Cubical Abstract Machine

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Stack frames and machine states A stack frame (K, \mathcal{E}, Ψ) represents a continuation whose whole binds Ψ dimensions. For a stack π , let $\vec{\Psi}$ be the aggregation of all the dimension bindings Ψ mentioned in π ; then the extension of a stack (K, \mathcal{E}, Φ) :: π is wellformed when \mathcal{E} is an environment for $\mathbf{FD}(K) \setminus \vec{\Psi}$; moreover, K is allowed to mention the dimensions in $\vec{\Psi}$.

A machine state $\langle M \triangleleft \mathcal{E} \parallel \pi \rangle$, where $\vec{\Psi}$ is the aggregation of dimension bindings in π , is wellformed when \mathcal{E} is an environment for $\mathbf{FD}(M) \setminus \vec{\Psi}$; moreover, M is allowed to mention the dimensions in $\vec{\Psi}$.

Selected transition rules We define a transition judgment $C \succ_{\mu} C'$ with μ ranging over the following possible modes:

- 1. Ø, denoting a cubically stable transition
- 2. •, denoting a non-cubically stable transition

Moreover, if $\mathcal{C} \succ_{\square} \mathcal{C}'$ then $\mathcal{C} \succ_{\bullet} \mathcal{C}'$.

$$\frac{\mathcal{E} \equiv (\mathbf{b}, \gamma, \psi)}{\langle a \triangleleft \mathcal{E} \parallel \pi \rangle } \searrow_{\square} \langle \gamma(a) \triangleleft \mathcal{E} \parallel \pi \rangle}$$

$$\underline{\mathcal{E} \equiv (\mathbf{b}, \gamma, \psi)} \qquad \underline{\mathbf{b}}(\mathbf{x}) \equiv \vec{x}; \vec{a}. N \qquad \gamma' \triangleq (\gamma, \overline{a \hookrightarrow M \triangleleft \mathcal{E}}) \qquad \psi' \triangleq (\psi, \overline{x \hookrightarrow r \psi})$$

$$\langle \mathbf{x} [\vec{r}; \vec{M}] \triangleleft \mathcal{E} \parallel \pi \rangle \searrow_{\square} \langle N \triangleleft [(\mathbf{b}, \gamma', \psi')] \parallel \pi \rangle$$

$$\overline{\langle M(N) \triangleleft \mathcal{E} \parallel \pi \rangle} \searrow_{\square} \langle M \triangleleft \mathcal{E} \parallel (\square(N), \mathcal{E}, \square) :: \pi \rangle$$

$$\overline{\langle \lambda a. M \triangleleft \mathcal{E} \parallel (\square(N), \mathcal{E}', \square) :: \pi \rangle} \qquad \overline{\langle \text{snd}(M) \triangleleft \mathcal{E} \parallel \pi \rangle} \searrow_{\square} \langle M \triangleleft \mathcal{E} \parallel (\text{snd}(\square), \mathcal{E}, \square) :: \pi \rangle}$$

$$\overline{\langle \text{fst}(M) \triangleleft \mathcal{E} \parallel \pi \rangle} \searrow_{\square} \langle M \triangleleft \mathcal{E} \parallel (\text{fst}(\square), \mathcal{E}, \square) :: \pi \rangle} \qquad \overline{\langle \text{snd}(M) \triangleleft \mathcal{E} \parallel \pi \rangle} \searrow_{\square} \langle M \triangleleft \mathcal{E} \parallel (\text{snd}(\square), \mathcal{E}, \square) :: \pi \rangle}$$

$$\overline{\langle \langle M, N \rangle \triangleleft \mathcal{E} \parallel (\text{fst}(\square), \mathcal{E}', \square) :: \pi \rangle} \searrow_{\square} \langle M \triangleleft \mathcal{E} \parallel \pi \rangle} \qquad \overline{\langle \langle M, N \rangle \triangleleft \mathcal{E} \parallel (\text{snd}(\square), \mathcal{E}', \square) :: \pi \rangle} \searrow_{\square} \langle N \triangleleft \mathcal{E} \parallel \pi \rangle}$$

$$\underline{r\psi} \equiv \epsilon} \qquad \overline{\langle \text{loop}_r \triangleleft (\mathfrak{g}, \gamma, \psi) \parallel \pi \rangle} \searrow_{\square} \langle \text{base} \triangleleft (\mathfrak{g}, \gamma, \psi) \parallel \pi \rangle}$$

$$\overline{\langle \text{coe}_{x.A}^{r \hookrightarrow r'}(M) \triangleleft \mathcal{E} \parallel \pi \rangle} \searrow_{\square} \langle A \triangleleft \mathcal{E} \parallel (\text{coe}_{\square}^{r \hookrightarrow r'}(M), \mathcal{E}', [x]) :: \pi \rangle}$$

$$\overline{\langle \text{bool} \triangleleft \mathcal{E} \parallel (\text{coe}_{\square}^{r \hookrightarrow r'}(M), \mathcal{E}', [x]) :: \pi \rangle} \searrow_{\square} \langle M \triangleleft \mathcal{E}' \parallel \pi \rangle}$$

The rules for coercions are more difficult, but using second-order substitutions we can account for them without needing to force closures.

$$\frac{\mathcal{E}''\triangleq [\mathcal{E}',\mathbf{x}\hookrightarrow x; a.B\lhd[\mathcal{E},a\hookrightarrow \mathsf{coe}^{r'\overset{\hookrightarrow}{\sim}x}_{x.\mathbf{z}[x]}(a)\lhd[\mathcal{E}',\mathbf{z}\hookrightarrow x.A\lhd\mathcal{E}]],\mathbf{y}\hookrightarrow x.A\lhd\mathcal{E}]}{\left\langle (a:A)\to B\lhd\mathcal{E}\, \middle\|\, (\mathsf{coe}^{r\overset{\hookrightarrow}{\sim}r'}_\square(M),\mathcal{E}',[x])::\pi\right\rangle \succ_{\square} \left\langle \lambda a.\mathsf{coe}^{r\overset{\hookrightarrow}{\sim}r'}_{x.\mathbf{x}[x;a]}(M(\mathsf{coe}^{r'\overset{\hookrightarrow}{\sim}r}_{x.\mathbf{y}[x]}(a)))\lhd\mathcal{E}''\, \middle\|\, \pi\right\rangle}}{\mathcal{E}''\triangleq [\mathcal{E}',\mathbf{x}\hookrightarrow x.A\lhd\mathcal{E},\mathbf{y}\hookrightarrow x.B\lhd[\mathcal{E},a\hookrightarrow \mathsf{coe}^{r\overset{\hookrightarrow}{\sim}r'}_{y.\mathbf{x}[y]}(\mathsf{snd}(M))\lhd[\mathcal{E},\mathbf{x}\hookrightarrow x.A\lhd\mathcal{E}]]]}{\left\langle (a:A)\times B\lhd\mathcal{E}\, \middle\|\, (\mathsf{coe}^{r\overset{\hookrightarrow}{\sim}r'}_\square(M),\mathcal{E}',[x])::\pi\right\rangle \succ_{\square} \left\langle \left\langle \mathsf{coe}^{r\overset{\hookrightarrow}{\sim}r'}_{x.\mathbf{x}[x]}(\mathsf{fst}(M)),\mathsf{coe}^{r\overset{\hookrightarrow}{\sim}r'}_{x.\mathbf{y}[x]}(\mathsf{snd}(M))\right\rangle \lhd\mathcal{E}''\, \middle\|\, \pi\right\rangle}$$

Unloading the machine We can unload the machine at any time; this is useful if we are computing an open term and hit a variable.

$$\frac{\langle K[\Psi.M[\mathcal{E}]] \triangleleft \mathcal{E}' \parallel \pi \rangle \Longrightarrow N}{\langle M \triangleleft \mathcal{E} \parallel \cdot \rangle \Longrightarrow M[\mathcal{E}]} \qquad \qquad \frac{\langle K[\Psi.M[\mathcal{E}]] \triangleleft \mathcal{E}' \parallel \pi \rangle \Longrightarrow N}{\langle M \triangleleft \mathcal{E} \parallel (K[\square], \mathcal{E}', \Psi) :: \pi \rangle \Longrightarrow N}$$