

# UniDyn--Demo-04.nb

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**Abstract:** Use the **UniDyn** Evolver function to calculate the evolution of the magnetization of a two coupled spin = 1/2 particles.

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## Set the path to the package

Tell *Mathematica* the path to the directory containing the package.

EDIT THE FOLLOWING PATH STRING:

```
In[1]:= $UniDynPath =
"/Users/jam99/Dropbox/MarohnGroup__Software_Library/UniDyn/
unidyn";
```

YOU SHOULD NOT NEED TO EDIT ANYTHING FROM HERE ONWARDS.

---

## Load the package

Append the package path to the system path. Before trying to load the package, ask *Mathematica* to find it. This is a test that we directed *Mathematica* to the correct directory. The output of this command should be the full system path to the UniDyn.m file.

```
In[2]:= $Path = AppendTo[$Path, $UniDynPath];
FindFile["UniDyn`"]
Out[3]= /Users/jam99/Dropbox/MarohnGroup__Software_Library/UniDyn/unidyn/UniDyn.m
```

Now that we are confident that the path is set correctly, load the package. Setting the global \$VerboseLoad variable to True will print out the help strings for key commands in the package.

```
In[4]:= $VerboseLoad = True;
```

```
Needs["UniDyn`"]
```

- **CreateOperator**: CreateOperator[] is used to batch–define a bunch of operators. Example: CreateOperator[{{Ix, Iy, Iz}, {Sx, Sy, Sz}}] will create six operators, where each of the operators in the first list will commute with each of the operators in the second list.
- **CreateScalar**: CreateScalar[list] is used to batch–define a bunch of scalars. The parameter list can be a single scalar or a list of scalars. Example: CreateScalar[{w1, w2}].
- **NCSort**: NCSort[list] sorts the operators in list into canonical order.
- **SortedMult**: SortedMult[list] returns Mult[list\$ordered], where list\$ordered are the elements of list sorted into canonical order.
- **MultSort**: MultSort[NonCommutativeMultiply[list]] returns NonCommutativeMultiply[list\$ordered], where list\$ordered are the elements of list sorted into canonical order.
- **Comm**: Comm[a,b] calculates the commutator of two operators.
- **Inv**: Inv[a] returns the inverse of the expression.
- **SpinSingle\$CreateOperators**: SpinSingle\$CreateOperators[Ix, ly, lz, L] creates Ix, ly, and lz angular momentum operators and defines their commutation relations. When the total angular momentum L = 1/2, additional rules are defined to simplify products of the angular momentum operators. When the total angular momentum L is unspecified, no such simplification rules are defined.
- **OscSingle\$CreateOperators**: OscSingle\$CreateOperators[aL, aR] creates a raising operator aR and a lowering operator aL for single harmonic oscillator and defines the operator commutation relations.
- **Evolve**: Evolve[H, t, ρ(0)] calculates  $\rho(t) = \text{Exp}[-iHt]\rho(0)\text{Exp}[+iHt]$ , assuming that H is time independent, according to the commutation rules followed by ρ(0) and H.
- **AllCommutingQ**: A test to see if all the terms in a sum commute with each other.
- **VisualComplexity**: A cost function to coax Mathematica into writing simpler-looking answers.
- **Evolver1**: Evolver1[H, t, ρ(0)] calculates  $\rho(t) = \text{Exp}[-iHt]\rho(0)\text{Exp}[+iHt]$ , assuming that H is time independent, according to the commutation rules followed by ρ(0) and H.
- **Evolver2**: Evolver2[H, t, ρ(0)] calculates  $\rho(t) = \text{Exp}[-iHt]\rho(0)\text{Exp}[+iHt]$ , assuming that H is time independent, according to the commutation rules followed by ρ(0) and H.
- **SpinBoson\$CreateOperators**: SpinBoson\$CreateOperators[Ix, ly, lz, lp, lm, aR, aL] creates Ix, ly, lz spin one half angular–momentum operators; the associated spin raising and lowering operators lp, lm; and harmonic–oscillator raising and lowering operators aR, aL.

## Evolver with simplification

```
In[6]:= SimplifyingEvolver[H_, t_, ρ$0_] :=
  Evolver2[H, t, ρ$0] // Simplify // ExpToTrig // FullSimplify
```

# Unitary evolution of a single spin 1/2

## Create a single spin

The assumptions define below are required for *Mathematica* to recognize  $\sqrt{-\Delta^2 - \omega^2} = i\sqrt{\Delta^2 + \omega^2}$  inside an exponential. One of the variables has to be defined to be  $> 0$  and not just  $\geq 0$ .

```
In[7]:= Clear[
  Δω, (* resonance offset frequency *)
  ω$1, (* Rabi frequency of the applied irradiation *)
  t, (* time *)
  Ix, Iy, Iz, (* spin angular momentum operators *)
  ρ$0, (* initial spin density operator *)
  H (* spin Hamiltonian *)]

CreateScalar[Δω, ω$1, t];
SpinSingle$CreateOperators[Ix, Iy, Iz, L = 1/2];

$Assumptions = {Element[Δω, Reals], Δω ≥ 0,
  Element[ω$1, Reals], ω$1 > 0, Element[t, Reals], t ≥ 0};

*** SpinSingle$CreateOperators: Creating spin operators.
*** SpinSingle$CreateOperators: Adding spin commutations relations.
*** SpinSingle$CreateOperators: Angular momentum L = 1/2. Adding operator simplification rules.
```

## On-resonance nutation

On-resonance irradiation Hamiltonian written in the interaction representation. The initial density operator is parallel to  $I_x$ .

```
In[11]:= H = ω$1 Ix;
ρ$0 = Iz;
```

Calculate the time-dependent density operator.

```
In[13]:= SimplifyingEvolver[H, t, ρ$0] /. {ω$1 → Subscript[ω, 1]}
Out[13]= Iz Cos[t ω1] - Iy Sin[t ω1]
```

## Free evolution

Zeeman-interaction Hamiltonian written in the interaction representation. The initial density operator is parallel to  $I_x$ .

```
In[14]:= H = Δω Iz;
```

```
ρ$0 = Ix;
```

Calculate the time-dependent density operator.

```
In[16]:= SimplifyingEvolver[H, t, ρ$0]
```

```
Out[16]= Ix Cos[t Δω] + Iy Sin[t Δω]
```

## Unitary evolution of two coupled spins

### Create a two spins

The assumptions define below are required for *Mathematica* to recognize

$\sqrt{-Δ^2 - ω^2} = i \sqrt{Δ^2 + ω^2}$  inside an exponential. One of the variables has to be defined to be  $> 0$  and not just  $\geq 0$ .

```
In[17]:= Clear[
Δ$I, (* resonance offset frequency *)
Δ$S, (* resonance offset frequency *)
J, (* spin-spin coupling *)
Ix, Iy, Iz, (* spin angular momentum operators *)
Sx, Sy, Sz, (* spin angular momentum operators *)
ρ, (* spin density operator *)
t, (* time *)
ρ$0, (* initial spin density operator *)
H (* spin Hamiltonian *)]

CreateScalar[Δ$I, Δ$S, J, t];
CreateOperator[{Ix, Iy, Iz}, {Sx, Sy, Sz}];
SpinSingle$CreateOperators[Ix, Iy, Iz, L = 1/2];
SpinSingle$CreateOperators[Sx, Sy, Sz, L = 1/2];

$Assumptions = {Element[Δ$I, Reals], Δ$I ≥ 0,
Element[Δ$S, Reals], Δ$S ≥ 0, Element[J, Reals], J > 0};

... SpinSingle$CreateOperators: Spin operators already exist.
... SpinSingle$CreateOperators: Adding spin commutations relations.
... SpinSingle$CreateOperators: Angular momentum L = 1/2. Adding operator simplification rules.
... SpinSingle$CreateOperators: Spin operators already exist.
... SpinSingle$CreateOperators: Adding spin commutations relations.
... SpinSingle$CreateOperators: Angular momentum L = 1/2. Adding operator simplification rules.
```

## Evolution under J coupling

On-resonance irradiation Hamiltonian written in the interaction representation. The initial density operator is parallel to  $I_x$ .

```
In[23]:= H$0 = Δ$I Iz + Δ$S Sz;
H$J = J Mult[Iz, Sz];
ρ$0 = Ix;
```

Try to calculate the density operator in a single step.

In[26]:= **SimplifyingEvolver**[H\$0 + H\$J, t, ρ\$0]

$$\begin{aligned} \text{Out}[26]= & \text{Mult}\left[\cosh\left[t \sqrt{\text{Inv}[Ix], \text{Inv}\left[\text{Inv}\left[-\frac{1}{4} Ix (J^2 + 4 \Delta\$I^2) - 2 J \Delta\$I \text{Mult}[Ix, Sz]\right]\right]\right]}, Ix\right] + \\ & \Delta\$I \text{Mult}\left[\sinh\left[t \sqrt{\text{Inv}[Ix], \text{Inv}\left[\text{Inv}\left[-\frac{1}{4} Ix (J^2 + 4 \Delta\$I^2) - 2 J \Delta\$I \text{Mult}[Ix, Sz]\right]\right]\right]}, \right. \\ & \left. Iy, \frac{1}{\sqrt{\text{Mult}\left[\text{Inv}[Ix], \text{Inv}\left[\text{Inv}\left[-\frac{1}{4} Ix (J^2 + 4 \Delta\$I^2) - 2 J \Delta\$I \text{Mult}[Ix, Sz]\right]\right]\right]}}\right] + \\ & J \text{Mult}\left[\sinh\left[t \sqrt{\text{Inv}[Ix], \text{Inv}\left[\text{Inv}\left[-\frac{1}{4} Ix (J^2 + 4 \Delta\$I^2) - 2 J \Delta\$I \text{Mult}[Ix, Sz]\right]\right]\right]}, \right. \\ & \left. Iy, Sz, \frac{1}{\sqrt{\text{Mult}\left[\text{Inv}[Ix], \text{Inv}\left[\text{Inv}\left[-\frac{1}{4} Ix (J^2 + 4 \Delta\$I^2) - 2 J \Delta\$I \text{Mult}[Ix, Sz]\right]\right]\right]}}\right] \end{aligned}$$

This single-step evolution fails. Instead, let us calculate the density operator in two steps. We can do this because the H\$J and H\$0 Hamiltonians commute. Evolve under the J-coupling first and under the chemical shift second.

In[27]:=  $\rho = \rho\$0 // \text{SimplifyingEvolver}[H\$J, t, \#] \& //$

**SimplifyingEvolver**[H\$0, t, \#] &

$$\begin{aligned} \text{Out}[27]= & \cos\left[\frac{Jt}{2}\right] (Ix \cos[t \Delta\$I] + Iy \sin[t \Delta\$I]) + \\ & 2 \sin\left[\frac{Jt}{2}\right] (\cos[t \Delta\$I] \text{Mult}[Iy, Sz] - \text{Mult}[Ix, Sz] \sin[t \Delta\$I]) \end{aligned}$$

Try another way -- evolve under the chemical shift first and under the J-coupling second.

In[28]:=  $\rho = \rho\$0 // \text{SimplifyingEvolver}[H\$0, t, \#] \& //$

**SimplifyingEvolver**[H\$J, t, \#] &

$$\begin{aligned} \text{Out}[28]= & \cos\left[\frac{Jt}{2}\right] (Ix \cos[t \Delta\$I] + Iy \sin[t \Delta\$I]) + \\ & 2 \sin\left[\frac{Jt}{2}\right] (\cos[t \Delta\$I] \text{Mult}[Iy, Sz] - \text{Mult}[Ix, Sz] \sin[t \Delta\$I]) \end{aligned}$$

We see by inspection that we get the same answer either way.

### Load spin 1/2 operator matrices

Load matrices representing two J=1/2 spins and give them a test drive.

In[29]:= << Matrices--two-spin-half.m

Loaded spin one-half matrices for mIx, mIy, mIz, mSx, mSy, mSz

Let's look at one of the matrices, the matrix for Iz.

In[30]:= mIz // MatrixForm  
mSz // MatrixForm

Out[30]//MatrixForm=

$$\begin{pmatrix} -\frac{1}{2} & 0 & 0 & 0 \\ 0 & -\frac{1}{2} & 0 & 0 \\ 0 & 0 & \frac{1}{2} & 0 \\ 0 & 0 & 0 & \frac{1}{2} \end{pmatrix}$$

Out[31]//MatrixForm=

$$\begin{pmatrix} -\frac{1}{2} & 0 & 0 & 0 \\ 0 & \frac{1}{2} & 0 & 0 \\ 0 & 0 & -\frac{1}{2} & 0 \\ 0 & 0 & 0 & \frac{1}{2} \end{pmatrix}$$

For a single spin 1/2 particle,  $I_z$  is a  $2 \times 2$  matrix. In the product space of two spin 1/2 particles, however, the  $I_z$  is now a  $4 \times 4$  matrix. All the operators are  $4 \times 4$  matrices. Check that the  $[I_x, I_y]/I = I_z$  and  $[S_x, S_y]/I = S_z$  commutation relations hold with the matrices.

In[32]:= m\$1 =  $\frac{mIx \cdot mIy - mIy \cdot mIx}{I}$  ;  
m\$2 =  $\frac{mSx \cdot mSy - mSy \cdot mSx}{I}$  ;

In[34]:= m\$1 // MatrixForm (\* should be Iz \*)  
m\$2 // MatrixForm (\* should be Sz \*)

Out[34]//MatrixForm=

$$\begin{pmatrix} -\frac{1}{2} & 0 & 0 & 0 \\ 0 & -\frac{1}{2} & 0 & 0 \\ 0 & 0 & \frac{1}{2} & 0 \\ 0 & 0 & 0 & \frac{1}{2} \end{pmatrix}$$

Out[35]//MatrixForm=

$$\begin{pmatrix} -\frac{1}{2} & 0 & 0 & 0 \\ 0 & \frac{1}{2} & 0 & 0 \\ 0 & 0 & -\frac{1}{2} & 0 \\ 0 & 0 & 0 & \frac{1}{2} \end{pmatrix}$$

Make the (matrix) operator  $I^{\text{total}} = I_z + S_z$ . We can see by inspection that this matrix looks correct. The  $I^{\text{total}}$  operator should have eigenvalues  $\{-1, 0, 0, +1\}$ . The matrix  $I^{\text{total}}$  is diagonal and we can read the eigenvalues off diagonal elements.

```
In[36]:= m$1 + m$2 // MatrixForm
```

Out[36]//MatrixForm=

$$\begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

## Calculate the analytical signal using the matrices

Write the matrix representation of  $I_y S_z$  using the matrices loaded above.

```
In[37]:= mIy . mSz // MatrixForm
```

Out[37]//MatrixForm=

$$\begin{pmatrix} 0 & 0 & -\frac{i}{4} & 0 \\ 0 & 0 & 0 & \frac{i}{4} \\ \frac{i}{4} & 0 & 0 & 0 \\ 0 & -\frac{i}{4} & 0 & 0 \end{pmatrix}$$

Show that we can obtain the same matrix from the symbolic product  $I_y^{**}S_z$  by (1) substituting the spin operators with matrices and (2) replacing the NonCommutative-Multiply operator with the Dot (i.e., matrix multiplication) operator.

```
In[38]:= Mult[Iy, Sz] // Mult → Dot /.
          {Ix → mIx, Iy → mIy, Iz → mIz, Sx → mSx, Sy → mSy, Sz → mSz} // MatrixForm
```

Out[38]//MatrixForm=

$$\begin{pmatrix} 0 & 0 & -\frac{i}{4} & 0 \\ 0 & 0 & 0 & \frac{i}{4} \\ \frac{i}{4} & 0 & 0 & 0 \\ 0 & -\frac{i}{4} & 0 & 0 \end{pmatrix}$$

```
In[39]:= ρ // Expand
```

```
Out[39]= Ix Cos[ $\frac{Jt}{2}$ ] Cos[t Δ$I] + 2 Cos[t Δ$I] Mult[Iy, Sz] Sin[ $\frac{Jt}{2}$ ] +
          Iy Cos[ $\frac{Jt}{2}$ ] Sin[t Δ$I] - 2 Mult[Ix, Sz] Sin[ $\frac{Jt}{2}$ ] Sin[t Δ$I]
```

Calculate the density operator matrix.

```
In[40]:= Clear[\rho$matrix, \rho$temp, t];
\rho$temp = Expand[\rho] /. Mult \rightarrow Dot;
\rho$matrix[t_] = \rho$temp /.
{Ix \rightarrow mIx, Iy \rightarrow mIy, Iz \rightarrow mIz, Sx \rightarrow mSx, Sy \rightarrow mSy, Sz \rightarrow mSz};
```

Check that the resulting object is a 4 x 4 matrix as expected.

```
In[43]:= Dimensions[\rho$matrix[t]]
```

```
Out[43]= {4, 4}
```

From this matrix we can calculate the I-spin signals as  $\text{Trace}[\rho \text{Ix}]$  and  $\text{Trace}[\rho \text{Iy}]$

```
In[44]:= Tr[\rho$matrix[t] . mIx] // Simplify
Tr[\rho$matrix[t] . mIy] // Simplify
```

```
Out[44]= Cos[\frac{Jt}{2}] \cos[t \Delta\$I]
```

```
Out[45]= Cos[\frac{Jt}{2}] \sin[t \Delta\$I]
```

We can mimic the complex signal collected by the NMR spectrometer by calculating the expectation value of the  $\text{Ix} + \text{Iy}$  operator:  $\text{Trace}[\rho (\text{Ix} + \text{Iy})]$ .

```
In[46]:= Clear[S];
S[t_] := Tr[\rho$matrix[t] . (mIx - I mIy)] // TrigToExp
```

### Plot the signal as a function of time

Create a more realistic experimental signal by multiplying the above-calculated signal by a decaying exponential.

```
In[48]:= Clear[S$expt];
S$expt[t_] := S[t] Exp[-t / T2]
```

To plot, set the total number of points (NN) and the total time (T). We set NN equal to a power of 2 in anticipation of taking a digital Fourier transform.

```
In[50]:= NN = 2^10;
T$final = 10.0;
```

From NN and T\$final we derive the time step (dt) and the frequency step (df). Now generate a list of data points based on the above function. At the same time, generate a list of time points (t) and frequency points (f) for plotting.

```
In[52]:= dt = T$final / (NN - 1);
df = 1 / dt;

f$table = Table[jj / T$final, {jj, -NN / 2, NN / 2 - 1}];
t$table = Table[ii * dt, {ii, 0, NN - 1}];

Print["The time step is ", dt, " s"]
Print["The Nyquist frequency is ", 1 / (2 * dt), " Hz"]

The time step is 0.00977517 s
The Nyquist frequency is 51.15 Hz

Give numbers for the chemical shift and the J-coupling.

In[58]:= Δ$I$value = 2 π × 10.0; (* chemical shift [rad/s] *)
J$value = 2 π × 8.0;
(* heteronuclear scalar coupling constant [rad/s] *)
T2$value = 1.0;          (* spin dephasing time [s] *)
```

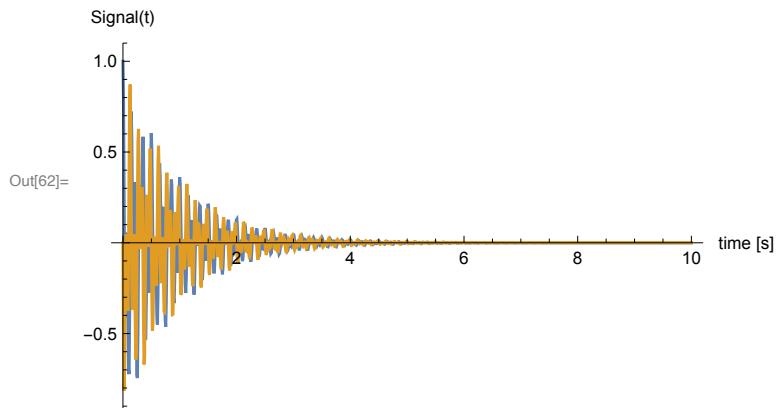
Create an array of complex signal.

```
In[61]:= S$table = S$expt[t] /.
{t → t$table, Δ$I → Δ$I$value, J → J$value, T2 → T2$value};
```

Plot the real and imaginary part of the signal

```
In[62]:= ListLinePlot[
{
  {t$table, Re[S$table]} // Transpose,
  {t$table, Im[S$table]} // Transpose},
 Joined → True,
 PlotRange → All,
 PlotLabel → "Free induction decay of 1H coupled to 13C\n",
 AxesLabel → {"time [s]", "Signal(t)"}]
```

Free induction decay of 1H coupled to 13C



A function to calculate the digital Fourier transform of signal.

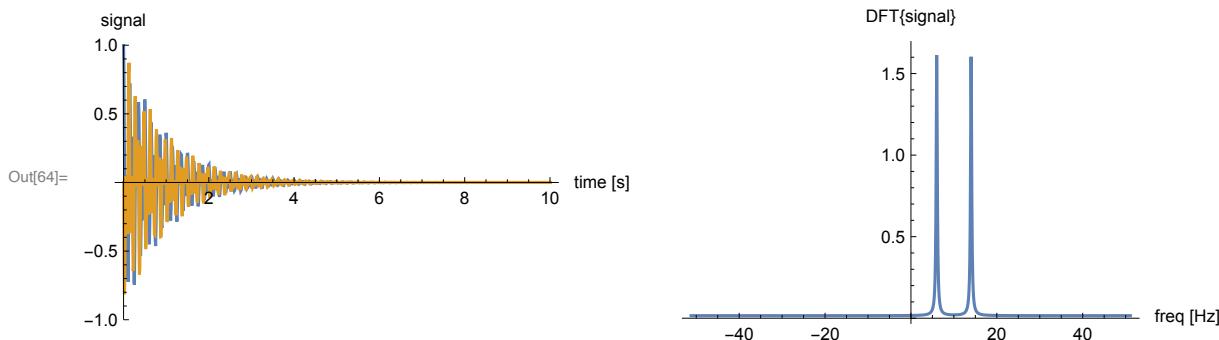
```

In[63]:= DFFT[signal$table_, time$table_, query_ : True] :=
  (* True to plot both real and imaginary FT parts *)
  Module[{NN, T$final},
    NN = Dimensions[signal$table][[1]];
    FFTS$table = RotateRight[Fourier[signal$table], NN / 2];
    T$final = time$table[[-1]];
    f$table = Table[jj / T$final, {jj, -NN / 2, NN / 2 - 1}];
    p1 = ListLinePlot[
      {Transpose[{time$table, Re[signal$table]}],
       Transpose[{time$table, Im[signal$table]}]},
      PlotRange → {-1, 1},
      AxesLabel → {"time [s]", "signal"}];
    If[query == True,
      p2 = ListLinePlot[
        {Transpose[{f$table, Re[FFTS$table]}],
         Transpose[{f$table, Im[FFTS$table]}]},
        PlotRange → All,
        AxesLabel → {"freq [Hz]", "DFT{signal}"}],
      p2 = ListLinePlot[
        Transpose[{f$table, Re[FFTS$table]}],
        PlotRange → All,
        AxesLabel → {"freq [Hz]", "DFT{signal}"}]
    ];
    Show[GraphicsGrid[{{p1, p2}}]]
  ]

```

Fourier transform the signal to obtain the spectrum.

In[64]:= **DFFT[S\$table, t\$table, False]**



Find peaks in the FT spectrum. Require the peaks to be larger than 1.0. See the FindPeaks function documentation here.

In[65]:= **peaks = FindPeaks[Re[FFTS\$table], 0, 0, 1.]**

Out[65]= { {573, 1.61249}, {653, 1.60297} }

Read out the frequencies at which the peaks are located. We see peaks at the expected frequencies of  $10-8/2 = 6$  Hz and  $10+8/2 = 14$  Hz.

In[66]:= **Part[f\$table, Transpose[peaks][1]]**

Out[66]= {6., 14.}

## Clean up