

Thermal phase transition of 2D classical Ising model

PHYS3999 Directed Studies in Physics

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- 3 Monte Carlo simulations
- 4 Physical observables and critical phenomena
- 5 Finite size scaling
- 6 References

Main objectives

- Investigate phase transition of 2D classical Ising model using Monte Carlo method

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 - Find the critical temperature T_c of 2D classical Ising model
- Explore finite scaling effect of physical observables
 - Find the critical exponents $\eta, \alpha, \beta, \gamma, \nu$ of 2D classical Ising model

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\rightarrow Solve complex models with large degrees of freedom

2D Ising Model

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Heisenberg model without external field

$$H = - \sum_{\langle i,j \rangle} J_{ij} \sigma_i \sigma_j \quad (1)$$

$\langle i, j \rangle$: nearest neighbors

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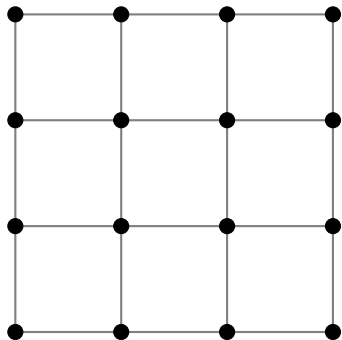
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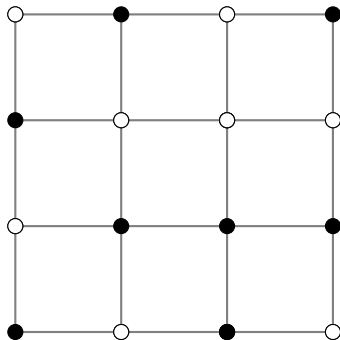
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Onsager (1944) exact solutions: $T_c \approx 2.269$

2D Ising Model



(a) Ordered phase $T < T_c$



(b) Disordered phase $T > T_c$

Figure: Ising model with system size $L = 4$

Solid node: up spin $+1$

Open node: down spin -1

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 - Swendsen-Wang algorithm - Cluster update (Swendsen & Wang, 1987)

Comparison between two algorithms

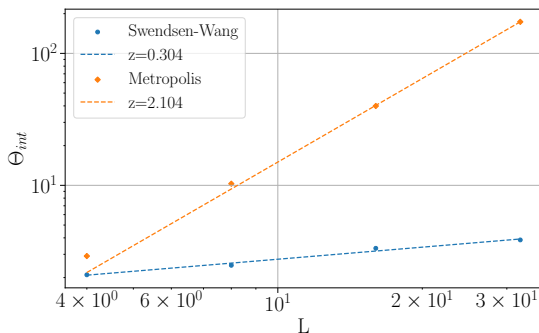
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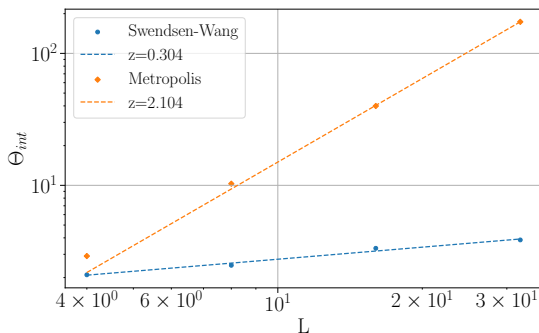
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- Metropolis algorithm takes longer steps to generate statistically independent configurations than Swendsen-Wang algorithm.

Physical observables and critical phenomena

Useful variables:

Reduced temperature:

$$t = \frac{|T - T_c|}{T_c} \quad (3)$$

Correlation length:

$$\xi \sim t^{-\nu} \quad (4)$$

ν : one of the critical exponents

Physical observables: spin-spin correlation

Spin-spin correlation:

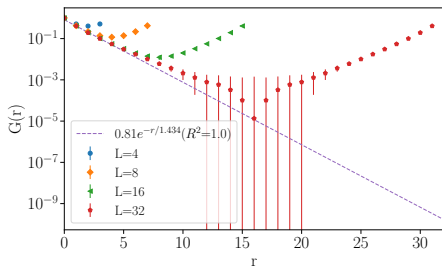
$$G(\vec{r}) \sim \begin{cases} \frac{e^{-r/\xi}}{r^{(d-2)/2}} & \text{when } T \neq T_c \\ \frac{1}{r^{d-2+\eta}} & \text{when } T = T_c, \xi \rightarrow \infty \end{cases} \quad (5)$$

d : dimension of the system

η : one of the critical exponents

Physical observables: spin-spin correlation

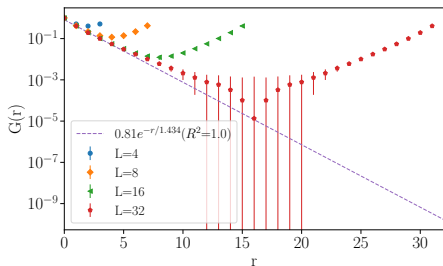
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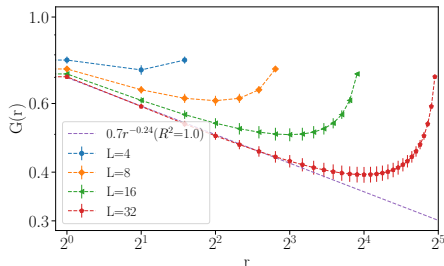
(a) Semi-log plot of $G(r)$ at $T=3$

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(a) Semi-log plot of $G(r)$ at $T=3$



(b) Log-log plot of $G(r)$ at $T=2.285$

Physical observables in infinite lattice

Net Magnetization:

$$\langle |m| \rangle \sim \begin{cases} (T_c - T)^\beta & \text{when } T < T_c \\ 0 & \text{when } T > T_c \end{cases} \quad (6)$$

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Specific Heat Capacity:

$$C \sim t^{-\alpha} \quad (8)$$

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Scaling hypotheses (Kadanoff, 1966):

$$|m|(T_c - T, L) = L^{-\beta/\nu} f_{|m|} \left((T_c - T) L^{1/\nu} \right) \quad (9)$$

$$\chi(t, L) = L^{\gamma/\nu} f_{\chi} \left(t L^{1/\nu} \right) \quad (10)$$

$$C(t, L) = \ln(L) f_C \left(t L^{1/\nu} \right) \quad (11)$$

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In equation (11), $\ln(L)$ is used because $\alpha = 0$

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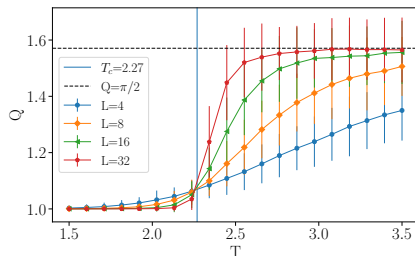
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Data collapse

Procedure of obtaining T_c and critical exponents with $|m|$ as an example:

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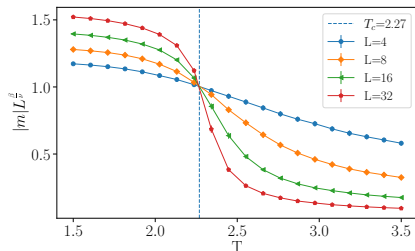
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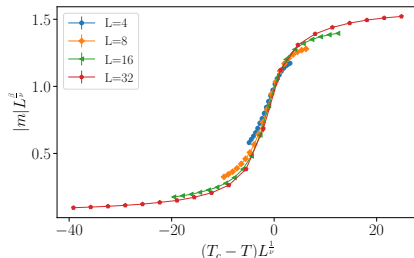
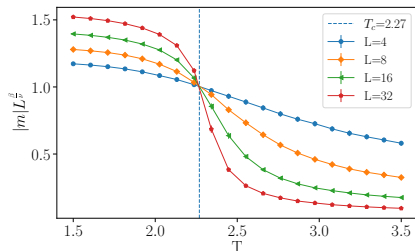
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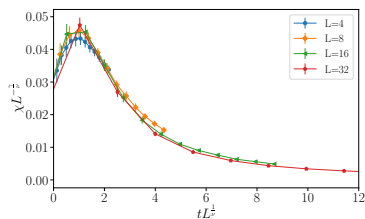
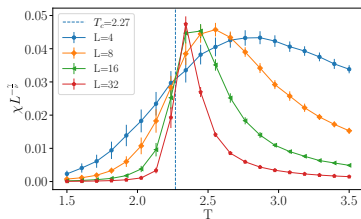
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- 2 Plot scaling function $f_m = L^{\beta/\nu}|m|$ against T to find out β/ν .
- 3 Plot scaling function f_m against its argument $(T_c - T)L^{1/\nu}$ to find out $1/\nu$ and hence β



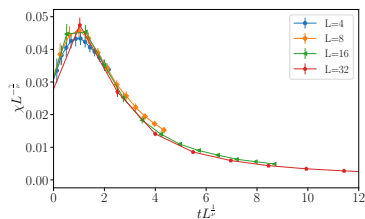
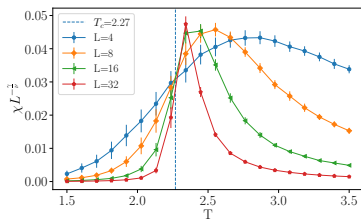
Finite size scaling

Suceptibility χ :

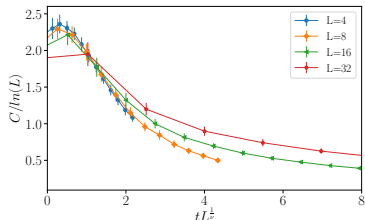
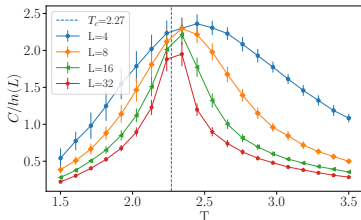


Finite size scaling

Suceptibility χ :



Specific heat capacity C :



Results

- Spin-spin correlation

- $\eta \approx -0.243$

- Binder ratio

- $T_c = 2.269$

- Data collapse

- $\alpha = 0$

- $\beta = 1/8$

- $\nu = 1$

- $\gamma = 7/4$

All results above are in reasonable accord with the exact solutions.

Conclusion

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- 4 Obtained critical exponents using data collapse

References

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