

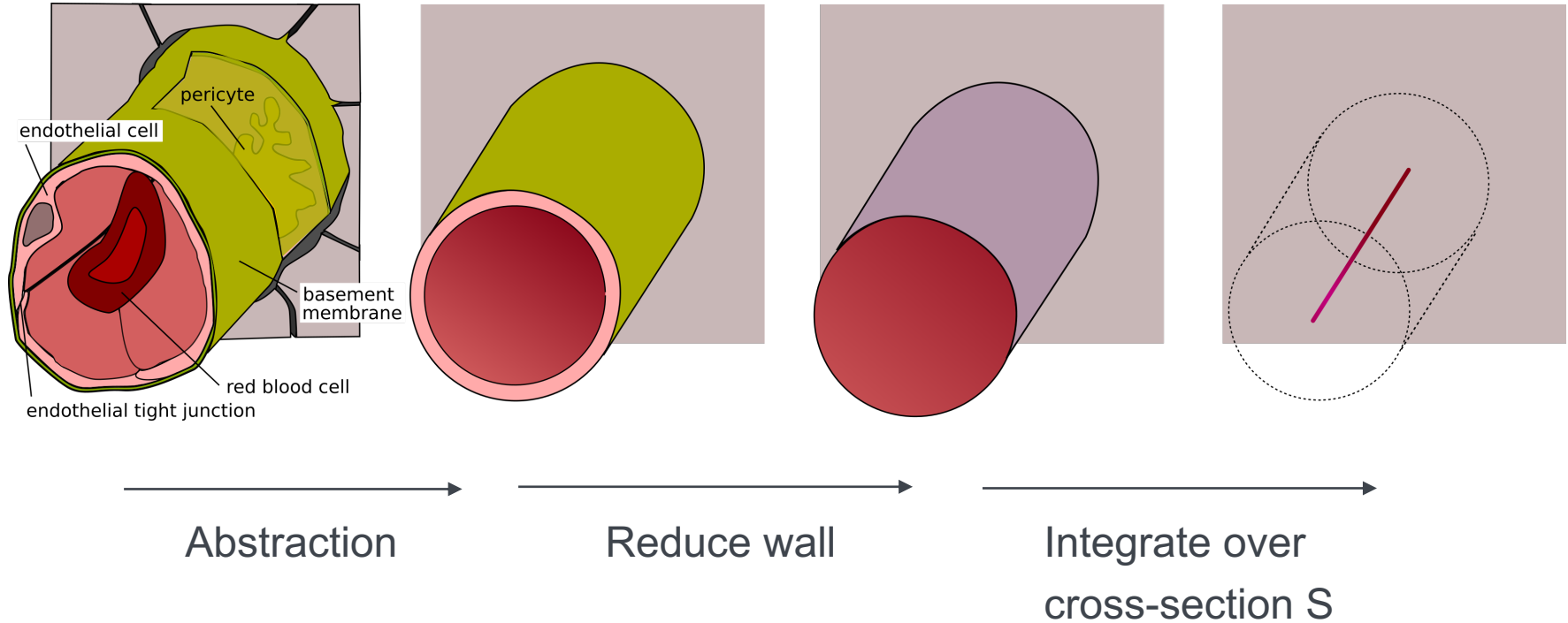
Chapter 4

Modeling flow in capillaries



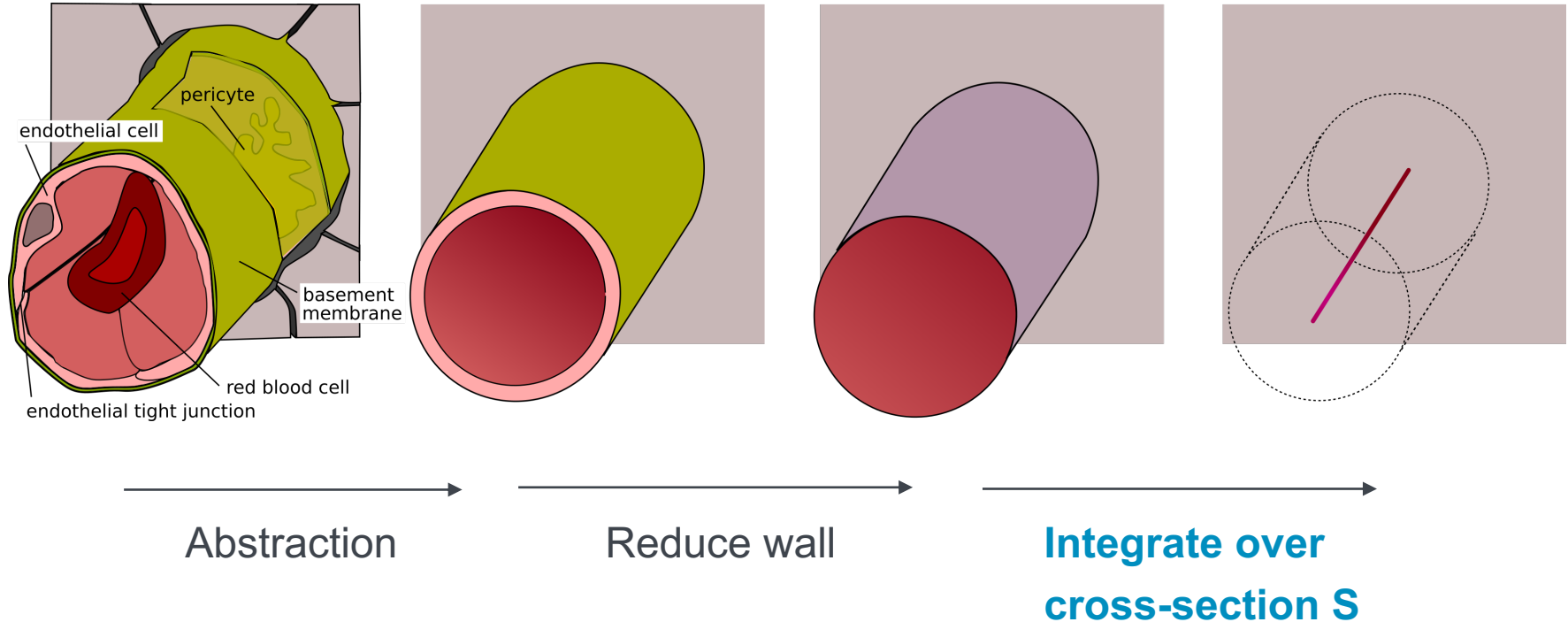
Modeling flow in capillaries

General idea



Modeling flow in capillaries

General idea



Modeling flow in capillaries

Assumptions

- Low Reynolds numbers ($Re \ll 1$)
 - Neglection of higher-order terms (quadratic in velocity)
 - Neglection of inertia terms (stationary flow)
- Rigid vessel wall

$$A(z, t) = |S(z, t)| = 2\pi R(z, t)$$

- Leaky vessel (fluid crosses the vessel wall)

$$v_r(z, R) := v_R(z)$$



Modeling flow in capillaries

Radially-symmetric Stokes equations

$$\frac{1}{\rho} \frac{\partial p}{\partial r} - \frac{\hat{\mu}_B}{\rho} \left(\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial v_r}{\partial r} \right) + \frac{\partial^2 v_r}{\partial z^2} \right) = 0$$

$$\frac{1}{\rho} \frac{\partial p}{\partial z} - \frac{\hat{\mu}_B}{\rho} \left(\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial v_z}{\partial r} \right) + \frac{\partial^2 v_z}{\partial z^2} \right) = 0$$

$$\frac{1}{r} \frac{\partial(rv_r)}{\partial r} + \frac{\partial v_z}{\partial z} = 0$$

mass balance equation



Modeling flow in capillaries

Mass balance equation – Integration over cross-section

$$\frac{1}{r} \frac{\partial(rv_r)}{\partial r} + \frac{\partial v_z}{\partial z} = 0,$$

$$A(z, t) = |S(z, t)| = 2\pi R(z, t)$$


$$\int_S \frac{1}{r} \frac{\partial(rv_r)}{\partial r} dA + \int_S \frac{\partial v_z}{\partial z} dA = 0$$



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Mass balance equation – Integration over cross-section

$$\boxed{\int_S \frac{1}{r} \frac{\partial(rv_r)}{\partial r} dA} + \int_S \frac{\partial v_z}{\partial z} dA = 0$$


$$\int_S \frac{1}{r} \frac{\partial(rv_r)}{\partial r} dA = \int_0^{2\pi} \int_0^R \frac{\partial(rv_r)}{\partial r} dr d\theta = 2\pi R v_r(z, R) := 2\pi R v_R(z).$$

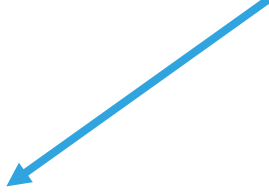
„Leaky vessel“



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Mass balance equation – Integration over cross-section

$$\int_S \frac{1}{r} \frac{\partial(rv_r)}{\partial r} dA + \int_S \frac{\partial v_z}{\partial z} dA = 0$$



$$\int_S \frac{\partial u_z}{\partial z} dA = \frac{\partial}{\partial z} \left(\int_S u_z dA \right) - \underbrace{2\pi R u_z(t, R)}_{=0} \frac{\partial R}{\partial z} = \frac{\partial Q}{\partial z}$$



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Mass balance equation – 1D model

$$\int_S \frac{1}{r} \frac{\partial(rv_r)}{\partial r} dA + \int_S \frac{\partial v_z}{\partial z} dA = 0$$


$$\frac{\partial Q(z)}{\partial z} = -2\pi R v_R(z)$$



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Momentum balance equation – 1D model

Same as for elastic tubes but without inertia / nonlinear term and constant A

$$\cancel{\frac{\partial Q}{\partial t}} + \cancel{\frac{\partial}{\partial z} \left(\cancel{\alpha} \frac{Q^2}{A} \right)} + \overset{A_0}{\frac{A}{\rho}} \frac{\partial p}{\partial z} + K_R \frac{Q}{A} = 0.$$

$$\longrightarrow \frac{A_0}{\rho} \frac{\partial p}{\partial z} + K_R \frac{Q}{A_0} = 0$$



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1D model

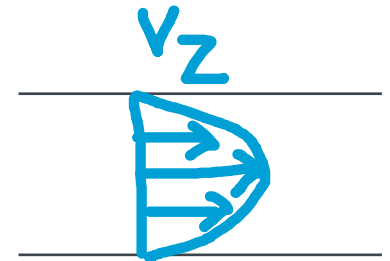
$$\frac{A_0}{\rho} \frac{\partial p}{\partial z} + K_R \frac{Q}{A_0} = 0$$

$$\frac{\partial Q(z)}{\partial z} = -2\pi R v_R(z)$$

→ insert momentum balance into mass balance + power-law velocity profile

$$-\frac{\partial}{\partial z} \left(\frac{A_0^2}{\rho K_R} \frac{\partial p}{\partial z} \right) = -2\pi R v_R(z)$$

$$-\frac{\partial}{\partial z} \left(\frac{\rho}{\mu_B} \frac{\pi R^4}{2(\gamma + 2)} \frac{\partial p}{\partial z} \right) = -\rho 2\pi R v_R(z)$$



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1D model

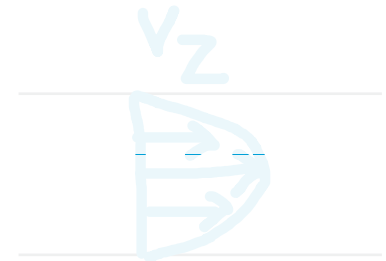
$$\frac{A_0}{\rho} \frac{\partial p}{\partial z} + K_R \frac{Q}{A_0} = 0$$

$$\frac{\partial Q(z)}{\partial z} = -2\pi R v_R(z)$$

→ insert momentum balance into mass balance + power-law velocity profile

$$-\frac{\partial}{\partial z} \left(\frac{A_0^2}{\rho K_R} \frac{\partial p}{\partial z} \right) = -2\pi R v_R(z)$$

$$-\frac{\partial}{\partial z} \left(\frac{\rho}{\mu_B} \frac{\pi R^4}{2(\gamma + 2)} \frac{\partial p}{\partial z} \right) = -\rho 2\pi R v_R(z)$$



Non-Newtonian behavior

Flow across vessel wall

→ need for closure models

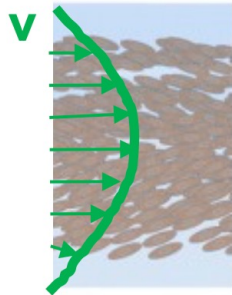


Modeling flow in capillaries

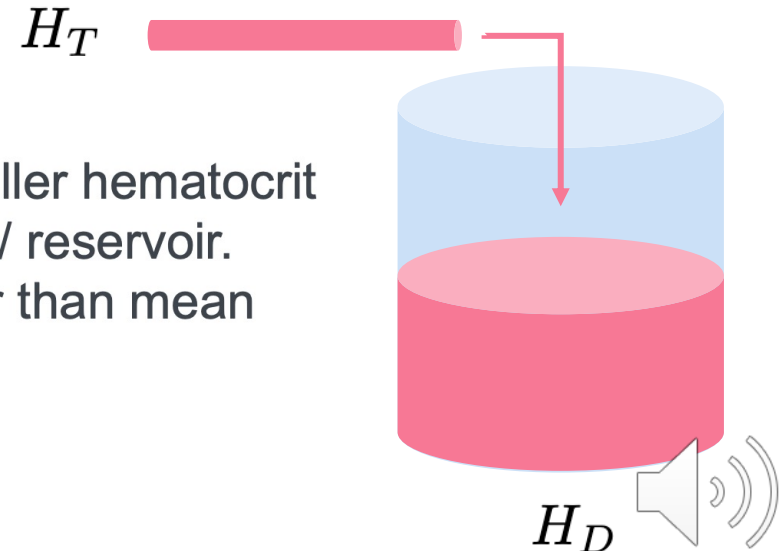
Non-Newtonian behavior

→ Fåhræus effect

$$\frac{H_T}{H_D} = H_D + (1 - H_D) (1 + 1.7e^{-0.415D} - 0.6e^{-0.011D})$$



Axial migration leads to smaller hematocrit than in feeding large vessel / reservoir.
Mean RBC velocity is higher than mean plasma velocity.



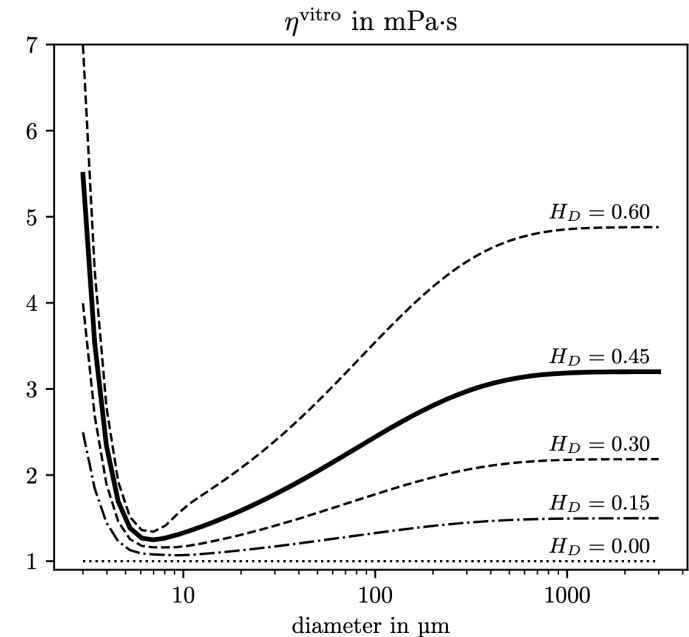
Modeling flow in capillaries

Non-Newtonian behavior

→ Empirical relationship for μ_B derived from experiments

$$\eta^{\text{vitro}} := \frac{\mu_B^{\text{vitro}}}{\mu_P} = 1 + (\eta_{0.45}^{\text{vitro}} - 1) \frac{(1 - H_D)^C - 1}{(1 - 0.45)^C - 1},$$
$$\eta_{0.45}^{\text{vitro}}(D) = 220e^{-1.3D} - 2.44e^{-0.06D^{0.645}} + 3.2,$$
$$C(D) = (0.8 + e^{-0.075D}) \left(\frac{1}{1 + 10^{-11}D^{12}} - 1 \right) + \frac{1}{1 + 10^{-11}D^{12}},$$

In-vitro: experiments with blood in glass tubes



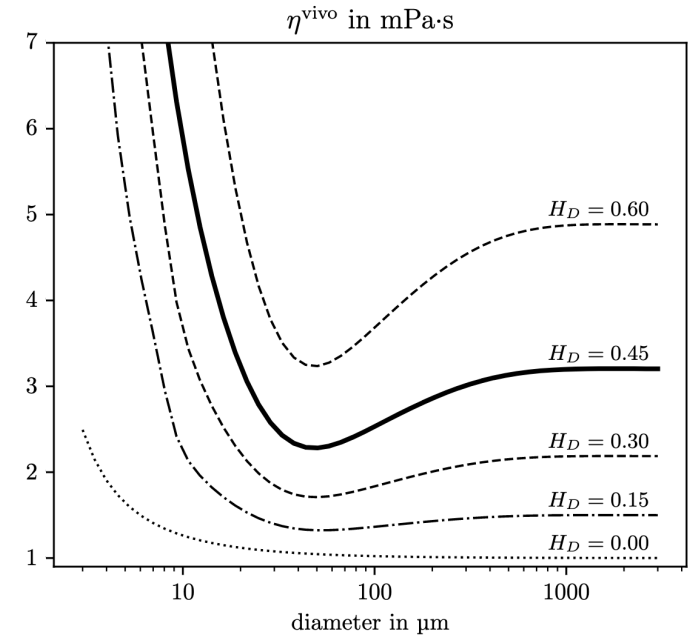
Modeling flow in capillaries

Non-Newtonian behavior

→ Empirical relationship for μ_B derived from experiments

$$\eta^{\text{vivo}} := \frac{\mu_B^{\text{vivo}}}{\mu_P} = \left[1 + (\eta_{0.45}^{\text{vivo}} - 1) \frac{(1 - H_D)^C - 1}{(1 - 0.45)^C - 1} \left(\frac{D}{D - 1.1} \right)^2 \right] \left(\frac{D}{D - 1.1} \right)^2,$$
$$\eta_{0.45}^{\text{vivo}}(D) = 6e^{-0.085D} - 2.44e^{-0.06D^{0.645}} + 3.2,$$
$$C(D) = (0.8 + e^{-0.075D}) \left(\frac{1}{1 + 10^{-11}D^{12}} - 1 \right) + \frac{1}{1 + 10^{-11}D^{12}},$$

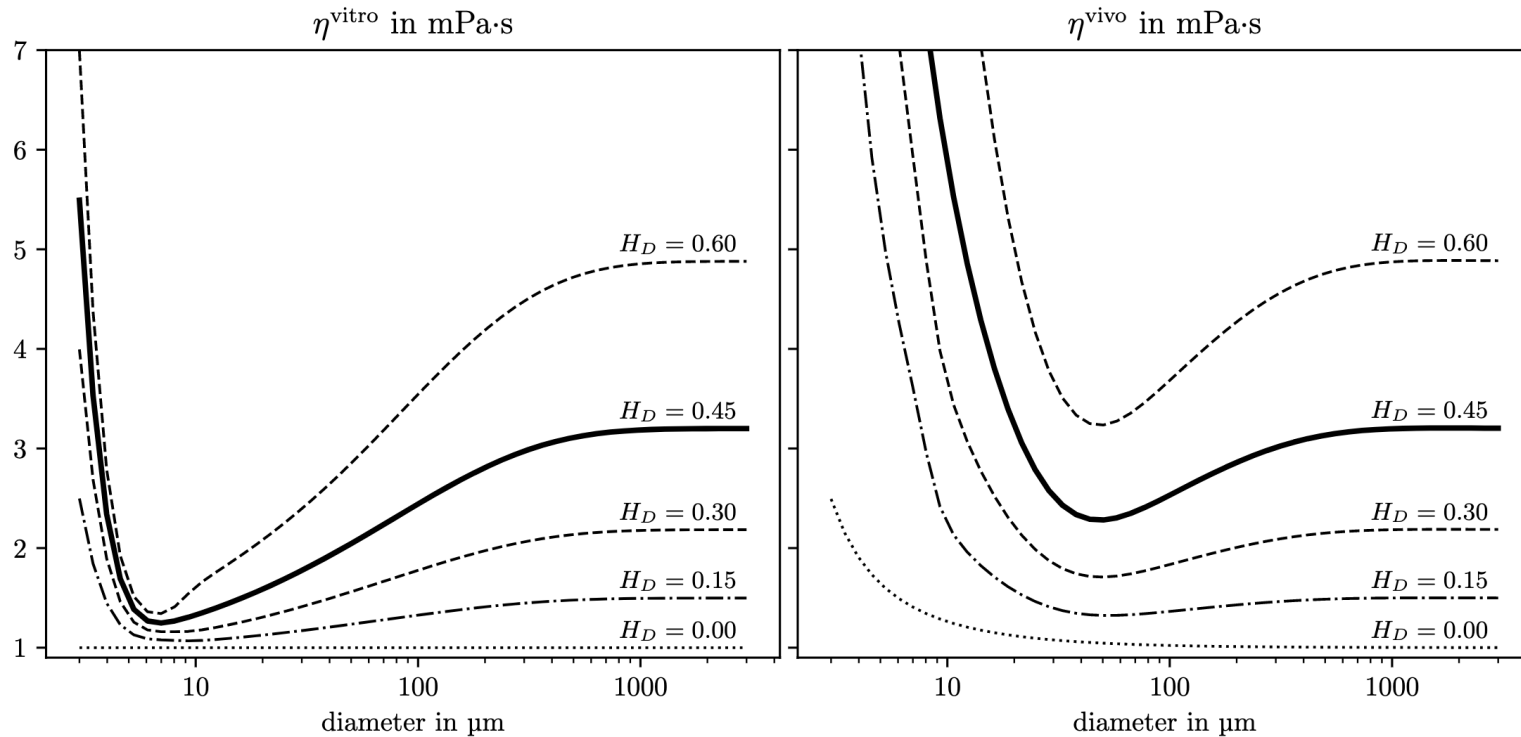
In-vivo: experiments with alive rats



Modeling flow in capillaries

Non-Newtonian behavior

→ Empirical relationship for μ_B derived from experiments



Modeling flow in capillaries

Transport across capillary wall

- Assumption: wall is porous
- Driving forces: hydrostatic pressure, osmotic pressure
 - Starling's law:

$$v_R(z) = L_p(p - P_{\text{ext}} - \sigma_w \Delta\Pi)$$

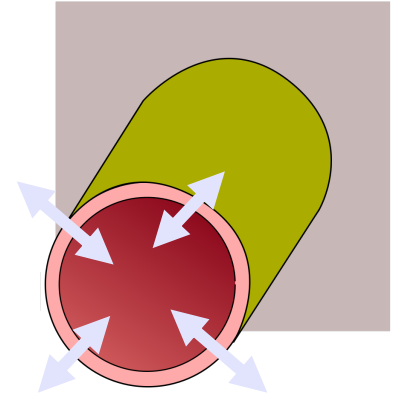
L_p Filtration coefficient (measurement data)

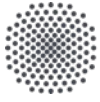
p Pressure in capillary (1d model)

P_{ext} Pressure outside wall (model?)

$\Delta\Pi$ Osmotic pressure difference (measurement data)

σ_w Osmotic reflection coefficient (measurement data)

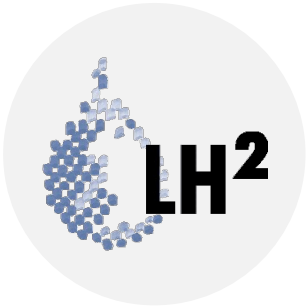




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Thank you!



<https://www.iws.uni-stuttgart.de/lh2/>

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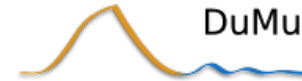
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Interface-Driven Multi-Field
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