

# UniDyn--Demo-02.nb

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**Abstract:** Use the **UniDyn** Evolver function to calculate the evolution of the magnetization of a single spin 1/2 particle under off-resonance, variable-phase irradiation. Plot the evolving magnetization for various combinations of resonance offset and irradiation phase.

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## Set the path to the package

Tell *Mathematica* the path to the directory containing the package.

EDIT THE FOLLOWING PATH STRINGS:

```
In[1]:= $NCPATH = "/Users/jam99/Dropbox";  
$UniDynPath =  
    "/Users/jam99/Dropbox/MarohnGroup__Software_Library/UniDyn/  
    unidyn";
```

YOU SHOULD NOT NEED TO EDIT ANYTHING FROM HERE ONWARDS.

---

## Load the package

Append the package path to the system path. Before trying to load the package, ask *Mathematica* to find it. This is a test that we directed *Mathematica* to the correct directory. The output of this command should be the full system path to the Uni-Dyn.m file.

```
In[3]:= $Path = AppendTo[$Path, $NCPath];
$Path = AppendTo[$Path, $UniDynPath];
FindFile["UniDyn`"]
FindFile["NC`"]
```

```
Out[5]= /Users/jam99/Dropbox/MarohnGroup__Software_Library/UniDyn/unidyn/UniDyn.m
```

```
Out[6]= /Users/jam99/Dropbox/NC/init.m
```

Now that we are confident that the path is set correctly, load the package. Setting the global `$VerboseLoad` variable to `True` will print out the help strings for key commands in the package.

```
In[7]:= $VerboseLoad = True;
Needs["UniDyn`"]
```

```
... NC: You are using the version of NCAIgebra which is found in: " /Users /jam99 /Dropbox /NC/".
```

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 NCAlgebra - Version 5.0.6  
 Compatible with Mathematica Version 10 and above

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 Helton 2002  
 Helton and Miller June 1991  
 All rights reserved.

The program was written by the authors and by:

David Hurst, Daniel Lamm, Orlando Merino, Robert Obar,  
 Henry Pfister, Mike Walker, John Wavrik, Lois Yu,  
 J. Camino, J. Griffin, J. Oval, T. Shaheen, John Shopple.  
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For NCAlgebra updates see:

[www.github.com/NCAlgebra/NC](http://www.github.com/NCAlgebra/NC)  
[www.math.ucsd.edu/~ncalg](http://www.math.ucsd.edu/~ncalg)

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 ... NCAlgebra : All lower cap single letter symbols (e.g. a,b,c,... ) were set as noncommutative.

... CreateOperator : CreateOperator [] is used to batch -define a bunch of operators. Example: CreateOperator [{{lx,  
 ly, lz},{Sx,Sy,Sz}}] will create six operators; each of the operators in the first list is meant to commute with  
 each of the operators in the second list.

- ... **CreateScalar** : CreateScalar [list] is used to batch —define a bunch of scalars. The parameter list can be a single scalar or a list of scalars. Example: CreateScalar [{w1,w2 }].
- ... **NCSort** : NCSort [list] sorts the operators in list into canonical order.
- ... **SortedMult** : SortedMult [list] returns Mult [list\$ordered ], where list\$ordered are the elements of list sorted into canonical order.
- ... **MultSort** : MultSort [NonCommutativeMultiply [list]] returns returns NonCommutativeMultiply [list\$ordered ], where list\$ordered are the elements of list sorted into canonical order.
- ... **Comm** : Comm [a,b] calculates the commutator of two operators.
- ... **SpinSingle\$CreateOperators** : SpinSingle\$CreateOperators [lx,ly,lz,L ] creates lx, ly, and lz angular momentum operators and defines their commutation relations. When the total angular momentum L = 1/2, additional rules are defined to simplify products of the angular momentum operators. When the total angular momentum L is unspecified, no such simplification rules are defined.
- ... **OscSingle\$CreateOperators** : OscSingle\$CreateOperators [aL,aR ] creates a raising operator aR and a lowering operator aL for single harmonic oscillator and defines the operator commutation relations.
- ... **Evolve** : Evolve [H, t,  $\rho$ ] represents unitary evolution of the density operator  $\rho$  for a time t under the Hamiltonian H. This function expands according to simplification rules but leaves the evolution unevaluated.
- ... **Evolver** : Evolver [H, t,  $\rho(0)$ ] calculates  $\rho(t) = \text{Exp}[-i H t] \rho(0) \text{Exp}[+i H t]$ , assuming that H is time independent, according to the commutation rules followed by  $\rho(0)$  and H.

---

## Function to help draw the magnetization

```
In[9]:= Clear[my$drawing];
SetAttributes[my$drawing, HoldAll];

my$drawing[func_[t_, a___], t$final_, N$step_] :=

Module[{ $\rho$ $vector$data,  $\rho$ $arrows, axes$arrows, big$plot},

(* Calculate a final time and a time step *)
(* The final time point should not be included in the plot *)

T$max = t$final * (N$step - 1) / N$step;
T$step = t$final / N$step;

(* Make a table of data of the form *)
(* {{0.,{0.,0.,1.}},{0.5,{0.,-0.9,-0.5}}} *)
```

```

ρ$vector$data =
  Table[{N[t/T$max], N[func[t, a]]}, {t, 0, T$max, T$step}];

(* Add arrows; the arrows grow from
   light to dark as time progresses in the plot*)

ρ$arrows =
  Graphics3D[{GrayLevel[1.0 - N[#1]], {Arrowheads[0.015],
    Arrow[Tube[{0, 0, 0}, #2]]}}] & /@ ρ$vector$data;

(* Add axes arrows. Here we make
   the assumption that the magnetization vector *)
(* has a magnitude of 1. *)

axes$arrows =
  Graphics3D[{Black, Arrow[Tube[{0, 0, 0}, #]]]} & /@
    {{0, 0, 1.25}, {0, 1.25, 0}, {1.25, 0, 0},
     {0, 0, -1.25}, {0, -1.25, 0}, {-1.25, 0, 0}};

big$plot = Flatten[Append[ρ$arrows, axes$arrows]];

(* Add axes labels. *)

big$plot = Flatten[Append[big$plot,
  Graphics3D[Text[Style[z, Large], {0, 0, 1.35}]]]];
big$plot = Flatten[Append[big$plot,
  Graphics3D[Text[Style[y, Large], {0, 1.35, 0}]]]];
big$plot = Flatten[Append[big$plot,
  Graphics3D[Text[Style[x, Large], {1.35, 0, 0}]]]];

(* Plot all the arrows. The neutral
   lighting helps making the rendering fast -- *)
(* the default Mathematica camera has three-
   colored lights which makes funny reflections *)
(* of off small objects like our arrows. *)

```

```
Show[big$plot, Boxed → False, ViewVertical → {0, 0, 1},
      ViewPoint → {2.0, -1.0, 1.0}, Lighting → "Neutral"]

]
```

## Examples of unitary evolution in a spin 1/2 system

### Create a single spin

The assumptions define below are required for *Mathematica* to recognize  $\sqrt{-\Delta^2 - \omega^2} = i \sqrt{\Delta^2 + \omega^2}$  inside an exponential. One of the variables has to be defined to be  $> 0$  and not just  $\geq 0$ .

```
In[12]:= Clear[

      Δ,          (* resonance offset frequency *)
      ω,          (* Rabi frequency of the applied irradiation *)
      ϕ,          (* phase of the applied irradiation *)
      t,          (* time *)
      Ix, Iy, Iz, (* spin angular momentum operators *)
      ρ,          (* spin density operator *)
      ρ$0,        (* initial spin density operator *)
      H           (* spin Hamiltonian *)]
```

```
CreateScalar[Δ, ω, ϕ, t];
SpinSingle$CreateOperators[Ix, Iy, Iz, L = 1/2];
```

```
$Assumptions = {Element[Δ, Reals], Δ ≥ 0,
                Element[ω, Reals], ω > 0, Element[t, Reals], t ≥ 0};
```

```
... SpinSingle$CreateOperators : Creating spin operators.
```

```
... SpinSingle$CreateOperators : Adding spin commutations relations.
```

```
... SpinSingle$CreateOperators : Angular momentum L = 1/2. Adding operator simplification rules.
```

### Off-resonance variable-phase nutation

Irradiation Hamiltonian written in the interaction representation. The initial density

operator is parallel to  $I_z$ .

```
In[16]:= H = Δ Iz + ω (Cos[φ] Ix + Sin[φ] Iy);
ρ$0 = Iz;
```

Calculating the time-dependent density operator might take as long as 10 to 15 seconds to complete.

```
In[18]:= ρ[t_, Δ_, ω_, φ_] = Collect[
    (Evolver[H, t, ρ$0] // Simplify // ExpToTrig //
    FullSimplify),
    {Ix, Iy, Iz}];

In[19]:= ρ[t, Δ, ω, φ] /. {Δ → Subscript[ω, 0], ω → Subscript[ω, 1]}
```

$$\begin{aligned} \text{Out[19]= } & \frac{Iz \left( \omega_0^2 + \cos \left[ t \sqrt{\omega_0^2 + \omega_1^2} \right] \omega_1^2 \right)}{\omega_0^2 + \omega_1^2} + \frac{1}{\omega_0^2 + \omega_1^2} \\ & Iy \left( \sin[\phi] \omega_0 \omega_1 - \cos \left[ t \sqrt{\omega_0^2 + \omega_1^2} \right] \sin[\phi] \omega_0 \omega_1 - \right. \\ & \quad \left. \cos[\phi] \sin \left[ t \sqrt{\omega_0^2 + \omega_1^2} \right] \omega_1 \sqrt{\omega_0^2 + \omega_1^2} \right) + \frac{1}{\omega_0^2 + \omega_1^2} \\ & Ix \left( \cos[\phi] \omega_0 \omega_1 - \cos[\phi] \cos \left[ t \sqrt{\omega_0^2 + \omega_1^2} \right] \omega_0 \omega_1 + \right. \\ & \quad \left. \sin[\phi] \sin \left[ t \sqrt{\omega_0^2 + \omega_1^2} \right] \omega_1 \sqrt{\omega_0^2 + \omega_1^2} \right) \end{aligned}$$

Below we want a function that returns a triple of numbers describing the magnetization vector. We turn the above expression for the density operator into a triple of numbers using the *Mathematica* function `Coefficient`. I tried using the `NCAgebra`'s `NCCoefficient` function but could not get it to work. The function below does what we want.

```
In[20]:=  $\rho$ $vector[t_, Δ_, ω_, ϕ_] =
Simplify[Coefficient[ρ[t, Δ, ω, ϕ], #, 1] & /@ {Ix, Iy, Iz}]
```

$$\text{Out[20]} = \left\{ \frac{1}{\Delta^2 + \omega^2} \omega \left( \cos[\phi] \left( \Delta - \Delta \cos\left[t \sqrt{\Delta^2 + \omega^2}\right] \right) + \sqrt{\Delta^2 + \omega^2} \sin[\phi] \sin\left[t \sqrt{\Delta^2 + \omega^2}\right] \right), \right. \\ \left. - \frac{1}{\Delta^2 + \omega^2} \omega \left( \Delta \left( -1 + \cos\left[t \sqrt{\Delta^2 + \omega^2}\right] \right) \sin[\phi] + \sqrt{\Delta^2 + \omega^2} \cos[\phi] \sin\left[t \sqrt{\Delta^2 + \omega^2}\right] \right), \frac{\Delta^2 + \omega^2 \cos\left[t \sqrt{\Delta^2 + \omega^2}\right]}{\Delta^2 + \omega^2} \right\}$$

### Check limiting cases

On resonance, the effective field is in the x-y plane. The z magnetization will oscillate co-sinusoidally while the magnetization in the x-y plane will oscillate sinusoidally.

```
In[21]:= ρ[t, 0, ω, ϕ] // PowerExpand // FullSimplify
```

```
Out[21]:= Iz Cos[t ω] + (-Iy Cos[ϕ] + Ix Sin[ϕ]) Sin[t ω]
```

Apply a  $\pi$  pulse. Observe that the magnetization is indeed inverted.

```
In[22]:= ρ[π/ω, 0, ω, 0] // PowerExpand
```

```
Out[22]:= -Iz
```

Now apply a  $\pi/2$  pulse. Applying an “x” pulse, one with a relative phase of  $\phi=0$ , places the magnetization along the -y axis. A “y” pulse, one with a relative phase of  $\phi = \pi/2$ , places the magnetization along the +x axis.

```
In[23]:= ρ[π/(2 ω), 0, ω, #] & /@ {0, π/2} // PowerExpand
```

```
Out[23]:= {-Iy, Ix}
```

### Draw the magnetization

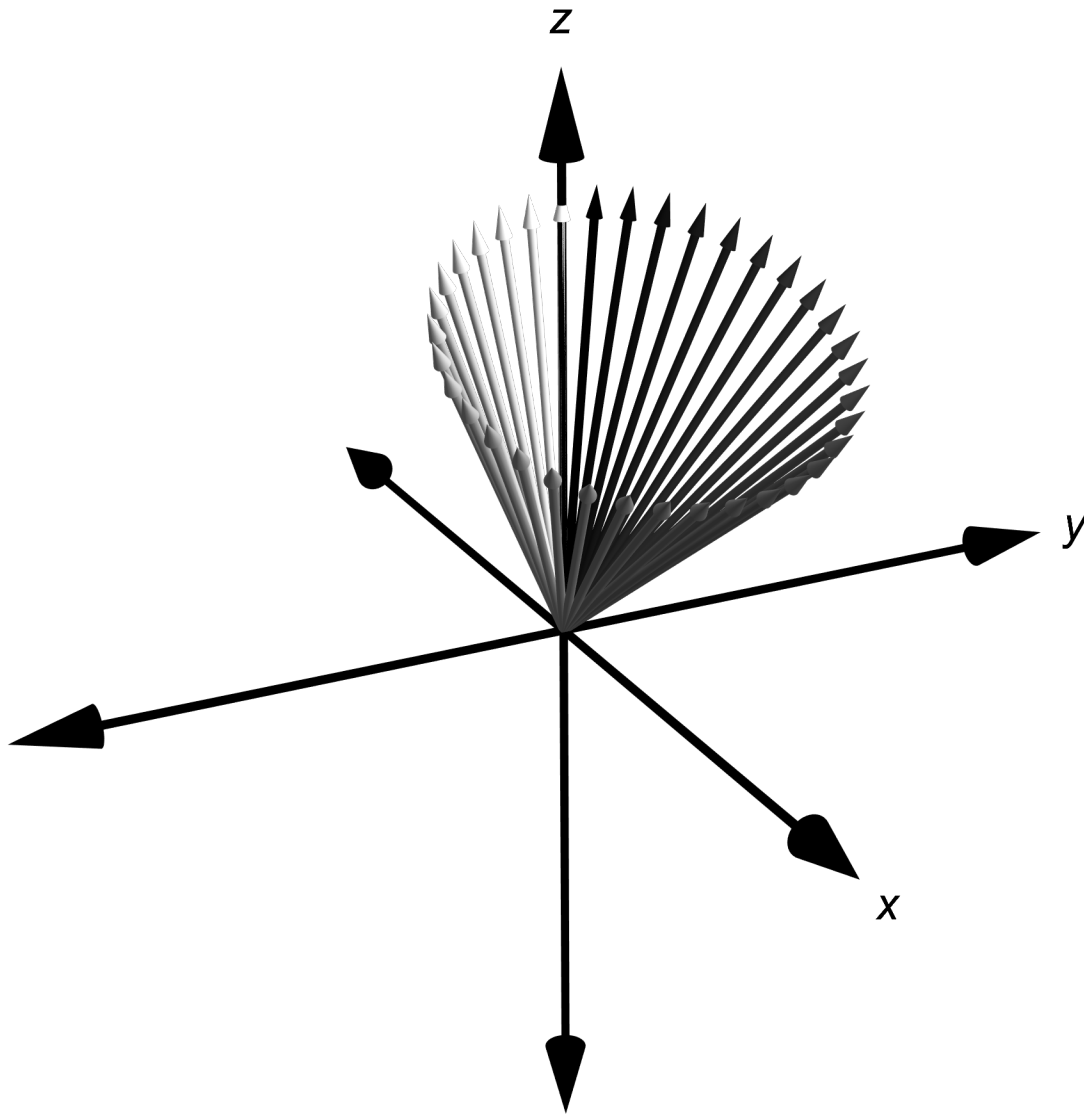
Set the rf phase to  $\phi=0$ , set the Rabi frequency to  $\omega=1$ , and set the resonance offset to  $\Delta=2$ .

The effective field has a magnitude of  $\sqrt{2^2 + 1^2} = \sqrt{5}$ , so we'll watch the magnetization out to a time of  $2\pi / \sqrt{5}$  in order to capture the magnetization orbiting once around the effective field.



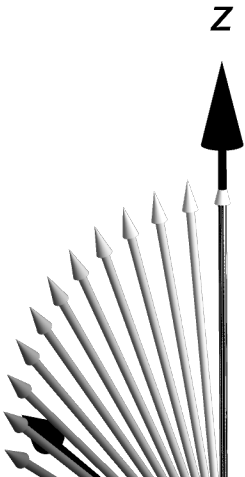
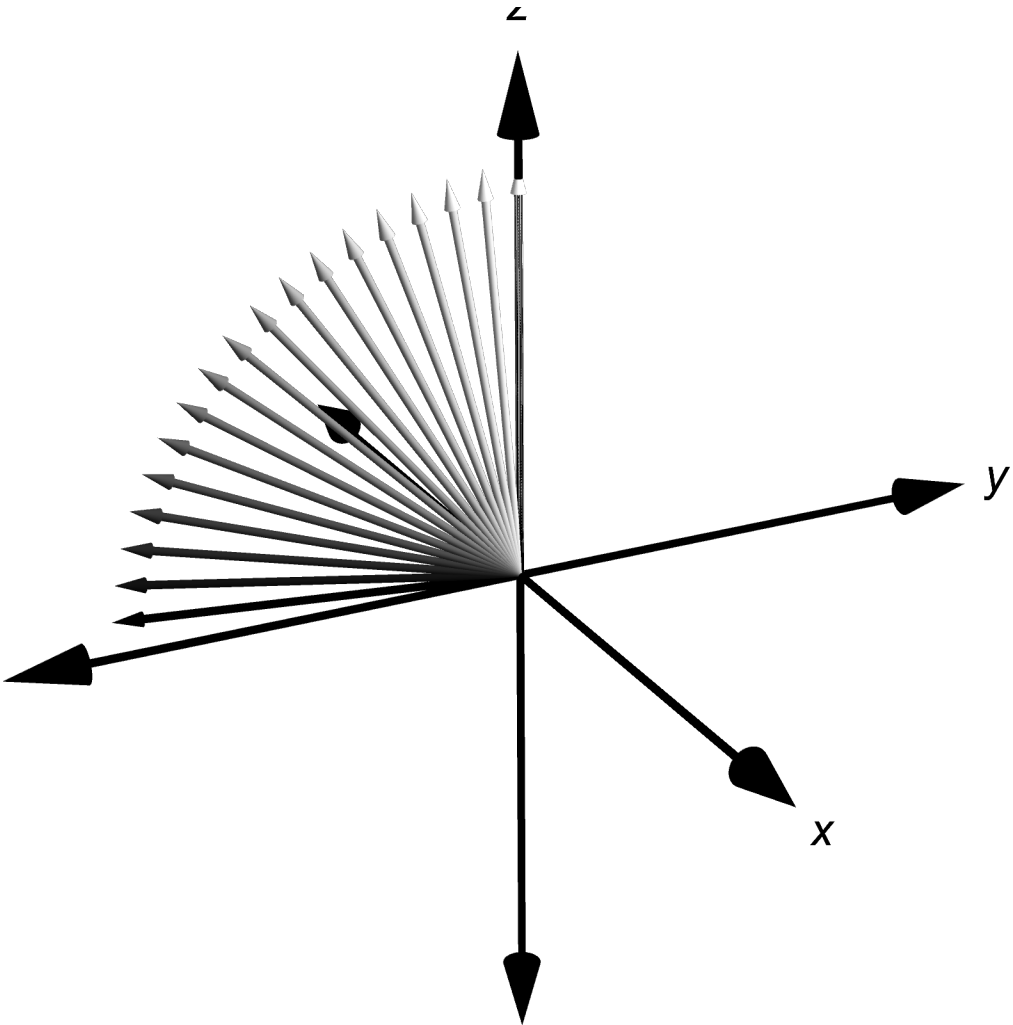
```
In[24]:= Show[my$drawing[ρ$vector[t, 2, 1, 0],  $\frac{2\pi}{\sqrt{5}}$ , 36], ImageSize → Full]
```

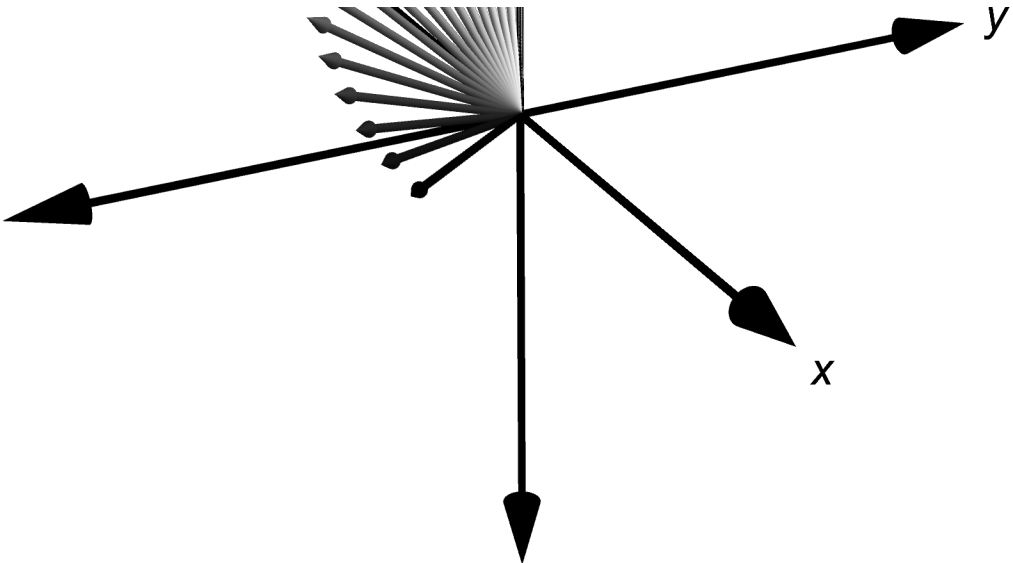
Out[24]=



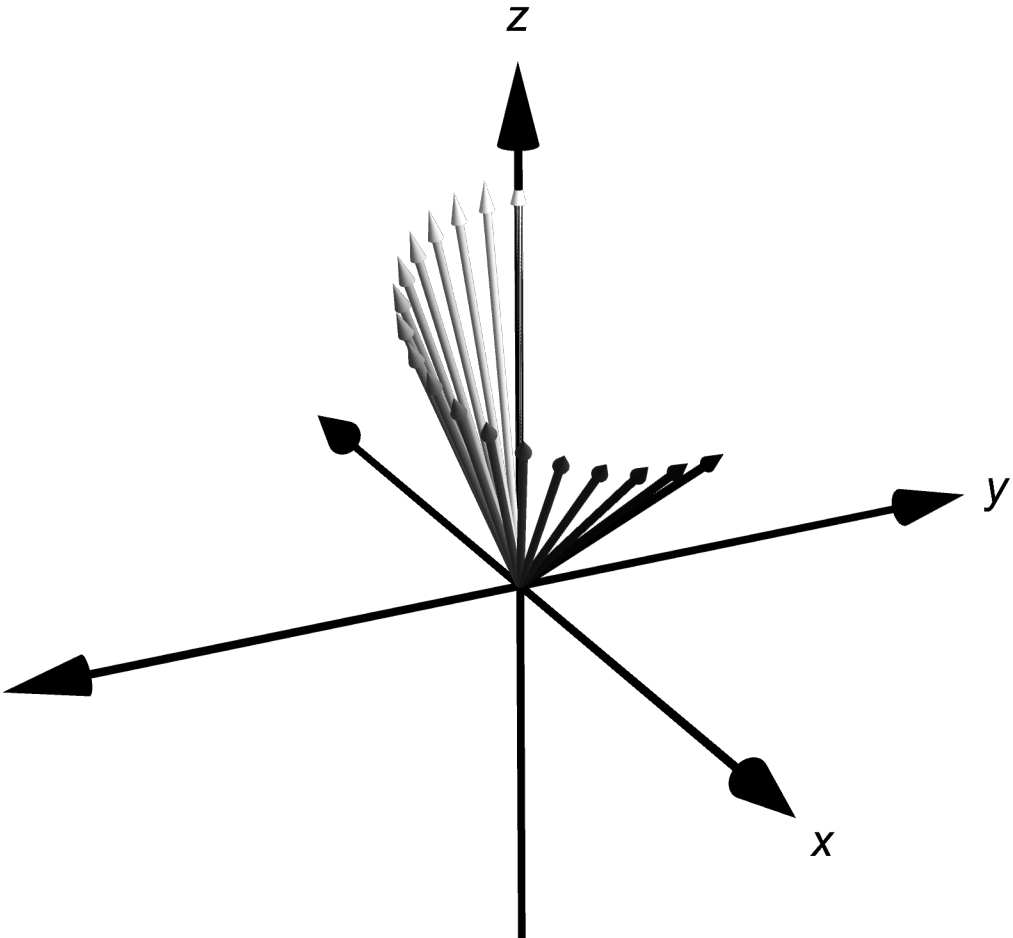
Set the rf phase to  $\phi = 0$ , set the Rabi frequency to  $\omega = 1$ , and look at magnetization out to times equal to  $\pi/2$ . Vary the resonance offset and plot the magnetization.

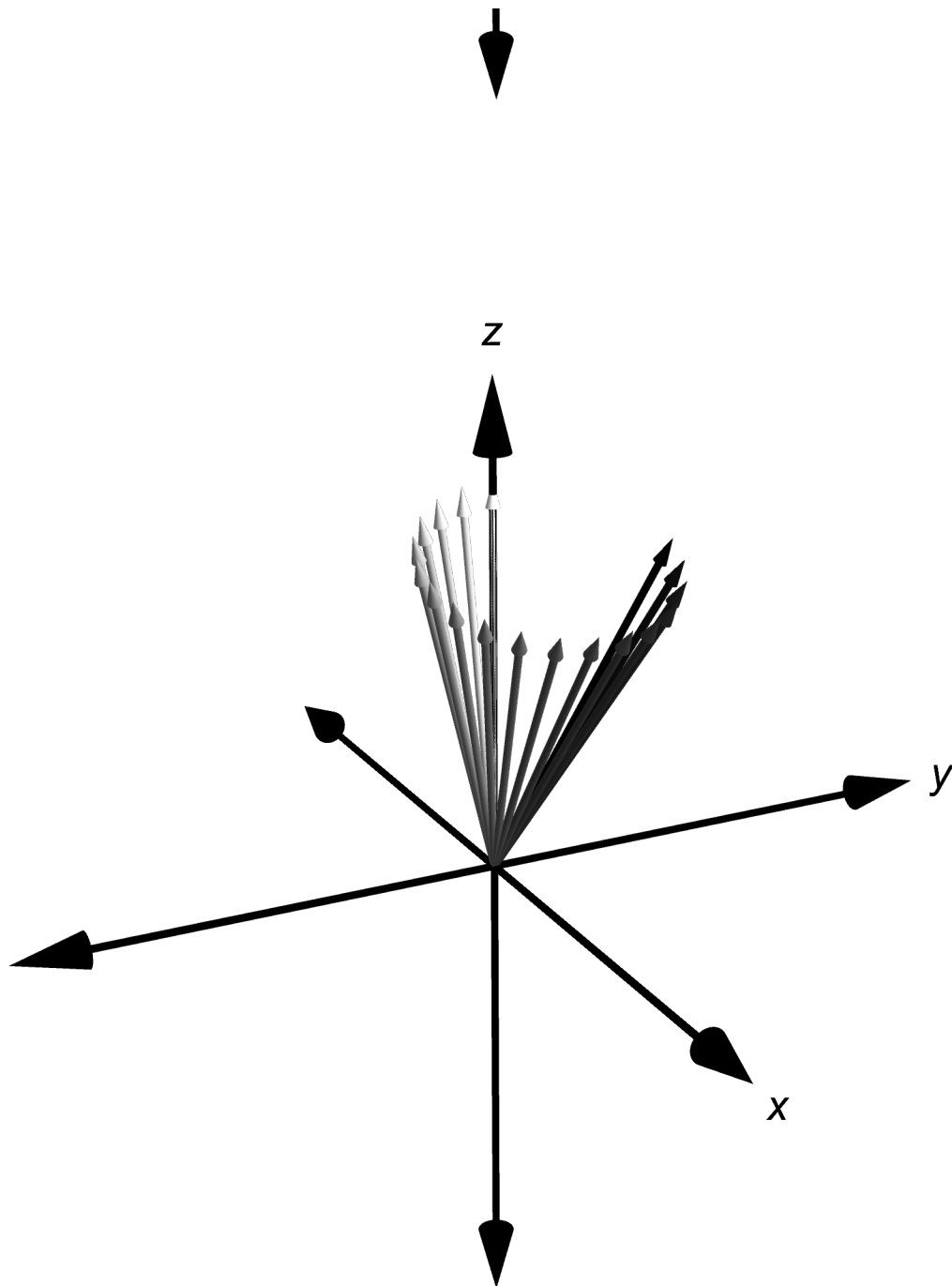
```
In[25]:= Show[GraphicsGrid @@  
  {{my$drawing[ρ$vector[t, #, 1, 0],  $\pi/2$ , 18]} &  
    /@ {0., 1., 2., 3.}}, ImageSize → Full]
```





Out[25]=



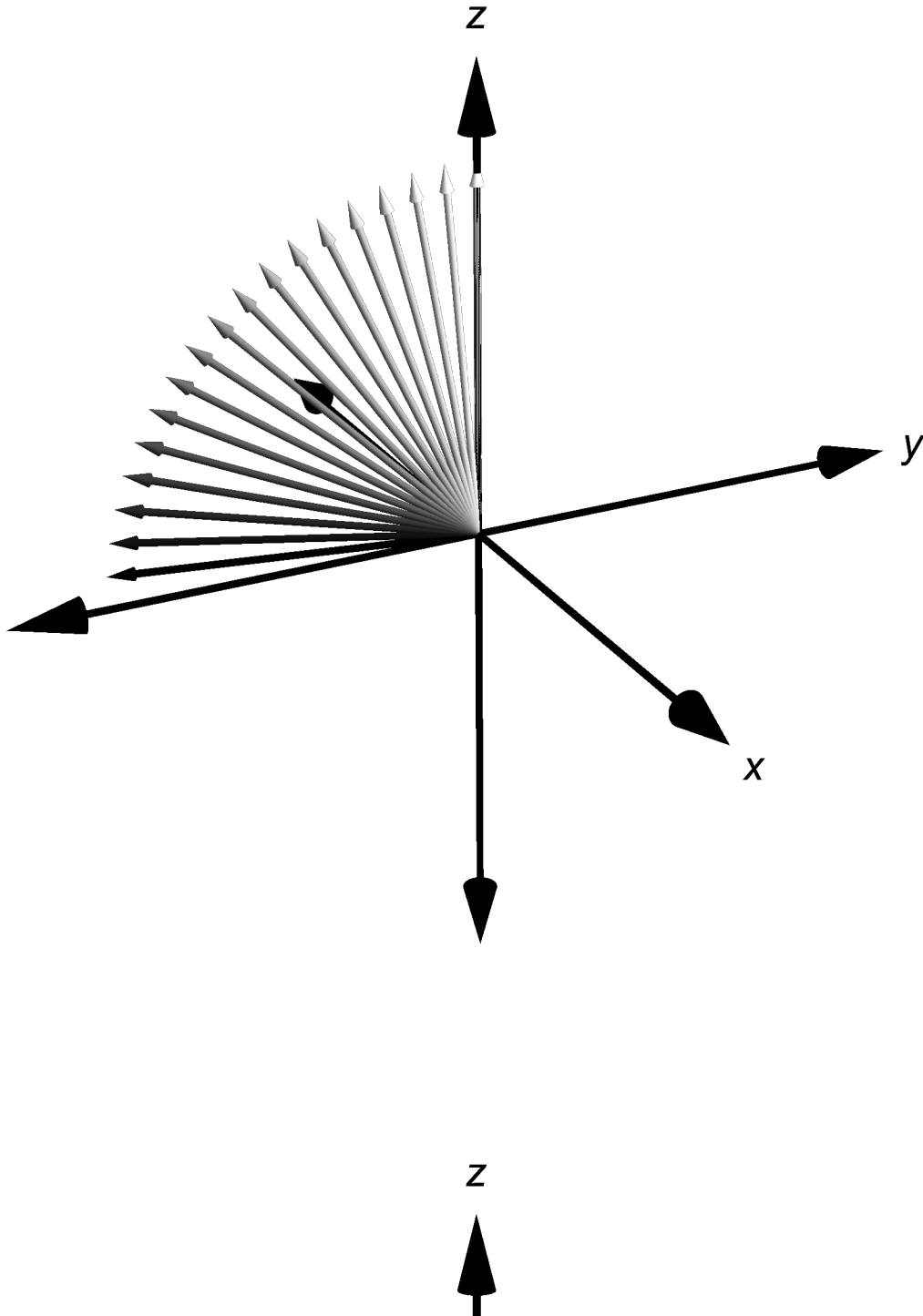


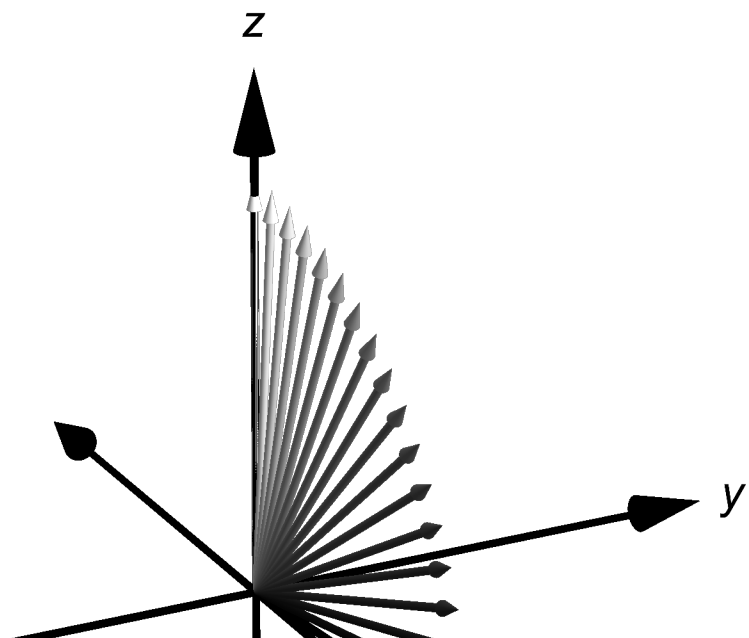
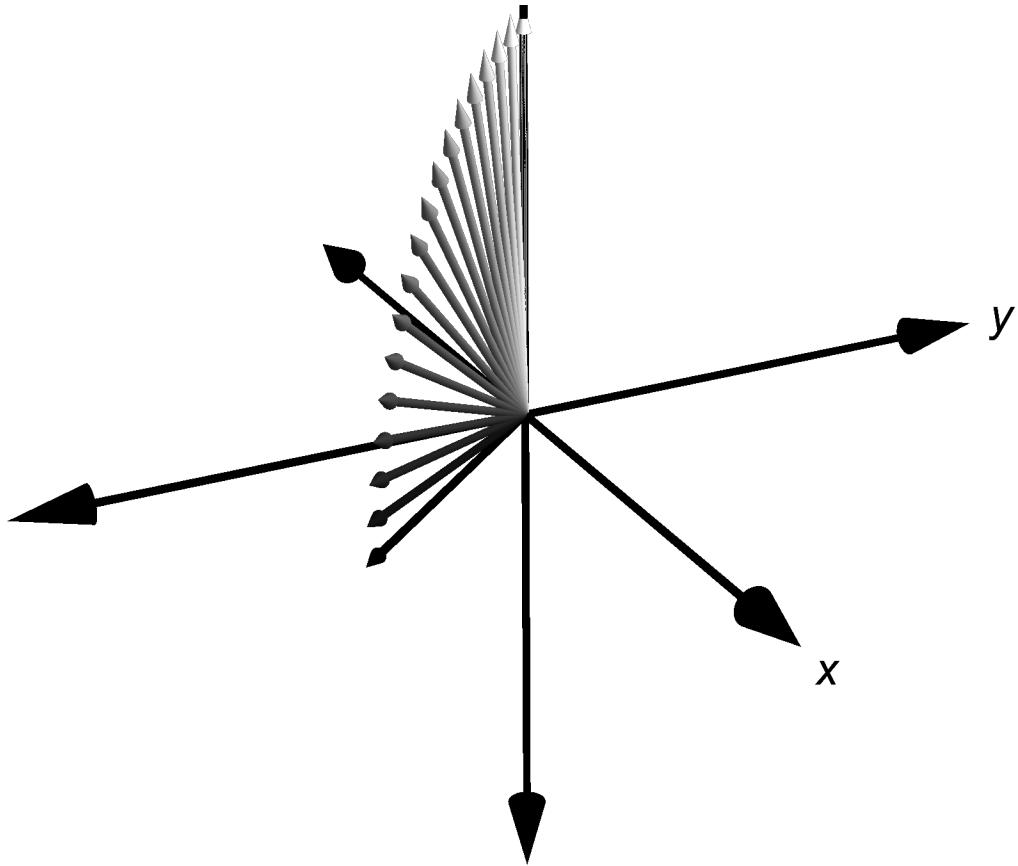
Set the resonance offset  $\Delta = 0$ , set the Rabi frequency to  $\omega = 1$ , and look at magnetization out to times equal to  $\pi/2$ . Vary the rf phase and plot the magnetization.

```

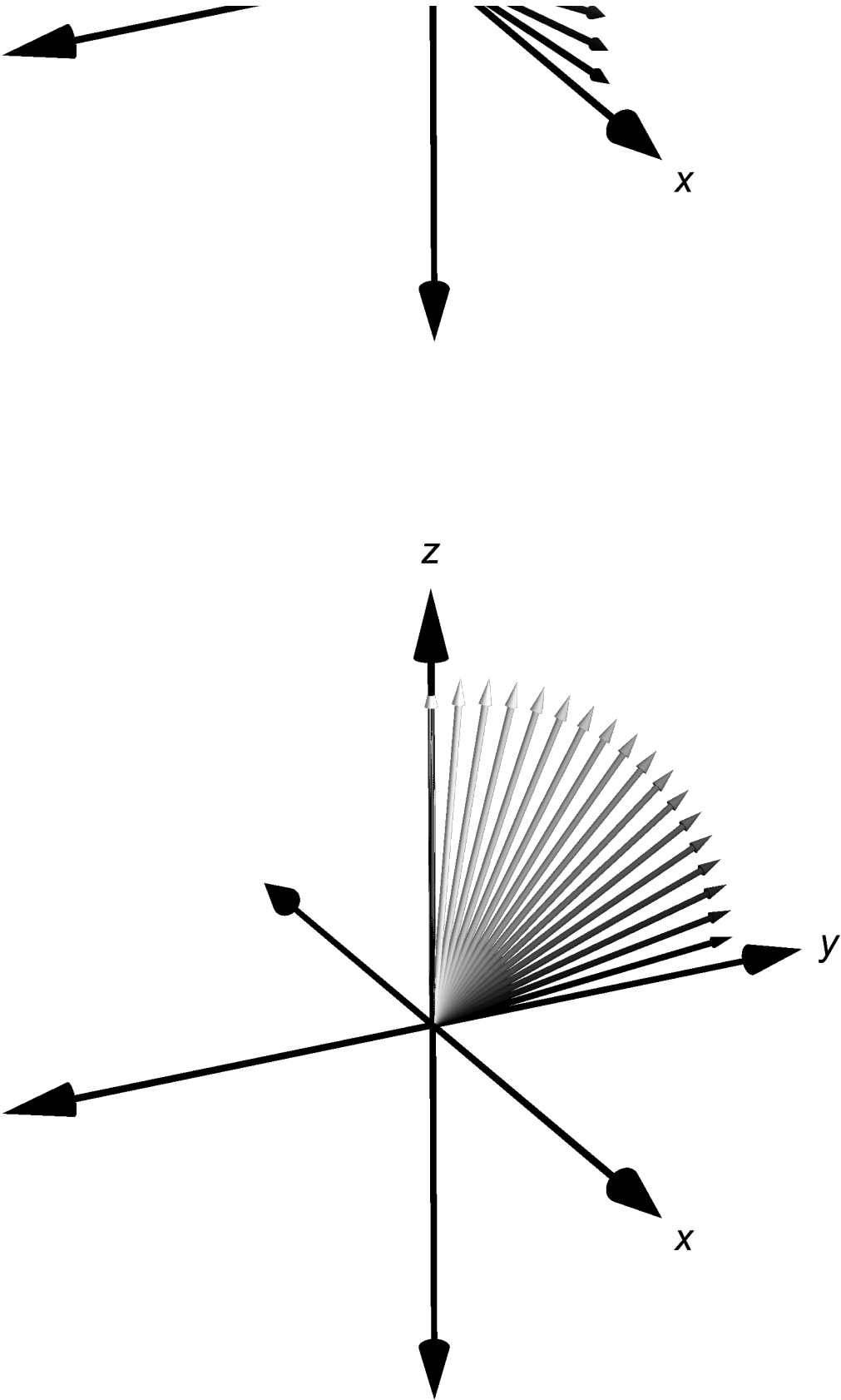
In[26]:= Show[GraphicsGrid @@
  { {my$drawing[ρ$vector[t, 0, 1, #], π/2, 18]} &
    /@ ({0., 45, 90, 180, 270}  $\frac{\pi}{180}$ ) }, ImageSize → Full]

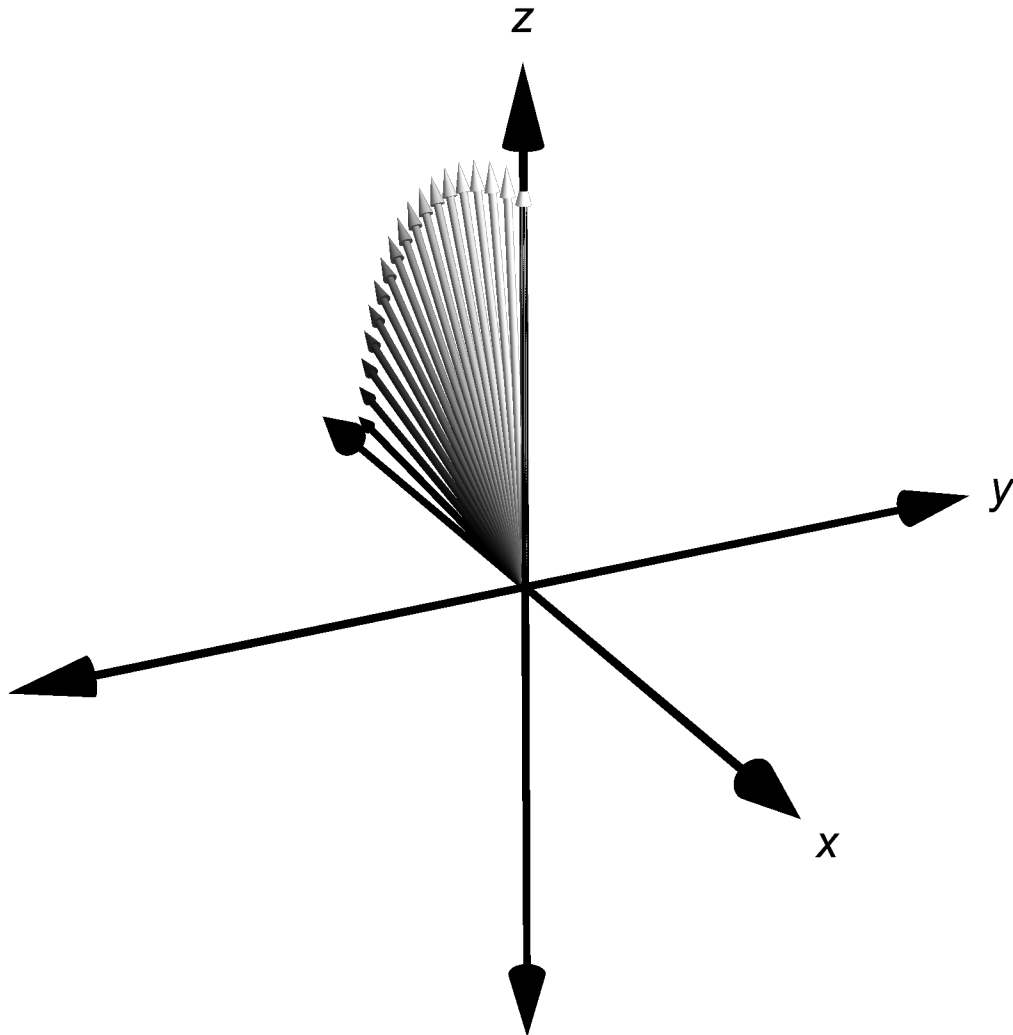
```





Out[26]=





## Clean up

```
In[27]:= (*
Clear[ $\omega$ ,  $\Delta$ ,  $\phi$ , t, Ix, Iy, Iz,  $\rho$ ,  $\rho_0$ , H]
*)
```