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STRONG GRAVITATIONAL LENSING BY WAVE DARK MATTER HALOS

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ABSTRACT

Wave Dark Matter (WaveDM) has recently gained attention as a viable candidate to account for the dark matter content of the Universe. In this paper we explore the extent to which, and under what conditions, dark matter halos in this model are able to reproduce strong lensing systems. First, we explore analytically the lensing properties of the model, finding that a pure WaveDM density profile, soliton profile, produces a weaker lensing effect than similar cored profiles. Then we analyze models with a soliton embedded within an NFW profile, as has been found in numerical simulations of structure formation. We use a benchmark model with a boson mass of $m_a = 10^{-22}\,\mathrm{eV}$, for which we see that there is a bi-modality in the contribution of the external NFW part of the profile, and some of the free parameters associated with it are not well constrained. We find that for configurations with boson masses 10^{-23} – $10^{-22}\,\mathrm{eV}$, a range of masses preferred by dwarf galaxy kinematics, the soliton profile alone can fit the data but its size is incompatible with the luminous extent of the lens galaxies. Likewise, boson masses of the order of $10^{-21}\,\mathrm{eV}$, which would be consistent with Lyman-α constraints and consist of more compact soliton configurations, necessarily require the NFW part in order to reproduce the observed Einstein radii. We then conclude that lens systems impose a conservative lower bound $m_a > 10^{-24}\,\mathrm{eV}$ and that the NFW envelope around the soliton must be present to satisfy the observational requirements.

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1. INTRODUCTION

The ΛCDM model is the most successful theoretical framework in modern cosmology to explain the process of structure formation in the Universe on large scales. This model requires the existence of a cold dark matter (CDM) component that comprises 26% of the total energy budget, which is best described by a non-relativistic (cold) and non-interacting fluid (see Ade et al. 2016).

One of the main predictions from only CDM simulations of structure formation is the appearance of universal cuspy density profiles for the galaxy halos, with the Navarro, Frenk and White (NFW) profile the one most used to describe CDM (see Navarro et al. 1997). Despite the successes of CDM at large scales, there are some open questions regarding the observations on galactic scales, such as: the "missing satellite problem", the "cusp core problem", and the "too-big-to-fail problem" (e.g. Burkert 1995; Maccio et al. 2012; de Blok 2010; Oh et al. 2011; Sawala et al. 2011; Klypin et al. 1999; Boylan-Kolchin et al. 2011), see also Bullock & Boylan-Kolchin (2017) for a recent review. These refer to both theoretical and observational questions on how CDM and barions interplay leads to the shapes, inner density profiles, and abundance of the DM structure at sub-galactic scales. This also opens the question of weather such observables, or others, can be used to learn more about the dark matter (DM) nature. This has lead to explore the viability of other DM candidates. Actually there is a wide range of DM proposals such as Self-Interacting Dark Matter (see Kaplinghat et al. 2016), Warm Dark Matter (see Gonzalez-Samaniego et al. 2016; Drewes et al. 2017), Axion/Scalar or Wave Dark Matter (e.g. Matos et al. 2000; Hu et al. 2000; Goodman 2000; Matos & Urena-Lopez 2001; Böhmer & Harko 2007; Schive Hsi-Yu et al. 2014; Robles & Matos 2013), and other specifications of the nature of dark matter particles, which can actually be described in a more general effective theory (e.g. Cyr-Racine et al. 2016).

In this paper, our approach is to describe the dark matter as an axion/scalar field that we will refer to as a Wave Dark Matter model (WaveDM, also referred sometimes to as scalar field DM or SFDM, ultralight axion-like DM, fuzzy DM, etc.). This type of model has been investigated by several other authors (e.g. Matos et al. 2000; Hu et al. 2000; Goodman 2000; Matos & Urena-Lopez 2001; Böhmer & Harko 2007; Lee 2017a; Suárez et al. 2014; Lee 2017b), and has been found to be able to reproduce the success of the ΛCDM model on cosmological scales, but it predicts a natural cut-off on the mass power spectrum of linear perturbations that could help to alleviate some of the low-scale issues of CDM (e.g. Ureña-López & Gonzalez-Morales 2016; Hlozek et al. 2015; Matos & Urena-Lopez 2001; Hu et al. 2000). Interestingly enough, all cosmological effects are directly related to a single parameter, which is the boson mass, m_a, of the scalar field particle, although extra observational effects may arise from quartic self-interactions (e.g. Schive & Chiueh 2017; Linares Cedeño et al. 2017; Zhang & Chiueh 2017a,b). Based on considering the cut-off of the mass power spectrum, the halo mass function, the reionization time or the Lyman-α forest, the most up-to-date constraints suggest that the boson mass must satisfy $m_a > 1 \times 10^{-21}$ eV (see Iršič et al. 2017; Armengaud et al. 2017).

However, the non-linear process of structure formation under the SFDM hypothesis does not depend on only a single parameter, but instead one requires to take into account at least a second parameter. This fact is indeed considered in many recent studies that try to put constraints on the WaveDM parameters with data coming from, for instance, satellite galaxies in the Milky Way (e.g. Bernal et al. 2017; González-Morales et al. 2016; Chen et al. 2016; Schive et al. 2014; Schive Hsi-Yu et al. 2014). The aforementioned studies consider that galaxies are described by a solitonic core with a negligible self-interaction. The soliton solution is just the ground state of the so-called Schrodinger-

Poisson system of equations (Ruffini & Bonazzola 1969; Guzman & Urena-Lopez 2004), and its wave-like properties provide stability against gravitational collapse – opening the possibility of naturally-supported, cored halos. The full prescription of the WaveDM profile requires specification of the boson mass m_a together with one of its structure parameters, which can be taken to be either the central density or the scale radius, while the other is determined by the relation,

$$\frac{\rho_s}{M_\odot \text{pc}^{-3}} = 2.4 \times 10^{12} \left(\frac{r_s}{\text{pc}}\right)^{-4} \left(\frac{m_a}{10^{-22}\text{eV}}\right)^{-2} \quad . \tag{1}$$

The boson mass m_a is expected to be a fundamental parameter with a single value for all galaxies, while the other two parameters may take values that differ from galaxy to galaxy. Hence, it is necessary to think more carefully if we are to obtain meaningful constraints on the boson mass. More specifically, if we consider the boson mass as an universal parameter, on the same footing as any other cosmological parameter, we should certainly be able to use statistical analysis of galaxy data to constrain which values are permitted, as has been proposed in González-Morales et al. (2013); Diez-Tejedor et al. (2014) and more recently carried out in Chen et al. (2016); González-Morales et al. (2016). However, in general, we may be unable to assert whether there is one single value of m_a that is suitable to satisfy all the possible constraints. For this purpose, in this paper we have selected gravitational lensing as a possible additional tool to study or assess the viability of the WaveDM profile whose parameters are subjected to the constraint in Eq. (1).

Gravitational lensing (strong and weak) has become a powerful astrophysical tool for the study of the background cosmology, structure and substructures of galactic halos. (see Schneider et al. 1999; Cao et al. 2015; Koopmans 2005; Cao et al. 2016; Bolton et al. 2006). In particular, it is possible to extract important information from the stellar kinematics and geometry of strong lensed systems. Furthermore, particular cosmologies where dark energy is not a cosmological constant can be tested by fitting the observed critical lines, Einstein angle and stellar dynamics, given a suitable lens model for the observed systems (see e.g. Futamase & Yoshida 2001; Grillo et al. 2008; Sereno 2002; Newton et al. 2011). As the geometry of the lensing system can be obtained from image astrometry, it also is a helpful probe for the Hubble parameter and the dark energy contents (Mitchell et al. 2005; Biesiada et al. 2010), whereas at the same time it gives information about the structure and formation of Early type galaxies (Koopmans et al. 2006; Gavazzi et al. 2007).

There are different procedures to extract particular parameters, like the Einstein angle, from the lensing geometry in combination with velocity dispersions from stellar dynamics (Bolton et al. 2008b; Auger et al. 2010), where the latter can be obtained along with redshifts from spectroscopic imaging (Ofek et al. 2003, 2006; Kochanek 1991; Bolton et al. 2006; Cao & Zhu 2012). When assuming a particular lens model it is usual, as a first approach, to describe strong lenses using an axially symmetric model, the most popular being the Singular Isothermal Sphere (SIS) (Mollerach & Roulet 2002; Schneider et al. 1999; Auger et al. 2009; Cao et al. 2012). Sometimes, however, there are deviations from the SIS model and the Singular Isothermal Ellipsoid (SIE) is used instead as a non-axially symmetric extension (Gavazzi et al. 2007; Cao et al. 2015). By assuming these models, which consider the total mass distribution of the lens, there is some freedom to test the cosmological parameters (Cao et al. 2016; Jie et al. 2016; Sereno 2002). On the other hand, using a fixed cosmology allows one to obtain information of the physical parameters regarding the structure of individual galaxies or clusters of galaxies (Futamase & Yoshida 2001; Grillo et al. 2008; Sereno 2002). Furthermore, using information obtained from the strong lensed galaxies is possible to test the gravitational weak field,

through the post-Newtonian parameter γ, where the parameterization of the profiles to describe are sensitive to the total mass distribution. This can be used to test the validity of General Relativity on large scales. Recently, it has been found that for at least ~ 2kpc, the lensing galaxies profiles agree with General Relativity (see e.g. Cao et al. 2017; Collett et al. 2018).

For the description of lensing systems, several mass distribution profiles exist which have been successful enough to represent observed data; the most popular, which include the SIS and SIE (see e.g. Keeton 2001), follow a power-law distribution. These are very successful to describe the observed data for lensing geometry and stellar dynamics when considering the total matter contents (luminous+dark matter) (Suyu et al. 2014; Bolton et al. 2008a,b; Auger et al. 2009; Cao et al. 2012; Cao et al. 2015). It is important to stress out that these modelling choices does not impact on the estimated Einstein angle of the lens. Nevertheless, these modelling choices do not explicitly give much information about the properties of the internal structure of the dark matter distribution in the galaxies. For this reason, other composite models where the baryonic matter and the dark matter are treated separately are also used (e.g. Suyu et al. 2014; Navarro et al. 1997; Park & Ferguson 2003). The luminous part, which contains the baryonic matter, is described by a SIS, but it is common to use Sersic or more general power-law luminosity profiles (Cardone 2004; Cao et al. 2016). The most popular to describe the DM halo of a galaxy is the NFW density profile (Wright & Brainerd 2000; Bartelmann 1996), which introduces a family of generalized NFW profiles (Wyithe et al. 2001). Other popular profiles are the Burkert profile (Park & Ferguson 2003; Burkert 1995), and cuspy halo models (Keeton 2001; Muñoz et al. 2001). It has been shown that the NFW profile correctly describes the observed lensing signal in large samples of galaxies (Gavazzi et al. 2007) and clusters of galaxies (Niikura et al. 2015).

Since the WaveDM model is considered a feasible candidate for DM, in this work we study the behavior of, and put constraints upon, a WaveDM type of profile acting as a single galactic gravitational lens, and we obtain the conditions under which the profile will be able to produce strong lensing. As we shall see, the WaveDM profile consists of a solitonic core plus a tail in the outer parts that follows the prescription of an NFW profile, whose properties are closely interlinked by their matching conditions such that a soliton is always present in the centre of the DM halo. In this respect, we are not interested in the individual properties of the single profiles (soliton vs NFW), but rather on the conditions under which the complete WaveDM profile could be consistent with lensing data.

The remainder of this paper is organized as follows. The basic lensing equations for any given density profile are described in Sec. 2, where we also introduce the explicit lensing expressions for the particular case of the WaveDM profile. In Sec. 3 we describe our statistical analysis and present the results arising from the comparison of the WaveDM model predictions with selected data from the SLACS catalog. Finally, the general conclusions are presented in Sec. 4. Some analytic solutions of the lens equations used in the text are shown in the appendix.

2. GRAVITATIONAL LENSING WITH A ψDM PROFILE

2.1. *General lensing equations*

One of the main predictions from Einstein's General Relativity is the bending of light as it passes close to a massive body. The deflection angle produced by this effect depends on the mass of the deflector, which then acts like a lens. This deflector may be approximated by a point-like mass,

like in the case of a star, but for more massive objects like galaxies it is better to represent them as extended masses which are described by their density profiles. For the purposes of this work, we shall consider galactic lenses, and therefore the density profile described in this section will be representing one galaxy acting as a single lens.

The simplest type of lens is a system with a point mass M located close to the line of sight to a luminous source S. Due to the gravitational field of the point mass, a light ray is deflected in its path to the observer; this is described by the lens equation in the thin-lens approximation. The same approximation also holds for a mass distribution, in which case the lens equation is (Mollerach & Roulet 2002),

$$\beta = \theta - \frac{m(\theta)}{\pi \Sigma_{cr} D_{OL}^2 \, \theta}\,, \tag{2}$$

that relates the (unobservable) angle between the line of sight and the path from the observer to the actual position of the source, β, and to the apparent position of the source (the image), θ, and the mass distribution that is causing the lensing $m(\theta)$. Also, D_{OL} is the angular distance from the observer to the lens, that we denote by the subindex (OL). Here we have assumed that $m(\theta)$ is the projected mass enclosed in a circle of radius $\xi \equiv D_{OL}\theta$; more explicitly, we can write

$$m(\xi) = 2\pi \int_0^\xi d\hat{\xi}\,\hat{\xi}\,\Sigma(\hat{\xi})\,. \tag{3a}$$

The projected surface mass density $\Sigma(\xi)$ can be calculated directly from the (spherically symmetric) density profile $\rho(r)$ of the lensing object as:

$$\Sigma(\xi) = 2 \int_0^{z_{\max}} dz\,\rho(z,\xi)\,, \tag{3b}$$

where $z \equiv \sqrt{r^2 - \xi^2}$ is a coordinate orthogonal to the line of sight, so that $0 \le \xi \le r$. If the lens system has a finite radius $r_{\max}$, then $z_{\max} = \sqrt{r_{\max}^2 - \xi}$; otherwise, we can put $z_{\max} \to \infty$ in the integral given by Eq. (3b).

Let us consider the case in which the density profile $\rho(r)$ has a characteristic density ρ_s, and a characteristic radius r_s, such that $\rho(r) = \rho_s f(r/r_s)$, where f is the function that accounts for the shape of the profile. We can then write Eq. (2) in the dimensionless form

$$\beta_*(\theta_*) = \theta_* - \lambda \frac{m_*(\theta_*)}{\theta_*}\,, \tag{4}$$

where the different distances are normalized in terms of r_s: $\beta_* = D_{OL}\beta/r_s$, $\theta_* = D_{OL}\theta/r_s$, and then $\xi_* = \xi/r_s = \theta_*$. The latter equation means that the normalized variables ξ_* and θ_* can be used interchangeably, and then hereafter we will use θ_* as our distance variable[1]. Likewise, the total mass, as given in Eq. (3a), is normalized as

$$m_*(\theta_*) = \frac{m(\theta_*)}{\rho_s r_s^3} = 2\pi \int_0^{\theta_*} d\hat{\theta}_*\,\hat{\theta}_*\,\Sigma_*(\hat{\theta}_*)\,. \tag{5a}$$

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[1] For the sake of simplicity in the notation, we are using the same same angular variables (together with an asterisk) to denote the new normalized distances.

Name	Density profile $f(r)$	λ_{cr}	References
NFW	$\left[(r/r_s)(1+r/r_s)^2\right]^{-1}$	0	Wright & Brainerd (2000)
Burkert	$\left[(1+r/r_s)(1+^2 r/r_s^2)\right]$	$2/\pi^2 \simeq 0.203$	Park & Ferguson (2003)
SFDM	$\sin(\pi r/r_s)/(\pi r/r_s)$	0.27	González-Morales et al. (2013)
WaveDM	$(1+r^2/r_s^2)^{-8}$	$\frac{2048}{429\pi^2} \simeq 0.484$	(Marsh & Pop 2015; Schive et al. 2014)

Table 1. The intrinsic value λ_{cr}, calculated from Eq. (8) for dark matter halos with different density profiles.

The normalized projected surface mass density, from Eq. (3b), is

$$\Sigma_*(\theta_*) = \frac{\Sigma(\theta_*)}{\rho_s r_s} = 2 \int_0^{z_{\mathrm{max}*}} dz\, f(z,\theta_*)\,, \tag{5b}$$

with $z = \sqrt{r_*^2 - \theta_*^2}$ and $r_* = r/r_s$. The new parameter λ in Eq. (4) is then given by

$$\lambda \equiv \frac{\rho_s r_s}{\pi \Sigma_{\mathrm{cr}}} = 10^{-3}\frac{0.57}{h}\left(\frac{\rho_s r_s}{M_\odot \mathrm{pc}^{-2}}\right)\frac{d_{\mathrm{OL}} d_{\mathrm{LS}}}{d_{\mathrm{OS}}}\,, \tag{6}$$

where we have defined the reduced (dimensionless) angular distances $d_A = D_A H_0/c$. The angular diameter distance D_A as a function of redshift is computed in the standard way (see e.g. Hogg 1999), assuming cosmological model parameters as given by the Planck 2015 results (Ade et al. 2016). Equation (6) contains information about the lensing properties of any given model, together with that of the different distances involved in the lens system, namely that between the observer and the lens (OL), the observer and the source (OS), and the lens and the source (LS).[2]

One particular case of interest is that of perfect alignment between the luminous source and the lens system for which $\beta_*(\theta_{*E}) = 0$. This in turn defines an Einstein ring with radius $R_E = D_{OL}\theta_E$, and an associated angular radius θ_E. In terms of our normalized variables, we see that the normalized angular Einstein radius θ_{*E} directly is the ratio of the Einstein radius to the (characteristic) scale radius of the density profile, $\theta_{*E} = R_E/r_s$. Moreover, the angular radius θ_{*E} must also be a solution of the equation [see Eq. (4)]

$$\lambda = \frac{\theta_{*E}^2}{m_*(\theta_{*E})}\,. \tag{7}$$

Interestingly enough, Eq. (7) shows that the lensing properties of a system with a density profile of the form $\rho(r) = \rho_s f(r/r_s)$ are independent of the density and distance scales, and are mostly sensitive to the particular shape of the density profile. The physical parameters of the system are then concentrated in the dimensionless parameter λ in Eq. (6), and the latter can be calculated from

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[2] This is the same parameter used in González-Morales et al. (2013), but also see Park & Ferguson (2003) in which the definition of λ differs by a factor of $1/4\pi$.
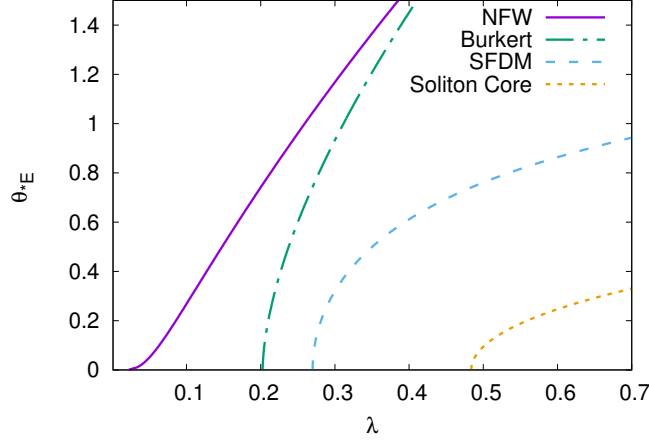
Figure 1. Normalized Einstein radius θ_{*E} as a function of λ for different density profiles, see Eq. (7). The point where each curve crosses the horizontal axis indicates the (intrinsic) critical value λ_{cr} for each profile, in agreement with the values calculated from Eq. (8) as shown in Table 1.

Eq. (7) without any prior knowledge of the given physical scales in the system, namely ρ_s and r_s, under only the assumption of perfect alignment (see Fig. 1 below for an example).

There is a critical value λ_{cr} that is the smallest value of λ for which an Einstein ring appears, which must correspond to the limit $\theta_{*E} \to 0$ in Eq. (7). As we shall show now, such a critical value can be calculated analytically in the general case. To avoid the divergence at $\theta_{*E} = 0$ (where $m_*(0) = 0$), we make use of the L'Hôpital rule in Eq. (7), and from Eq. (3a) we finally obtain

$$\lambda_{\mathrm{cr}}^{-1} = \pi\Sigma_*(0) = 2\pi \int_0^{r_{\mathrm{max}*}} d\hat{r}_* \, f(\hat{r}_*) \,, \tag{8}$$

where $\Sigma(0)$ is the central value of the projected surface mass density given by Eq. (3b). Eq. (8) is quite a simple formula for the calculation of λ_{cr} for any given density profile $\rho(r)$.[3]

As said before, Eq. (8) suggests that the critical value λ_{cr} just depends on the particular shape of the given density profile and no information is necessary about its other physical parameters. The values of λ_{crit}, calculated from Eq. (8) for density profiles that are well-known in the literature, are shown in Table 1. For these profiles we also show in Fig. 1 the Einstein angle θ_{*E} as calculated from Eq. (7). As expected, the Einstein angle is the smallest for the WaveDM profile (Eq. (10)) alone, which also means that it is the one with the weakest lensing signal.

We should mention here an additional use of Eq. (7) to constrain the free parameters of a given density profile. It relates to the fact that any DM halo characterized by a particular density profile needs to satisfy the constraint $\lambda \geq \lambda_{\mathrm{cr}}$ if it is to produce a lensing signal. Using Eqs. (6) and (8), the latter statement can be re-written as

$$\frac{\rho_s r_s}{M_\odot \mathrm{pc}^{-2}} \geq 10^3 \frac{h}{0.57} \frac{d_{\mathrm{OS}}}{d_{\mathrm{OL}} d_{\mathrm{LS}}} \lambda_{\mathrm{cr}} \,. \tag{9}$$

Equation (9) establishes a minimum value for the (structural) surface density $\rho_s r_s$ of any given DM profile in terms of the measured quantities of a lens system. Although the constraint Eq. (9) is

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[3] It should be noted that the definition of λ_{cr} depends on the chosen scale radius for normalization r_s, so that the value obtained from Eq. (8) in our case is considering that r_s coincides with the intrinsic distance scale in the density profile $\rho(r)$.

satisfied automatically by the NFW profile, for which $\lambda_{\mathrm{crit}} = 0$, this is not the case for the other profiles listed in Table 1.

2.2. *Combined density profile of WaveDM*

For the density profile of WaveDM halos we will consider the model described in Schive Hsi-Yu et al. (2014); Schive et al. (2014), which arises from the study of extensive N-body simulations. The profile consists basically of two parts: one part describing a core sustained by the quantum pressure of the boson particles, also known as the soliton profile, and another part that resembles a NFW-like profile in the outer parts of the halo. As argued in Marsh & Pop (2015), the transition at some radius to a NFW profile must be expected from the change of behavior to CDM on scales larger than the natural length of coherence, which should be proportional to the associated Compton length of the boson particles (in full units, the Compton length is $L_C = \hbar/(mc)$, where $\hbar$ is the (reduced) Planck's constant and c is the speed of light).

The soliton profile is given by

$$\rho_{\mathrm{sol}}(r) = \frac{\rho_s}{(1 + r^2/r_s^2)^8}\,, \tag{10}$$

where r_s and ρ_s are its characteristic radius and central density contrast, respectively. This profile was first studied in detail in Schive et al. (2014), although here we are following the nomenclature adopted in Marsh & Pop (2015), where it is also shown that the profile fits well the ground-state solution of the so-called Schrödinger-Poisson (SP) system of equations (Ruffini & Bonazzola 1969; Guzman & Urena-Lopez 2004). In this respect, the soliton profile is strongly related to the wave properties (via the Schrödinger equation) of the boson particles.

One important property of the profile given in Eq. (10) is that it must also obey the intrinsic scaling symmetry of the SP system (Guzman & Urena-Lopez 2004). If $0 < \hat{\lambda} \ll 1$ is a constant parameter, it can be shown that the central density and radius in the soliton profile are given by

$$\rho_s = \hat{\lambda}^4 m_a^2\, m_{\mathrm{Pl}}^2/4\pi\,, \quad r_s = (0.23\,\hat{\lambda}\,m_a)^{-1}\,, \tag{11}$$

This equation suggests that the intrinsic, physical, quantities of the soliton profile in Eq. (10) are related as shown in Eq. (1). This relation will be important later when we discuss the constraints on the boson mass m_a.

For the NFW profile at the outskirts of the galaxy halo we adopt the following parametrization

$$\rho_{\mathrm{NFW}}(r) = \frac{\rho_s\, \rho_{\mathrm{NFW}*}}{\alpha_{\mathrm{NFW}}\,(r/r_s)\,(1 + \alpha_{\mathrm{NFW}}\, r/r_s)^2}\,. \tag{12}$$

Notice that in writing Eq. (12) we are assuming the following implicit definitions for the scale radius and density, respectively, of the NFW profile: $r_{\mathrm{NFW}} = r_s/\alpha_{\mathrm{NFW}}$ and $\rho_{\mathrm{NFW}} = \rho_s\, \rho_{\mathrm{NFW}*}$, where both α_{NFW} and $\rho_{\mathrm{NFW}*}$ are dimensionless numbers.

Unfortunately, there is not precise information in Schive Hsi-Yu et al. (2014) about the transition in a galaxy halo from the soliton profile of Eq. (10) to the NFW profile of Eq. (12) in the general case. Hence, for the present work we adopt the convention for a combined profile as suggested in Marsh & Pop (2015)

$$\rho(r) = \Theta(r_\epsilon - r)\rho_{\mathrm{sol}}(r) + \Theta(r - r_\epsilon)\rho_{\mathrm{NFW}}(r)\,. \tag{13}$$

where $\Theta(r_\epsilon - r)$ is the Heaviside step function. Here, r_ϵ is the matching radius where the transition between the individual profiles occurs, and which satisfies the condition $\rho(r_\epsilon) = \epsilon\rho_s$. Notice that $0 < \epsilon < 1$ if the transition between the profiles is to occur at the outskirts of the galaxy halo.

In general terms, and under our parametrization, there are six free parameters in the combined profile (Eq. (13)): $(\rho_s, r_s, \rho_{\mathrm{NFW}*}, \epsilon, r_\epsilon, \alpha_{\mathrm{NFW}})$. We will now derive two new constraints that arise from the continuity of the combined density profile at the matching radius which will help us to reduce the number of free parameters.

For a continuous density function, we must impose the condition

$$\rho_{\mathrm{sol}}(r_\epsilon) = \epsilon\rho_s = \rho_{\mathrm{NFW}}(r_\epsilon) \,. \tag{14}$$

When Eq. (14) is applied to the soliton profile of Eq. (10), we obtain

$$r_{\epsilon*} = r_\epsilon/r_s = \left(\epsilon^{-1/8} - 1\right)^{1/2} \,, \tag{15a}$$

which basically establishes the interchangeability of the (dimensionless) matching radius $r_{\epsilon*}$ and ϵ. In the case of the NFW profile (Eq. (12)), the continuity condition (Eq. (14)) establishes that

$$\epsilon^{-1}\rho_{\mathrm{NFW}*} = \alpha_{\mathrm{NFW}}\, r_{\epsilon*}\left(1 + \alpha_{\mathrm{NFW}}r_{\epsilon*}\right)^2 \,, \tag{15b}$$

which, taking into account Eq. (15a), can be written as

$$\rho_{\mathrm{NFW}*} = \frac{\alpha_{\mathrm{NFW}}\, r_{\epsilon*}\left(1 + \alpha_{\mathrm{NFW}}\, r_{\epsilon*}\right)^2}{\left(1 + r_{\epsilon*}^2\right)^8} \,. \tag{15c}$$

Equation (15c) indicates the (normalized) density $\rho_{\mathrm{NFW}*}$ that is required for a correct matching between the soliton and NFW profiles, for given values of α_{NFW} and $r_{\epsilon*}$.

However, one can see that the continuity constraint (Eq. (15c)) actually shows a hidden degeneracy: once the values of α_{NFW} and $\rho_{\mathrm{NFW}*}$ are fixed, there can be up to two solutions for the matching radius $r_{\epsilon*}$. This is a direct consequence of the fact that the crossing of the density profiles (Eqs. (10) and (12)) can occur at most at two different points, as illustrated in the left-hand panel of Fig. 2, which shows normalized density profiles for $\alpha_{NFW} = 0.1$ and different values of the normalized density $\rho_{\mathrm{NFW}*}$.

Fig. 2 (left panel) also shows that there exists a maximum value of $\rho_{\mathrm{NFW}*}$ beyond which the profiles do not cross each other. This fact can be understood in terms of Eq. (15c), which we evaluate for different values of α_{NFW} in the right-hand panel of Fig. 2. Here we can see that, for each α_{NFW}, there is always a maximum value of $\rho_{\mathrm{NFW}*}$, as a function of the matching radius $r_{\epsilon*}$, that corresponds to the case in which the soliton and NFW density profiles barely touch, as seen in the left-hand panel of Fig. 2.

To avoid the hidden degeneracy, and to select always a combined profile with an interior soliton shape, we will choose those cases for which $r_{\epsilon*} \geq r_{\epsilon*,\mathrm{max}}$, where $r_{\epsilon*,\mathrm{max}}$ is the matching radius corresponding to the maximum value of $\rho_{\mathrm{NFW}*}$. A straightforward calculation from Eq. (15c) shows that $r_{\epsilon*,\mathrm{max}}$ is a root of the cubic equation

$$13\alpha_{\mathrm{NFW}}r_{\epsilon*,\mathrm{max}}^3 + 15r_{\epsilon*,\mathrm{max}}^2 - 3\alpha_{\mathrm{NFW}}r_{\epsilon*,\mathrm{max}} = 1 \,. \tag{16}$$
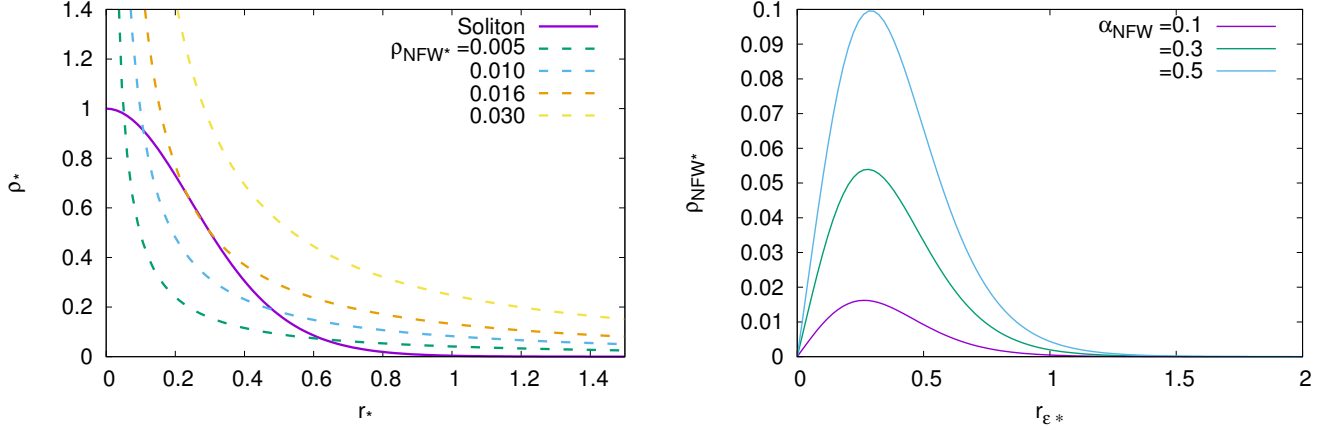
Figure 2. Determination of the matching radius, $r_{\epsilon*}$, for the density profile in Eq. (14). (Left) Normalized soliton and NFW density profiles showing there are at most two possible for $r_{\epsilon*}$ for each combination of α_{NFW} (set to 0.1 only for illustrative purposes) and $\rho_{\mathrm{NFW}*}$. (Right) NFW normalizing density factor $\rho_{\mathrm{NFW}*}$ as a function of $r_{\epsilon*}$ for different values of α_{NFW}, as obtained from Eq. (15c).

Although there is a general solution to this equation, it can be shown that the limits for small and large values of α_{NFW} are

$$\lim_{\alpha_{\mathrm{NFW}} \to 0} r_{\epsilon*,\mathrm{max}} = \left(1/\sqrt{15}\right), \tag{17a}$$

$$\lim_{\alpha_{\mathrm{NFW}} \to \infty} r_{\epsilon*,\mathrm{max}} = \left(\sqrt{3/13}\right). \tag{17b}$$

This means that in absolute terms the maximum value of $\rho_{\mathrm{NFW}*}$ must be located in the range $0.25 < r_{\epsilon*,\mathrm{max}} < 0.48$, which is in agreement with the values observed in the right-hand panel of Fig. 2. After all this, it is possible to reduce the number of free parameters that describe the combined profile (13) to only four: ρ_s, r_s, r_ϵ and α_{NFW}. By means of these parameters and the constraints discussed above, the other parameters are fully specified.

Notice that our chosen normalization is such that the physical parameters in the NFW profile (Eq. (12)) are given in terms of those in the soliton profile (Eq. (10)). This means, for instance, that $\rho_{\mathrm{NFW}*} > 1$ ($\rho_{\mathrm{NFW}*} < 1$) is equivalent to $\rho_{\mathrm{NFW}} > \rho_s$ ($\rho_{\mathrm{NFW}} < \rho_s$), whatever the physical value of ρ_s is. Likewise, we find that $\alpha_{\mathrm{NFW}} < 1$ ($\alpha_{\mathrm{NFW}} > 1$) corresponds to $r_{\mathrm{NFW}} > r_s$ ($r_{\mathrm{NFW}} < r_s$), even if the physical value of r_s is not known beforehand. The same will apply for the matching radius, since $r_{\epsilon*} > 1$ ($r_{\epsilon*} < 1$) means that matching occurs beyond the soliton radius and then $r_\epsilon > r_s$ (before the soliton radius, and then $r_\epsilon < r_s$).

It must be noticed also that the prescription above for the matching of the density profiles in Eq. (13) means that the NFW part is always subjected to the presence of the central soliton. For instance, the NFW profile can be diluted away if the matching radius $r_{\epsilon*} \to \infty$ (which also means that $\rho_{\mathrm{NFW}*} \to 0$), and then the density profile becomes the soliton one alone, $\rho(r) \simeq \rho_{\mathrm{sol}}(r)$. On the other hand, if $r_s \to 0$ the central soliton becomes small but much more massive and more dense, because of the scaling symmetry shown in Eq. (1), so that it dominates the matter contents over that of the NFW profile. The conclusion here is that under our parametrization the density profile (Eq. (13)) can become the soliton profile only if $r_{\epsilon*} \to \infty$, but it is not possible to do the same for the

NFW part; in this sense, the complete profile (Eq. (13)) should always be seen as that of a central soliton with a subdominant NFW tail.

We want to stress that the complete profile (13) should not be confused with the so-called cored NFW profile that exists already in the literature. The latter is of the form $\rho(r) \sim (r/r_s)^{-\beta}/(1+r/r_s)^2$ with $0 < \beta < 1$, and whose lensing properties have been analyzed in Wyithe et al. (2001). A comparison of the lenses produced by the cored-NFW profile and the WaveDM one is beyond the purpose of this work, as our primary intention is to constrain the parameters in Eq. (13) and to obtain from them credible bounds on the mass of the boson particles.

As a final note, we emphasize the convenience of the chosen parametrization in terms of the soliton characteristic quantities, as the soliton and NFW parameters must follow well defined scaling constraints that are intrinsic to the WaveDM. These scaling properties will then be already explicit in the complete profile (13) when making a comparison of the model with lensing data.

2.3. Gravitational Lensing

To obtain the lensing properties of the combined profile given by Eq. (13), we follow the recipe described in Sec. 2.1. We first need to compute the projected surface mass density (Eq. (3b)). Because of the presence of the step functions in Eq. (13), the integral in Eq. (3b) naturally separates as

$$\Sigma_*(\theta_*, r_{\epsilon*}, \alpha_{\mathrm{NFW}}) = 2 \begin{cases} \displaystyle\int_0^{\sqrt{r_{\epsilon*}^2 - \theta_*^2}} \frac{dz}{(1+\hat{r}^2)^8} + \frac{r_{\epsilon*}(1+\alpha_{\mathrm{NFW}}\,r_{\epsilon*})^2}{(1+r_{\epsilon*}^2)^8} \displaystyle\int_{\sqrt{r_{\epsilon*}^2 - \theta_*^2}}^{\infty} \frac{dz}{\hat{r}(1+\alpha_{\mathrm{NFW}}\,\hat{r})^2}\,, & \theta_* < r_{\epsilon*}\,, \\[20pt] \displaystyle\frac{r_{\epsilon*}(1+\alpha_{\mathrm{NFW}}\,r_{\epsilon*})^2}{(1+r_{\epsilon*}^2)^8} \displaystyle\int_0^{\infty} \frac{dz}{\hat{r}(1+\alpha_{\mathrm{NFW}}\,\hat{r})^2}\,, & \theta_* \geq r_{\epsilon*}\,. \end{cases} \tag{18}$$

It should be understood that the integrals in Eq. (18) are done along the line of sight. Notice also that we are following our convention in Sec. 2 for normalized quantities, namely $\Sigma_* = \Sigma/(\rho_s r_s)$, $\theta_* = \xi/r_s$ and $z = \sqrt{r_*^2 - \theta_*^2}$. The analytical expressions for the integrals in Eq. (18) can be found in appendix A.

Equation (18) shows that the projected surface mass density only depends upon the characteristic radii and densities of the combined density profile (Eq. (13)). For instance, if we keep r_s fixed, it can be shown that

$$\lim_{r_{\epsilon*} \to \infty} \Sigma_*(\theta_*, r_{\epsilon*}, \alpha_{\mathrm{NFW}}) = 0.658\,(1 + \theta_*^2)^{-15/2}\,, \tag{19}$$

a result that is obtained from the first branch in Eq. (18). Notice that Eq. (19) is exactly the result for the soliton profile (Eq. (10)) alone. Also, as we have mentioned before, it is not possible to recover the standard result of the surface density for the NFW profile by letting $r_s \to 0$, as in this case the matter content is still dominated by the central soliton.

Going back to the complete profile (Eq. (13)), we start with the calculation of the critical value λ_{crit} from the analytical formula in Eq. (8). The (total) projected surface mass density for the special value $\theta_* = 0$ is obtained from the first branch, in Eq. (18), as

$$\Sigma_*(0, r_{\epsilon*}, \alpha_{\mathrm{NFW}}) = 2 \left[\int_0^{r_{\epsilon*}} \frac{dz}{(1+z^2)^8} + \frac{r_{\epsilon*}(1+\alpha_{\mathrm{NFW}}\,r_{\epsilon*})^2}{(1+r_{\epsilon*}^2)^8} \int_{r_{\epsilon*}}^{\infty} \frac{dz}{z(1+\alpha_{\mathrm{NFW}}\,z)^2} \right]\,,$$
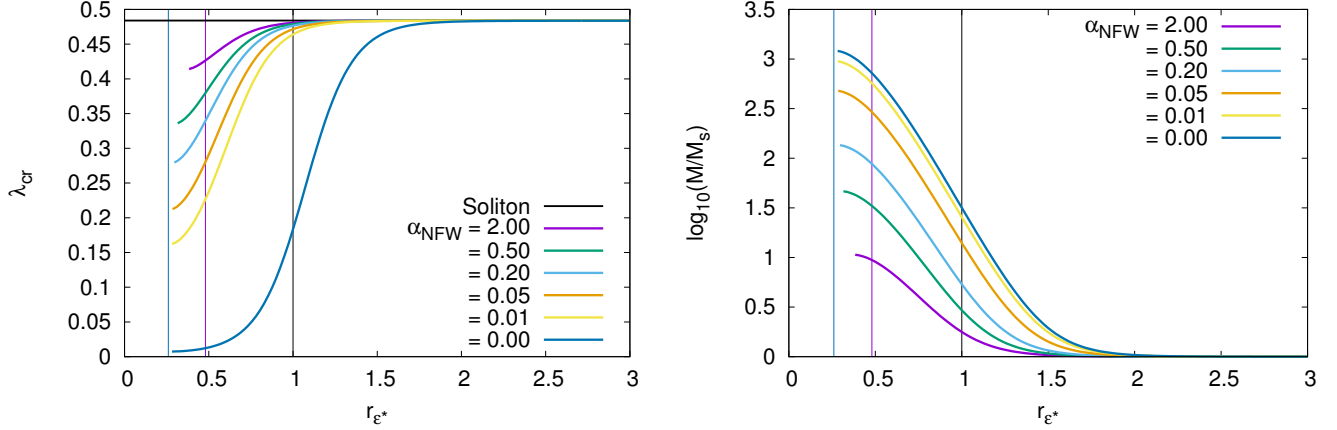
Figure 3. Critical value $\lambda_{\rm cr}$ to produce strong lensing (left). Total lens mass M normalized by the soliton mass M_s (right). Both as a function of the matching radius $r_{\epsilon*}$ for different density profiles characterized by $\alpha_{\rm NFW}$. The pure soliton case, $\lambda_{\rm cr} \simeq 0.48$ is recovered asymptotically in the limit $r_{\epsilon*} \to \infty$. Vertical lines marks the position of the normalized soliton radius at $r_{\epsilon*} = 1, 0.5, 0.25$ (gray,purple and blue). The presence of the NFW part ease the formation of multiple images. See the text for more details.

which indicates, together with Eq. (7), that the critical value $\lambda_{\rm cr}$ of the combined profile (Eq. (13)) is a function of $r_{\epsilon*}$ and $\alpha_{\rm NFW}$, and its behavior for different combinations of these parameters is shown in the left panel of Fig. 3. Notice that we have taken into account the constraint $r_{\epsilon*} \geq r_{\epsilon*,\rm max}$, see Eq. (16). Moreover, it can be seen that the lowest value of $\lambda_{\rm cr}$, for any given value of $\alpha_{\rm NFW}$, is indeed attained at $r_{\epsilon*,\rm max}$ as indicated by the vertical lines with the corresponding colors. Not surprisingly, the addition of the NFW outer part helps the soliton profile to achieve small values of $\lambda_{\rm crit}$, which in turn eases the accomplishment of the inequality in Eq. (9). In particular, Fig. 3 shows that $\lambda_{\rm crit} \to 0$ as $\alpha_{\rm NFW} \to 0$, which means that the combined profile (13) will be able to produce a lensing signal for any non-trivial combination of its parameters ρ_s and r_s.

In the case of the combined profile the total mass $M(r)$ inside a sphere of any given radius $r > r_\epsilon$ is simply given by the integral

$$\frac{M(r)}{10^{13} M_\odot} = 3 \left(\frac{m_a}{10^{-22}\,{\rm eV}}\right)^{-2} \left(\frac{r_s}{\rm pc}\right)^{-1} \times \left[\int_0^{r_{\epsilon*}} \frac{dx\,x^2}{(1+x^2)^8} + \frac{r_{\epsilon*}(1+\alpha_{\rm NFW}\,r_{\epsilon*})^2}{(1+r_{\epsilon*}^2)^8} \int_{r_{\epsilon*}}^{r_*} \frac{dx\,x}{(1+\alpha_{\rm NFW}\,x)^2}\right]. \quad (20)$$

In the general case the total mass diverges as $r \to \infty$, whereas for the soliton profile only (which requires $r_{\epsilon*} \to \infty$) we simply obtain that its total mass M_s is (Guzman & Urena-Lopez 2004; Marsh & Pop 2015; Chen et al. 2016)

$$\frac{M_s}{10^{11} M_\odot} = 7.7 \left(\frac{m_a}{10^{-22}\,{\rm eV}}\right)^{-2} \left(\frac{r_s}{\rm pc}\right)^{-1}. \quad (21)$$

In general, we expect from Eq. (20) the total mass in the combined profile to be larger than the soliton alone, that is $M(r) \geq M_s$. However, the value of the total mass M will depend on the upper limit of integration r_*, and the largest values for any given r_* will be obtained for the case where $\alpha_{\rm NFW} \to 0$, similar to the case of the critical value $\lambda_{\rm crit}$. The aforementioned general behaviour of

the total mass M as a function of the free parameters $r_{\epsilon*}$ and α_{NFW} is shown in the right-hand panel of Fig. 3. For the numerical examples we considered the upper limit of integration $r_* = 20$, for which we then see that that the difference between M and M_s can be as large as three orders of magnitude in the case $\alpha_{\mathrm{NFW}} = 0$. In other words, the total mass in the WaveDM profile is intrinsically attached to that of its soliton, and then the latter should be large enough if we are going to get the right mass scales in galaxies. This is a non-trivial property, as it shows that any non-zero value of the parameter α_{NFW} could point out the existence of a soliton core with a non-negligible mass contribution to the lens system (see for instance Fig. (5) below).

3. DATA ANALYSIS

In this section we will use our theoretical results to infer information about the WaveDM profile from observations of specific lens systems. We recall from Sec. 2.2 that there are four free parameters that are needed to describe the lensing properties of the combined density profile, Eq. (13). However, the lens equation, discussed in Sec. 2.3, is not explicitly dependent on two of them, namely ρ_s and r_s, and depends only on the free parameters of the NFW outer profile $r_{\epsilon*}$ and α_{NFW}. Therefore we could use the right-hand side of the lens equation (Eq. (4)) to put constraints on the surface density through the combination of parameters $\rho_s r_s$ – see also the discussion in Sec. 2.1.

However, the special properties of the WaveDM profile, as represented by Eq. (1), suggest that the lens equation could be written in a more convenient form. Using the fact that the (normalized) angular Einstein radius is $\theta_{*E} = R_E/r_s$, Eq. (6) can be re-cast in the form

$$m_{a22}^{-2}\,\theta_{*E}\,m_*(\theta_{*E}, \alpha_{\mathrm{NFW}}, r_{\epsilon*}) = \frac{1}{2.4}\frac{d_{OS}}{d_{OL}d_{LS}}\frac{h}{0.57}\left(\frac{R_E}{\mathrm{kpc}}\right)^3, \tag{22}$$

where we have set $m_{a22} \equiv m_a/10^{-22}\mathrm{eV}$. Equation (22) then defines a different observable, which results solely from the combination of the distances involved in the measurement of the lens system, so that we can put constraints directly on the boson mass m_a rather than on the energy density ρ_s, but in any case in combination with the rest of parameters, namely θ_{*E}, α_{NFW}, and $r_{\epsilon*}$.

In general, we expect that, given the data from a single galaxy, there will always be a region in the parameter space that will satisfy Eq. (22). Thus, for a given sample of galaxies, we could in principle determine the range of possible values of m_a that is consistent with the observed data. However, we must recall that the boson mass m_a is a fundamental physical parameter of the model which in principle should have a unique value. This means that the boson mass should be treated differently from other parameters in the model and should *not* be given the freedom to vary from galaxy to galaxy.

Our proposal, therefore, is to study the lensing properties of the WaveDM profile by fixing the value of the boson mass and finding, via statistical analysis, the best-fit values of the remaining free parameters θ_{*E}, α_{NFW} and $r_{\epsilon*}$. As we are interested in the properties of the WaveDM profile alone, we select a particular sub-sample of early-type galaxies, we will focus in particular on those lensed systems in which the galaxy is known to have a relatively high dark matter fraction. We have set a threshold at a fraction of luminous matter of 50% or less, i.e. reducing as much as possible its effects. Also, due to the consideration of the lens model to be comprised of a single galaxy, samples where contribution of more than one component is known were excluded, e.g. MG 2016+112 (see Nair & Garrett 1996).

Name	$f_{*,\mathrm{Ein}}^{Salp}$	z_{lens}	z_{source}	$d_{OS}/(d_{OL}d_{LS})$	R_E (kpc)
SLACS					
J0008-0004	0.50 ± 0.16	0.44	1.192	6.6855	6.7965
J0935-0003	0.35 ± 0.05	0.347	0.467	18.2172	4.4063
J0946+1006	0.46 ± 0.13	0.222	0.609	9.7613	5.0934
J1143-0144	0.46 ± 0.10	0.106	0.402	14.9617	3.3683
J1306+0600	0.47 ± 0.08	0.173	0.472	11.7208	4.0050
J1318-0313	0.42 ± 0.08	0.24	1.3	7.2634	6.1840
LSD					
CFRS03.1077	0.46 ± 0.15	0.94	2.94	5.3188	10.0470
HST1417+5226	0.38 ± 0.11	0.81	3.40	4.7801	10.9360
SL2S					
J220329+020518	0.24 ± 0.06	0.40	2.150	5.4526	10.8130

Table 2. Selected galaxies from SLACS, LSD and SL2S. Columns correspond to: label within the SDSS catalog (Name), fraction of luminous matter within the Einstein radius ($f_{*,\mathrm{Ein}}^{Salp}$), redshift of the lens (z_{lens}) and the source (z_{source}), distance factor $d_{OS}/(d_{OL}d_{LS})$, and measured Einstein radius (R_E). Selection was based on the condition $f_{*,\mathrm{Ein}}^{Salp} \leq 0.5$.

We use strongly-lensed galactic-scale systems observed by the Sloan data from the Sloan Lens ACS (SLAC) survey, which is comprised of nearly 100 likely and confirmed lensed systems (see Gavazzi et al. 2007; Auger et al. 2009). Our criteria reduces this sample to a sub-sample of only 6 galaxies. We also include samples from Lens Structure and Dynamics(LSD), and Strong Lensing Legacy Survey(LS2S) (see e.g. Cao et al. 2015; Sonnenfeld et al. 2013a; Treu & Koopmans 2004), but after applying the same criteria we ended selecting only three galaxies. The names of the nine chosen galaxies are shown in Table 2 together with the values of their lens parameters.

Another advantage from the SLAC survey is it was analized previously with other, similar, scalar field dark matter models (see e.g. Gonzalez-Morales et al. 2013; Robles & Matos 2013). Due to the

lack of studies of lensing for this kind of model, we therefore adopt SLACS as the main sub-sample for consistency with previous studies and as a proof of concept for the possible use of the methodology.

The Einstein radius R_E is obtained using $R_E = D_{OL}\theta_E$, where θ_E is as given in (see Bolton et al. 2008a; Auger et al. 2009; Sonnenfeld et al. 2013b; Treu & Koopmans 2004). It has been shown that θ_E is nearly model independent and well constrained, and have been used before to determine the lens mass where only big asymmetric arcs in the images could produce some difference[4] (see e.g. Kochanek 1991; Kochanek et al. 2001; Trick et al. 2016; Tortora et al. 2018; Lyskova et al. 2018; Cardone et al. 2009); this justifies its direct use as a reliable observable in our analysis. The distances and R_E values in Table 2 are obtained considering a cosmology with matter-density parameter $\Omega_M = 0.3089$, vacuum energy-density parameter $\Omega_\Lambda = 0.6911$, and Hubble parameter $H_0 = 67.74\,\mathrm{km\,s^{-1}\,Mpc^{-1}}$ from the Planck Collaboration et al. (2015).

3.1. *Soliton core profile*

As a first case of study, let us consider the soliton core profile without the external NFW part. There are in this case only two free parameters: m_{a22} and θ_{*E}. In section 3.2 a Bayesian analysis will be carried out, taking into account the results from this section. The projected mass surface density given by Eq. (5a), with the help of Eq. (19), has in this case an analytical expression,

$$m_*(\theta_{*E}) = \frac{2}{13\lambda_{\mathrm{crit}}} \frac{(1+\theta_{*E}^2)^{13/2} - 1}{(1+\theta_{*E}^2)^{13/2}} , \tag{23}$$

where $\lambda_{\mathrm{crit}} \simeq 0.484$ is the critical value calculated from Eq. (8); see also Table 1. Notice that $m_*(0) = 0$, whereas its asymptotic limit is $m_*(\infty) = 2/(13\lambda_{\mathrm{crit}})$.

To obtain a basic understanding of the solutions that will be found for the physical parameters, we show in the left panel of Fig. 4 the expected behavior of the left-hand side of Eq. (22) as a function of the Einstein angle θ_{*E}. We also show, as the series of horizontal lines, the values of the right-hand side of Eq. (22) obtained from the observed data for the galaxies listed in Table 2.

Figure 4 shows that it will always be possible to identify a value of the Einstein angle θ_{*E} for which the left-hand and right-hand sides of Eq. (22) are in agreement, irrespective of the value of the boson mass – although as the boson mass increases the agreement occurs at increasingly large values of θ_{*E}. For the examples shown in Fig. 4, a boson mass of order $m_{a22} \simeq 0.02$ seems to fit well the SLACS galaxies listed in Table 2 – corresponding to an allowed range for the angular Einstein radius of $5 < \theta_{*E} < 10$. This latter range can also be translated into an allowed range for the soliton radius, and suggests that $r_s \sim$ kpc for the given example galaxies. However, note that it is always possible to find a solution that matches the left-hand and right-hand side of Eq. (22), for any given value of the boson mass m_a, by a suitably large choice of Einstein angle θ_{*E}, that is, by choosing $r_s \to 0$. We must recall that the latter condition means that the density profile is dominated by a very massive and compact soliton, but this can be in disagreement with other indications about the actual size of the dark matter halo in the lens galaxies.

To summarize, given that we have only one observable constraint, the most we can do is first to fix the value of the boson mass m_a and from this to obtain constraints on the remaining free parameters that are consistent with that boson mass. Specifically, by adopting a proposed value for the boson

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[4] In particular for SLACS, it is important to remark that they adjusted different mass models and found that the images where visually indistinguishable and the Einstein angles where the same within error. This is stated in section section 5.2 and Table 5 of Bolton et al. (2008a).
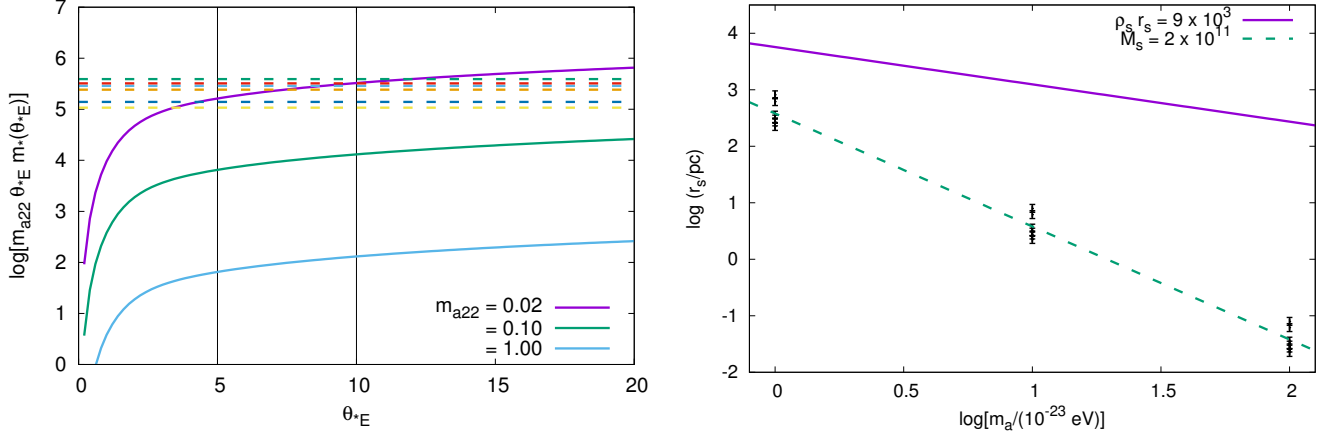
Figure 4. (Left) Left hand side of Eq. (22) as function of the normalized Einstein angle, $\theta_{*E} = R_E/r_s$, according to Eq. (23). The intersection with the dashed horizontal lines, r.h.s of Eq. (22) for each galaxy in our sample defines the value of θ_{*E}, i.e the scale radius. (Right) Fitted soliton radius, r_s, as a function of the boson mass m_a. The resultant r_s lies along the line of constant soliton mass M_s. This example is for $M_s \simeq 2 \times 10^{11} M_\odot$. See the text for more details.

mass m_a in Eqs. (22) and (23), we can obtain for each galaxy the corresponding best-fit value for θ_{*E}, and from that the best-fit value for r_s.

The results obtained for our selected sample of galaxies are shown in Table 3, and also plotted in the right panel of Fig. 4. The latter figure speaks for itself, and shows that the data points for all galaxies lie along the line with a constant soliton mass $M_s \simeq 10^{11} M_\odot$ (see Eq. (21)), and (as required) all lie below the line that represents the inequality, Eq. (9), for the galaxy in Table 2 (J0935-0003) with the most extreme value for the ratio of distances on the right-hand side of Eq. (22). The different values obtained for the characteristic radius r_s give an enclosed mass which corresponds closely to the values reported in Auger et al. (2009). Nevertheless, these models are found to be considerably too compact when the characteristic radius and corresponding enclosed mass are considered together. For example, galaxy J0008-0004 has a value for $M_{\rm Eins} = 3.1 \times 10^{11} M_\odot$ which is comparable with the value of $M_s = 3.4 \times 10^{11} M_\odot$ obtained using the best fit parameters of the soliton model. Notwithstanding that the soliton model gives an enclosed mass that is adequate and realistic, we think that the characteristic radius is most definitely not so. This is by taking into consideration the results from rotation curves where the effects of dark matter are expected to be at larger radii than the luminous part of the galaxy, and this contrasts with the values obtained for the soliton alone where the mean effective radius for galaxy J0008-0004 is observed to be $r_e \approx 9.6\,{\rm kpc}$, which is several orders of magnitude larger than the characteristic radius r_s obtained for any of the different boson masses presented in Table 3, including the samples from the other surveys. Therefore we think the soliton profile alone is actually not helping to explain the distribution of dark matter around the selected galaxies in a consistent way.

There are two valuable lessons from the above exercise. The first one is that the soliton core profile alone will always be able to fulfill the lensing constraints even without the consideration of the NFW contribution given the Einstein-radius as the only measurement to satisfy. This is not surprising, as the lensing equations can be solved even if we consider a point particle with the required total mass (which formally corresponds to the soliton core profile with $m_a \to \infty$). The second lesson is that

	$m_{a22} = 10$	$m_{a22} = 1$	$m_{a22} = 0.1$
Galaxy	$\log_{10}(r_s/\mathrm{pc})$		
SLACS			
J0008-0004	$-1.67^{+0.07}_{-0.06}$	$0.33^{+0.07}_{-0.06}$	$2.33^{+0.07}_{-0.06}$
J0935-0003	$-1.73^{+0.07}_{-0.06}$	$0.27^{+0.07}_{-0.06}$	$2.27^{+0.07}_{-0.06}$
J0946+1006	$-1.59^{+0.07}_{-0.06}$	$0.41^{+0.07}_{-0.06}$	$2.41^{+0.07}_{-0.06}$
J1143-0144	$-1.41^{+0.07}_{-0.06}$	$0.59^{+0.07}_{-0.06}$	$2.59^{+0.07}_{-0.06}$
J1306+0600	$-1.46^{+0.07}_{-0.06}$	$0.54^{+0.07}_{-0.06}$	$2.54^{+0.07}_{-0.06}$
J1318-0313	$-1.62^{+0.07}_{-0.06}$	$0.37^{+0.07}_{-0.06}$	$2.37^{+0.07}_{-0.06}$
LSD			
CFRS03.1077	$-1.58^{+0.17}_{-0.12}$	$0.42^{+0.17}_{-0.12}$	$2.42^{+0.17}_{-0.12}$
HST1417+5226	$-1.7^{+0.18}_{-0.12}$	$0.30^{+0.18}_{-0.12}$	$2.30^{+0.17}_{-0.12}$
SL2S			
J220329+020518	$-1.90^{+0.06}_{-0.05}$	$0.10^{+0.06}_{-0.05}$	$2.10^{+0.06}_{-0.05}$

Table 3. Soliton radius, r_s, obtained from the fits to each galaxy and for three different values of the boson mass m_a. Note that for SLACS samples, all combinations have a total soliton mass mass contained within the Einstein radius of $M_s \simeq 10^{11.5}\,M_\odot$. For the LSD and SL2S samples have a total mass of $M_s \simeq 10^{11.8}\,M_\odot$

even though the soliton profile may be adequate, formally speaking, to explain the lensing properties of the galaxies in Table 2, we will, in any case, have to consider the NFW outskirts in the complete profile (Eq. (13)) in order to satisfy other constraints that suggest that the boson mass should be in the range $m_{a22} = 1 - 10$ (see Hui et al. 2017).

3.2. *Complete profile*

Taking into account the above experience gained with the soliton profile alone, we will now consider the following procedure for the complete WaveDM profile. Since the total mass inside the Einstein radius is the only constraint provided by the lens systems, we will fix the values of the boson mass

m_a and soliton mass M_s. This approach is considered due to the results from the soliton analysis where the boson mass can satisfy different values for the Einstein radius, and other studies have found that m_a needs to be in certain range. For this, we take following values of the boson mass $m_{a22} = 0.1, 1, 10$, and for the soliton mass $\log_{10}(M_s/M_\odot) = 11.5, 10.5, 9.5, 8.5, 7.5$, from which we will calculate the values of r_s by means of Eq. (21), which allows to avoid a possible overcompensation of the soliton mass.

We will adopt a uniform prior for the other parameters over the following ranges: $\alpha_{\mathrm{NFW}} = [0 : 10]$, and $r_{\epsilon\star} = [r_{\epsilon\star,max} : 10]$. Here $r_{\epsilon\star,\max}$ is found from the cubic equation (16) for a given value of α_{NFW}, and the extreme values $\alpha_{\mathrm{NFW}} = 10$ and $r_{\epsilon\star} = 10$ are suggested by Figs. 3 and 4.

We will obtain the values of $\theta_{\ast E}$ by sampling from a Gaussian distribution, using the relation

$$\theta_{\ast E}(p) = \theta_{\ast Em} + \sigma\sqrt{2}\mathrm{erf}^{-1}(2p - 1), \quad p \in (0, 1). \tag{24}$$

The value for $\theta_{\ast Em} = R_E/r_s$ is the mean of the distribution using the observed value for the Einstein radius, $\sigma = 0.05 \ast \chi$ the error assigned and p is a random number sampled from a uniform distribution on the interval [0,1]. The inverse error function is approximated as described in Winitzki (2008). In this way, $\theta_{\ast E}$ will not enter into the fitting analysis as an extra variable.

Once the soliton mass is fixed, the rest of the mass that is included within the Einstein radius must be completed by the NFW profile. Because this requires a huge contribution, up to three orders of magnitude more, one sensible consideration is to set the total mass of the lens as composed by a simple representation of luminous matter and the selected model of dark matter. In a first approximation, the mass corresponding to the baryonic matter is simply a constant value modeled as a point particle. This is done from Eq. (2), and then the projected mass for the lens is composed of two parts,

$$m'(\theta) = m(\theta) + M', \tag{25}$$

where $m(\theta)$ is the mass from the dark matter component given by the profile in Eq. (13), and $M' = f_{\ast,Ein}M_{Ein}$ is the stellar mass contribution as described in Table 2. These values are normalized accordingly and then the dimensionless total mass m' is

$$m'_\ast(\theta_{\ast E}, \alpha_{NFW}, r_{\epsilon\ast}) = m_\ast(\theta_{\ast E}, \alpha_{NFW}, r_{\epsilon\ast}) + M'_\ast, \tag{26a}$$

where

$$M'_\ast = 0.3208 f_{\ast,Ein}\left(\frac{M_{Ein}}{M_s}\right). \tag{26b}$$

Eq. (26a) is combined with Eq. (22) to produce a modified observable which uses the soliton mass directly,

$$\frac{M_s}{M_\odot} m'_\ast(\theta_{\ast E}, \alpha_{\mathrm{NFW}}, r_{\epsilon\ast}) = \frac{7.7 \times 10^8}{2.4}\frac{d_{OS}}{d_{OL}d_{LS}}\frac{h}{0.57}\left(\frac{R_E}{\mathrm{kpc}}\right)^2. \tag{27}$$

3.3. General results

Using the samples mentioned in section 3 we will try to constrain the free parameters that will satisfy Eq. (27). As said before, the information available from the data is the Einstein radius, R_E, the lens distances (d_{OL}, d_{LS}, d_{OS}), the lens redshift and the source redshift. This information is used in the Multinest code (Feroz et al. 2009) to carry out a parameter search for each individual galaxy.
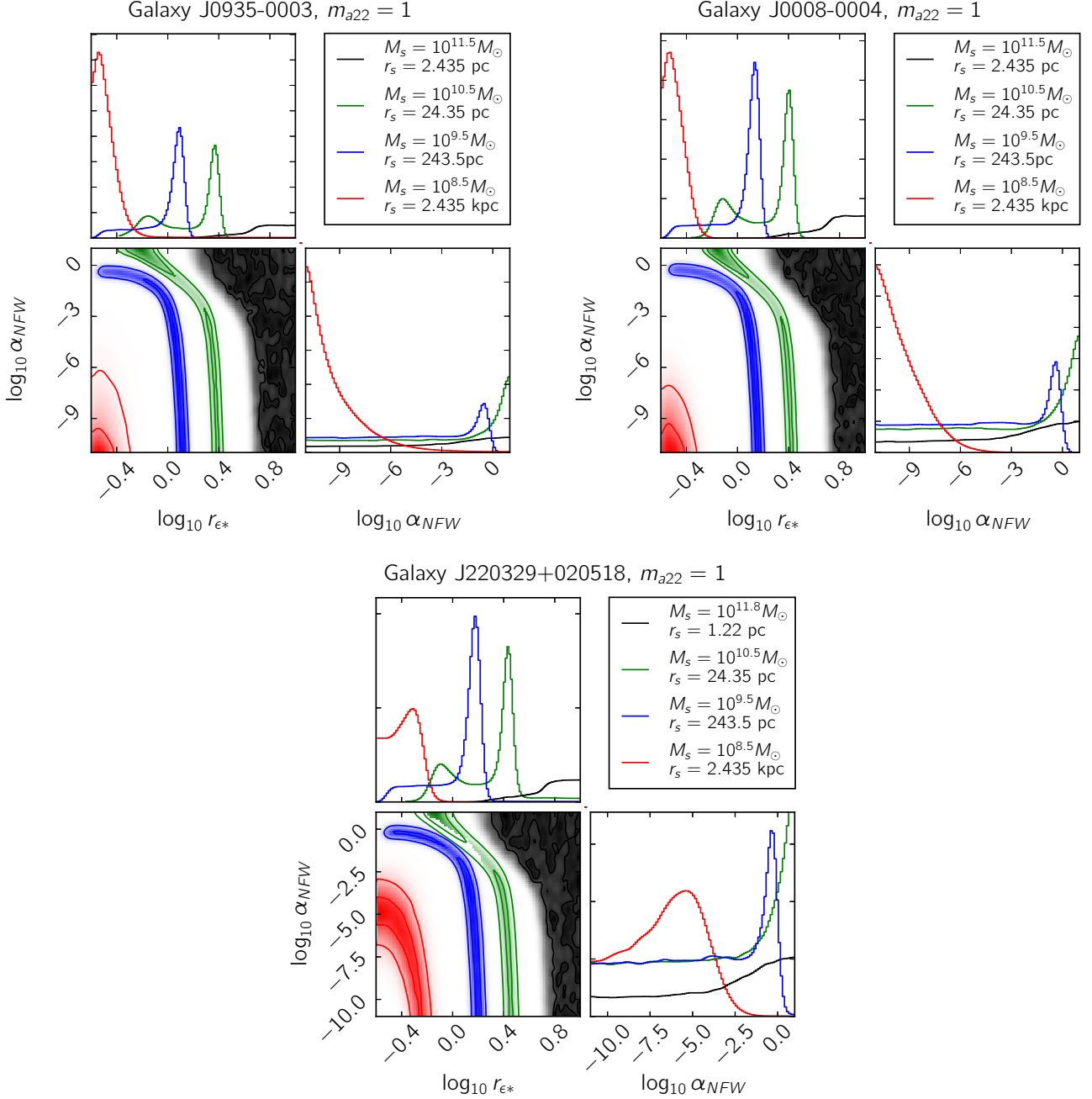
Figure 5. Posterior distribution for the parameters fitted to galaxies J0935-0003 (left), J0008-0004 (right) and J220329+020518(bottom); the contribution of the luminous matter are 35%, 50% and 24%, respectively, of the total reduced mass inside the Einstein radius.The colors indicate different combinations of the soliton mass, M_s, and scale radius, r_s, computed with a fixed (normalized) boson mass $m_{a22} = 1$.

We carried out the analysis on the nine galaxies of our sub-sample of the surveys. Nevertheless, they showed a similar behaviour for the range of values of f_*.

For brevity only representative results are shown, as in Fig. 5, for the individual cases of galaxies J0935-0003, J0008-0004 and J220329+020518; these cases include the contribution of the luminous matter to the total mass of the lens as in Eq. (27). For the purposes of clarity, in each figure we
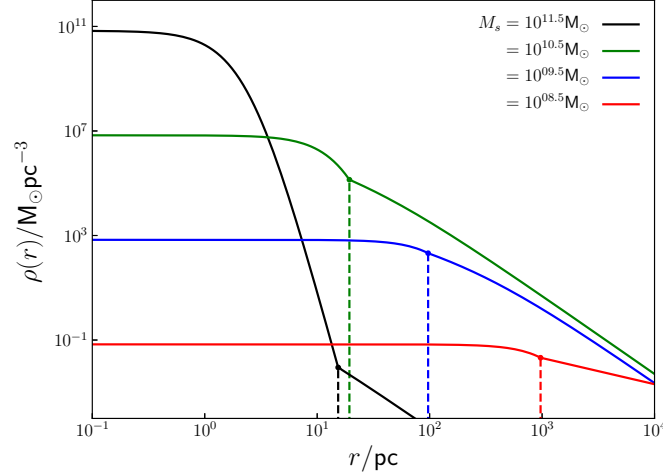
Figure 6. Examples of the density profiles from some of the selected configurations for Galaxy J0008-0004 obtained from the constraints in Fig. 5. The core region is indicated by the plateau in the curves. Note that the transition to a NFW-like profile happens at larger radii for smaller central density. Shown are cases with the value $\alpha_{NFW} = 1$, except in the case $M_s = 10^{8.5} M_\odot$ (red line) for which $\log(\alpha_{NFW}) = -7$.

indicate the radius r_s and total mass M_s of the soliton profile. We note that the free parameters $r_{\epsilon*}$ and α_{NFW} appear well constrained if the soliton mass cannot provide the total mass required by the lens system; in the examples shown, this happens if $M_s < 10^{11.5} M_\odot$. Secondly, the credible regions for the parameters in Fig. 5 are in agreement with the theoretical expectations discussed in Sec. 2.2: there is a minimum value for $r_{\epsilon*}$ due to the constraint imposed by Eq. (16), and a maximum value of α_{NFW} appears due to the maximal contribution of the NFW part of the profile to the total mass in the lens, see also the right panel in Fig. 3. Likewise, notice that as $\alpha_{\mathrm{NFW}} \to 0$ the value of the matching radius $r_{\epsilon*}$ is very well constrained, and this is easily understood from Eq. (20): it is $r_{\epsilon*}$ which determines alone the contribution of the NFW part of the profile to the total mass. Indeed, according to our parametrization in Sec. 2.2 the NFW part of the density profile, under the limit $\alpha_{\mathrm{NFW}} \to 0$, becomes

$$\rho_{\mathrm{NFW}}(r) = \frac{\rho_s \, r_{\epsilon*}}{(r/r_s)\left(1 + r_{\epsilon*}^2\right)^8} \, . \tag{28}$$

Apart from the presence of $r_{\epsilon*}$ (which in this case is bounded from above $r_{\epsilon*} \leq 1/\sqrt{15}$, see Eq. (17a)), we also see that the behavior $1/r$ is the only one that survives from the NFW functional form, and then our results indicate that the outermost behavior $1/r^3$ is left unconstrained.

Finally, observe that the value $\log_{10}(M_s/M_\odot) = 7.5$ is excluded because the soliton mass M_s is so small that the NFW part cannot compensate the required mass for the lens. Recall that there is a matching (continuity) condition for the density profile in which the NFW density ρ_{NFW} is always smaller than ρ_s, and this condition makes the NFW part of the profile unable to account for the total mass of the lens even in the limit $\alpha_{\mathrm{NFW}} \to 0$.

In summary, if the soliton is allowed to provide enough mass to fulfill the matter contribution in the lens, say $M_s \sim 10^{11.5} M_\odot$, the analysis will select large values for $r_{\epsilon*}$ so that the NFW tail contribution to the total matter is minimal, see Eq. (20). In contrast, if the soliton mass is not large enough, $M_s < 10^{11.5} M_\odot$, it is then possible to find appropriate pairs $(\alpha_{\mathrm{NFW}}, r_{\epsilon*})$ for the NFW part of the profile to provide the needed mass for the lens. In this respect, the striped credible regions in

Fig. 5 represent the degeneracy regions in the plane $(\alpha_{\mathrm{NFW}}, r_{\epsilon*})$ for the same mass contribution of the NFW tail to the lens system. Thus, we can see that distinct credible regions can be found for the NFW parameters if the soliton mass is $10^{8.5} < M_s/M_\odot < 10^{11.5}$, and that the constraints are in agreement with the semi-analytic analysis in Sec. 2.

Another quantity of interest is the resultant density profile of DM in the lens system. Fig. 6 shows examples of the density profiles inferred from the posteriors of galaxy J0008-0004 in Fig. 5 for a boson mass $m_{a22} = 1$. The soliton core is clearly seen in all curves, and so too is the transition to the NFW part of the profile. The corresponding matching radius r_ϵ, in full units, is selected to be at 15.36 pc, 19.34 pc, 96.94 pc and 969.4 pc for the soliton masses $10^{11.5} M_\odot$, $10^{10.5} M_\odot$, $10^{9.5} M_\odot$ and $10^{8.5} M_\odot$, respectively. Not surprisingly, the largest core corresponds to the configuration with the lowest soliton mass for which the matching radius is close to the lower bound suggested in Eq. (17a).

We also report in Fig. 7 the results obtained for the lens system J0008-0004 and J220329+020518, for larger or smaller values of the boson mass. For the boson mass $m_{a22} = 10$ we obtain good constraints on the NFW parameters, but the soliton core is very compact in all cases, although a constraint cannot be found if $M_s = 10^{6.5} M_\odot$. For the boson mass of $m_{a22} = 0.1$, we can only obtain well defined constraints on the NFW parameters when the soliton mass is $M_s = 10^{10.5} M_\odot$, but not for larger or smaller values. Low values of M_s imply values of the soliton radius r_s that are larger than the Einstein radius, and this kind of cases are unable to satisfy the lensing constraints. Hence, for a mass of $m_{a22} = 10$, the soliton is much more compact, and it is not by itself adequate to describe a galaxy. But given the fact that the parameters α_{NFW} and $r_{\epsilon*}$ are also well constrained we conclude that the lensing effect must be mostly attributed to the NFW part. This is not surprising, as we had already indicated in Sec. 2.3 that strong lensing could be achieved if $\alpha_{\mathrm{NFW}} \ll 1$. Moreover, a larger boson mass is also in better agreement with recent cosmological constraints (see Iršič et al. 2017) and with estimations based upon satellite galaxies of the Milky Way and Andromeda (see Ureña-López et al. 2017).

In contrast, we can see that the constraints become more diffuse if we consider a smaller boson mass of $m_{a22} = 0.1$, although there seems to be some preference for the case in which $M_s = 10^{10.5}$, that also corresponds to a larger soliton radius. This time the resultant configuration would be in agreement with those found in the statistical analysis carried out in González-Morales et al. (2016), which suggests that satellite galaxies put an upper bound on the boson mass that takes the form $m_{a22} < 0.4$.

4. CONCLUSIONS

We have studied the properties of the so-called WaveDM density profile, assuming that it comprises the total DM contribution in galaxies for which a gravitational lens has been detected and measured. In doing so we have adapted the standard lens equations to the particular features of the WaveDM, in that we took into account its soliton core together with its NFW envelope, which is the complete form suggested by numerical simulations of cosmological structure under the WaveDM hypothesis.

We then used the lens equations to make a comparison with actual observations of some lens systems that seem to be DM dominated, although we took into account their baryonic components in a simplified manner. In carrying out the statistical analysis we considered carefully the role of the
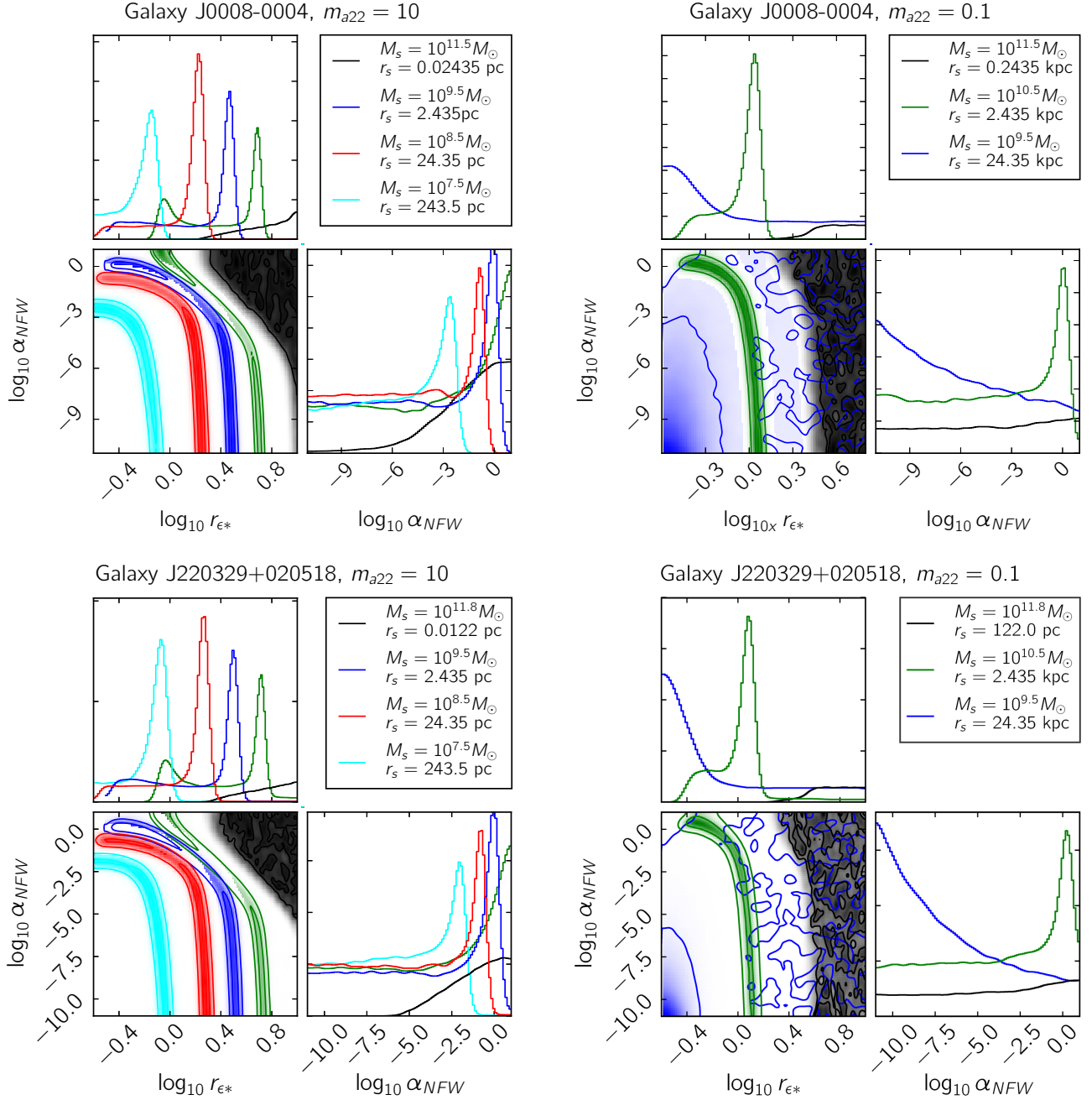
Figure 7. Same as figure 5 but for galaxies J0008-0004 and J220329+020518. Here we show different boson masses: $m_{a22} = 10$ (left panel) and $m_{a22} = 0.1$ (right panel). See the text for more details.

different free parameters of the WaveDM profile, and in particular the boson mass m_a which has to be regarded as a fundamental parameter that should not vary from one galaxy to another.

The overall procedure was then to fix the value of the boson mass and the total mass within the soliton core in the configuration. In consequence, the soliton radius was fixed and the only free parameters were those of the NFW part of the density profile. In general terms, for large or small values of the boson mass, our results indicate that the soliton structure, if it is as massive as $10^{11.5} M_\odot$, is able to fit the measured Einstein radius in the lens systems studied, although this also

requires the soliton structure to be extremely compact when compared to the measured scales of the lensing galaxies. This result then indicates that galaxies in general cannot be explained by the soliton structure alone.

Because of the above, we had to consider the complete WaveDM density profile and constrain the NFW free parameters. Generically, and so far for the cases we explored, our analyses suggest that the matching radius for the soliton and NFW parts of the profile is of the same order of magnitude as the soliton radius, $r_\epsilon \sim r_s$, which is in agreement with the expectation from numerical simulations (e.g. Schwabe et al. 2016; Veltmaat & Niemeyer 2016; Mocz et al. 2017). In addition, the second free parameter is in general bounded from above as $\alpha_{\mathrm{NFW}} < 1$, which just means that the characteristic NFW radius is larger than the soliton radius, $r_{\mathrm{NFW}} > r_s$. Moreover, our results also suggest that the case $\alpha_{\mathrm{NFW}} \to 0$ is also possible, which in turn means that the density profile decays as $\rho \sim r^{-1}$ at large radii.

On the other hand, for any given value of the boson mass, it was not possible to constrain the NFW parameters in the case where the soliton radius was larger than the Einstein radius, as in such cases the soliton mass is insufficient to produce the required lensing signal. Together with the aforementioned difficulty that the soliton should not provide the whole mass of the lens, we can summarize our results as $M_s/M_\odot < 10^{11.5}$ and $r_s < 6\,\mathrm{kpc}$. Given the similar masses and values of the Einstein radii in the selected sample of galaxies, these constraints can be taken as characteristic of the WaveDM model if the latter is considered to be the dark matter in them.

By means of Eq. (21), the above inequalities can be combined in the following lower bound on the boson mass $m_a > 10^{-24}\mathrm{eV}$. Notice that this lower bound is in agreement with previous constraints from cosmological and galactic scales, (see for instance Hlozek et al. 2015; Chen et al. 2016; González-Morales et al. 2016; Ureña-López et al. 2017). Although the lens systems we considered are not able to put strong bounds on the boson mass, they certainly indicate that most likely a complete WaveDM profile (i.e. comprising a soliton core + NFW tail) is necessary to account for all the diverse observations at galaxy scales.

As a final note, the lens systems studied here have a subdominant, although non-negligible, baryonic contribution. We expect to extend our analysis to a larger sample considering other surveys with a more detailed and specific description of the baryonic matter contained that could give us better constraints on the soliton features. This is ongoing work that will be presented elsewhere.

Fronteras 281; Fundación Marcos Moshinsky; and the Instituto Avanzado de Cosmología Collaboration.

APPENDIX

A. INTEGRAL SOLUTIONS

Some useful analytical solutions are given here for the integrals in Eq. (18). For the first branch $\theta_* < r_{\epsilon*}$ the formula for the first integral is

$$\int_0^{\sqrt{r_{\epsilon*}^2 - \xi_*^2}} \frac{dz}{[1 + \alpha_{\rm sol}^2 \hat{r}^2]^8} = \frac{\alpha_{\rm sol}^{-1}}{(1 + \alpha_{\rm sol}^2 \xi_*^2)^{15/2}} \int_0^x \cos^{14} u \, du \,, \quad \tan x = \alpha_{\rm sol} \left(\frac{r_{\epsilon*}^2 - \xi_*^2}{1 + \alpha_{\rm sol}^2 \xi_*^2} \right)^{1/2} \,, \tag{A1}$$

where

$$\int_0^x \cos^{14} u \, du = \frac{429}{2048} x + \frac{1001}{16384} \left[3\sin(2x) + \sin(4x) + \frac{1}{3}\sin(6x) + \frac{1}{11}\sin(8x) + \frac{1}{55}\sin(10x) \right.$$
$$\left. + \frac{1}{429}\sin(12x) + \frac{1}{7007}\sin(14x) \right] \,, \tag{A2}$$

whereas for the second integral we obtain

$$\int_{\sqrt{r_{\epsilon*}^2 - \xi_*^2}}^{\infty} \frac{dz}{\hat{r}\,(1 + \alpha_{\rm NFW}\,\hat{r})^2} = \begin{cases} \dfrac{1}{x^2 - 1}\left(1 - \dfrac{\sqrt{y^2 - x^2}}{1 + y} - \dfrac{2\operatorname{arctanh}}{\sqrt{1 - x^2}}\left[\dfrac{\sqrt{1 - x^2}}{1 + y + \sqrt{y^2 - x^2}}\right]\right) & x < 1\,, \\[2ex] \dfrac{1}{3}\left(1 - \dfrac{y+2}{y+1}\sqrt{\dfrac{y-1}{y+1}}\right) & x = 1\,, \,, \\[2ex] \dfrac{1}{x^2 - 1}\left(1 - \dfrac{\sqrt{y^2 - x^2}}{1 + y} - \dfrac{2\arctan}{\sqrt{x^2 - 1}}\left[\dfrac{\sqrt{x^2 - 1}}{1 + y + \sqrt{y^2 - x^2}}\right]\right) & x > 1\,. \end{cases} \tag{A3}$$

where $x = \alpha_{\rm NFW}\xi_*$ and $y = \alpha_{\rm NFW} r_{\epsilon*}$. By setting $y = x$, which is equivalent to $r_{\epsilon*} = \xi_*$, in Eq. (A3) we obtain the solution for the second branch in Eq. (18). For the case $\xi_* = 0$, which is used in Eq. (3), the integral result simply is

$$\int_{r_{\epsilon*}}^{\infty} \frac{dz}{z\,(1 + \alpha_{\rm NFW}\,z)^2} = \ln\frac{(1 + \alpha_{\rm NFW}\,r_{\epsilon*})}{\alpha_{\rm NFW} r_{\epsilon*}} - \frac{1}{(1 + \alpha_{\rm NFW}\,r_{\epsilon*})} \,. \tag{A4}$$

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