\$p]$$

Compare to

$$\text{In}[140]:= \text{Evolver}[g \text{Mult}[Iz, aL], t, I\$p]$$

**Evolver**: Unrecognized evolution

$$\text{Out}[140]= \{I\$p, -\text{i} g \text{Mult}[I\$p, aL], -g^2 \text{Mult}[I\$p, aL, aL], \\ \text{i} g^3 \text{Mult}[I\$p, aL, aL, aL], g^4 \text{Mult}[I\$p, aL, aL, aL, aL]\}$$

The 2nd case is when the second commutator commutes with the initial condition.

$$\text{In}[141]:= \text{Evolver2}[\Delta Iz, t, Ix, \text{quiet} \rightarrow \text{False}]$$

$$\rho^{(n)} = \{Ix, Iy \Delta, -Ix \Delta^2, -Iy \Delta^3\}$$

$$\text{commutators} = \{Iy \Delta, \text{i} Iz \Delta, 0, 0\}$$

$$\text{divisions} = \{\Delta \text{Mult}[Iz, Iy], -\Delta^2, -\Delta^2\}$$

Case 2E

$$w = \Delta$$

$$\text{Out}[141]= Ix \cos[t \Delta] + Iy \sin[t \Delta]$$

Compare to

$$\text{In}[142]:= \text{Evolver}[\Delta Iz, t, Ix]$$

$$\text{Out}[142]= Ix \cos[t \Delta] + Iy \sin[t \Delta]$$

The 3rd case is when the third commutator commutes with the first commutator.

$$\text{In}[143]:= \text{Evolver2}[\Delta Iz + \omega1 Ix, t, Iz, \text{quiet} \rightarrow \text{False}] // \text{Collect}[\#, \{Ix, Iy, Iz\}, \text{Simplify}] &$$

$$\rho^{(n)} = \{Iz, -Iy \omega1, \omega1 (Ix \Delta - Iz \omega1), Iy \omega1 (\Delta^2 + \omega1^2)\}$$

$$\text{commutators} = \{-Iy \omega1, \text{i} Ix \omega1, \text{i} Iy \Delta \omega1, 0\}$$

$$\text{divisions} = \{-\omega1 \text{Mult}[Iz, Iy], \omega1 (-\omega1 + \Delta \text{Mult}[Iz, Ix]), -\Delta^2 - \omega1^2\}$$

Case 3E

$$w = \sqrt{\Delta^2 + \omega1^2}$$

$$\text{Out}[143]= \frac{Iz (\Delta^2 + \omega1^2 \cos[t \sqrt{\Delta^2 + \omega1^2}])}{\Delta^2 + \omega1^2} + \frac{2 Ix \Delta \omega1 \sin[\frac{1}{2} t \sqrt{\Delta^2 + \omega1^2}]^2}{\Delta^2 + \omega1^2} - \frac{Iy \omega1 \sin[t \sqrt{\Delta^2 + \omega1^2}]}{\sqrt{\Delta^2 + \omega1^2}}$$

Compare to

```
In[144]:= Evolver[Δ Iz + ω1 Ix, t, Iz] // Collect[#, {Ix, Iy, Iz}, Simplify] &
Out[144]= 
$$\frac{Iz (\Delta^2 + \omega1^2 \cos[t \sqrt{\Delta^2 + \omega1^2}])}{\Delta^2 + \omega1^2} + \frac{2 Ix \Delta \omega1 \sin[\frac{1}{2} t \sqrt{\Delta^2 + \omega1^2}]^2}{\Delta^2 + \omega1^2} - \frac{Iy \omega1 \sin[t \sqrt{\Delta^2 + \omega1^2}]}{\sqrt{\Delta^2 + \omega1^2}}$$

```

## Speed test

```
In[145]:= timing1 = Timing[
  Evolver[Δ Iz + ω1 Ix, t, Iz]
]

Out[145]= {0.026488, 
$$\frac{\Delta (Iz \Delta + Ix \omega1) + \omega1 (-Ix \Delta + Iz \omega1) \cos[t \sqrt{\Delta^2 + \omega1^2}] - Iy \omega1 \sqrt{\Delta^2 + \omega1^2} \sin[t \sqrt{\Delta^2 + \omega1^2}]}{\Delta^2 + \omega1^2}$$
}
```

  

```
In[146]:= timing2 = Timing[
  Evolver2[Δ Iz + ω1 Ix, t, Iz]
]

Out[146]= {0.00099, Iz + 
$$\frac{\omega1 (Ix \Delta - Iz \omega1)}{\Delta^2 + \omega1^2} - \frac{\omega1 (Ix \Delta - Iz \omega1) \cos[t \sqrt{\Delta^2 + \omega1^2}]}{\Delta^2 + \omega1^2} - \frac{Iy \omega1 \sin[t \sqrt{\Delta^2 + \omega1^2}]}{\sqrt{\Delta^2 + \omega1^2}}$$
}
```

The ratio of the execution times is

```
In[147]:= Round[timing1[[1]] / timing2[[1]], 1]
Out[147]= 27
```

For this example of **Case 3E** evolution, **Evolver2[]** is approximately **28x faster** than **Evolver[]**.

## Electron transfer test cases

This first case my old code handled .

```
In[148]:= { Evolver[-2 I g Mult[Iz, aR - aL], 1, aL],
  Evolver2[-2 I g Mult[Iz, aR - aL], 1, aL]}
Out[148]= {aL + 2 g Iz, aL + 2 g Iz}
```

These next two more complicated cases, **Evolver[]** failed to handle.

**Evolver2[]** handles these two cases of non-trivial spin-boson evolution!

```
In[149]:= { Evolver[I g (Mult[Iz, aR] - Mult[Iz, aL]), t, I$p],  
Evolver2[I g (Mult[Iz, aR] - Mult[Iz, aL]), t, I$p]}
```

... Evolver: Unrecognized evolution

```
Out[149]= { {I$p, -g (Mult[I$p, aL] - Mult[I$p, aR]),  
-g2 (I$p - Mult[I$p, aL, aL] + 2 Mult[I$p, aR, aL] - Mult[I$p, aR, aR]),  
g3 (3 Mult[I$p, aL] - 3 Mult[I$p, aR] - Mult[I$p, aL, aL, aL] +  
3 (Mult[I$p, aR, aL, aL] - Mult[I$p, aR, aR, aL]) + Mult[I$p, aR, aR, aR]),  
g4 (3 I$p - 6 Mult[I$p, aL, aL] + 12 Mult[I$p, aR, aL] - 6 Mult[I$p, aR, aR] + Mult[  
I$p, aL, aL, aL] - 4 Mult[I$p, aR, aL, aL] + 6 Mult[I$p, aR, aR, aL, aL] -  
4 Mult[I$p, aR, aR, aR, aL] + Mult[I$p, aR, aR, aR, aR])}, Mult[e-(aL-aR) g t, I$p]}
```

```
In[150]:= { Evolver[-2 I g Mult[Iz, aR - aL], t, I$p],  
Evolver2[-2 I g Mult[Iz, aR - aL], t, I$p]}
```

... Evolver: Unrecognized evolution

```
Out[150]= { {I$p, 2 g (Mult[I$p, aL] - Mult[I$p, aR]),  
-4 g2 (I$p - Mult[I$p, aL, aL] + 2 Mult[I$p, aR, aL] - Mult[I$p, aR, aR]),  
-8 g3 (3 Mult[I$p, aL] - 3 Mult[I$p, aR] - Mult[I$p, aL, aL, aL] +  
3 (Mult[I$p, aR, aL, aL] - Mult[I$p, aR, aR, aL]) + Mult[I$p, aR, aR, aR]),  
16 g4 (3 I$p - 6 Mult[I$p, aL, aL] + 12 Mult[I$p, aR, aL] - 6 Mult[I$p, aR, aR] + Mult[  
I$p, aL, aL, aL] - 4 Mult[I$p, aR, aL, aL] + 6 Mult[I$p, aR, aR, aL, aL] -  
4 Mult[I$p, aR, aR, aR, aL] + Mult[I$p, aR, aR, aR, aR])}, Mult[e2 (aL-aR) g t, I$p]}
```

## Create two spin operators (both $I = 1/2$ )

```
In[151]:= Clear[Ix, Iy, Iz, Sx, Sy, Sz];  
CreateOperator[{{Ix, Iy, Iz}, {Sx, Sy, Sz}}];  
SpinSingle$CreateOperators[Ix, Iy, Iz, 1/2];  
SpinSingle$CreateOperators[Sx, Sy, Sz, 1/2];
```

... SpinSingle\$CreateOperators: Spin operators already exist.

... SpinSingle\$CreateOperators: Adding spin commutations relations.

... SpinSingle\$CreateOperators: Angular momentum L = 1/2. Adding operator simplification rules.

... SpinSingle\$CreateOperators: Spin operators already exist.

... SpinSingle\$CreateOperators: Adding spin commutations relations.

... SpinSingle\$CreateOperators: Angular momentum L = 1/2. Adding operator simplification rules.

```
In[155]:= Clear[Δ, J];  
CreateScalar[{Δ, J}];  
$Assumptions = {Δ ∈ Reals, J ∈ Reals};
```

This rotation shows up when calculating evolution of an  $I_x$  under a heteronuclear  $J$  coupling.

For the  $I = 1/2$  case, both **Evolver[]** and **Evolver2[]** know what to do.

```
In[158]:= { Evolver[J Mult[Sz, Iz], t, Ix],
    Evolver2[J Mult[Sz, Iz], t, Ix]}

Out[158]= {Ix Cos[J t/2] + 2 Mult[Iy, Sz] Sin[J t/2], Ix Cos[J t/2] + 2 Mult[Iy, Sz] Sin[J t/2]}
```

---

## Create two spin operators (one $I = 1/2$ )

```
In[159]:= Clear[Ix, Iy, Iz, Sx, Sy, Sz];
CreateOperator[{{Ix, Iy, Iz}, {Sx, Sy, Sz}}];
SpinSingle$CreateOperators[Ix, Iy, Iz, 1/2];
SpinSingle$CreateOperators[Sx, Sy, Sz];

... SpinSingle$CreateOperators: Spin operators already exist.

... SpinSingle$CreateOperators: Adding spin commutations relations.

... SpinSingle$CreateOperators: Angular momentum L = 1/2. Adding operator simplification rules.

... SpinSingle$CreateOperators: Spin operators already exist.

... SpinSingle$CreateOperators: Adding spin commutations relations.

... SpinSingle$CreateOperators: No angular momentum L defined.
```

```
In[163]:= Clear[Δ, J];
CreateScalar[{Δ, J}];
$Assumptions = {Δ ∈ Reals, J ∈ Reals};
```

For the general- $I$  case, **Evolver[]** cannot handle the following rotation.

```
In[166]:= Evolver[J Mult[Sz, Iz], t, Ix]

... Evolver: Unrecognized evolution

Out[166]= {Ix, J Mult[Iy, Sz], -J^2 Mult[Ix, Sz, Sz],
-J^3 Mult[Iy, Sz, Sz, Sz], J^4 Mult[Ix, Sz, Sz, Sz, Sz]}
```

**Evolver2[]** handles the general  $I$ -spin heteronuclear  $J$ -coupling evolution nicely.

```
In[167]:= Evolver2[J Mult[Sz, Iz], t, Ix, quiet → False]
```

```

 $\rho^{(n)} = \{Ix, J \text{Mult}[Iy, Sz], -J^2 \text{Mult}[Ix, Sz, Sz], -J^3 \text{Mult}[Iy, Sz, Sz, Sz]\}$ 
commutators = {J Mult[Iy, Sz], I J Mult[Iz, Sz], 0, 0}
divisions = {J Mult[Inv[Ix], Iy, Sz], -J^2 Mult[Sz, Sz], -J^2 Mult[Sz, Sz]}
Case 2E

```

$w = J \sqrt{\text{Mult}[Sz, Sz]}$

```
Out[167]= Mult[Cos[J t \sqrt{\text{Mult}[Sz, Sz]}], Ix] + Mult[\sin[J t \sqrt{\text{Mult}[Sz, Sz]}], \frac{\text{Mult}[Iy, Sz]}{\sqrt{\text{Mult}[Sz, Sz]}}]
```

We could now think about using the Lowdin projection operator theorem to expand functions like  $\cos[J t \sqrt{\text{Mult}[Sz, Sz]}]$  and  $\sin[J t \sqrt{\text{Mult}[Sz, Sz]}]$ .

Neato!

## Create two spin operators (both general I) [unrecognized evolution]

```

In[168]:= Clear[Ix, Iy, Iz, Sx, Sy, Sz];
CreateOperator[{{Ix, Iy, Iz}, {Sx, Sy, Sz}}];
SpinSingle$CreateOperators[Ix, Iy, Iz];
SpinSingle$CreateOperators[Sx, Sy, Sz];

... SpinSingle$CreateOperators: Spin operators already exist.
... SpinSingle$CreateOperators: Adding spin commutations relations.
... SpinSingle$CreateOperators: No angular momentum L defined.
... SpinSingle$CreateOperators: Spin operators already exist.
... SpinSingle$CreateOperators: Adding spin commutations relations.
... SpinSingle$CreateOperators: No angular momentum L defined.

```

```

In[172]:= Clear[Δ, J];
CreateScalar[{Δ, J}];
$Assumptions = {Δ ∈ Reals, J ∈ Reals};

```

For the general-I case, **Evolver2[]** cannot handle the following rotation.  
We can see that the commutators quickly get out of hand.

```

In[175]:= Evolver2[J Mult[Sz, Iz, Iy], t, Ix]
... Evolver2: Unrecognized evolution
Out[175]= {Ix, J (Mult[Iy, Iy, Sz] - Mult[Iz, Iz, Sz]), 
-J^2 (Mult[Ix, Iy, Iy, Sz, Sz] + Mult[Iy, Ix, Iy, Sz, Sz] + 
Mult[Iz, Ix, Iz, Sz, Sz] + Mult[Iz, Iz, Ix, Sz, Sz]), 
J^3 (2 Mult[Ix, Ix, Iy, Iy, Sz, Sz] + Mult[Ix, Iy, Ix, Iy, Sz, Sz] + 
Mult[Ix, Iy, Ix, Iy, Sz, Sz] + Mult[Ix, Iy, Ix, Iy, Sz, Sz])}

```

$$\begin{aligned}
& \text{Mult[Iy, Ix, Ix, Iy, Sz, Sz, Sz]} - 2 \text{Mult[Iy, Iy, Iy, Iy, Sz, Sz, Sz]} - \\
& \text{Mult[Iz, Ix, Ix, Iz, Sz, Sz, Sz]} - \text{Mult[Iz, Iz, Iz, Ix, Sz, Sz, Sz]} - \\
& 2 (\text{Mult[Iz, Iz, Ix, Ix, Sz, Sz, Sz]} - \text{Mult[Iz, Iz, Iz, Iz, Sz, Sz, Sz]})) \}, \\
\{ & J (\text{Mult[Iy, Iy, Sz]} - \text{Mult[Iz, Iz, Sz]}), J (2 i \text{Mult[Iy, Iz, Sz]} + 2 i \text{Mult[Iz, Iy, Sz]}), \\
& - J^2 (i \text{Mult[Ix, Iy, Iz, Sz, Sz]} + i \text{Mult[Ix, Iz, Iy, Sz, Sz]} - \\
& i \text{Mult[Iy, Iz, Ix, Sz, Sz]} - i \text{Mult[Iz, Iy, Ix, Sz, Sz]}), \\
J & (J^3 (-i \text{Mult[Ix, Ix, Ix, Iz, Sz, Sz, Sz]} - i \text{Mult[Ix, Ix, Iz, Ix, Sz, Sz, Sz]}, \\
& Sz, Iy, Sz] - i \text{Mult[Ix, Iy, Iz, Iy, Sz, Sz, Sz]} - i \text{Mult[Iy, Ix, \\
& Ix, Ix, Iz, Sz, Sz, Sz]} - i \text{Mult[Iy, Ix, Ix, Iz, Ix, Sz, Sz, Sz]} - \\
& i \text{Mult[Iy, Ix, Iy, Iz, Iy, Sz, Sz, Sz]} - i \text{Mult[Iy, Iy, Ix, Ix, Iy, \\
& Sz, Sz, Iy, Sz] - i \text{Mult[Iy, Iy, Ix, Ix, Iy, Sz, Sz, Sz]} - \\
& i \text{Mult[Iy, Iy, Iz, Ix, Iy, Sz, Sz, Sz]} - 2 i \text{Mult[Iy, Iz, Ix, Ix, \\
& Ix, Sz, Sz, Sz]} - i \text{Mult[Iy, Iz, Ix, Iy, Sz, Sz, Sz]} + \\
& 2 i \text{Mult[Iy, Iz, Ix, Iz, Sz, Sz, Sz]} - i \text{Mult[Iy, Iz, Iy, Ix, \\
& Iy, Sz, Sz, Sz]} + i \text{Mult[Iy, Iz, Iz, Ix, Ix, Ix, Sz, Sz, Sz]} + \\
& i \text{Mult[Iy, Iz, Iz, Ix, Sz, Sz, Sz]} - 2 i \text{Mult[Iz, Ix, Ix, Ix, \\
& Sz, Sz, Sz, Iy, Sz]} + 2 (-i \text{Mult[Ix, Iz, Iy, Iy, Sz, Sz, Sz, Iy, Sz]} - \\
& i \text{Mult[Iy, Ix, Iz, Iy, Sz, Sz, Sz]} - i \text{Mult[Iy, Iz, Ix, Iy, \\
& Iy, Sz, Sz, Sz]} - i \text{Mult[Iz, Ix, Iy, Sz, Sz, Sz, Iy, Sz]})) + \\
& 2 i \text{Mult[Iz, Ix, Iz, Iz, Sz, Sz, Sz, Iy, Sz]} - i \text{Mult[Iz, Iy, Ix, Iy, \\
& Sz, Sz, Sz, Sz]} + i \text{Mult[Iz, Iz, Ix, Ix, Ix, Sz, Sz, Sz, Iy, Sz]} - \\
& 2 (i \text{Mult[Ix, Iz, Ix, Ix, Sz, Sz, Sz, Iy, Sz]} - i \text{Mult[Ix, Iz, Iz, Ix, Sz, \\
& Sz, Sz, Iy, Sz]} + i \text{Mult[Iy, Ix, Iz, Ix, Ix, Sz, Sz, Sz, Sz]} - i \text{Mult[Iy, Ix, \\
& Iz, Iz, Iz, Sz, Sz, Sz]} + i \text{Mult[Iy, Iz, Ix, Ix, Ix, Sz, Sz, Sz, Sz]} - \\
& i \text{Mult[Iy, Iz, Ix, Iz, Iz, Sz, Sz, Sz, Sz]} - 2 i \text{Mult[Iy, Iz, Iz, Ix, \\
& Iz, Sz, Sz, Sz]} - 2 i \text{Mult[Iy, Iz, Iz, Iz, Ix, Sz, Sz, Sz, Sz]} + \\
& i \text{Mult[Iz, Ix, Ix, Ix, Sz, Sz, Sz, Iy, Sz]} - i \text{Mult[Iz, Ix, Iz, Ix, Sz, Sz, \\
& Sz, Iy, Sz]} - 2 i \text{Mult[Iz, Iz, Ix, Iz, Sz, Sz, Iy, Sz]} - 2 i \text{Mult[Iz, Iz, \\
& Iz, Ix, Sz, Sz, Sz, Iy, Sz]} + i \text{Mult[Iz, Iz, Iz, Ix, Sz, Sz, Sz, Iy, Sz]})) - \\
J^3 & (-2 i \text{Mult[Ix, Ix, Ix, Iy, Sz, Sz, Sz, Iz, Sz]} - i \text{Mult[Ix, Iy, Ix, Ix, \\
& Sz, Sz, Sz, Ix, Sz, Sz]} + i \text{Mult[Ix, Iy, Iy, Iy, Sz, Sz, Sz, Ix, Sz]} - \\
& i \text{Mult[Iy, Ix, Ix, Ix, Sz, Sz, Sz, Ix, Sz]} + i \text{Mult[Iy, Ix, Iy, Iy, \\
& Sz, Sz, Sz, Ix, Sz]} + 2 i \text{Mult[Iy, Iy, Ix, Iy, Sz, Sz, Sz, Ix, Sz]} - \\
& 2 i \text{Mult[Iz, Ix, Ix, Ix, Iy, Sz, Sz, Sz, Sz]} - i \text{Mult[Iz, Ix, Iy, Ix, \\
& Ix, Sz, Sz, Sz, Sz]} + i \text{Mult[Iz, Ix, Ix, Iy, Iy, Iy, Sz, Sz, Sz, Sz]} - \\
& i \text{Mult[Iz, Ix, Ix, Iy, Ix, Sz, Sz, Sz, Sz]} - i \text{Mult[Iz, Ix, Ix, Ix, Iy, \\
& Sz, Sz, Sz, Ix, Sz, Sz, Sz]} + i \text{Mult[Iz, Ix, Ix, Ix, Ix, Sz, Sz, Sz, Sz]} + \\
& 2 (-i \text{Mult[Ix, Ix, Ix, Iy, Sz, Sz, Sz, Iz, Sz]} - i \text{Mult[Ix, Ix, Iy, Ix, \\
& Sz, Sz, Sz, Sz]} + i \text{Mult[Ix, Iy, Iy, Iy, Sz, Sz, Sz, Sz]} + i \text{Mult[Iy, Ix, \\
& Iy, Iy, Sz, Sz, Sz, Iz, Sz]} - i \text{Mult[Iz, Ix, Ix, Ix, Iy, Sz, Sz, Sz, Sz]} - \\
& i \text{Mult[Iz, Ix, Ix, Iy, Ix, Sz, Sz, Sz, Sz]} + i \text{Mult[Iz, Ix, Iy, Iy, \\
& Iy, Sz, Sz, Sz, Sz]} + i \text{Mult[Iz, Iy, Ix, Iy, Iy, Sz, Sz, Sz, Sz]} + \\
& i \text{Mult[Iz, Iy, Ix, Iy, Iy, Sz, Sz, Sz, Sz]} - i \text{Mult[Iz, Iy, Ix, Ix, Ix, Iy, \\
& Iy, Sz, Sz, Sz, Sz]} + i \text{Mult[Iz, Iy, Ix, Ix, Iy, Iy, Sz, Sz, Sz, Sz])) +
\end{aligned}$$

$$\begin{aligned}
& \text{Sz, Sz, Sz, Iz, Sz] + 2 i \text{Mult}[Iz, Iy, Iy, Ix, Iy, Sz, Sz, Sz, Sz] -} \\
& 2 (-i \text{Mult}[Ix, Iy, Iy, Iy, Sz, Sz, Sz, Iz, Sz] - i \text{Mult}[Iy, Ix, Iy, Iy, Sz, \\
& \quad \text{Sz, Sz, Iz, Sz}] - i \text{Mult}[Iy, Iy, Ix, Iy, Sz, Sz, Sz, Iz, Sz] - i \text{Mult}[Iy, Iy, \\
& \quad \text{Iy, Ix, Sz, Sz, Sz, Iz, Sz}] - i \text{Mult}[Iz, Ix, Iy, Iy, Iy, Sz, Sz, Sz, Sz] - \\
& \quad i \text{Mult}[Iz, Iy, Ix, Iy, Iy, Sz, Sz, Sz] - i \text{Mult}[Iz, Iy, Iy, Ix, \\
& \quad \text{Iy, Sz, Sz, Sz}] - i \text{Mult}[Iz, Iy, Iy, Iy, Ix, Sz, Sz, Sz, Sz]) - \\
& i \text{Mult}[Iz, Iy, Iz, Ix, Sz, Sz, Sz, Iz, Sz] - i \text{Mult}[Iz, Iz, Ix, Iy, \\
& \quad \text{Iz, Sz, Sz, Sz}] - i \text{Mult}[Iz, Iz, Ix, Iy, Sz, Sz, Sz, Sz] - \\
& i \text{Mult}[Iz, Iz, Iy, Iz, Ix, Sz, Sz, Sz, Sz] - \\
& 2 (i \text{Mult}[Iz, Iz, Ix, Iy, Sz, Sz, Sz, Iz, Sz] + \\
& \quad i \text{Mult}[Iz, Iz, Iy, Ix, Sz, Sz, Sz, Iz, Sz] + i \text{Mult}[Iz, Iz, Iz, Ix, Iy, \\
& \quad \text{Sz, Sz, Sz, Sz}] + i \text{Mult}[Iz, Iz, Iz, Iy, Ix, Sz, Sz, Sz, Sz])) \}, \\
& \{ J (\text{Mult}[\text{Inv}[Ix], Iy, Iy, Sz] - \text{Mult}[\text{Inv}[Ix], Iz, Iz, Sz]), \\
& - J^2 \\
& (\text{Mult}[Iy, Iy, Sz, Sz] + \\
& \quad \text{Mult}[\text{Inv}[Ix], Iy, Ix, Iy, Sz, Sz] + \\
& \quad \text{Mult}[\text{Inv}[Ix], Iz, Ix, Iz, Sz, Sz] + \\
& \quad \text{Mult}[\text{Inv}[Ix], Iz, Iz, Ix, Sz, Sz]), \\
J^2 & (2 \text{Mult}[\text{Inv}[\text{Mult}[Iy, Iy, Sz] - \text{Mult}[Iz, Iz, Sz]], Ix, Ix, Iy, Iy, Sz, Sz, Sz] + \\
& \quad \text{Mult}[\text{Inv}[\text{Mult}[Iy, Iy, Sz] - \text{Mult}[Iz, Iz, Sz]], \\
& \quad Ix, Iy, Ix, Iy, Sz, Sz, Sz] + \\
& \quad \text{Mult}[\text{Inv}[\text{Mult}[Iy, Iy, Sz] - \text{Mult}[Iz, Iz, Sz]], \\
& \quad Iy, Ix, Ix, Iy, Sz, Sz, Sz] - \\
& 2 \text{Mult}[\text{Inv}[\text{Mult}[Iy, Iy, Sz] - \text{Mult}[Iz, Iz, Sz]], Iy, Iy, Iy, Iy, Sz, Sz, Sz] - \\
& \quad \text{Mult}[\text{Inv}[\text{Mult}[Iy, Iy, Sz] - \text{Mult}[Iz, Iz, Sz]], \\
& \quad Iz, Ix, Ix, Iz, Sz, Sz, Sz] - \\
& \quad \text{Mult}[\text{Inv}[\text{Mult}[Iy, Iy, Sz] - \text{Mult}[Iz, Iz, Sz]], \\
& \quad Iz, Ix, Iz, Ix, Sz, Sz, Sz] - \\
& 2 (\text{Mult}[\text{Inv}[\text{Mult}[Iy, Iy, Sz] - \text{Mult}[Iz, Iz, Sz]], Iz, Iz, Ix, Ix, Sz, Sz, Sz] - \\
& \quad \text{Mult}[\text{Inv}[\text{Mult}[Iy, Iy, Sz] - \text{Mult}[Iz, Iz, Sz]], Iz, Iz, Iz, Iz, Sz, Sz, Sz])) \}
\end{aligned}$$

## Harmonic oscillator test case

```

In[176]:= Clear[aL$sym, aR$sym, ω, Q$sym, P$sym, H$sym, Q, P];
CreateScalar[ω, delta$x$sym, delta$p$sym];
OscSingle$CreateOperators[aL$sym, aR$sym];

... OscSingle$CreateOperators: Creating oscillator operators.
... OscSingle$CreateOperators: Adding oscillator commutations relations.

In[179]:= H$sym = ω (Mult[aL$sym, aR$sym] + Mult[aR$sym, aL$sym]) / 2;

```

```
In[180]:= CreateOperator[{{Q, P}}];
{Q$sym, P$sym} = {(aR$sym + aL$sym) / Sqrt[2], I (aR$sym - aL$sym) / Sqrt[2]};
QP$rules = {aR$sym → (Q - I P) / Sqrt[2], aL$sym → (Q + I P) / Sqrt[2]};
```

Check that we recover the expected Hamiltonian written in terms of (P,Q).

```
In[183]:= H$sym /. QP$rules // Simplify
Out[183]=  $\frac{1}{2} \omega (\text{Mult}[P, P] + \text{Mult}[Q, Q])$ 
```

Here is a test of **Evolver[]** from Evolve-tests.m. It gives the expected result.

```
In[184]:= ExpToTrig[Expand[Evolver[H$sym, t$sym, Q$sym]]] /. QP$rules // Simplify
Out[184]= Q Cos[t$sym  $\omega$ ] - P Sin[t$sym  $\omega$ ]
```

Here is the same test applied to **Evolver2[]**. It also gives the expected result.

```
In[185]:= ExpToTrig[Expand[Evolver2[H$sym, t$sym, Q$sym]]] /. QP$rules // Simplify
Out[185]= Q Cos[t$sym  $\omega$ ] - P Sin[t$sym  $\omega$ ]
```

Check that Evolver2[] evolves the raising and lower operators correctly. It does.

```
In[186]:= {Evolver2[H$sym, t$sym, aR$sym], Evolver2[H$sym, t$sym, aL$sym]}
Out[186]= {aR$sym  $e^{-i t$sym \omega}$ , aL$sym  $e^{i t$sym \omega}$ }
```

In the following cell we use Evolve[] expand Q\$sym into a sum of two operators, evolve these operators separately under the Hamiltonian using Evolver2[], add the results together, and convert raising and lower operators into (Q,P) operators. This procedure also gives the correct result.

```
In[187]:= Evolve[H$sym, t$sym, Q$sym] /. Evolve → Evolver2 /. QP$rules // FullSimplify
Out[187]= Q Cos[t$sym  $\omega$ ] - P Sin[t$sym  $\omega$ ]
```

Let's dig a little deeper into how Evolver2[] gets the right answer. We see that the 0th and 2nd commutators are proportional. The algorithm correctly identifies this as a 2E case. However, when it tries to "divide" the operators, it cannot simplify the resulting expressions.

```
In[188]:= Evolver2[H$sym, t$sym, Q$sym, quiet → False]
```

$$\rho^{(n)} = \left\{ \frac{aL\$sym + aR\$sym}{\sqrt{2}}, \frac{\pm (aL\$sym - aR\$sym) \omega}{\sqrt{2}}, -\frac{(aL\$sym + aR\$sym) \omega^2}{\sqrt{2}}, -\frac{\pm (aL\$sym - aR\$sym) \omega^3}{\sqrt{2}} \right\}$$

$$\text{commutators} = \left\{ \frac{\pm (aL\$sym - aR\$sym) \omega}{\sqrt{2}}, -\pm \omega, 0, 0 \right\}$$

$$\text{divisions} = \{ \pm \omega (\text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aL\$sym] - \text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aR\$sym]), -\omega^2 (\text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aL\$sym] + \text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aR\$sym]), -\omega^2 (\text{Mult}[\text{Inv}[aL\$sym - aR\$sym], aL\$sym] - \text{Mult}[\text{Inv}[aL\$sym - aR\$sym], aR\$sym]) \}$$

Case 2E

$$w = \omega \sqrt{\text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aL\$sym] + \text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aR\$sym]}$$

$$\text{Out[188]} = \frac{1}{\sqrt{2}} \left( \text{Mult}[\text{Cos}[t\$sym \omega \sqrt{\text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aL\$sym] + \text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aR\$sym]}], aL\$sym] + \text{Mult}[\text{Cos}[t\$sym \omega \sqrt{\text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aL\$sym] + \text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aR\$sym]}], aR\$sym] \right) + \frac{1}{\sqrt{2}} \pm \left( \text{Mult}[\text{Sin}[t\$sym \omega \sqrt{\text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aL\$sym] + \text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aR\$sym]}], aL\$sym], \frac{1}{\sqrt{2}} \text{Mult}[\text{Sin}[t\$sym \omega \sqrt{\text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aL\$sym] + \text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aR\$sym]}], aR\$sym] \right) - \frac{1}{\sqrt{2}} \left( \text{Mult}[\text{Sin}[t\$sym \omega \sqrt{\text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aL\$sym] + \text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aR\$sym]}], aL\$sym], \frac{1}{\sqrt{2}} \text{Mult}[\text{Sin}[t\$sym \omega \sqrt{\text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aL\$sym] + \text{Mult}[\text{Inv}[aL\$sym + aR\$sym], aR\$sym]}], aR\$sym] \right)$$

In the next step we write the raising and lowering operators in terms of  $(Q, P)$ .

In[189]:= Evolver2[H\\$sym, t\\$sym, Q\\$sym] /. QP\$rules

$$\text{Out[189]} = \frac{1}{\sqrt{2}} \left( \frac{1}{\sqrt{2}} \left( -\pm \text{Mult}[\text{Cos}[t\$sym \omega \sqrt{\frac{-\pm P+Q}{\sqrt{2}} + \frac{\pm P+Q}{\sqrt{2}}}], P] + \text{Mult}[\text{Inv}[\frac{-\pm P+Q}{\sqrt{2}} + \frac{\pm P+Q}{\sqrt{2}}], Q] \right) + \frac{\pm \text{Mult}[\text{Inv}[\frac{-\pm P+Q}{\sqrt{2}} + \frac{\pm P+Q}{\sqrt{2}}], P] + \text{Mult}[\text{Inv}[\frac{-\pm P+Q}{\sqrt{2}} + \frac{\pm P+Q}{\sqrt{2}}], Q]}{\sqrt{2}} \right), P \right) +$$



$$\begin{aligned}
& \frac{\frac{i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}\Bigg], Q, \\
& \frac{1}{\sqrt{\frac{-i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}+\frac{i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}}} \\
& \Bigg]+\frac{1}{\sqrt{2}}\left(\frac{i \operatorname{Mult}\left[\sin [\operatorname{t\$sym} \omega\right.}{\sqrt{\left(\frac{-i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}+\right.}\right.} \\
& \left.\left.\frac{i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}\right)\Bigg], P, \\
& \frac{1}{\sqrt{\frac{-i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}+\frac{i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}}} \\
& \Bigg)+\operatorname{Mult}\left[\sin [\operatorname{t\$sym} \omega\right. \\
& \left.\left.\frac{-i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}+\right.}\right.} \\
& \left.\left.\frac{i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}\right)\Bigg], Q, \\
& \frac{1}{\sqrt{\frac{-i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}+\frac{i \operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], P\right]+\operatorname{Mult}\left[\operatorname{Inv}\left[\frac{-i P+Q}{\sqrt{2}}+\frac{i P+Q}{\sqrt{2}}\right], Q\right]}{\sqrt{2}}}} \\
& \Bigg)
\end{aligned}$$

Amazingly, this simplifies to the expected result.

```
In[190]:= Evolver2[H$sym, t$sym, Q$sym] /. QP$rules // Simplify
Out[190]= Q Cos[t$sym \[Omega]] - P Sin[t$sym \[Omega]]
```