Geometric Algebra for Special and General Relativity

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Part I.

Special Relativity and Geometric Algebra

Chapter 1.

Introduction

The Special Theory of Relativity is a model of *spacetime* — the geometry in which physical events take place. Spacetime comprises the Euclidean dimensions of space and time, but only in a way relative to each observer moving through it: there exists no single 'universal' ruler or clock. Instead, two observers in relative motion define different decompositions of spacetime, and their respective clocks and rulers are found to disagree according to the Lorentz transformation laws. The insight of special relativity is that one should focus not on the observer-dependent notions of space and time, but on the Lorentzian geometry of spacetime itself.

Seven years after Albert Einstein introduced this theory,¹ he succeeded in formulating a relativistic picture which included gravity. In this General Theory of Relativity, gravitation is identified with the curvature of spacetime over astronomical distances. Both theories coincide locally when confined to sufficiently small extents of spacetime, over which the effects of curvature are negligible. In part I, we will focus on special relativity, leaving gravity and curvature to part II.

In the wake of the Erlangen programme,² the study of local spacetime geometry amounts to the study of its intrinsic symmetries. These symmetries form the Poincaré group, and consist of spacetime translations and Lorentz transformations, the latter being the extension of the group of rotations of Euclidean space to the relativistic rotations of spacetime. The standard matrix representation of the Lorentz group, SO⁺(1, 3), is the connected component of the orthogonal group

$$O(1,3) = \left\{ \Lambda \in GL(\mathbb{R}^4) \, \middle| \, \Lambda^\mathsf{T} \eta \Lambda = \eta \right\}$$

with respect to the bilinear form $\eta = \pm \text{diag}(-1, +1, +1, +1)$. The rudimentary tools of matrix algebra are sufficient for an analys the Lorentz group, and are familiar to any physicist.

However, the last century has seen many other mathematical tools be applied to the study of generalised rotation groups such as $SO^+(1,3)$ or the rotation group SO(3) of \mathbb{R}^3 . Among these tools is the *geometric algebra*, invented by William Clifford in 1878 [4]. Geometric algebra remains largely unknown in the physics community, despite arguably being far superior for the analysis

- ¹ Einstein's paper [1] was published in 1905, the so-called *Annus Mirabilis* or "miracle year" during which he also published on the photoelectric effect, Brownian motion and the massenergy equivalence. Each of the four papers was a monumental contribution to modern physics.
- ² Introduced by Felix Klein in 1872 [2], the Erlangen program is the characterisation of geometries (Euclidean, hyperbolic, projective, etc.) by their symmetry groups and the properties invariant under those groups. E.g., Euclidean geometry studies the invariants of rigid transformations.
- ³ Clifford algebra was independently discovered by Rudolf Lipschitz two years later [3]. He was the first to use them to the study the orthogonal groups.

of rotations than traditional matrix techniques. It is informative to glean some of the history that led to this (perhaps unfortunate) state of affairs.

The quest for an optimal formalism for rotations

Mathematics has seen the invention of a variety of vector formalisms since the 1800s, and the question of which is best suited to physics has a long contentious history.

The vector algebra "war" of 1890–1945 saw William Hamilton's prized quaternion algebra \mathbb{H} , hailed as the optimal tool for describing rotations in \mathbb{R}^3 , struggle for popularity before being eventually left to gather dust as an old-fashioned curiosity.

Chapter 2.

Preliminary Theory

Many of the tools we will develop for the study of spacetime share the property of being associative algebras. As well as the geometric algebra of spacetime, we will encounter tensors, exterior forms, quaternions, and other structures in this category. Instead of defining each algebra axiomatically as needed, it is easier to develop the general theory and then define each algebra succinctly as a particular quotient of the free algebra. This enables the use of the same tools and the same terminology for the analysis of different algebras.

Therefore, this section is an overview of the abstract theory of associative algebras, which more generally belongs to *ring theory*. Algebras, quotients, gradings, homogeneous and inhomogeneous multivectors are defined. Throughout, $\mathbb F$ denotes the underlying field of some vector space. (Eventually, $\mathbb F$ will always be taken to be $\mathbb R$, but we may begin in generality.) Most definitions in this chapter can be readily generalised by replacing the field $\mathbb F$ with a ring.

2.1. Associative Algebras

Definition 1. An ASSOCIATIVE ALGEBRA A is a vector space equipped with a product $\otimes : A \times A \rightarrow A$ which is associative and bilinear.

Associativity means $(u \otimes v) \otimes w = u \otimes (v \otimes w)$, while bilinearity means the product is:

- compatible with scalars: $(\lambda u) \otimes v = u \otimes (\lambda v) = \lambda(u \otimes v)$ for $\lambda \in \mathbb{F}$; and
- distributive over addition: $(u + v) \otimes w = u \otimes w + v \otimes w$, and similarly for $u \otimes (v + w)$.

This definition can be generalised by relaxing associativity or by letting \mathbb{F} be a ring. However, we will use "algebra" exclusively to mean an associative algebra over a field (usually \mathbb{R}).

Any ring forms an associative algebra when considered as a one-dimensional vector space. The complex numbers can be viewed as a real 2-dimensional algebra by defining \otimes to be complex multiplication; $(x_1, y_1) \otimes (x_2, y_2) := (x_1x_2 - y_1y_2, x_1y_2 + y_1x_2)$.

⁴ A RING is a field without the requirement that multiplicative inverses exist or that multiplication commutes; a field is a commutative ring in which non-zero elements are invertible.

The free tensor algebra

The most general (associative) algebra containing a given vector space V is the Tensor algebra V^{\otimes} . The tensor product \otimes satisfies exactly the relations of definition 1 with no others. Thus, the tensor algebra associative, bilinear and *free* in the sense that no further information is required in its definition.

As a vector space, the tensor algebra is equal to the infinite direct sum

$$V^{\otimes} \cong \bigoplus_{k=0}^{\infty} V^{\otimes k} \equiv \mathbb{F} \oplus V \oplus (V \otimes V) \oplus (V \otimes V \otimes V) \oplus \cdots$$
 (2.1)

where each $V^{\otimes k}$ is the subspace of Tensors of Grade k.

2.1.1. Quotient algebras

Owing to the maximal generality of the free tensor algebra, any other associative algebras may be constructed as a *quotient* of V^{\otimes} . In order for a quotient V^{\otimes}/\sim by an equivalence relation \sim to itself form an algebra, the relation must preserve the associative algebra structure:

Definition 2. A CONGRUENCE on an algebra A is an equivalence relation \sim which is compatible with the algebraic relations, so that if $a \sim a'$ and $b \sim b'$ then $a+b\sim a'+b'$ and $a\otimes b\sim a'\otimes b'$.

The quotient of an algebra by a congruence naturally has the structure of an algebra, and so is called a QUOTIENT ALGEBRA.

Lemma 1. The QUOTIENT A/\sim of an algebra A by a congruence \sim , consisting of equivalence classes $[a] \in A/\sim$ as elements, forms an algebra with the naturally inherited operations [a] + [b] := [a + b] and $[a] \otimes [b] := [a \otimes b]$.

Proof. The fact that the operations + and \otimes of the quotient are well-defined follows from the structure-preserving properties of the congruence. Addition is well-defined if [a] + [b] does not depend on the choice of representatives: if $a' \in [a]$ then [a'] + [b] should be [a] + [b]. By congruence, we have from $a \sim a'$ so that [a + b] = [a' + b] and indeed [a] + [b] = [a'] + [b]. Likewise for \otimes . \square

Instead of presenting an equivalence relation, it is often easier to define a congruence by specifying the set of elements which are equivalent to zero, from which all other equivalences follow from the algebra axioms. Such a set of all 'zeroed' elements is called an ideal.

Definition 3. A (TWO-SIDED) IDEAL of an algebra A is a subset $I \subseteq A$ which is closed under addition and invariant under multiplication, so that

• if $a, b \in I$ then $a + b \in I$; and

• if $r \in A$ and $a \in I$ then $r \otimes a \in I \ni a \otimes r$.

Lemma 2. An ideal uniquely defines a congruence, and vice versa, by the identification of I as the set of elements equivalent to zero; $a \sim 0 \iff a \in I$.

Proof. The set $I := \{a \mid a \sim 0\}$ is indeed an ideal because it is closed under addition (for $a, b \in I$ we have $\implies a+b \sim 0+0=0$ so $a+b \in I$) and invariant under multiplication (for any $a \in I$ and $r \in A$, we have $r \otimes a \sim r \otimes 0 = 0 = 0 \otimes r \sim a \otimes r$). Conversely, let $a \sim a'$ and $b \sim b'$. Since \sim respects addition:

$$\begin{vmatrix} a-a' \in I \\ b-b' \in I \end{vmatrix} \implies (a+b)-(a'+b') \in I \iff a+b \sim a'+b',$$

and multiplication:

$$\left. \begin{array}{l} (a-a') \otimes b \in I \\ a' \otimes (b-b') \in I \end{array} \right\} \implies a \otimes b - a' \otimes b' \in I \iff a \otimes b \sim a' \otimes b',$$

the equivalence defined by $a \sim b \iff a - b \in I$ is a congruence.

The equivalence of ideals and congruences is a general feature of abstract algebra.⁵ Furthermore, both can be given in terms of a homomorphism between algebras,⁶ and this is often the most convenient way to define a quotient.

Theorem 1 (first isomorphism theorem). If $\Psi: A \to B$ is a homomorphism, between algebras, then

- 1. the relation $a \sim b$ defined by $\Psi(a) = \Psi(b)$ is a congruence;
- 2. the kernel $I := \ker \Psi$ is an ideal; and
- 3. the quotients $A/\sim \equiv A/I \cong \Psi(A)$ are all isomorphic.

Proof. We assume A and B associative algebras. (For a proof in universal algebra, see $[5, \S 15]$.)

To verify item 1, suppose that $\Psi(a) = \Psi(a')$ and $\Psi(b) = \Psi(b')$ and note that $\Psi(a+a') = \Psi(b+b')$ by linearity and $\Psi(a \otimes b) = \Psi(a' \otimes b')$ from $\Psi(a \otimes b) = \Psi(a) \otimes \Psi(b)$, so the congruence properties of definition 2 are satisfied.

For item 2, note that $\ker \Psi$ is a vector subspace, and that $a \in \ker \Psi$ implies $a \otimes r \in \ker \Psi$ for any $r \in A$ since $\Psi(a \otimes r) = \Psi(a) \otimes \Psi(r) = 0$. Thus, $\ker \Psi$ is an ideal by definition 3.

The first equivalence in item 3 follows from lemma 2. For an isomorphism $\Phi: A/\ker \Psi \to \Omega(A)$, pick $\Phi([a]) = \Psi(a)$. This is well-defined because the choice of representative of the equivalence class [a] does not matter; $a \sim a'$ if and only if $\Psi(a) = \Psi(a')$ by definition of \sim , which simultaneously shows that Φ is injective. Surjectivity follows since any element of $\Psi(A)$ is of the form $\Psi(a)$ which is the image of [a].

- ⁵ For example, in group theory, ideals are *normal* subgroups and a congruence is an equivalence relation satisfying $gag^{-1} \sim id$ whenever $a \sim id$. A group modulo a normal subgroup forms a quotient group.
- ⁶ A homomorphism is a structure-preserving map; in the case of algebras, a linear map $\Psi: A \to A'$ which satisfies $\Psi(a \circledast b) = \Psi(a) \circledast' \Psi(b)$.

With the free tensor algebra and theorem 1 in hand, we are able to describe any associative algebra as a quotient of the form V^{\otimes}/I .

Definition 4. The dimension dim A of a quotient algebra $A = V^{\otimes}/I$ is its dimension as a vector space. The base dimension of A is the dimension of the underlying vector space V.

Algebras may be infinite-dimensional, as is the case for the tensor algebra itself (which is a quotient by the trivial ideal).

2.1.2. Graded algebras

Associative algebras may possess another layer of useful structure: a grading. The grading of the tensor algebra has already been exhibited in eq. (2.1). A grading is a generalisation of the degree or rank of tensors or forms, and of the notion of parity for objects functions or polynomials.

Definition 5. An algebra A is R-GRADED for (R, +) a monoid 7 if there exists a decomposition

$$A = \bigoplus_{k \in R} A_k$$

such that $A_i \otimes A_j \subseteq A_{i+j}$, i.e., $a \in A_i, b \in A_j \Longrightarrow a \otimes b \in A_{i+j}$.

The monoid is usually taken to be $\mathbb N$ or $\mathbb Z$ with addition, possibly modulo some integer. The tensor algebra V^\otimes is $\mathbb N$ -graded, since if $a\in V^{\otimes p}$ and $b\in V^{\otimes q}$ then $a\otimes b\in V^{\otimes p+q}$. Indeed, V^\otimes is also $\mathbb Z$ -graded if for k<0 we understand $V^{\otimes k}:=\{\mathbf 0\}$ to be the trivial vector space. The tensor algebra is also $\mathbb Z_p$ -graded, where $\mathbb Z_p\equiv \mathbb Z/p\mathbb Z$ is addition modulo any p>0, since the decomposition

$$V^{\otimes} = \bigoplus_{k=0}^{p-1} Z_k \quad \text{where} \quad Z_k = \bigoplus_{n=0}^{\infty} V^{\otimes k + np} = V^{\otimes k} \oplus V^{\otimes (k+p)} \oplus \cdots$$

satisfies $Z_i \otimes Z_j \subseteq Z_k$ when $k \equiv i+j \mod p$. In particular, V^{\otimes} is \mathbb{Z}_2 -graded, its elements admit a notion of *parity*: elements of $Z_0 = \mathbb{F} \otimes V^{\otimes 2} \otimes \cdots$ are even, while elements of $Z_1 = V \otimes V^{\otimes 3} \otimes \cdots$ are odd, and parity is respected by \otimes as it is for integers.

Importantly, not all elements of a \mathbb{Z}_2 -graded algebra are either even or odd, and more generally not all elements of a graded algebra belong to a single graded subspace.

Definition 6. If $A = \bigoplus_{k \in R} A_k$ is an R-graded algerba, then an element is HO-MOGENEOUS if it belongs to some A_k , in which case it is said to be a k-VECTOR. If $a \in A_{k_1} \otimes \cdots \otimes A_{k_n}$ is inhomogeneous, we may call it $a\{k_1, ..., k_n\}$ -multivector.

⁷ A MONOID is a group without the requirement of inverses; i.e., a set with an associative binary operation for which there is an identity element.

⁸ Algebras which are \mathbb{Z}_2 -graded are sometimes called *superalgebras*, with the prefix 'super-' originating from supersymmetry theory.

A grading structure may or may not be inherited by quotients. When reasoning about quotients of a graded algebras, the following fact is useful.

Lemma 3. Quotients commute with direct sums, so if

$$A = \bigoplus_{k \in R} A_k$$
 then $A/\sim = \bigoplus_{k \in R} (A_k/\sim)$.

Therefore, a quotient algebra of V^{\otimes} will, as a vector space, admit a similar decomposition into subspaces of fixed grade as in eq. (2.1). However, this does *not* imply that the quotient algebra will itself be graded, in the sense of definition 5. This motivates a stronger notion of congruence:

Definition 7. A congruence \sim on an R-graded algebra $A = \bigoplus_{k \in R} A_k$ is homogeneous (with respect to R) if all equivalent pairs $a \sim b$ of homogeneous non-zero elements $a \in A_i$ and $b \in A_j$ are of the same grade $a, b \in A_k$.

The additional requirement that a congruence be homogeneous ensures that the equivalence relation, as well as respecting the basic algebraic relations, also preserves the grading structure. And so, we have an analogue to lemma 1:

Theorem 2. If an algebra A is R-graded and \sim is a homogeneous congruence on A, then the quotient A/\sim is also R-graded.

Proof. By lemma 3, the quotient $A/\sim = \bigoplus_{k\in R} (A_k/\sim)$ is a direct sum. If \sim is homogeneous, then the A_k/\sim together respect the *R*-grading, meaning

$$(A_i/\sim) \otimes (A_j/\sim) \subseteq A_{i+j}/\sim.$$

Hence, A/\sim is itself *R*-graded.

Part II.

General Relativity and Manifold Geometry

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