

# Lightlike form factors and OPE

Gang Yang

ITP, CAS

Based on the work with Yuanhong Guo, Lei Wang  
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Outline:

I. Introduction

II. Bootstrap  $\sim$  2-loop  $F_4^{(2)}$

III. FFOPE  $\sim$  Non-perturbative

IV. Summary & Outlook.

# I. Introduction and some history.

What are form factors?

EM vertex:

$$j^\mu = \bar{\psi} \gamma^\mu \psi$$

$$=$$

$$+ \dots$$

$$= \underbrace{\gamma^\mu F_1(q^2)}_{\substack{q=p-p' \\ \downarrow}} + i \frac{\sigma^{\mu\nu} q_\nu}{2m} \underbrace{F_2(q^2)}_{\substack{\downarrow \\ \text{Form factor.}}} + \dots$$

$$g-2 = 2F_2(q^2=0)$$

$$= \frac{\alpha}{\pi} + O(\alpha^2)$$

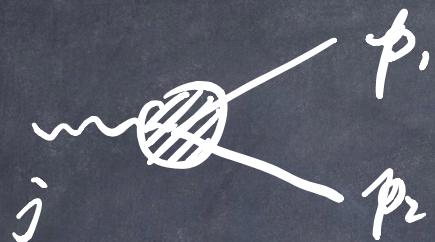


$$\sigma \sim \text{depend on } \{F_1, F_2\}$$

Nuclear form factor:  $F_1, F_2$

↳ characterize shape of particle  
deviation from point particle.

# Sudakov form factor.



- Sudakov (1954) sum over leading log.  
(off-shell)

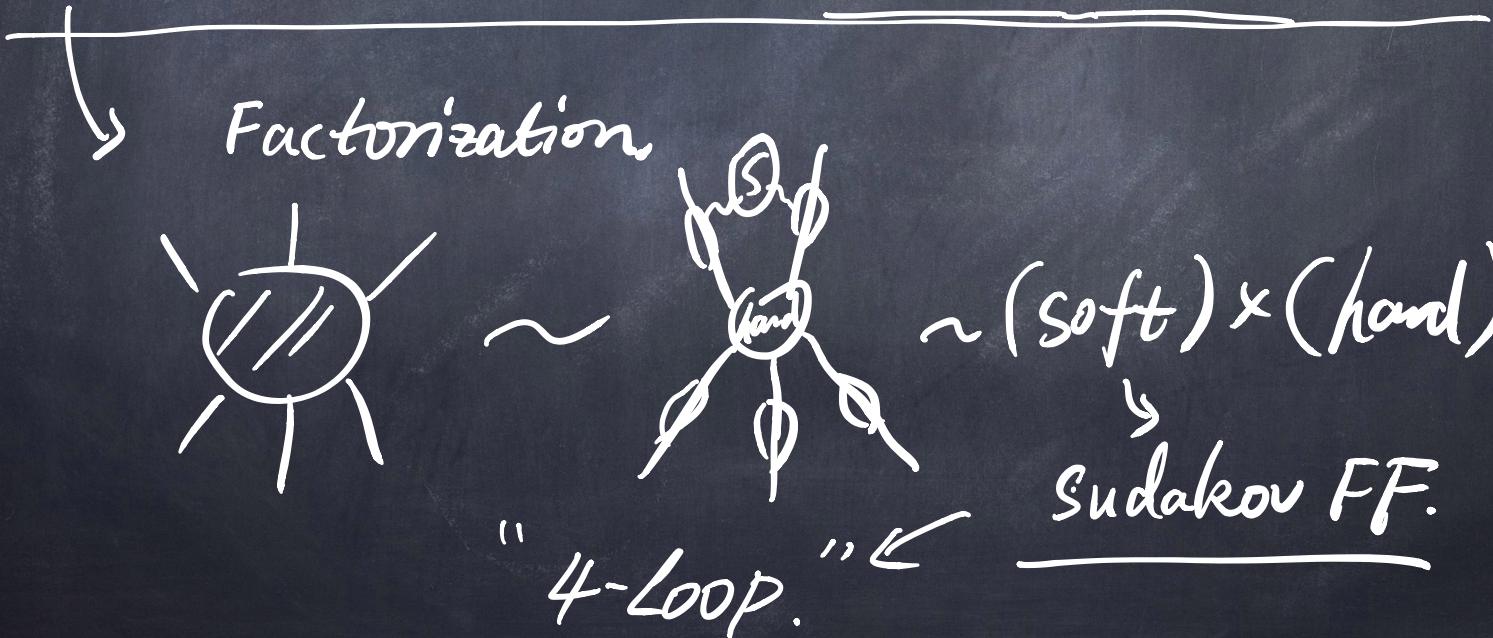
$$\Gamma(p_1, p_2, q) \sim \exp \left[ -e^2 \log \frac{q^2}{p_1^2} / \log \frac{q^2}{p_2^2} \right]$$
$$q^2 \gg p_1^2, p_2^2 \gg m^2$$

- Jackiw (1966) → PhD thesis high-energy form factor  
(on-shell)

$$\Gamma(q) \sim \exp \left[ -e^2 \log \frac{q^2}{m^2} \right]$$

Further development  $\sim 1980$ .

- A. Mueller (79')  $\leadsto$  sub-leading log. resum.  
J. Collins (80') in QED.
- A. Sen (81')  $\rightarrow$  Non-abelian.



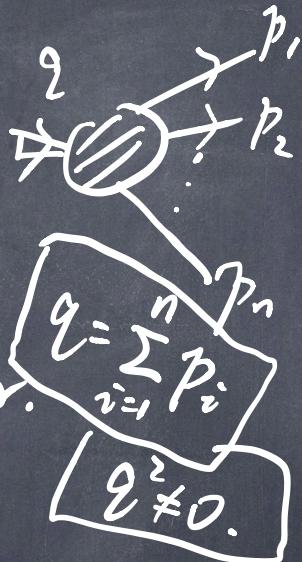
General high-point FF.

F.T.

$$\frac{\int d^Dx e^{-iQx} \langle p_1, p_2 \dots p_n | V(x) | 0 \rangle}{F_{n,V}(p_1 \dots p_n)} \quad \text{Asmpt. on-shell states.}$$

$$\langle p_1 \dots p_n | 0 \rangle$$

Amplitudes.



$$\langle V, O_2 \dots O_n \rangle$$

correlation funct.

2010 draw attentions.

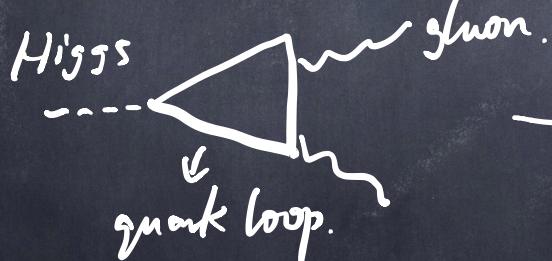
- Maldacena, Zwiebach . 2010.  
using AdS/CFT study strong coupling FF.
- Brandhuber et.al. 2010.      } weak coupling.  
Bork et.al. 2010.

MHV FF:

$$F_{\text{tr}(\phi^2)}^{(0), \text{MHV}} (i^g \dots i^\phi \dots j^\phi \dots n^g) = \frac{\langle ij \rangle^2 \delta^4(\sum_i p_i - q)}{\langle 12 \rangle \langle 23 \rangle \dots \langle n1 \rangle}$$
$$\underline{\text{tr}(F^2)} (i^{g^+} \quad i^{g^-} \quad j^{g^-} \quad n^{g^+}) \rightarrow \text{apply to } N=4 \text{ QCD}$$

# Applications of FF:

- Sudakov FF (IR divergence)  
→ in general gauge theory.
- Anomalous dimensions (UV renormalization)
- EFT Amplitudes.



→ integrate  
over quarks. effective vertices.



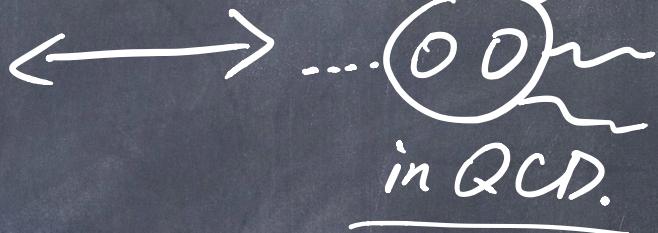
$H \text{ tr}(F_{\mu\nu} F^{\mu\nu})$  + higher  
dim-vert.

Higg + gluons Amp  $\leftrightarrow$  Formfactor with  $\text{tr}(f^2)$   
 $Q^2 = m_H^2$



$$F_{\text{tr}(f^2)}^{(2)}(123) \text{ in } N=4$$

Brandhuber et.al. 2012.



in QCD.

Gehmann et.al.

R 2011.  
Maximal Transcendental

Principle of Max. Transcendentality.

Lipatov et.al. 2001.

$$\gamma^{N=4} = \gamma^{\text{QCD}} / \text{L.T.}$$

This talk we focus on

"Lightlike" form factor. in  $N=4$  SYM.

$$F_4^{LL} = \langle 1^g 2^s 3^s 4^g | \mathcal{L}(q) | 0 \rangle$$

$\xrightarrow{\text{D}(F^2) + \text{super comp.}}$

• Bootstrap to 2-loops.

$$q^2 = 0$$

• FF OPE.

$$q = \sum_i p_i$$

# Bootstrap Method.

$$\text{Loop Amp}_{\text{FF}} = \sum_i \frac{\text{Coeff}_i \times \text{Master Integral}_i}{\text{Physical information}}$$

Physical information.

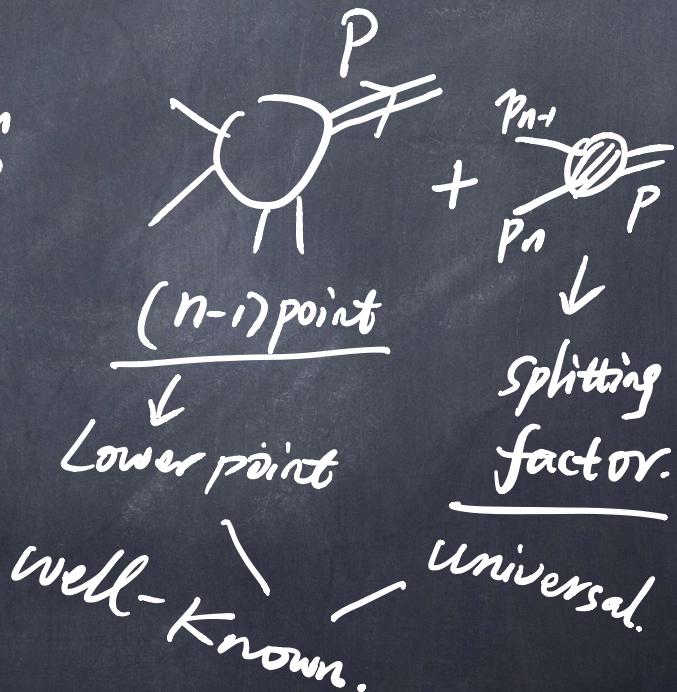
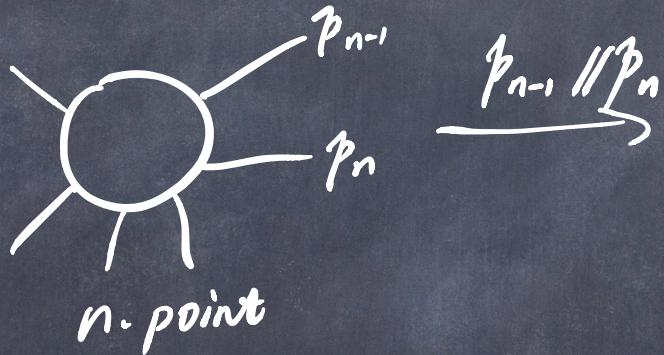
Theory-independent  
Known in many  
(once for all cases).

- Traditional method. "Top-down".
- Bootstrap "bottom-up".

Use physical constraints directly to fix final result.

Constraints :

- IR divergence:  $\frac{1}{\epsilon^\#}$  divergent terms are well-known.
- Collinear limit:



BDS subtraction : Bern, Dixon, Smirnov 2005

$$\underline{F}^{(1)}, \underline{F}^{(2)}$$

define Remainder function:

$$R^{(2)} = \underline{F_{(\epsilon)}^{(2)}} - \frac{1}{2} \left( F_{(\epsilon)}^{(1)} \right)^2 - \underline{f_{(\epsilon)}^{(2)} \times F_{(2\epsilon)}^{(1)}}$$

$$\underline{f^{(2)}(\epsilon)} = -2 \zeta_2 - 2 \zeta_3 \epsilon - 2 \zeta_4 \epsilon^2 \rightarrow \underline{\text{contain splitting info.}}$$

$R^{(2)}$  ① is finite.

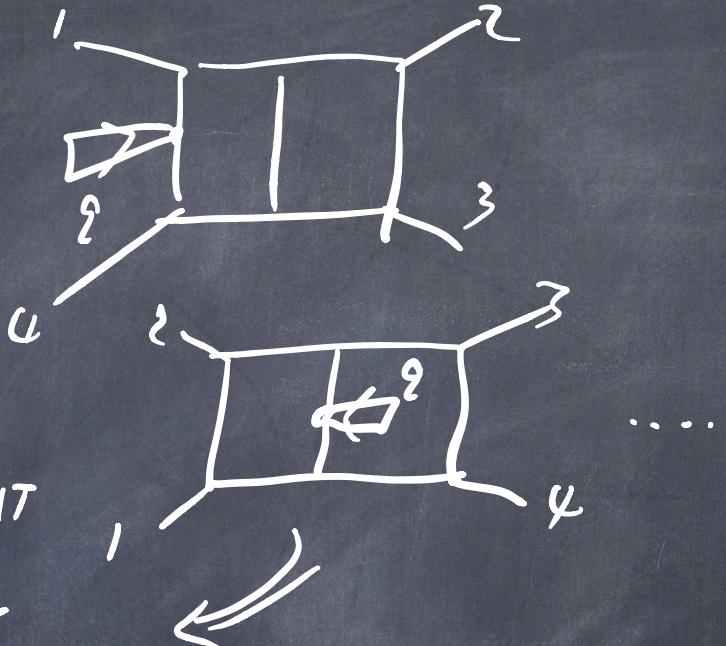
② has trivial collinear limit  $R_n^{(2)} \xrightarrow{P_n \parallel P_{n-1}} R_{n-1}^{(2)}$



① Ansatz

$$F_4^{(2)} = \sum_{i=1}^{590} C_i I_i^{(2), \text{UT}}$$

↓  
Un-known.



$q^2 = 0$  all master integral  
are known.

② One-loop is known

③ Apply constraints.



$$\frac{\#}{C^4} + \frac{\#}{C^3} + \dots$$

\* Symmetry  $D_4$

$$F_4^{(\ell)} = F_4^{(\ell)} \Big/ p_i \rightarrow p_{i+1} = F_4^{(\ell)} \Big/ p_i \rightarrow 5 - p_i$$

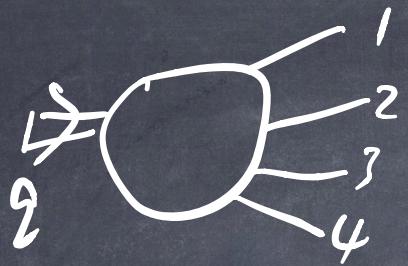
\* IR :  $F_4^{(2)} \xrightarrow{IR} \frac{1}{2}(F_4^{(1)})^2 + f^{(2)}(\epsilon) F_{(2\epsilon)}^{(1)}$

\* Collinear.

$$R_4^{(2)} \xrightarrow{p_i \parallel p_{i+1}} R_3^{(2)} = \text{const.}$$

—————  
known.

$\Rightarrow$  We can fix  $R_4^{(2)}$  uniquely! 



$$P_1, P_2, P_3, P_4, \quad Q = \sum_{i=1}^4 P_i$$

$$S_{ij} : S_{12}, S_{13}, S_{14}, S_{23}, S_{24}, S_{34}$$

↓

- $q^2 = 0, \Rightarrow \sum_{i < j} S_{ij} = 0 \Rightarrow 5\text{-independent}$

↓

- scale invariance.  $u \sim \frac{S_{ij}}{S_{12}} \Rightarrow 4\text{-independent}$

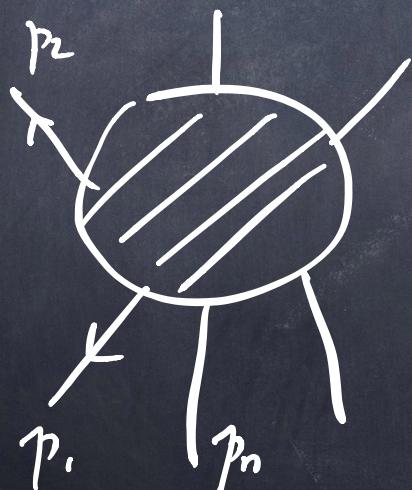
Result shows that there are only

3 independent ratio variables!

→ imply there is a hidden symmetry!

Duality between Amp/FF and Wilson loop.

Alday - Maldacena  $\stackrel{07}{=}$  ( strong coupling )



Amp.



$$x_{i+1} - x_i = p$$

Wilson-loop

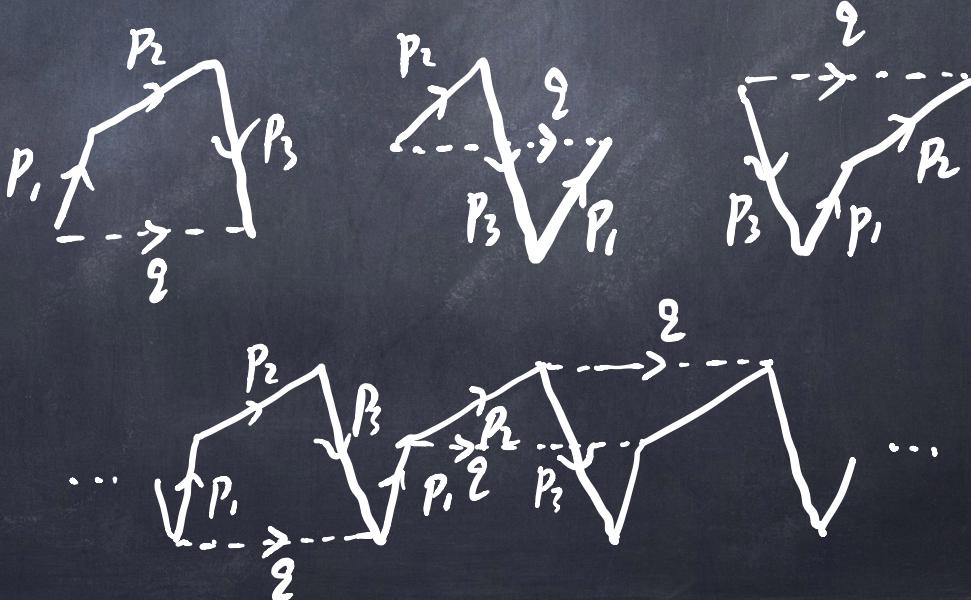
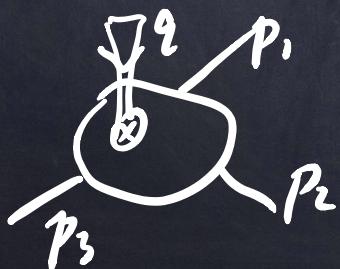
Dual  
 conf sym.  
 Hidden sym!?

$$kw(c) = \text{tr} P e^{\oint A_u(x) dx}$$

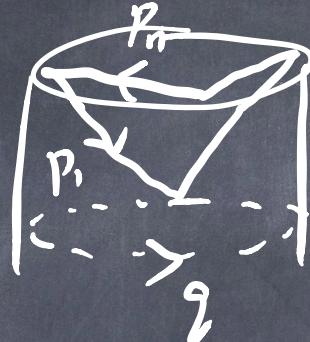
   
 conf Sym.

Form factor:

3-pt example:



form factor  $\leftrightarrow$  periodic Wilson loop



one  
 $q \sim$  period



$q^2 \neq 0$       non-lightlike      C.T.      curved line

straightline

$q^2 = 0$       Lightlike      C.T.      Lightlike

Lightlike

↙ A dual conformal sym along  $q$ -direction.

$$\delta_q x_i^{\mu} = \frac{1}{2} x_i^{\lambda} q^{\mu} - (x_i \cdot p) x_i^{\mu}$$

$$\delta_q \langle W.L. \rangle = 0 = \underline{\underline{\delta_q R^{(L)}}}$$

↳ one sym  $\rightarrow$  eliminate one variable  
as mentioned above.

$$A_n, \text{ } \begin{array}{c} \text{pentagon} \\ \text{W}_n \end{array} : \underline{3n-15} \quad F_{n, q^2=0} : 3n-9$$

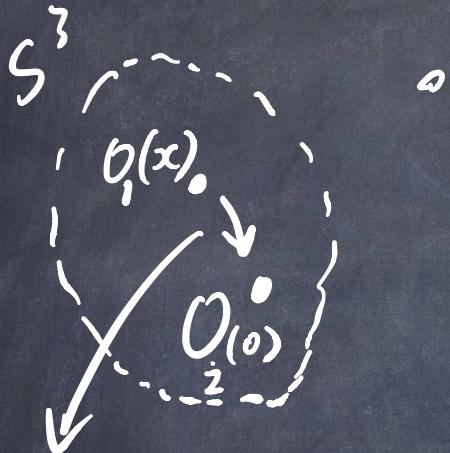
$$F_{n, q^2 \neq 0} \checkmark : 3n-7 \quad n=4 : 3 \times 4-7 = 5$$
$$= \underline{\underline{3}}$$

$$\{p_1, \dots, p_n\} \quad \sum_i p_i = 0 \quad 3n - \underbrace{15}_{\text{conf group}} \\ p_i^2 = 0 \quad 10 + 4 + 1 \\ K^n$$

$$\{p_1, \dots, p_n, q\} \quad p_i^2 = 0 \\ \sum_i p_i = q \quad 3n - (15 - \frac{4}{p^u} - \frac{4}{K^u}) \\ = 3n - 7$$

OPE

Analogy with usual OPE in CFT.



Dilatation transformation.

Hamiltonian.

$$O_i(x) O_j(o) \sim \sum_{k=1}^n f(x, \partial_x) O_k$$

structure  
const.

determine  
by C.S.

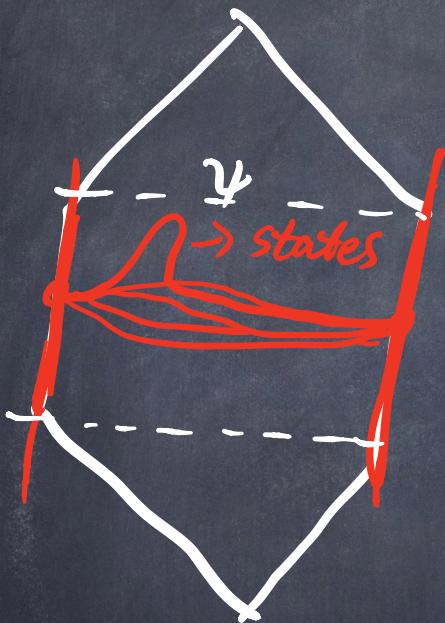
up to  $\Delta, l.$

$O_i$  characterized by  $\Delta, l.$

$$D O_i = \Delta_i D$$

$$\begin{aligned}
 &\text{5 pt correlator} \quad \langle U_1 U_2 \dots U_5 \rangle \\
 &\sim \sum_{\psi_1, \psi_2} \frac{C_{12\psi_1}}{\psi_1} \frac{C_{45\psi_2}}{\psi_2} \cdot \text{conf. P.W.}(\sigma_i) \\
 &\{ C_{12\psi}, \Delta_i \} \text{ CFT data.}
 \end{aligned}$$

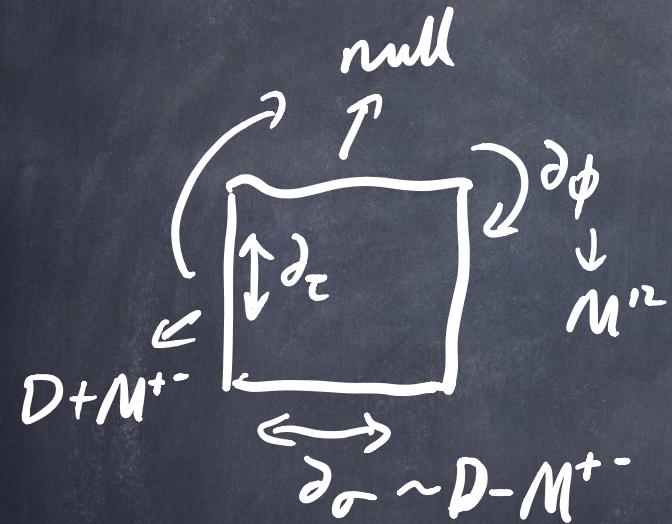
# Wilson Loop OPE.



$$W(\square) = \sum_{\psi} e^{-E_\psi \tau} C_n$$

Alday, Gaiotto, Maldacena.  
Swer. Vieira 2010.

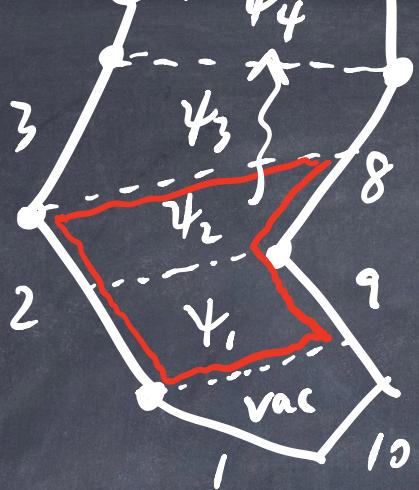
The expansion states are labeled  
by symmetries of a square  
null.



$$\begin{array}{c} \partial\tau, \partial\sigma, \partial\phi \text{ 对易} \\ \downarrow \quad \downarrow \quad \sim_{U(1)} \\ E \quad p \quad m \end{array} .$$



$$\begin{array}{c} \gamma_i \sim \{\sigma_i, \tau_i, \phi_i\} \\ \downarrow \\ \# = 5 \end{array} .$$



$$3n - 15 \Big|_{n=10} = 15 = 3 \times 5.$$

$$W_{10} = \sum_{\{k_i\}} \prod_{i=1}^5 e^{-\tau_i E_i + i p_i \theta_i + i m_i \phi_i}$$

$$\frac{P(\text{vac} / \psi_1) P(\psi_1 / \psi_2)}{P(\psi_5 / \text{vac})}$$

# Analogy :

$\langle \cdot, \cdot \rangle$   $\Delta$   $\longleftrightarrow$



null box.

$$\langle \dots \rangle_{ijk} \longleftrightarrow \{$$


Pentagon.

$P(q/q')$



hexagon

$F(q)$

$q^2 = 0$

$$\langle ::: \rangle G_4 \longleftrightarrow$$

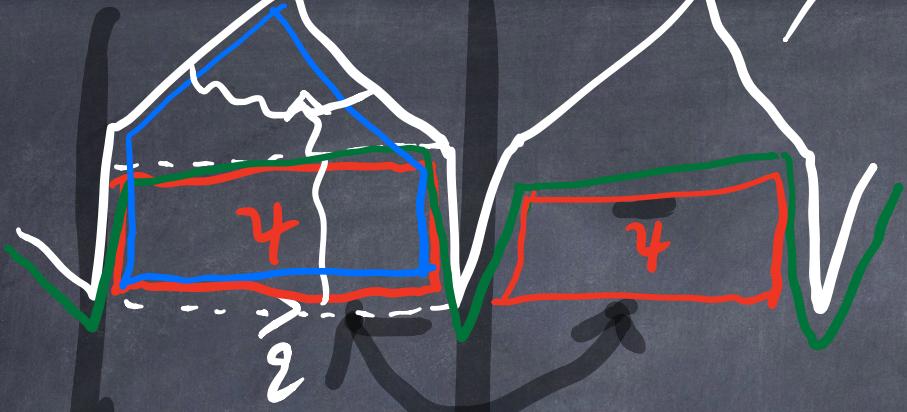
Form factor.

$$\underbrace{3n-9}_{n=3 \text{ const.}}$$



based on  
Sver. Tumanov.  
Wilhelmi 2020.

$$q^2 \neq 0,$$



$\gamma \sim (\tau, \phi, \sigma)$ .

$$\triangle F_4^{\text{LL.}} = \sum e^{-\tau E + i p \sigma + i m \phi} \frac{P(\text{vac}|\gamma) F(\gamma)}{\text{_____}}$$

New building block  
for Lightlike FF.

# Regularization :

$$W_4^{\text{Reg}} = F_4^{\text{Reg}} = \sqrt{4pt} \times \frac{\sqrt{3pt} \times \text{pen}}{\text{box}}$$

red  
box

blue.

Function of  $\{\sigma_i, \tau_i, \phi_i\}$

$$W_4^{\text{Reg}} = -\frac{1}{2} \log \left( \frac{(1-u_1)(1-u_2)}{1-u_3} \right) \log \left( \frac{(1-u_1)(1-u_2)}{1-u_3} \frac{u_3}{u_1} \right)$$

$$U_1 = \frac{s_{12}}{s_{34}} = \frac{x_{4\bar{3}}^2}{x_{3\bar{1}}^2} = e^{-2\sigma}$$

$$U_2 = \frac{s_{23}}{s_{41}} = \frac{x_{2\bar{4}}^2}{x_{4\bar{2}}^2} = e^{-2\tau}$$

$$U_3 = \frac{s_{123}s_{341}}{s_{234}s_{412}} = \frac{x_{14}^2 x_{3\bar{2}}^2}{x_{2\bar{1}}^2 x_{4\bar{3}}^2} = \frac{\cosh(\sigma - \tau) + \cos\phi}{\cosh(\sigma + \tau) + \cos\phi}$$



$$\delta_q U_i = 0$$

From perturbative : Large  $\tau$  expansion

$$W_4^{(1)} = \frac{-2 e^{-2\tau} \cos^2 \phi}{D(e^{-3\tau})}$$

$$W_4^{(2)} = 2 e^{-\tau} \cos \phi h^{(2)}(\sigma)$$

↓ Map to OPE .  $\tau$ -dependence.

$$W_4^{(1)} = 1 + \cos \phi \int \frac{du}{2\pi} e^{-E(u)\tau + i P(u)\sigma} \times \underline{P(0/u) F_F(u)}$$

(first excitation).

$$E(u), P(u), \overline{P}(0/u), M_F(u).$$

△ Known in previous work !

→ From perturbative  
we can determine.

$$F_F(u) = g^0 \times 0 + g^2 \times (*) + V(g^4)$$

$$\stackrel{\equiv}{\downarrow} -2M_F(u) |_{g^2}$$

$$E(u) = \underline{1} + g^{(1)} \cancel{g^2} + \dots$$

$$W_4^{(l)} = \underline{e}^{-\tau} \left( \tau^0 \bar{X}_0 + \tau^1 \bar{X}_1 + \dots + \underline{\tau^{l-2}} \bar{X}_{l-2} \right)$$

$$\bar{X}_{l-1} = 0 \quad (\text{from 1-loop})$$

$$\bar{X}_{l-2} = \int \frac{du}{2\pi} e^{2iu\sigma} \left( \frac{(j^{(1)})^{l-2}}{(l-2)!} \right) \times \left[ \frac{m_p(u)}{g^2} \right]^2$$

$\downarrow$   
 2 loop  
 data.

$$\underline{e}^{-j^{(1)} g^2} (g^2)^{l-2}$$

This is for any loop order!

# Summary & Outlook.

- Amp/FF  $\longleftrightarrow$  WL ( $\square/\sqrt{n}$ )
  - momenta.
- Bootstrap + OPE  $\Rightarrow$  2D.
  - Geometrical

