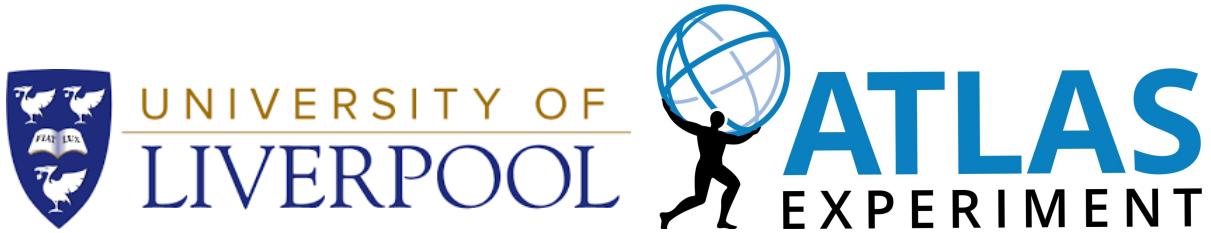


Search for Higgs boson pair-production in the $b\bar{b}\tau\tau$
final state using proton-proton collisions at $\sqrt{s} = 13$
TeV data with the ATLAS detector

Zhiyuan Li

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1 Introduction

2 Theory and Motivation

2.1 The Standard Model and the Higgs boson

2.2 Beyond the Standard Model

3 The ATLAS experiment at the Large Hadron Collider

3.1 The Large Hadron Collider

The Large Hadron Collider [1] is the world's largest and most powerful particle accelerator. It started in 2008 and retains its crucial role in the many accelerators at CERN and in the world. The main body of the collider consists of a ring tunnel of circumference 26.7 km, which lies beneath the French-Swiss border near Geneva, with superconducting magnets along the tunnel to keep the particle beam in direction and a large number of accelerating structures to boost the beam to the desired energy. Inside the tunnel, *bunches* of up to 1.15×10^{11} protons travelling at close to the speed of light in opposite direction are collided 40 million times per second at four cross points, around which are positioned four main detectors: ATLAS (A Toroidal LHC ApparatuS) [2], CMS [3] (Compact Muon Solenoid), ALICE (A Large Ion Collider Experiment) [4] and LHCb (b stands for beauty) [5].

3.1.1 The LHC Accelerator Complex

The proton beams colliding in the LHC are designed to carry energy of the order of a few TeV. To reach such high energy, a series of acceleration is required for the beams before entering the LHC ring. The protons are supplied by the injector chain Linac 2 — Proton Synchrotron Booster (PSB) — Proton Synchrotron (PS) — Super Proton Synchrotron (SPS), as shown in Figure 1.

The protons are produced by stripping off the electrons from hydrogen gas in an electric field. Linac 2, the first accelerator in the chain, is a linear accelerator which accelerates the protons to an energy of 50 MeV. They then enter the PSB, which accelerates the protons to 1.4 GeV, followed by the PS, where they reach 25 GeV. The series of radio frequency cavities in the PS splits the beam into discrete bunches of protons of 25 ns spacing. These bunches are then accelerated to 450 GeV in the SPS, from which they are finally injected into the beam pipes of the LHC. Each proton beam contains 2808 bunches, arranged in “trains” with 72 bunches each “carriage” of a gap of around 320 ns between each carriage. The beams are required to have well defined transverse and longitudinal *pittance*.

The beam pipes are kept at ultra-high vacuum, a vacuum thinner than interstellar void, maintained for 24 km of low-temperature section and 3 km of room-temperature section. For the low temperature section, the vacuum is achieved by pumping in 9000 m^3 of cryogenic gas, which later will be condensed and adhered to the surface of the beampipe. For the room temperature section, the vacuum is achieved by use of non-evaporable getter (NEG) that absorbs residue gas particles when heated. More residue is absorbed by an ion pumper.

The beam pipes are installed in the existing tunnel that was constructed between 1984 and 1989 for the Large Electron-Positron Collider (LEP), which lies between 45 m and 170

The CERN accelerator complex Complexe des accélérateurs du CERN

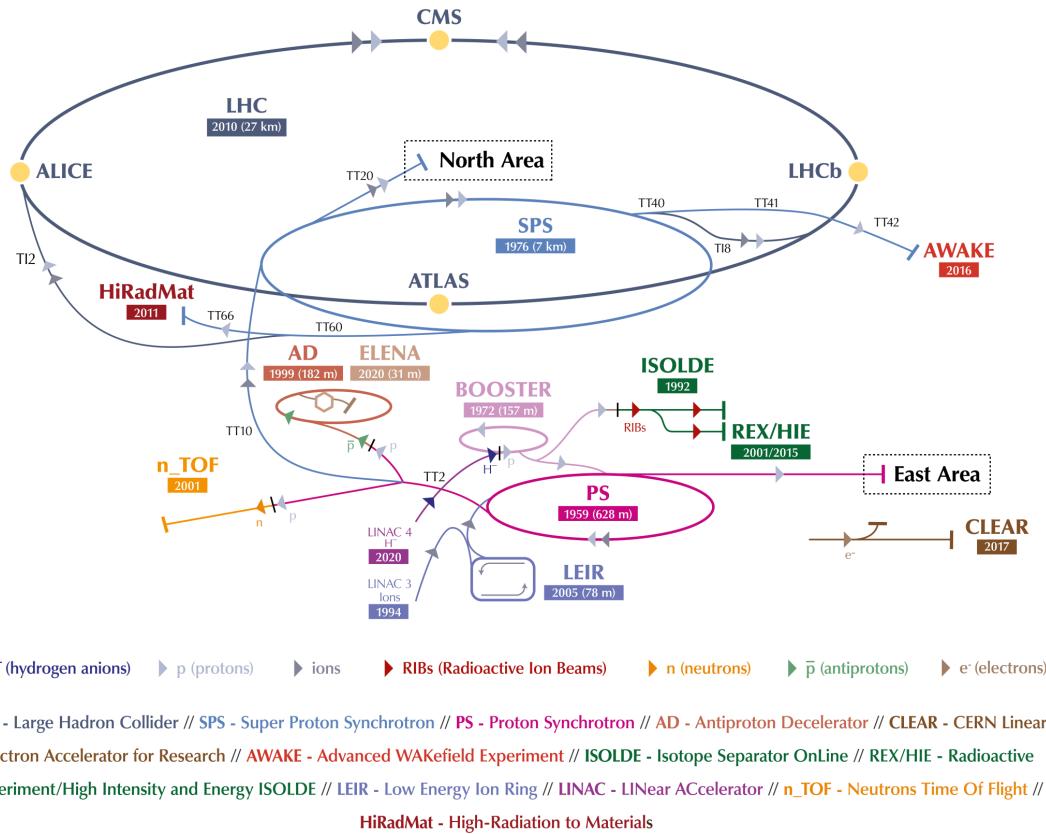


figure 1: The LHC is the last ring (dark blue line) in a complex chain of particle accelerators. The smaller machines are used in a chain to help boost the particles to their final energies and provide beams to a whole set of smaller experiments.

m below the surface on a plane inclined at 1.4% sloping towards the Léman lake. There are advantages and disadvantages for a proton-proton collider, compared to the electron-positron collider like LEP or proton-anti-proton collider. For the LHC, two rings are needed to accommodate the two counter-rotating beams, unlike particle-antiparticle colliders that can have both beams sharing the same phase space in a single ring, and therefore sharing the same magnets, same radio frequency cavities, etc. However, LHC is able to achieve very high collision energy which is not possible using electron-positron colliders, neither linear nor circular ones. This is due to the fact that in a circular electron collider, synchrotron radiation lost is proportional to the Lorentz factor $\gamma = E/m$ to the power of four, where E, m is the energy, mass of the particle. Since electrons are about 2000 times lighter than protons, synchrotron radiation lost is at the order of 10^{13} faster for electrons than for protons. For linear electron-positron collider, an extremely long acceleration section is required with

current technologies, which makes it an impractical option. As for the proton-anti-proton collider, it would not be possible to achieve such high luminosity using anti-proton beams, since it's much more difficult to produce anti-protons than to produce protons. In addition, at high energies the proton anti-proton collider starts losing one of its advantages of having higher cross section, which is due to the quark sea and anti-quarks in protons become more “visible” at high energies (more details in the parton model section).

As explained above, two separate rings are required to accomodate the two beam pipes, while the internal diameter of the tunnel is only about 3.8 m. It's technically challenging to install them in such small space. LHC therefore adopted the twin-bore magnet design [6], as shown in Figure 2. It was first proposed by John Blewett at the Brookhaven laboratory in 1971 due to cost considerations [7], but in the case of the LHC the overriding reason for adopting this solution is the lack of space in the tunnel.

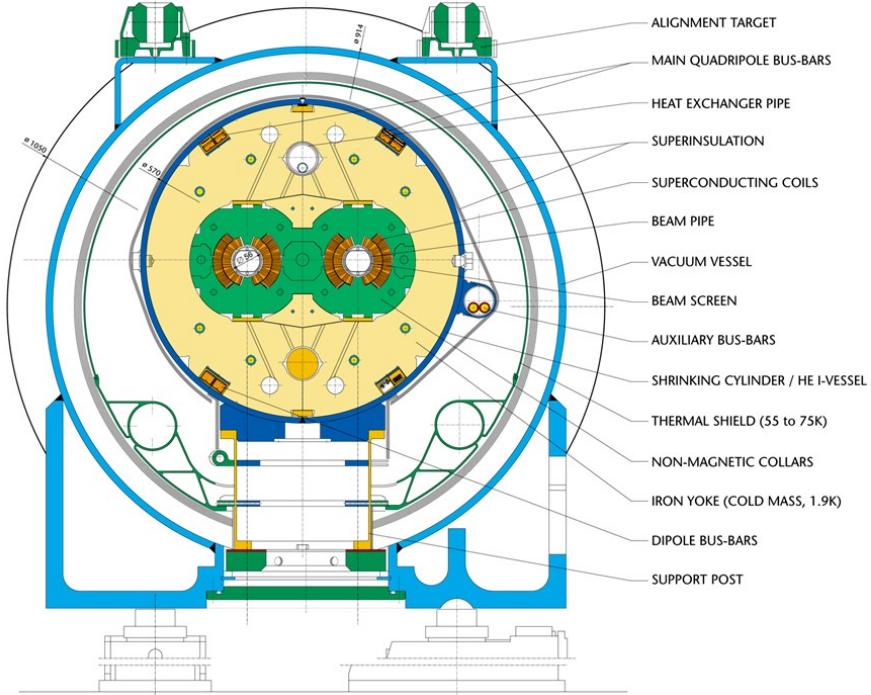


Figure 2: Double-bore magnet configuration of the LHC superconducting magnets [6].

3.1.2 Luminosity and pileup

Luminosity is an important measure of a collider's performance. It comes closely to the number of events generated per second, given by:

$$N_{event} = L\sigma_{event},$$

where $\sigma_{e\gamma}$ is the cross section for the event under study and L is the machine luminosity. For the cross-section, it is more common to use barn as the unit, where $1\text{b} = 10^{-28} \text{ m}^2 = 10^{-24} \text{ cm}^2$, since particle interactions usually have very small cross-sections. The machine luminosity depends on the beam parameters and can be written for a Gaussian beam distribution as:

$$\mathcal{L} = \frac{N_b^2 n_b f_{rev} \gamma_r}{4\pi \epsilon_n \beta^*} F,$$

where N_b refers to the number of protons per bunch, n_b is number of bunches per beam, f_{rev} is the revolution frequency, γ_r is the relativistic gamma factor, ϵ_r is the normalised transverse beam emittance, β^* is the beta function at the collision point which describes the size of the beam, and F refers to the geometric luminosity reduction factor due to the crossing angle at the interaction point (IP). While the instantaneous luminosity measures the rate of collisions (per unit cross-section), the total number of collisions (per unit cross-section) is measured by the integrated luminosity L , given by:

$$L = \int \mathcal{L} dt,$$

which is the integral of the instantaneous luminosity over time.

The two high luminosity experiments, ATLAS and CMS are both aiming at a peak luminosity of $L = 34 \text{ cm}^{-2}\text{s}^{-1}$ for proton operation, which corresponds to about one billion pp collisions per second. The instantaneous luminosity was much improved at real-time operations reaching about twice the nominal (from 2015 to 2018) thanks to the effort of the LHC experts.

Another important parameter for the LHC experiments is the pileup, which is a measure of the number of inelastic pp interactions that occur per bunch crossing. Higher pileup gives more luminosity (for a fixed number of bunches) but makes physics analysis more difficult due to the signals in the detector from the additional interactions. The distribution of the recorded luminosity over the pileup is shown in Figure 3 for operations from 2015 to 2018 (Run 2), where the $\langle \mu \rangle$ stands for the mean number of interactions per bunch crossing. There are two main sources of pileup: in-time pileup and out-of-time pileup. The former refers to the additional proton-proton collisions occurring in the same bunch-crossing as the collision of interest. The latter refers to the additional proton-proton collisions occurring in bunch-crossings just before and after the collision of interest. The out-of-time pileup is a challenge for the detector subsystems that have sensitivity windows longer than 25 ns, the interval between bunch crossings.

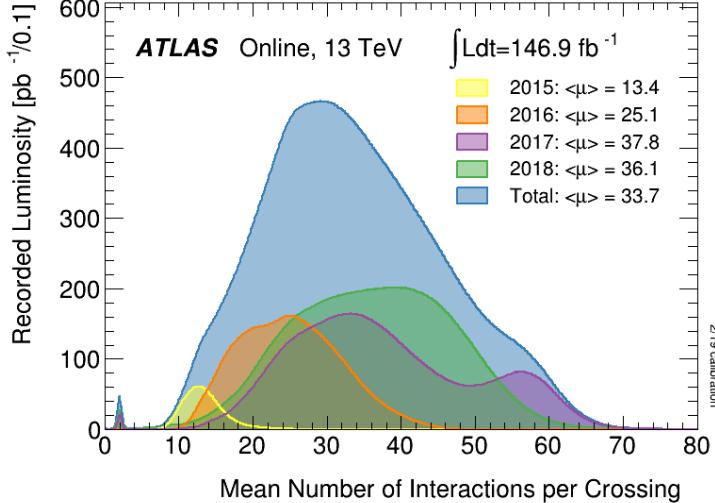


Figure 3: Shown is the luminosity-weighted distribution of the mean number of interactions per crossing for the data collected by the ATLAS from 2015 to 2018. All data recorded by ATLAS during stable beams is shown, and the integrated luminosity and the mean μ value are given in the figure.

3.1.3 Operation schedule

The LHC operation and shutdown schedule is shown in Figure 4 and Figure 5. Following the downtime after an incident in one of the main dipole circuits during the first commissioning in 2008 [8], the operation restarted at lower beam energy to minimise the risk. Therefore, the first proton run (2010-2013) [9] was carried out at 3.5–4 TeV per beam (centre of mass energy 7–8 TeV). Furthermore, a bunch spacing of 50 ns was used instead of the nominal 25 ns, with peak luminosity of $0.8 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$, and pileup larger than nominal.

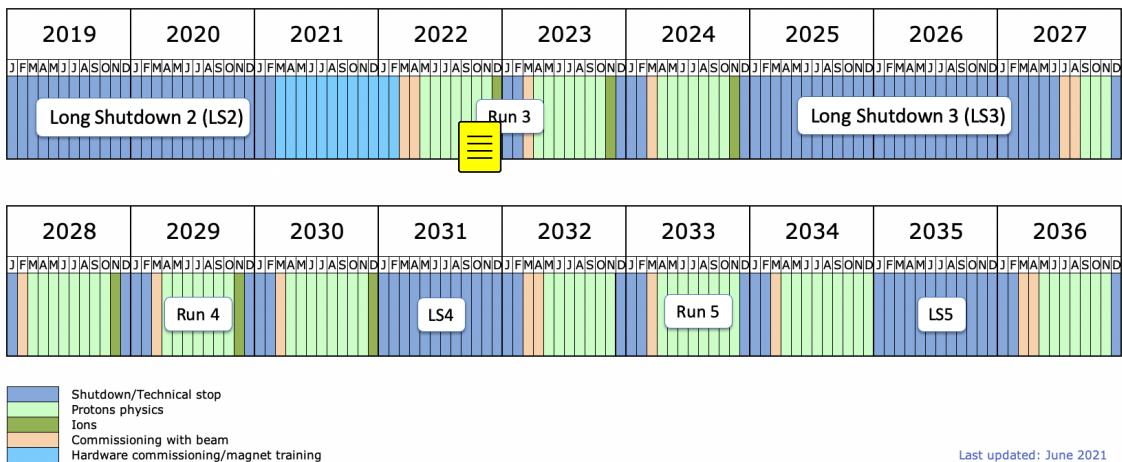


Figure 4: Longer term LHC operation schedule.



Figure 5: LHC operation schedule and luminosity targets.

In Run 1, the LHC delivered about 30 fb^{-1} of proton data and important physics results, most notably the discovery of the Higgs boson [10, 11]. Run 1 was followed by a long shutdown (LS1, 2013–2014) with a large number of consolidation and upgrade activities [12]. The bus-bar splices between the superconducting magnets were improved, in order to make sure that the LHC could operate at higher energy without risk of repeating the 2008 incident.

Run 2 (2016–2018) was carried out at 6.5 TeV per beam (center of mass energy 13 TeV) [13]. As shown in Figure 6, out of the 156 fb^{-1} of data LHC has delivered at 13 TeV centre of mass energy, the ATLAS detector has recorded 147 fb^{-1} and 139 fb^{-1} of data is certified to be good quality data. The delivered luminosity accounts for the luminosity delivered from the start of stable beams until the LHC requests ATLAS to put the detector in a safe standby mode to allow a beam dump or beam studies. The recorded luminosity is slightly smaller than the delivered luminosity, due to the inefficiency of the Data Acquisition (DAQ) and the so called “warm start”: when the stable beam flag is raised, the tracking detectors undergo a ramp of the high-voltage and, for the pixel system, turning on the pre-amplifiers. More details of the ATLAS detector can be found in the following sections. The recorded data is checked carefully to exclude possible hardware or software issues. This is achieved by monitoring detector-level quantities and reconstructed collision event characteristics at key stages of the data processing chain. This procedure led to high efficiency of good quality data: 95.6% [14].

In this thesis the 139 fb^{-1} data recorded by the ATLAS detector of Run 2 is used. The nominal bunch spacing of 25 ns was used, with slightly less bunches (2500) per beam. The

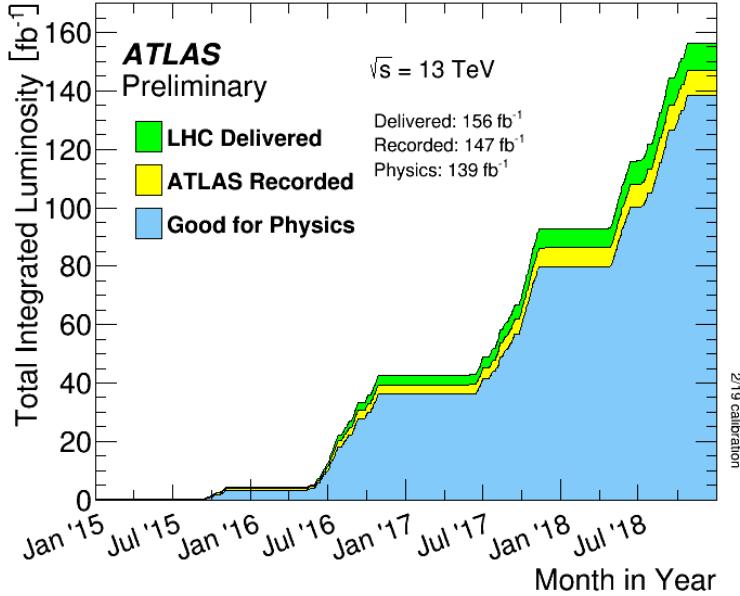


Figure 6: Cumulative luminosity versus time delivered to ATLAS (green), recorded by ATLAS (yellow), and certified to be good quality data (blue) during stable beams for pp collisions at 13 TeV centre-of-mass energy in Run 2.

LHC experts have continually improved the running scenario to increase the luminosity, and during Run 2 the luminosity surpassed the designed luminosity by a factor of 2. As well as improving the instantaneous luminosity, the availability of the machine was dramatically improved during Run 2 which is an important factor enabling the high efficiency of good quality data as mentioned above. During Run 2, the machine was providing physics collisions during 50% of the allocated physics time, which is very impressive for a super conducting collider.

The operation of CERN’s accelerators is subject to scheduled shutdowns to allow important repair and upgrade work to take place. The present shutdown, LS2, is devoted to preparations for Run 3 of the LHC, which will have an integrated luminosity equal to the two previous runs combined, and for the High-Luminosity LHC (HL-LHC), the successor to the LHC, which will begin operation at the end of 2027, and eventually generate 10 times the integrated luminosity of all Run 1, 2 and 3 combined!

The LS2 schedule has had to be modified due to the COVID-19 pandemic, which the new schedule anticipates that the first test beams will circulate in the LHC at the end of September 2021, four months later than the date planned before the COVID-19 crisis, to give the LHC’s main experiments – ATLAS, CMS, ALICE and LHCb – time to prepare their own upgrade. Run 3 of the LHC will begin at the start of March 2022. The third long shutdown (LS3) will begin at the start of 2025 and end in mid-2027. This is when the equipment for the HL-LHC and its experiments will be installed.

3.2 The ATLAS Detector

ATLAS is one of the two general purpose detectors built for probing proton-proton collision. This detector represents the work of a large collaboration of several thousand physicists, engineers, technicians, and students over a period of fifteen years of dedicated design, development, fabrication, and installation. The overall layout of the detector is shown in Figure 7. It has the dimensions of a cylinder, 46 m long, 25 m in diameter, and sits in a cavern 100 m below ground. The ATLAS detector weighs 7000 tonnes, similar to the weight of the Eiffel Tower. The detector itself is a many-layered instrument designed to detect some of the tiniest yet most energetic particles ever created on earth. It consists of six different detecting subsystems wrapped concentrically in layers around the collision point of nearly 4π solid angle coverage to record the trajectory, momentum, and energy of particles, allowing them to be individually identified and measured. These six subsystems are the pixel detector [15], the semiconductor tracker (SCT) [16], the transition radiation tracker (TRT) [16], the electromagnetic (EM) calorimeter [17], the hadronic calorimeter [18] and the muon spectrometer (MS) [19]. The first three sub-detectors are collectively known as the inner detector (ID), described in section 3.2.3, and it is used for tracking charged particles. The electromagnetic and the hadronic calorimeter, described in section 3.2.4, are responsible for measuring the energies of the electromagnetic and hadronic particles respectively. The MS, described in section 3.2.5, is a unique sub-detector used for measuring the momentum of muons leaving the calorimeters.

A huge magnet system bends the paths of the charged particles so that their momenta can be measured as precisely as possible.

The high interaction rates, radiation doses, particle multiplicities and energies, as well as the requirements for precision measurements have set stringent standards on the design of the ATLAS detector. Therefore the ATLAS detector is designed to fulfill the following requirements:

- Fast, radiation-resistant electronics and sensor elements and high detector granularity. This is due to the high frequency of collisions, high particle fluxes and high radiation environment of the detector.
- Large acceptance in polar angle with almost full azimuthal angle coverage, due to the geometry of the detector (more details in section 3.2.1).
- Good energy resolution calorimetry, as required to enable accurate physical object reconstruction. The high resolution of energy can be obtained with very good electromagnetic calorimetry for electron and photon identification and measurements, complemented by full-coverage hadronic calorimetry for accurate jet and missing transverse energy measurements.

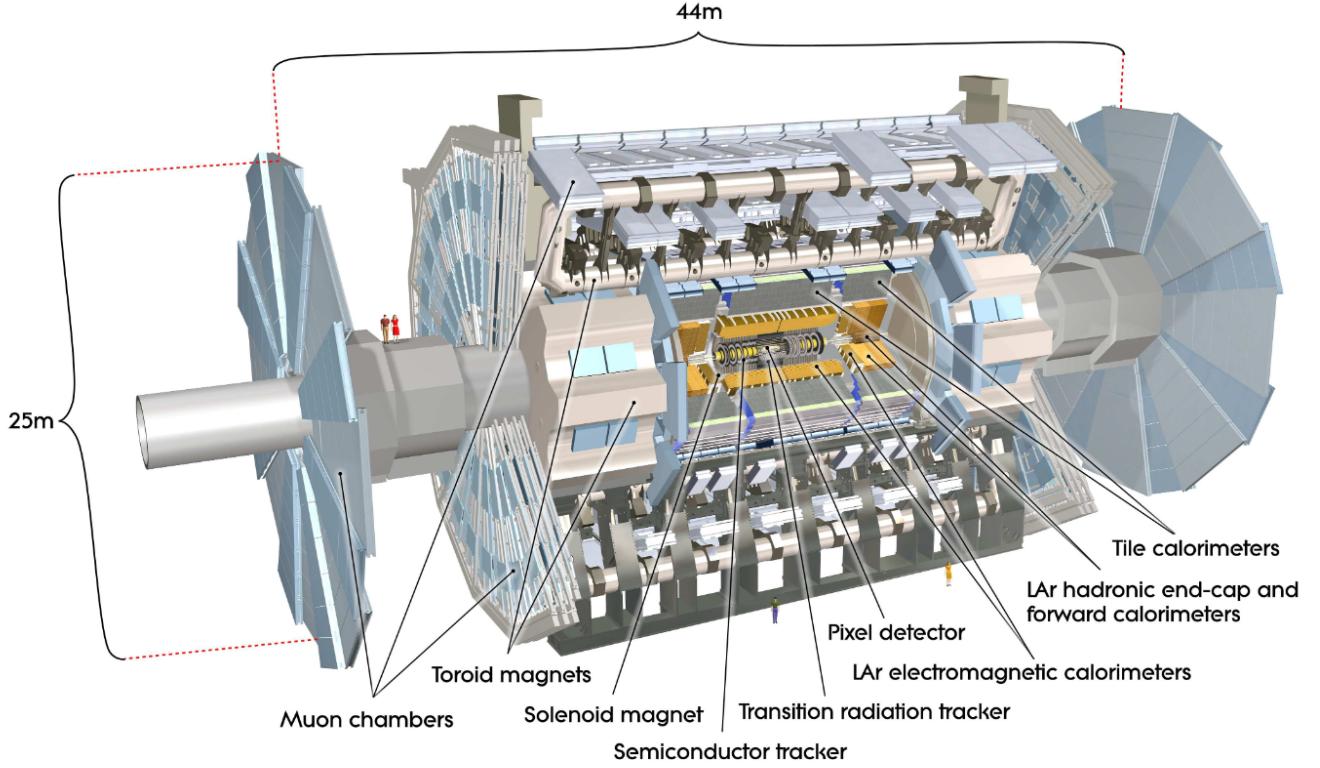


Figure 7: Cut-away view of the ATLAS detector. The dimensions of the detector are 25 m in height and 44 m in length. The overall weight of the detector is approximately 7000 tonnes.

- Tracking of precision in the ID, as required to provide high momentum resolution and to allow the reconstruction of secondary vertices to identify b -hadrons and τ -leptons.
- Good muon identification and momentum resolution over a wide range of momenta and the ability to determine unambiguously the charge of high- p_T muons in the muon spectrometer.
- Trigger system with high efficiency on low p_T objects with sufficient background rejection, which is a prerequisite to achieve an acceptable trigger rate for most physics processes of interest.

The main performance goals of the detector are listed in Table 1.

3.2.1 Coordinate system

The ATLAS coordinate system is a right-handed Cartesian system with the nominal interaction point defined as the origin of the coordinate system, while the beam direction defines the z -axis and the x - y plane is transverse to the beam direction. The positive x -axis is

Detector component	Required resolution	η coverage	
		Measurement	Trigger
Tracking	$\sigma_{p_T}/p_T = 0.05\% p_T \oplus 1\%$	± 2.5	None
EM calorimetry	$\sigma_E/E = 10\% /E \oplus 0.7\%$	± 3.2	± 2.5
dronic calorimetry barrel and end-cap forward	$\sigma_E/E = 50\% /E \oplus 3\%$ $\sigma_E/E = 100\% /E \oplus 10\%$	± 3.2 $3.1 < \eta < 4.9$	± 3.2 $3.1 < \eta < 4.9$
Muon spectrometer	$\sigma_{p_T}/p_T = 10\% \text{ at } p_T = 1 \text{ TeV}$	± 2.7	± 2.4

Table 1: General performance goals of the ATLAS detector. The units for energy of the particle, E and transverse momentum, p_T (detailed definition in section 3.2.1) are in GeV [2]. Note that, for high- p_T muons, the muon-spectrometer performance is independent of the inner-detector system.

defined as pointing from the interaction point to the centre of the LHC ring and the positive y -axis is defined as pointing upwards, as shown in Figure 8.

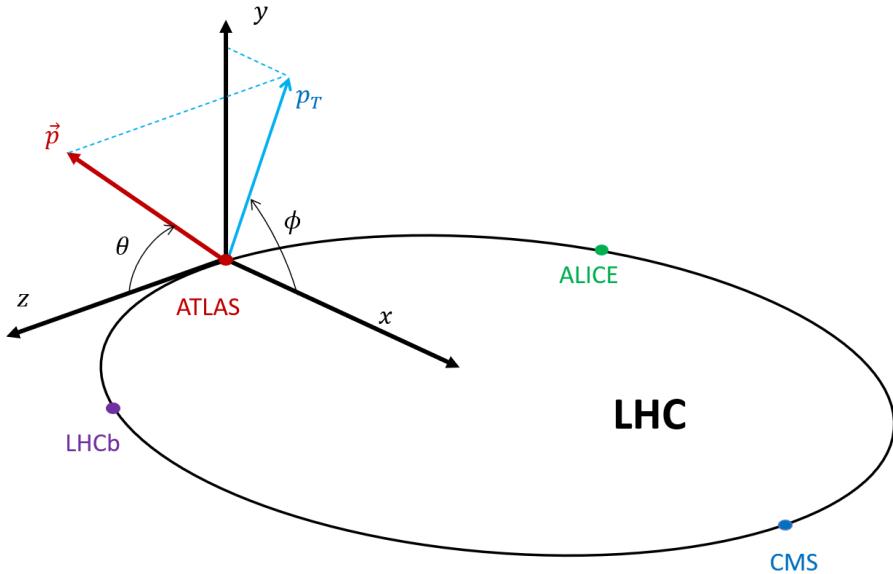


Figure 8: Illustration of the coordinate system used at the ATLAS experiment in the geographical context of the LHC.

The azimuthal angle ϕ is measured as usual around the beam axis, and the polar angle θ is the angle from the beam axis. In high energy physics, it's more common to use the pseudorapidity instead of the polar angle θ , defined as:

$$\eta = -\ln \tan(\theta/2).$$

In the case of highly relativistic particles (which is the common case in high energy physics),

the pseudorapidity approaches the rapidity,

$$y = 1/2 \ln[(E + p_z)/(E - p_z)],$$

where E is the energy of the particle, m is its mass and p_z is the momentum along the z -axis. There are two main reasons for using pseudorapidity rather than the polar angle θ nor the rapidity. The reason for using pseudorapidity but not θ is that the rapidity is invariant under Lorentz transformation, while capturing the characteristic of the particle direction of travel: $y \rightarrow \pm\infty$ when the particle is travelling close to the beam pipe (positive for along the beam pipe, negative for the opposite direction) and $y \rightarrow 0$ when p_z is small. The reason for using pseudorapidity but not rapidity is that due to the limited angle coverage of the detector, it's usually hard to determine the total energy and the momentum along the z -axis, especially when the direction of the particles are close to the beam pipe. While the pseudorapidity is determined only by the polar angle, which is much easier and faster to compute. Another commonly used variable, transverse momentum p_T , is defined as the momentum of a particle transverse to the beam direction (z -direction):

$$\vec{p}_T = (p_x, p_y).$$



The reason for using transverse momentum is that, because the partons that make up a proton share the momentum, the initial longitudinal momentum is unknown; we do know, however, that the initial transverse momentum was zero. And hence we can look for the missing transverse momentum, defined as



$$E_T^{miss} = - \sum_i p_{Ti}$$

for visible particles i (Confusingly E_T^{miss} is commonly called missing transverse energy or MET. Missing transverse energy is equivalent to missing transverse momentum only if the missing particle(s) were massless.). Finally, the distance ΔR in the pseudorapidity-azimuthal angle space is defined as:

$$\Delta R = \sqrt{\Delta\eta^2 + \Delta\phi^2},$$

where R is the radial distance from the particle position to the interaction vertex.

3.2.2 Magnets

ATLAS has a unique hybrid system of four large superconducting magnets [20]. This magnetic system is 22 m in diameter and 26 m in length, with a stored energy of 1.6 GJ. Figure 9 shows the real scale of the magnets system compared to a person. Figure 7 shows the general

layout, the four main layers of detectors and the four superconducting magnets which provide the magnetic field over a volume of approximately 12000 m³.



Figure 9: A picture showing the real size magnets compared to a person.

The spatial arrangement of the coil windings is shown in Figure 10. The ATLAS magnet system consists of two parts:

- a solenoid, which is aligned on the beam axis and provides a 2 T axial magnetic field in the z -direction for the ID. Because the magnet is located in front of the EM calorimeter, it is imperative to minimise possible interactions between the magnet and the particles being studied. This is achieved by embedding over 9 km of niobium-titanium superconductor wires into strengthened, pure aluminum strips, which is capable to provide such powerful magnetic field in just 4.5 cm thickness.
- A barrel toroid and two end-cap toroids, which produce a toroidal magnetic field of approximately 0.5 T and 1 T for the muon detectors in the central and end-cap regions, respectively. The barrel toroid generates the magnetic field in the central zone of the muon spectrometer, along the tangential direction of the circumferences centered on the z -axis (ϕ direction). The end-cap toroids are two smaller toroids designed to provide the magnetic field in the forward areas of the muon spectrometer. This magnet configuration provides a field mostly orthogonal to the muon trajectories.

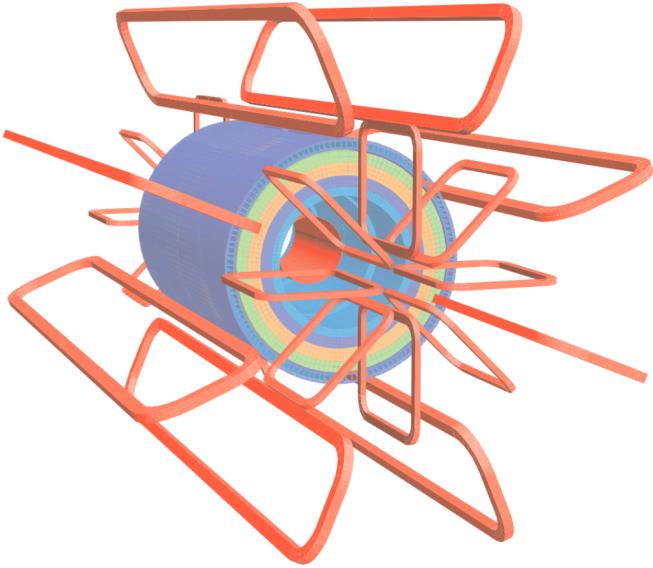


Figure 10: Geometry of magnet windings and tile calorimeter steel.

3.2.3 Inner detector

The inner detector (ID) is the closest sub-detector to the beam line, designed to track the early trajectories of charged particles for momentum calculations and locate their primary and secondary vertices with extremely high precision. The ID is required to deal with large numbers of tracks promptly, where 1000 particle collisions are taking place every 25 ns. The ID has full coverage in azimuthal angle ϕ and $|\eta| < 2.5$ acceptance in pseudorapidity. As mentioned briefly above, it consists of three parts: the pixel detector and the insertable B-Layer (IBL) [21] (as one part), the semiconductor tracker and the transition radiation tracker. The layout of the ID is shown in figure 11, with a charge track (in red) traversing the sensors and structural elements.

3.2.3.1 Pixel detector and IBL

The silicon pixel detector is the closest ATLAS component to the collision. It is composed of layers of silicon pixels and designed to have a very high granularity for reconstructing primary and secondary interaction vertices. The detector layers are formed of silicon sensor modules and in total there are approximately 92 million pixels (consequently, 92 million readout channels) in the system. It consists of three cylindrical layers in the barrel region, positioned at the radial distances of 50.5, 88.5 and 122.5 mm, and of disks perpendicular to the beams in the end-caps at the longitudinal distances of 49.5, 58.0 and 65.0 mm. In 2014, during the first LHC long shutdown, a fourth pixel layer was installed inside the existing detector, the insertable B-Layer (IBL) at a radius of 33 mm from the beam axis. The new

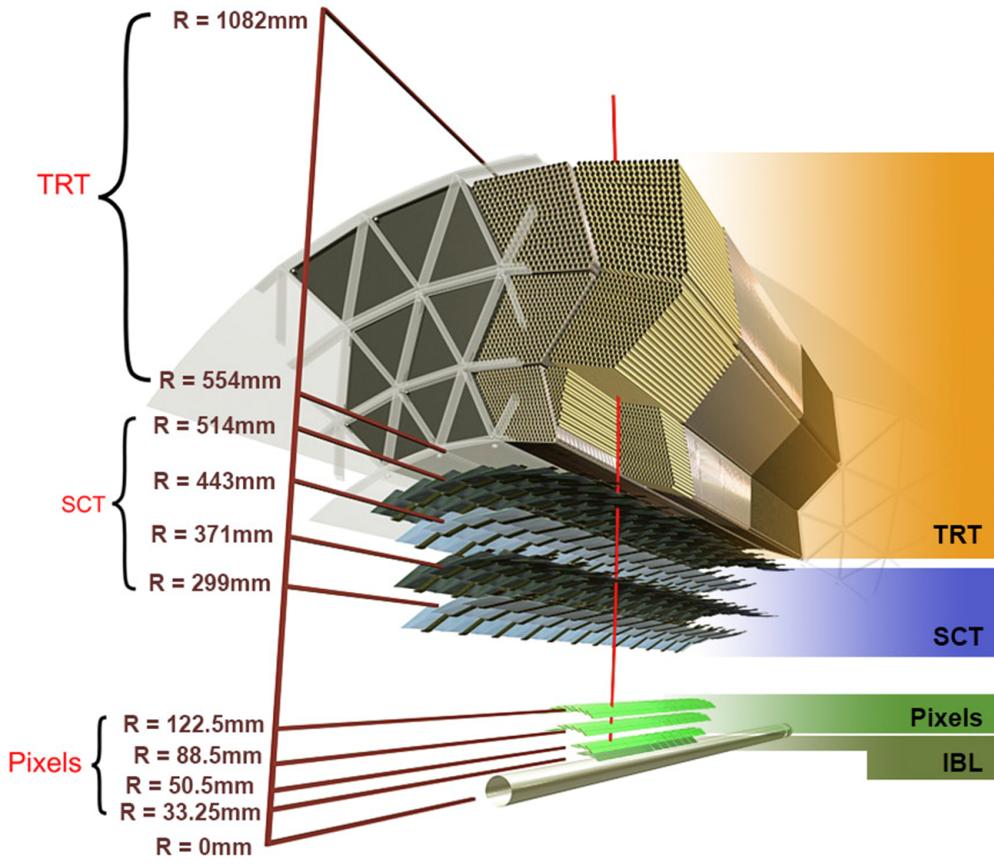


Figure 11: Cut-away view of the inner detector.

pixel layer provides an additional space point very close to the interaction point, which significantly improves the identification of jets coming from b -quark hadronisation (b jets). Particles with $|\eta| < 2.5$ traverses the four layers usually produce four space-points. The pixel detector provides a resolution of $\sigma_\phi = 10 \mu\text{m}$ in the bending direction ($R - \phi$), and a resolution of $\sigma = 115 \mu\text{m}$ in the $z(R)$ direction in the barrel (end-cap) region.

3.2.3.2 Semiconductor tracker

The next constituent of the inner detector is the SCT. It is a silicon microstrip detector with over six million readout channels, which surrounds the pixel detector and covers the region of radius between 299 mm and 560 mm. It consists of four layers of strips located axially on the beam direction in the barrel region and placed along the z -direction in the end-cap region. This configuration allows the particles along the beam pipe to be reconstructed. Each layer of strips is glued back to back with an angle of 40 mrad to form a two-sided module and make possible the measurement of the second coordinate. The sensors are 285 μm thick and

are constructed of high-resistivity n-type bulk silicon with p-type implants. Readout strips are positioned every $80\ \mu\text{m}$, providing a spatial resolution of $\sigma_\phi = 17\ \mu\text{m}$ in the bending direction ($R - \phi$) and $\sigma_\phi = 580\ \mu\text{m}$ in the z (barrel) and R (end-cap) direction.

3.2.3.3 Transition radiation tracker

The outermost part of the inner detector is the TRT, which covers the radial region between 563 mm and 1066 mm. It is a straw drift tube tracker, which consists of modules of 4 mm diameter polyimide straws, filled with a mixture of gas of 70% Xe, 27% CO₂ and 3% O₂ and a gold-plated tungsten wire in the centre. The straws are interleaved with propylene fibres (foils) in the barrel (end-cap) region. With a spatial resolution of $\sigma_\phi = 130\ \mu\text{m}$, the TRT measures the track position only in the bending direction ($R - \phi$). This is because when a charged particle passes through a straw tube, electrons from the gas are liberated through ionisation; under high voltage, these electrons then drift toward the wire in the centre, where a current flow is created and registered as a hit. Since a hit can happen on any location along the wire, the information of the z position of the particle is lost. In addition, the TRT provides capability of distinguishing electrons from other charged particles. When a highly-relativistic charged particle traverses the polymer straws interface, the particle emits transition radiation which is then absorbed by the Xeon gas. The intensity of the radiation depends on the gamma factor of the particle (strongest for lighter particles), hence this information can be exploited for electron identification.

3.2.4 Calorimeter system

Calorimeters are used to measure the energy of both charged and neutral particles. The ATLAS calorimeters [22], as shown in Figure 12, are consists of three major components, the Electromagnetic calorimeter, the Hadronic calorimeter and the Forward calorimeter (FCal). The fine granularity of the EM calorimeter is ideal for precision measurements of electrons and photons; the coarser granularity of hadronic calorimeter is sufficient for the hadronic jet reconstruction; the FCal provides coverage of large pseudorapidity region: these calorimeters cover the range $|\eta| < 4.9$. All the three calorimeters are sampling calorimeters. Sampling calorimeters use different material the absorber and the active part: the absorber (or passive material) is responsible for producing particle showers where the active part then measures their energy. Note well the a fraction of total particles energy deposited in the passive material and it's not measured; the overall energy must be deduced from the definite measurements taken in the active detector layers.

The electromagnetic and hadronic showers must be contained in the calorimeter to ensure precise measurement of the total energy of the particle and to avoid punch-through into the muon system. The calorimeter depth is hence an important design consideration. The

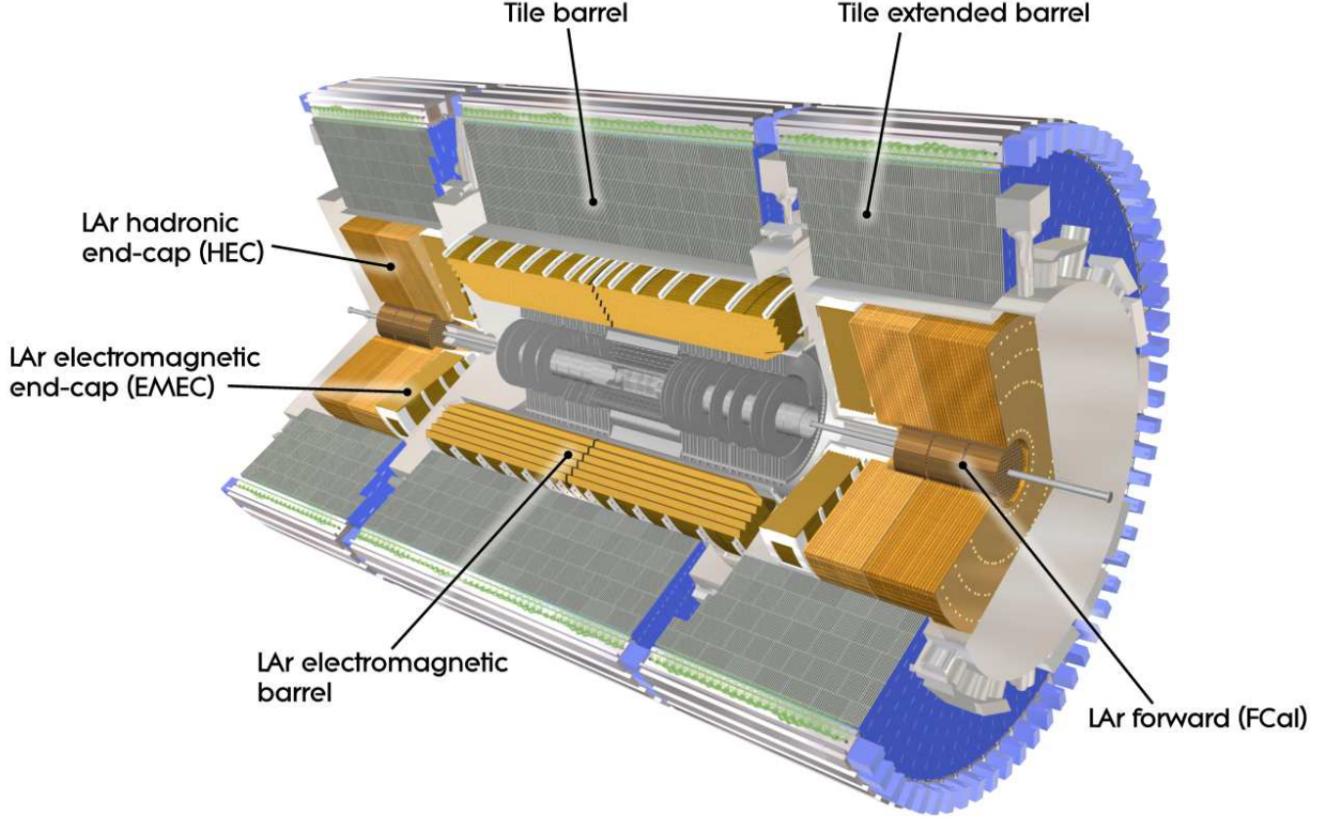


Figure 12: Cut-away view of the ATLAS calorimeter system.

thickness of the calorimeter is measured in radiation length X_0 , which is the mean length of a material over which an electron will lose all but $1/e$ of its initial energy through radiative processes; and nuclear interaction length λ , which is the mean distance travelled by a hadronic particle before undergoing an inelastic nuclear interaction. The total thickness of the EM calorimeter is greater than 22 radiation lengths (X_0) in the barrel and greater than $24 X_0$ in the end-caps. The total thickness of the calorimeters, including 1.3λ from the outer support, is 11λ at $\eta = 0$ and has been shown both by measurements and simulations to be sufficient to reduce punch-through well below the irreducible level of prompt or decay muons.

3.2.4.1 Electromagnetic calorimeter

The EM calorimeter is a lead-Liquid argon (lead-LAr) detector [17] with accordion-shaped (as shown in Figure 13) kapton electrodes and lead absorber plates over its full coverage, while using liquid argon (LAr) as the active material. The lead thickness in the absorber plates has been optimised as a function of η in terms of EM calorimeter performance in energy resolution. The EM calorimeter is divided into a barrel part ($|\eta| < 1.475$) and two end-cap components ($1.375 < |\eta| < 3.2$), each housed in their own cryostat. Additional material

needed to instrument and cool the detector creates a “crack” region at $1.375 < |\eta| < 1.52$, where the energy resolution is significantly degraded.



Figure 13: A figure of the accordion-shaped electrodes.

The barrel calorimeter consists of two identical half-barrels, separated by a small gap (4 mm) at $z = 0$. Each end-cap calorimeter is mechanically divided into two coaxial wheels: an outer wheel covering the region $1.375 < |\eta| < 2.5$, and an inner wheel covering the region $2.5 < |\eta| < 3.2$.

The calorimeter has three layers along the transverse direction: a pre-sampler with very high granularity in η , in order to reconstruct the neutral pions decaying to two photons and particles which already starts showering in the inner detector. The pre-sampler is followed by longer towers of relatively high granularity, which is the major part of detecting EM showers, and responsible for measuring the η and ϕ coordinates of the particles. The last layer detects showers generated from particles other than electrons or photons that start showering inside the EM calorimeter before leaving it.

3.2.4.2 Hadronic calorimeter

The Hadronic calorimeter is comprised of the Tile Hadronic calorimeter (HCAL) and the LAr hadronic end-cap calorimeter (HEC). The HCAL is placed directly outside the EM calorimeter envelope. Its barrel covers the region $|\eta| < 1.0$, and its two extended barrels the range $0.8 < |\eta| < 1.7$. It is using steel as the absorber and scintillating tiles as the active material. It is segmented in depth in three layers, approximately 1.5, 4.1 and 1.8 λ thick for the barrel and 1.5, 2.6, and 3.3 λ for the extended barrel. The total detector thickness at the outer edge of the tile-instrumented region is 9.7 λ at $\eta = 0$.

The HEC is similar to the construction to the ECAL, using LAr as the active material, but instead of using lead it uses copper as the absorber. It consists of two independent wheels

per end-cap, located directly behind the end-cap electromagnetic calorimeter and sharing the same LAr cryostats. The HEC covers the range of $1.5 < |\eta| < 3.2$, slightly overlapping with the forward calorimeter which will be described in the following paragraph (around $|\eta|=3.1$) and the tile calorimeter ($|\eta| < 1.7$). This overlap is to reduce the drop in material density at the transition between the different calorimeters.

3.2.4.3 Forward calorimeter

The Forward Calorimeter (FCal) covers $3.1 < |\eta| < 4.9$ and is approximately 10 interaction lengths deep. It consists of three modules in each end-cap: the first, made of copper, is optimised for electromagnetic measurements, while the other two, made of tungsten, measure predominantly the energy of hadronic interactions. All three modules use LAr as active material. Due to high particle fluxes and energies in the forward region, the calorimeter must contain relatively long showers in the small volume allowed by design constraints, and thus must be very dense.

3.2.5 Muon Spectrometer

The muon spectrometer is the outermost and largest sub-detector of ATLAS. A cut-away view of the MS is illustrated in Figure 14. It fully covers the calorimeter system and occupies a large part of the ATLAS cavern. It is based on the magnetic deflection of muon tracks in the large superconducting air-core toroid magnets, instrumented with separate trigger and high-precision tracking chambers. Over the range $|\eta| < 1.4$, magnetic bending is provided by the large barrel toroid. For $1.6 < |\eta| < 2.7$, muon tracks are bent by two smaller end-cap magnets inserted into both ends of the barrel toroid. Over $1.4 < |\eta| < 1.6$, usually referred to as the transition region, magnetic deflection is provided by a combination of barrel and end-cap fields. The configuration of magnets provides a field mostly orthogonal to the muon trajectories, hence minimising the degradation of resolution due to multiple scattering.

The MS consists of four subsystems which rely on four different gas detector technologies. Two of them, the resistive plate chambers (RPC) in the barrel region and the thin gap chambers (TGC) in the end-cap region, provide trigger signals, while the other two, the monitored drift tubes (MDT) in the barrel and the cathode strip chambers (CSC) in the end-cap region provide the momentum measurement. The MDT chambers provide high precision measurements in the bending direction over most of the detector acceptance while the CSC are used in the forward region where the particle flux is too high for the MDT chambers. The muon chambers are arranged in the barrel ($|\eta| < 1.05$) in three cylindrical layers around the beam axis, while in the end-cap regions ($1.05 < |\eta| < 2.7$) they are placed in three wheels. The resolution of muons tracks momentum measurement varies from typically 2–3% over most of the kinematic range, to about 10% at $p_T = 1$ TeV.

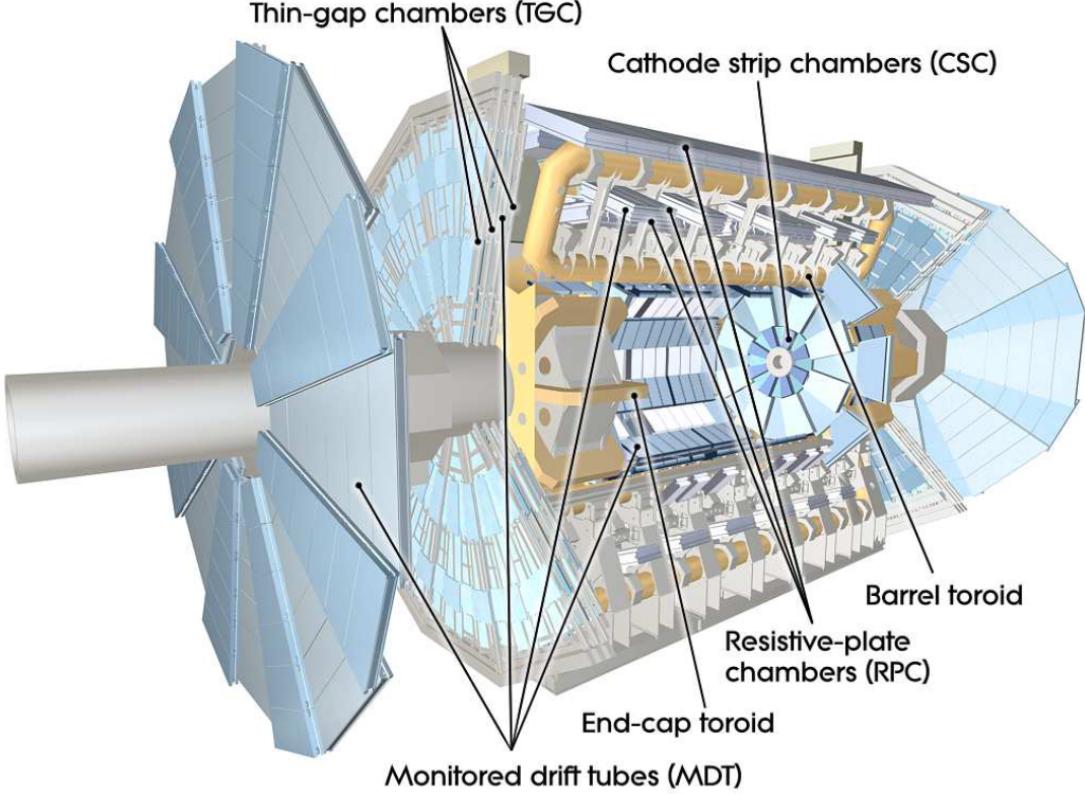


Figure 14: Cut-away view of the ATLAS muon spectrometer system.

3.2.6 Trigger system

As mentioned in section 3.1.3, the spacing of each bunch is 25 ns, which translates to a 40 MHz of bunch-crossing frequency, with up to 80 collisions per bunch crossing. This is far beyond the data collection bandwidth and storage capacity of ATLAS. Therefore, it's necessary to adopt a trigger system that make fast decisions whether an event is high quality, rare or "interesting" and to save the event or not. The ATLAS trigger system consists of two consecutive parts: the Level 1 trigger [23] which is hardware-based, followed by the software-based High Level Trigger (HLT) [24].

The L1 trigger searches for signatures from high- p_T muons, electrons/photons, jets, and τ -leptons decaying into hadrons. It also selects events with large MET and large total transverse energy. The L1 trigger uses reduced-granularity information from a subset of detectors: the RPC and TGC for high- p_T muons, and all the calorimeter sub-systems for electromagnetic clusters, jets, τ -leptons, E_{miss}^T , and large total transverse energy. As a result, the L1 trigger reduces the event rate from 40 MHz to a maximum of 100 kHz. The decision is made by Central Trigger Processor (CTP), which operates on signals from dedicated hardware in the calorimeter and muon detector systems. The decision time, at under 2.5 μs , is faster than the ID can process events so ID information is omitted. For each data-taking

period, the L1 trigger is loaded with a trigger menu, a list of up to 256 criteria used to determine whether an event is accepted. The trigger menus are designed to accomodate a broad physics programme, with high acceptance for both BSM searches and SM precision measurements. The L1 trigger also uses detector information with reduced granularity to identify Regions of Interest (RoI) [25] in ϕ and η . The ROI information with full granularity and precision and all the available detector data (including the ID information) within the RoI's are provided to the HLT. This trigger level reduces the rate of events by two orders of magnitude, reaching an average of 1 kHz with a latency of 0.2 μ s. These events are passed on to a data storage system for offline analysis.



4 Data and Monte Carlo samples

5 Physics Objects Reconstruction

5.1 Track and vertex

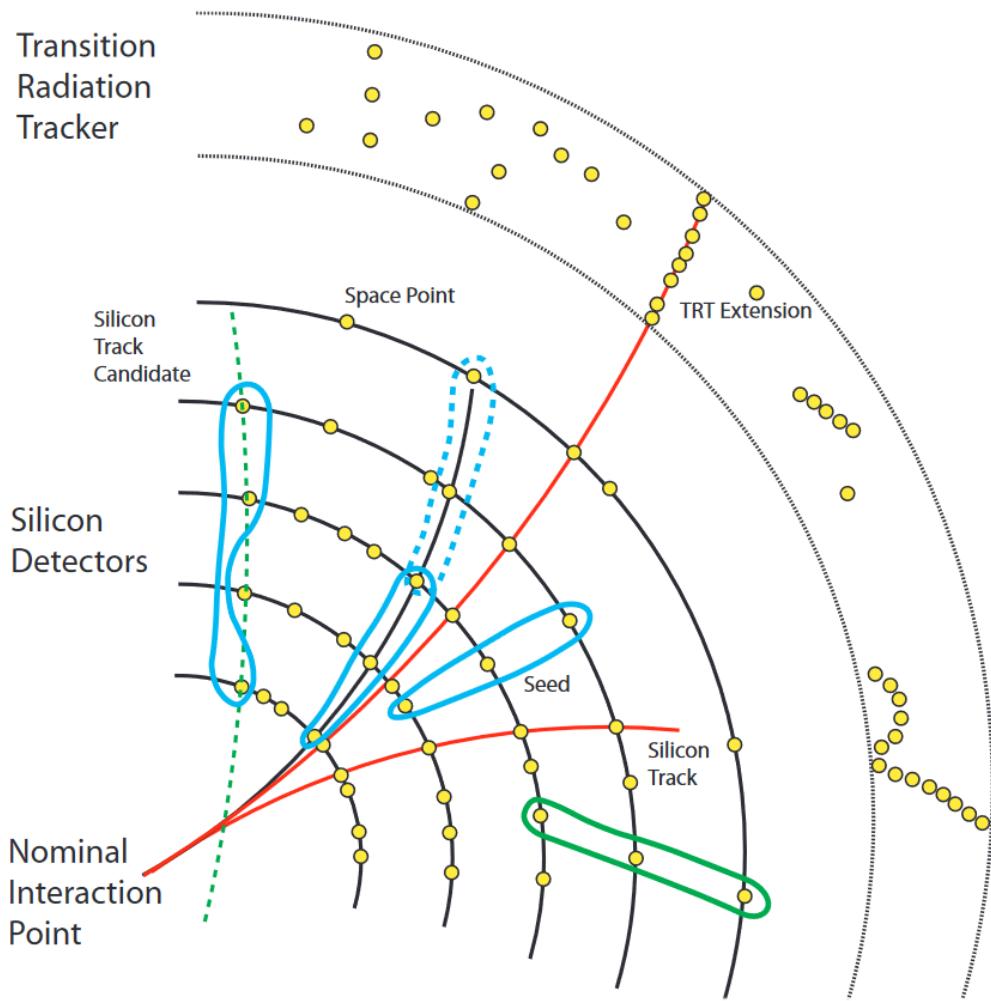


Figure 15: An example of track reconstruction [26].

Track reconstruction is a starting point of physics objects reconstruction, as it is not only used in reconstruction of charged particles but also almost every element of reconstruction, such as vertices, pile-up removal, missing energy and flavour tagging. Therefore it is crucial to understand how tracks are reconstructed in ATLAS. There are 3 stages of reconstruction [27], starting with the “inside-outside” algorithm, then the “outside-in” and finally reconstructing the TRT-standalone tracks. Figure 15 shows an illustrative view of elements used in tracks reconstruction.

The inside-out algorithm starts from assembling clusters from the raw measurements, from which three-dimensional measurements referred to as space points are created. They

represent the point where the charged particle traversed the active material of the ID. In the pixel detector, each cluster equates to one space-point, while in the SCT, clusters from both sides of a strip layer must be combined to obtain a three-dimensional measurement. Once the space-points are created, three sets of space points are combined to seeds. A combinatorial Kalman filter [28] is used to build track candidates from the chosen seeds by incorporating additional space-points from the remaining layers of the pixel and SCT detectors which are compatible with the preliminary trajectory. A track score is assigned to each track, which is computed by the χ^2 of fitting the track, intrinsic resolution, expected clusters multiplicity, number of holes (missing hits) and track p_T . Ambiguity is solved by selecting the track with the best score. The final step of the “inside-out” algorithm is to extend the silicon tracks into the TRT. The tracks are fitted again to estimate the final track parameter (IP, transverse distance from the interaction point) d_0 , distance from the interaction point along the z axis z_0 , azimuthal angle ϕ , polar angle θ and charge momentum ratio q/p_T . The “outside-in” algorithm starts with searching for tracks with segments reconstructed in the TRT, and extend the tracks inwards by adding silicon hits. Finally, tracks with a TRT segment but no extension into the silicon detectors are referred as TRT-standalone tracks. Vertices are matched to interactions by calculating the sum of the weights of the tracks in a vertex matched to each interaction.

Primary vertices are reconstructed using an iterative vertex finding algorithm [27]. Vertex seeds are obtained from the z -position at the beamline of the reconstructed tracks. An iterative χ^2 fit is made using the seed and nearby tracks. Tracks displaced by more than 7σ are used to seed a new vertex and the procedure is repeated until no additional vertices can be found. The beam spot position is used as a constraint. Vertices are required to contain at least two tracks.

starts from 3-point seeds in the silicon detectors and adds hits moving away from the interaction point using a combinatorial Kalman filter (Iterative algorithm that provides best estimate of the state based on projection of earlier measurements and current measurement). Ambiguities in the track candidates are resolved, and tracks are extended into the TRT. This is the baseline algorithm designed for the efficient reconstruction of primary charged particles. In a second stage, a track search starts from segments reconstructed in the TRT and extends them inwards by adding silicon hits. This back-tracking is mainly designed to reconstruct secondaries. Finally tracks with a TRT segment but no extension into the silicon detectors are referred as TRT-standalone tracks.

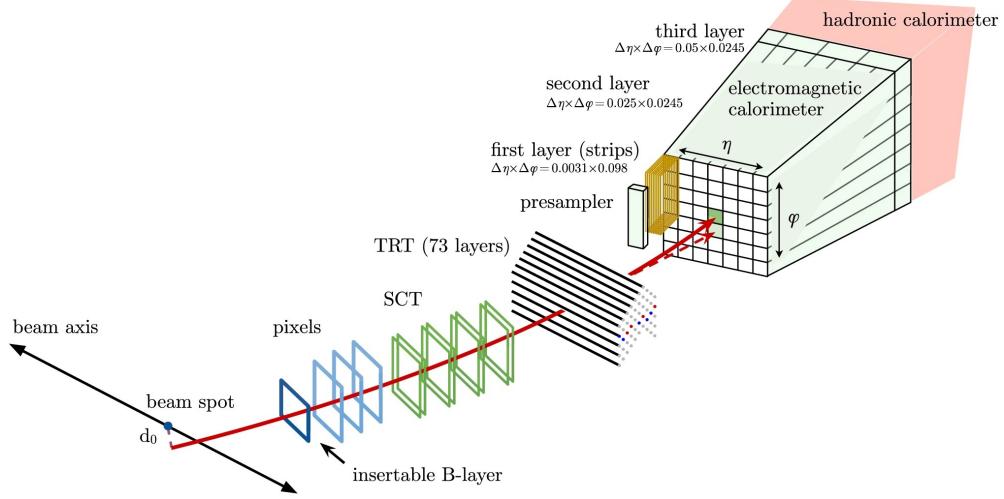


Figure 16: A schematic illustration of the path of an electron through the detector. The red trajectory shows the hypothetical path of an electron, which first traverses the tracking system (pixel detectors, then silicon-strip detectors and lastly the TRT) and then enters the electromagnetic calorimeter. The dashed red trajectory indicates the path of a photon produced by the interaction of the electron with the material in the tracking system[29].

5.2 Electron

5.3 Muon

5.4 Jet

5.5 b jets

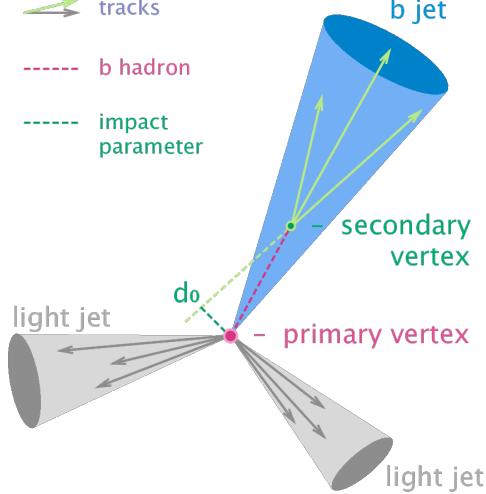


Figure 17: A diagram showing the b hadron decay initiated jets.

5.5.1 Flavour tagging

The identification of jets containing b -hadrons (b jets) against the large background of jets containing c -hadrons (c jets) or coming from the hadronization of light (u,d,s) quarks or gluons is of major importance in many areas of the physics programme of the ATLAS experiment at the LHC. It is crucial in a large number of Standard Model (SM) precision measurements, studies of the Higgs boson properties, and searches for new phenomena [30, 31, 32]. It also plays an important role in the $HH \rightarrow bb\tau\tau$ searches presented in Chapter 7.

The ATLAS Collaboration uses various algorithms to identify b jets [33], referred to as b -tagging algorithms, when analysing data recorded during Run 2 of the LHC. These algorithms exploit the long lifetime, high mass and high decay multiplicity of b -hadrons, as well as the properties of the b quark fragmentation. Given a lifetime of the order of 1.5 ps, b -hadrons have a significant mean flight length ($\langle c\tau \rangle \approx 450 \mu\text{m}$), in the detector before decaying, generally leading to at least one vertex displaced from the hard-scatter collision point, as illustrated in Figure 17.

The strategy developed by the ATLAS Collaboration is based on a two-stage approach. Firstly, low-level algorithms reconstruct the characteristic features of the b jets via two complementary approaches, one that uses the individual properties of charged-particle tracks, later referred to as tracks, associated with a hadronic jet, and a second which combines the tracks to explicitly reconstruct displaced vertices. These algorithms, first introduced during Run 1 [33], have been improved and retuned for Run 2 [34]. Secondly, in order to maximise the b -tagging performance, the results of the low-level b -tagging algorithms are combined into high-level algorithms via multivariate classifiers.

The most performant algorithms presently in use in physics analyses at ATLAS are based on multivariate combinations of the available information (MV2) or additionally using a deep

feed-forward neural network (DL1) [35, 36], as shown in Figure 18, where the performance is characterised by the probability of tagging a b jet (b jet tagging efficiency, ϵ_b) and the probability of mistakenly identifying a c jet or a light-flavour jet as a b jet labelled $\epsilon_c(\epsilon_l)$. In addition, the distribution of the output discriminant of the MV2 and DL1 tagger for b jet, c jet, and light-flavour jets in the $t\bar{t}$ simulated events are shown in Figure 19. Depending on the low-level algorithm, the DL1 tagger can be further separated into two taggers: DL1 and DL1r, where the DL1 tagger uses traditional track-based impact parameter taggers IP2D and IP3D [37] and the DL1r tagger uses a Recurrent Neural Network Impact Parameter tagger (RNNIP) [36]. The DL1r tagger is now the default b -tagging algorithm used for flavour tagging in ATLAS. The performance of the algorithms is quantified in terms of c jet (light jet) rejections, defined as $1/\epsilon_c$ and $1/\epsilon_l$.

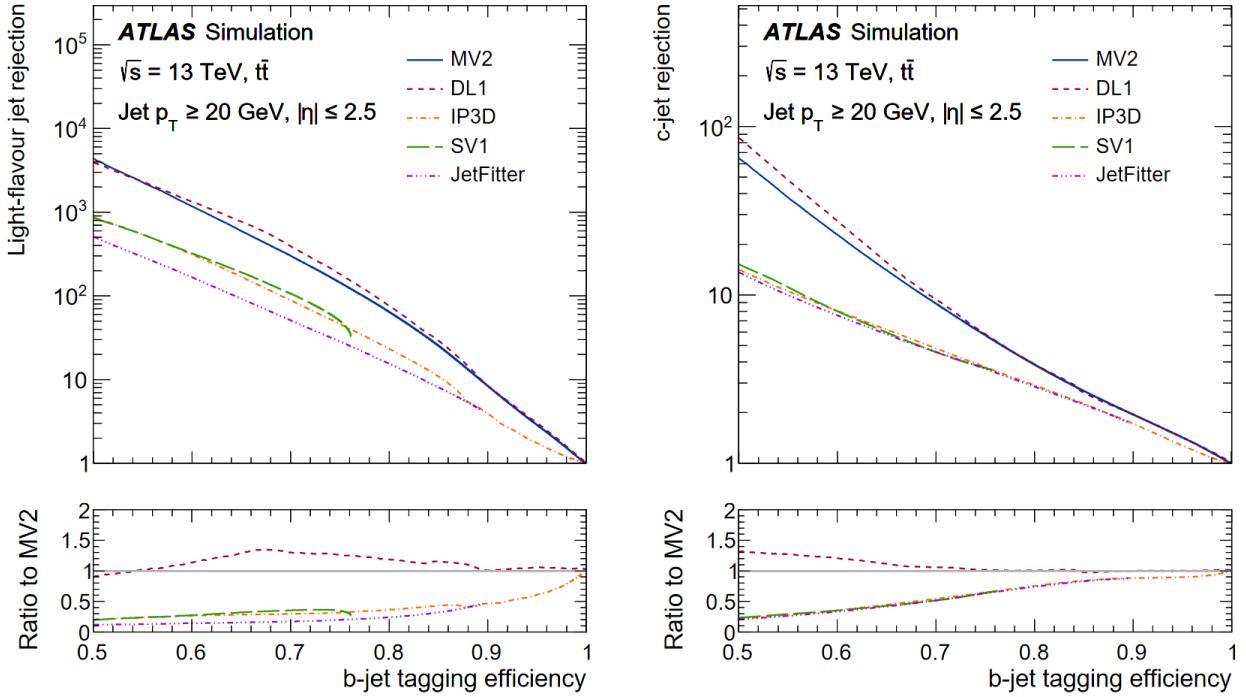


Figure 18: The light-flavour jet (left) and c jet (right) rejections versus the b jet tagging efficiency for the IP3D, SV1, JetFitter, MV2 and DL1 b -tagging algorithms evaluated on the $t\bar{t}$ events [34].

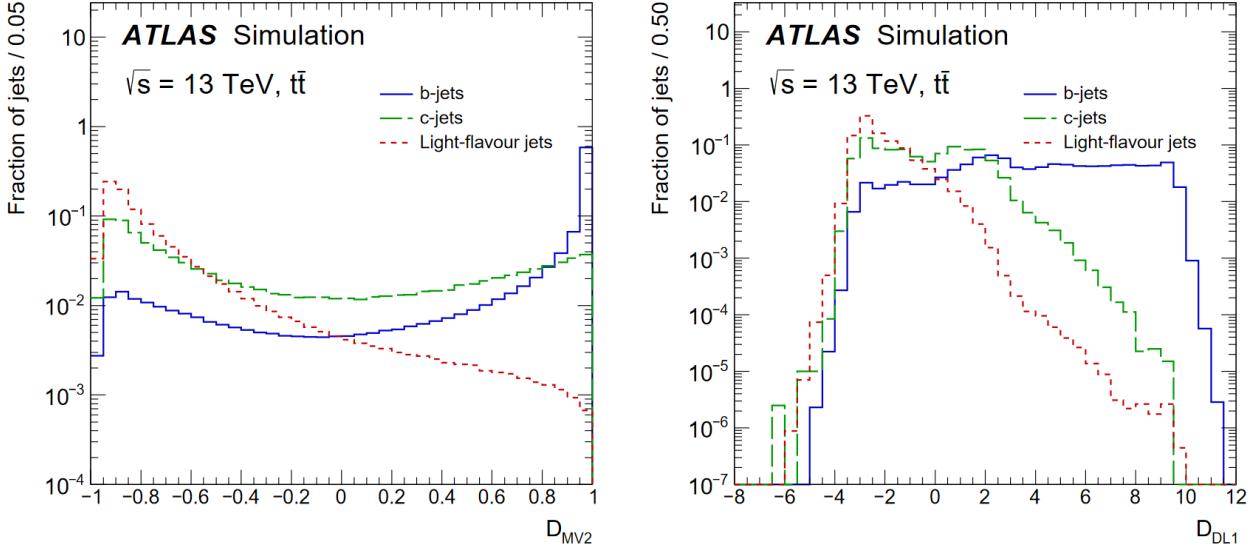


Figure 19: The fraction of light-flavour jets and c jets versus the b jets in the MV2 (left) and DL1 (right) b -tagging algorithms output distribution evaluated on the $t\bar{t}$ events [34].

5.6 Missing transverse energy

5.7 Hadronically decaying τ lepton

6 Charm jet mis-tagging calibration

MC simulations are not able to model exactly the performance of the b -tagging algorithms in data. For this reason calibration is required, i.e. correcting MC to recover the data in terms of b -tagging efficiency, charm jet mis-tagging and light jet mis-tagging rates [34]. The calibration is performed for all supported jet collections(TODO: refer back to the object definition chapter) and working points, which are cuts in the b -tagging algorithm output identifying the different tagging efficiencies and corresponding light jet and c jet rejection rate. In general, the efficiency is calculated with data and simulations, and scale factors are then calculated to match the efficiency extracted from simulations to the data.

6.1 Calibration methods for b jet and light jet

In general, the efficiency is calculated with data and simulations, and scale factors are then calculated to recover the efficiency extracted from simulations to the data. The production of $t\bar{t}$ pairs at the LHC provides an abundant source of b jets by virtue of the high cross-section and the $t \rightarrow Wb$ branching ratio being close to 100%. A very pure sample of $t\bar{t}$ events can be selected by requiring that both W bosons decay leptonically, referred to as di-leptonic $t\bar{t}$ decays in the following. For the b jet calibration, the performance of the b tagging algorithms

is evaluated in the simulation and the efficiency with which these algorithms identify jets containing b -hadrons is measured in collision data. The measurement uses a likelihood-based method in the di-leptonic $t\bar{t}$ sample, where events with exactly 2 jets and 2 opposite-sign leptons are selected. The data b jet efficiency is then extracted from a combined likelihood fit, and subsequently compared with that predicted by the simulation. Scale factors are then calculated to emulate the performance of the algorithms to the data [34].

For the light jet mis-tagging calibration, two methods are used to measure the mis-tagging rate from the data [38]. The first is the negative tag method, which uses a high statistics data sample enriched in light jets with the application of a modified algorithm which reverses some of the criteria used in the nominal identification algorithm. The second is the adjusted Monte Carlo (adjusted-MC) method, which adjusts the characteristic track observables in the simulation to im the data, and then compares the adjusted simulation to the “standard” simulation. The scale factors are then calculated using the these two methods. The scale factors of the two different methods are in good agreement within the systematic uncertainties.

6.2 Calibration method for charm jet

It is worth mentioning that the author’s qualification task to become an ATLAS author is to calibrate the rate of a charm jet being mis-identified as a b jet which is a part of the calibration of the b -tagging algorithm. During the task the calibration range has been extended down to 20 GeV (previously 25 GeV) in jet p_T and a new selection category has been developed to increase the data statistics of the scale factors in the high- p_T ($p_T^{jet} > 70$ GeV) region. The calibration is performed on the PFlow jets and the VR-Track jets (TODO: refer back to definition in Object reconstruction section).

As determined by the CKM matrix [39, 40], the W boson decays dominantly to a pair of light quarks (u quark and d quark) or to a s quark and a c quark. The W boson decays very rarely to pairs containing a b quark. More specifically, the branching ratio of a W boson decays to a u quark and d quark pair or a s quark and c quark pair is 33.1%, and to pairs containing a b quark is only 0.057% [41]. Therefore, b -tagged jets from the W decay are most likely to be mis-tagged c jets or light jets.

Furthermore, given the ratio between the DL1 light jet rejection and the corresponding charm jet rejection ranges from 10 to 40 (Figure 18), the c jet is much more likely to be mis-tagged than the light jet. This allows for a source of mis-tagged c jets to be obtained in the $t\bar{t}$ events, requiring that one W boson decays leptonically and the other decay hadronically (referred to as semi-leptonic $t\bar{t}$ decay in the following), where the b -tagged jets from the W decay are candidates of mis-tagged c jets. Requiring a W boson decaying leptonically reduces the number of combinations of jets of different flavour, and allows triggering with the lepton.

The events kinematics are shown by the diagram in Figure 20, where the $t\bar{t}$ pair decays to a b and a \bar{b} quark, circled in red. One of the W bosons, circled in blue, decays hadronically to quarks, and the other W boson decays leptonically to either an electron or a muon and the corresponding neutrinos, circled in green and purple, respectively. The lepton in the final state is used for triggering. The following notation will be used: the jets that are the decay products of the W boson are referred to as W jets and the remaining two jets are referred to as top-jets.

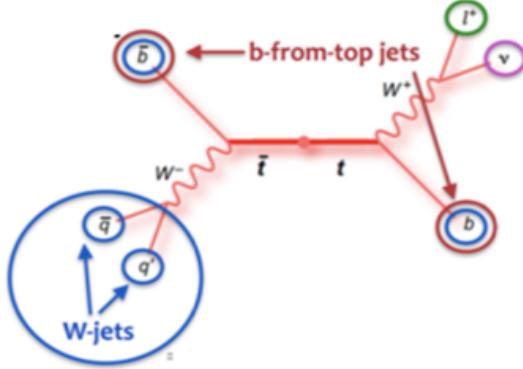


Figure 20: Feynman diagram of the semi-leptonic $t\bar{t}$ events.

A kinematic likelihood technique, referred to as KLFitter [42], is used to assign jets to the proper $t\bar{t}$ decay product (more details in Section 6.4). The following notation will be used: the jets that are assigned as the decay products of the W boson are referred to as W jets and the remaining two jets are referred to as top jets.

The charm jet efficiency is defined as the ratio of events with either of the W jet is tagged. The efficiency is evaluated in four p_T intervals, with boundaries of 20, 40, 65, 140 and 250 GeV for the PFlow jets and 15, 20, 40, 140 GeV for the VR-Track jets; and for four tagging intervals with the boundaries of 85%, 77%, 70% and 60%.

The choice of the bin boundaries ensures enough statistics for each bin and hence relatively flat statistical uncertainty, given the underlying charm-jet p_T spectrum as shown in Figure 26. The boundaries for the VR-Track jets are lower than for PFlow jets, since the track jets miss the neutral particles the reconstructed energy is significantly below the true jet energy.

The main method described in the chapter is for the “fixed-cut” calibration where the efficiency is defined as the fraction of b jets passing the tagger. Jets are said to be tagged (untagged) at particular working point if they have DL1r scores greater (less) than the DL1r score of that working point. The events with both W jet are discarded to simplify the fit described in the following.

To extract the scale factors of the charm jet mis-tagging, a fit is performed by minimising

the χ^2 defined as:

$$\begin{aligned} \chi^2 = & \sum_{t=1}^4 \sum_{i=1}^4 \sum_{j=1}^4 (N_{\text{data}}^t(i,j) - p(i,j)[c^t(i)N_C^t(i,j) + N_J^t(i,j) + \sum_k c^4(k)N_X^t(i,j,k)])^2 / N_{\text{data}}^t(i,j) \\ & + \sum_{i=1}^4 \sum_{j=i}^4 [N_{\text{data}}^{\text{untag}}(i,j) - p(i,j)N_{\text{MC}}^{\text{untag}}(i,j)]^2 / N_{\text{data}}^{\text{untag}}(i,j). \end{aligned} \quad (1)$$

The $c^t(i)$ is the main floating parameter in the fit, which is the charm jet mis-tagging scale factor at working point t of p_T bin labelled i . The other main floating parameter is $p(i,j)$ that is the normalisation factor scaling the MC to data. The $N_{\text{data}}^t(i,j)$ is the number of data events with a tagged W jet in the p_T bin labelled i and the other (un>tagged) W jet in the p_T bin labelled j . Similarly the $N_C^t(i,j)$ is the number of MC events with a tagged W jet while the tagged W jet is indeed a c jet which can be seen as “signal”. In contrast the $N_J^t(i,j)$ is the number of events with neither the tagged W jet nor the top jets are c jets and the $N_X^t(i,j,k)$ is the number of events with one of the top jets is a c jet. These two types of events can be seen as “background”. The later case is slightly more complicated, as the c jet lies in a different p_T bin to the tagged jet, denoted as k , and it only depends on the $c^4(k)$ (which is the scale factor of the 4th working point i.e. 60%) as the top jets are tagged at 60% working point. The calibration is then given as the scale factors of the four working points in bins of p_T defined in the above text.

6.3 Data and Monte Carlo samples

TODO: remove the overlap between this section and the Data MC chapter in the thesis
Dedicated MC are used to model SM processes. The data analysed in this study correspond to 139 fb^{-1} [43, 44, 45, 46], of pp collision data collected by the ATLAS detector between 2015 and 2018 with a centre-of-mass energy of 13 TeV. The data sample was collected using a set of single-muon [47] and single-electron triggers [48]. The single-muon triggers had p_T thresholds in the range 20–26 GeV for isolated muons and 50 GeV for muons without any isolation requirement. The single-electron triggers employed a range of p_T thresholds varying between 24–300 GeV and a combination of quality and isolation requirements depending on the data-taking period and the p_T threshold. All detector subsystems were required to be operational during data taking and to fulfil data quality requirements.

All samples were produced using the ATLAS simulation infrastructure [49] and GEANT4 [50]. A subset of samples use a faster simulation based on a parameterisation of the calorimeter response and GEANT4 for the other detector systems [49]. The simulated events are reconstructed with the same algorithms as used for data, and contain a realistic modelling of

pile-up interactions. The pile-up profiles in the simulation match those of each dataset between 2015 and 2018, and are obtained by overlaying minimum-bias events, simulated using the soft QCD processes of PYTHIA 8 [51] using the NNPDF2.3LO set of PDFs [52] and a set of tuned parameters called the A3 tune [53].

The events that are used in this study originate mostly due to $t\bar{t}$ production. This process is modelled using the POWHEGBOX v2 [54, 55, 56, 57] generator at NLO with the NNPDF3.0nlo parton distribution function (PDF) set and the h_{damp} parameter¹ set to 1.5 m_{top} [58]. The events were interfaced to PYTHIA 8.230 to model the parton shower, hadronisation, and underlying event, with parameters set according to the A14 tune and using the NNPDF2.3l0 set of PDFs. The decays of bottom and charm hadrons were performed by EVTGEN v1.6.0 [59]. The simulated $t\bar{t}$ events are split according to the origin of W jets. The notation “ $t\bar{t}, \text{ll}$ ” denotes that both W jets are light flavour jets. Similarly, “ $t\bar{t}, \text{cl}$ ” (“ $t\bar{t}, \text{bl}$ ”) indicates that one of the W jets is a c jet (b jet) whereas the other is a light flavour jet. W jets with origin other than what is discussed above fall into the category denoted by “ $t\bar{t}, \text{other}$ ”. This category includes events in which at least one of the W jets comes from a hadronically decaying τ -lepton.

In addition to $t\bar{t}$ production, there are some minor backgrounds that contribute to the final event sample that is used for the calibration. These backgrounds consist mostly of single-top and diboson production, the production of $t\bar{t}$ in association with a vector boson and the production of a vector boson in association with jets. The details of the modeling of these samples are given in the following.

Single-top s -channel production is modelled using the POWHEGBOX v2 generator at NLO in QCD in the five-flavour scheme with the NNPDF3.0nlo [60] parton distribution function (PDF) set. The associated production of top quarks with W bosons (tW) is modelled using the POWHEGBOX v2 [61, 55, 56, 57] generator at NLO in QCD using the five-flavour scheme and the NNPDF3.0nlo set of PDFs [60]. The diagram removal scheme [62] is used to remove interference and overlap with $t\bar{t}$ production. The events for both single-top s -channel and tW production are interfaced to PYTHIA8.230 using the A14 tune and the NNPDF2.3l0 set of PDFs.

The production of $Z+\text{jets}$ and $W+\text{jets}$ is simulated with the SHERPA v2.2.1 [63] generator using next-to-leading order (NLO) matrix elements (ME) for up to two partons, and leading order (LO) matrix elements for up to four partons calculated with the Comix [64] and OPENLOOPS [65, 66, 67] libraries. They are matched with the SHERPA parton shower [68] using the MEPS@NLO prescription [69, 70, 71, 72] using the set of tuned parameters developed by the SHERPA authors. The NNPDF3.0nnlo set of PDFs [60] is used and the samples

¹The h_{damp} parameter is a resummation damping factor and one of the parameters that controls the matching of POWHEG matrix elements to the parton shower and thus effectively regulates the high- p_T radiation against which the $t\bar{t}$ system recoils.

are normalised to a next-to-next-to-leading order (NNLO) prediction [73].

Samples of diboson final states (VV) are simulated with the SHERPA v2.2.1 or v2.2.2 [63] generator depending on the process, including off-shell effects and Higgs-boson contributions, where appropriate. Fully leptonic final states and semileptonic final states, where one boson decays leptonically and the other hadronically, are generated using matrix elements at NLO accuracy in QCD for up to one additional parton and at LO accuracy for up to three additional parton emissions. Samples for the loop-induced processes $gg \rightarrow VV$ are generated using LO-accurate matrix elements for up to one additional parton emission for both cases of fully leptonic and semileptonic final states. The matrix element calculations are matched and merged with the SHERPA parton shower based on Catani-Seymour dipole factorisation [64, 68] using the MEPS@NLO prescription [69, 70, 71, 72]. The virtual QCD correction are provided by the OPENLOOP library [65, 66, 67]. The NNPDF3.0nnlo set of PDFs is used, along with the dedicated set of tuned parton-shower parameters developed by the SHERPA authors.

The production of $t\bar{t}$ in association with a vector boson is modelled using the MADGRAPH5_aMC@NLO v2.3.3 [74] generator at NLO with the NNPDF3.0nlo [60] parton distribution function (PDF). The events are interfaced to PYTHIA8.210 [51] using the A14 tune [75] and the NNPDF2.3llo [60] PDF set. The decays of bottom and charm hadrons are simulated using the EVTGEN v1.2.0 program [76].

6.4 Kinematic Likelihood Fitter

The four-vectors of the four highest p_T jets, the lepton and the event E_T^{miss} are used as inputs to a likelihood-based $t\bar{t}$ event reconstruction algorithm, which is described in more detail in Ref. [42]. This algorithm uses a likelihood function to assign the four jets to the $t\bar{t}$ decay topology. In particular, the algorithm assigns one jet to be the b jet from the leptonically decaying top quark ($t \rightarrow Wb \rightarrow \ell\nu b$), another to the b jet from the hadronically decaying top quark ($t \rightarrow Wb \rightarrow qq'b$, where qq' are the quarks in which the W boson decays) and the remaining two jets to the jets that come from the hadronic W boson decay. The jet assignment does not use any b -tagging information to avoid bias.

6.5 Maximising likelihood

Taking only four jets in the event limits the total number of possible jet orderings (permutations) in the event. In the semi-leptonic channel, four jets can be permuted a total number of times equal to $4! = 24$. However, the two W jets are kinematically indistinguishable. This reduces the possible number of permutations to 12. For every combination of jet ordering, the likelihood is maximised over its free parameters, the energy of the four jets, the lepton energy and the three components of the momentum of the neutrino, and provides a value

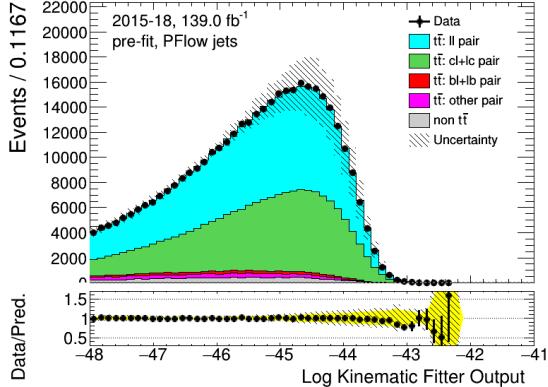


Figure 21: Distribution of the negative logarithm of the likelihood that is used to reconstruct the $t\bar{t}$ decay.

based on how closely the kinematic information from the reconstructed objects for a specific jet ordering resembles the expected kinematic behaviour of the decay of a Standard Model semi-leptonic $t\bar{t}$ event. The likelihood therefore distinguishes the possible permutations on an event-by-event basis. The best permutation, given by the largest log-likelihood value, is adopted as the jet ordering for the event. An additional requirement of $\log\text{-likelihood} > -48$ is placed on the output of the likelihood value for the chosen event permutation. An example of the distribution of log-likelihood of the best permutations is shown in Figure 21. In this figure, the data events are compared against the simulation. The majority of the events come from $t\bar{t}$ production. There is only a very small fraction of events, which is denoted as “non $t\bar{t}$ ” on the figure, that come from other processes like W or Z production in association with jets or single-top production.

6.6 Event selection

6.6.1 Standard selection

	PFlow jets	Track jets
Data	227118	218351
$t\bar{t}$	235670 ± 200	223770 ± 180
Non $t\bar{t}$	7610 ± 120	7280 ± 100
Data/MC	0.934 ± 0.002	0.945 ± 0.002

Table 2: Standard selection: prefit comparison of the number of events in data and in simulation considering the PFlow jets and the VR-Track jets for events with exactly 4 jets.

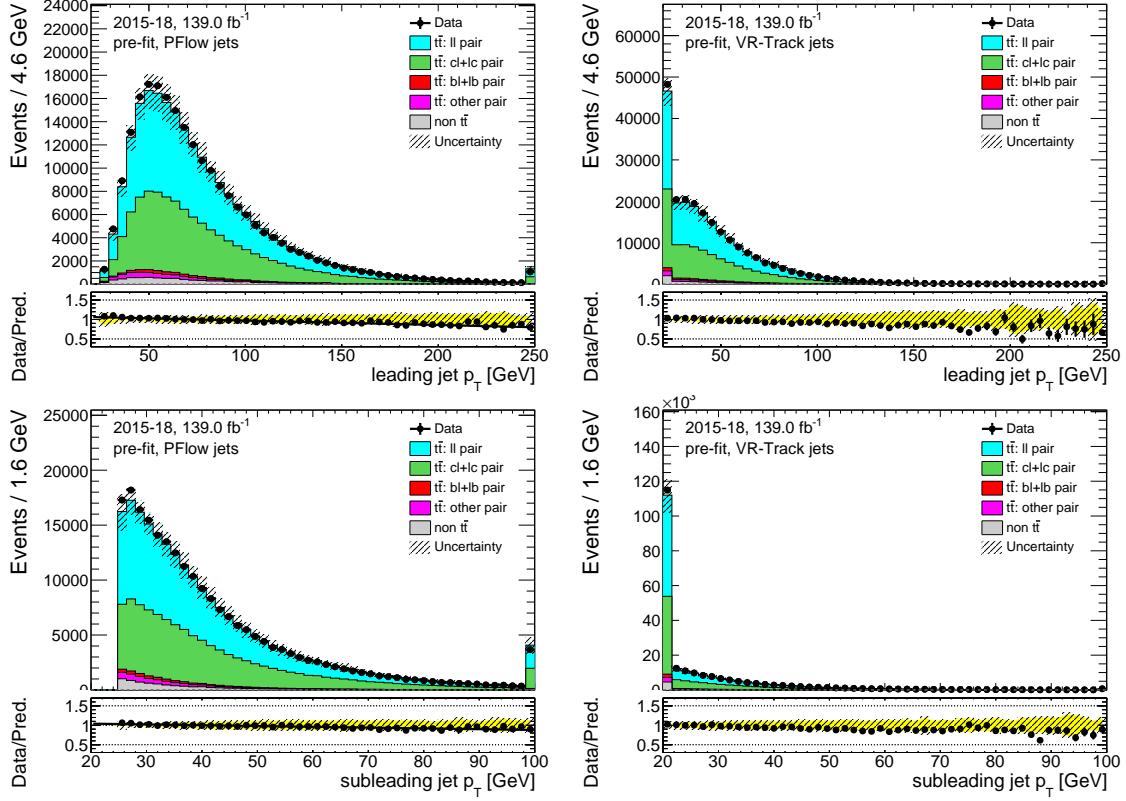


Figure 22: Standard selection: data versus simulation of the leading and sub-leading W jet p_T for the PFlow jets in the left column and for VR-Track jets in the right column. The leading jet and sub-leading jet refer to the highest p_T W jet and the second highest p_T jet, respectively. The 'non $t\bar{t}$ ' background indicates background comes from non- $t\bar{t}$ processes like W or Z production in association with jets or single-top production. The error in the table (and the following yields tables for different selection) is stats-only.

Events are required to contain exactly one trigger-matched lepton with $p_T > 27$ GeV and exactly four jets with $p_T > 25$ GeV. Leptons are required to have p_T above 27 GeV in order to avoid the turn-on curve for the single lepton triggers. Events which contain an additional lepton with $p_T > 27$ GeV are rejected. The events are also required to have $E_T^{\text{miss}} > 20$ GeV, which is assumed to be the result of the neutrino from the leptonically decaying W boson. The transverse mass m_T between the lepton and the E_T^{miss} , is constrained as follows:

$$m_T = \sqrt{2p_T^\ell E_T^{\text{miss}}(1 - \cos \Delta\phi)} > 40 \text{ GeV},$$

where $\Delta\phi = \phi(E_T^{\text{miss}}) - \phi(\ell)$ is the azimuthal difference between the lepton and E_T^{miss} . The yields of the data and the MC are given in Table 2. An example of the p_T distributions before any tagging or fitting and after the standard selection is shown in Figure 22. More plots can be found in Appendix A.1.1. The yellow band in the lower pad shows the overall systematic

uncertainties, combining the experimental uncertainties and the $t\bar{t}$ modelling uncertainties, as described in Section 6.7. The data/MC ratio shows good agreement within the systematic uncertainties.

6.6.2 Low- p_T selection

The author has developed an orthogonal selection to extend the calibration in the low- p_T region so that the calibration can be applied to PFlow jets with $20 < p_T < 25$ GeV. The p_T threshold of the VR-Track jets is 10 GeV therefore the low- p_T selection is not needed. Instead of requiring events to have exactly 4 jets $p_T > 25$ GeV, events are required to have exactly 3 jets with $p_T > 25$ GeV and exactly 1 jet with $25 \text{ GeV} > p_T > 20$ GeV. Other than that, all requirements for the selection are the same. This additional cut provides candidates for the PFlow W jet that is used for calibration in the $20 - 25$ GeV region. The inclusive yields of the low- p_T selection of the data and the MC are given in Table 3, and the p_T distributions of the W jets are shown in Figure 23. More plots of the kinematic distributions are shown in Appendix A.1.2. Good agreement between MC and data is shown in these distributions, and the p_T range of the sub-leading has gone down to 20 GeV.

	PFlow jets
Data	59987
$t\bar{t}$	56530 ± 90
Non $t\bar{t}$	3340 ± 60
Data/MC	1.002 ± 0.004

Table 3: Low- p_T selection: prefit comparison of the number of events in data and MC for the PFlow W jets. Events are required to have exactly 3 jets with $p_T > 25$ GeV and one jet with $20 < p_T < 25$ GeV.

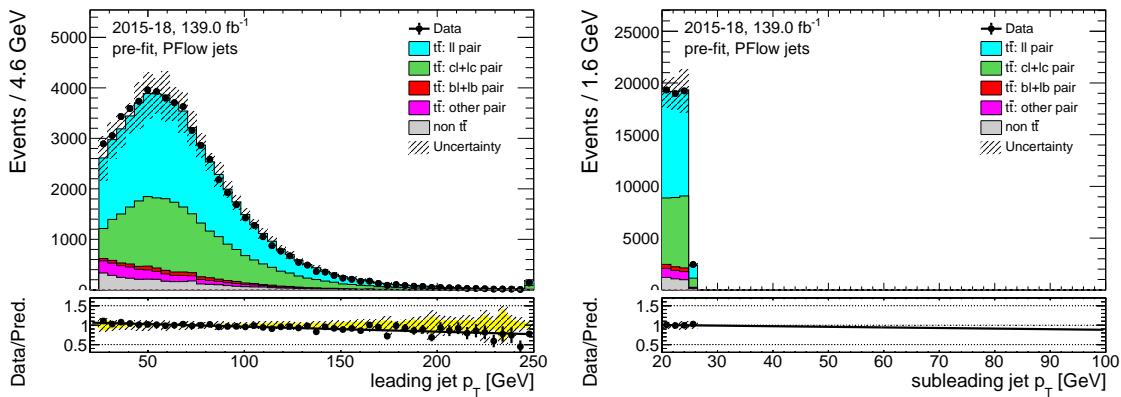


Figure 23: Low- p_T selection: data versus simulation of the PFlow W jets p_T .

6.6.3 High- p_T selection

It has been observed that in the previous calibrations that the statistics are relatively low for the high- p_T region (e.g. jet $p_T > 100$ GeV). Therefore, the author has worked on an orthogonal selection to improve this situation. Instead of requiring events to have exactly 4 jets, events are required to have at least 5 jets with $p_T > 25$ GeV, in which at least 1 jet with $p_T > 70$ GeV. Other than that, all requirements for the selection remain the same.

The choice of cut value at 70 GeV is based on the study shown in the following. The effect on the c jet purity and the potential statistical gain is investigated, where the c jet purity is defined as:

$$c \text{ jet purity} = \frac{N_{\text{true } c \text{ jet}}}{N_{\text{all}}}, \quad (2)$$

where $N_{\text{true } c \text{ jet}}$ stands for the number of events with a true c jet from the W decay, and N_{all} stands for the number of all events. The ideal situation is the high- p_T selection will maximally increase the statistics while minimally decreasing the c jet purity, therefore a figure of merit P^{Cut} is defined as:

$$P^{\text{Cut}} = \frac{\sum_i \text{Gain in stats}_i^2}{\sum_i c \text{ jet purity}_i^2},$$

where i stands for the number of bins. The “Gain in stats” stands for increase in statistics and it’s summed over all bins in Figure 24. The c jet purity and the statistical gain are calculated for 4 different cut values as shown in Figure 24, comparing with the cut value of 0. The value of 70 GeV is chosen as it gives the highest value of P^{Cut} .

	PFlow jets	Track jets
Data	98273	83957
$t\bar{t}$	99430 ± 120	87476 ± 110
Non $t\bar{t}$	1842 ± 21	1570 ± 20
Data/MC	0.97 ± 0.003	0.94 ± 0.003

Table 4: High- p_T selection: prefit comparison of the number of events in data and in simulation considering the PFlow W jets and the VR-Track jets.

The yields of the data and the MC are given in Table 4. An example of the p_T distributions before any tagging or fitting, applying the high- p_T selection is shown in Figure 25. In general the event statistics improve about 80% in region with $p_T > 70$ GeV as desired. More plots can be found in Appendix A.2.

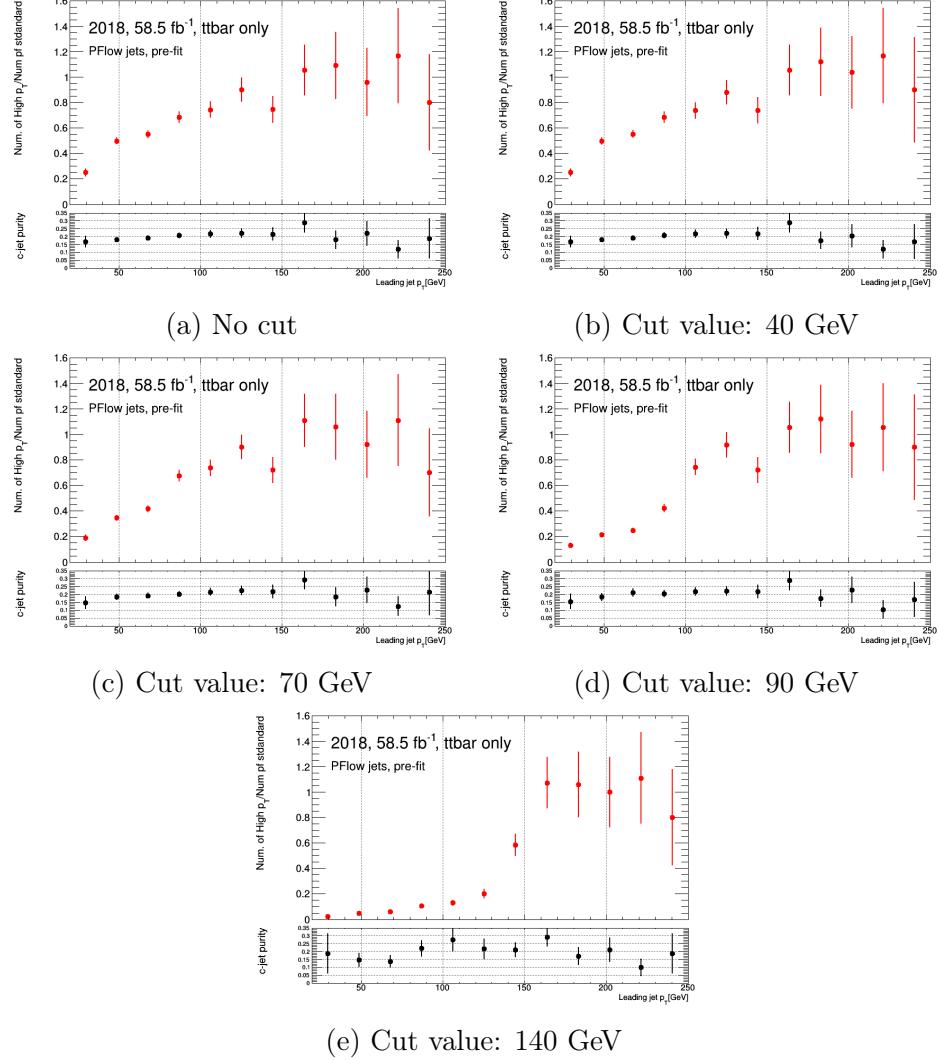


Figure 24: Comparison of different cut values in terms of gain in stats and c jet purity.

6.6.4 Combined selection

As the standard selections, low- p_T selection and high- p_T selection are orthogonal to each other, all the selections are combined to provide the maximum range and statistics for the calibration. The yields of the data and the MC are given in Table 5, an example of the p_T distributions before any tagging or fitting and after the combined selection is shown in Figure 26. More plots can be found in Appendix A.3.

6.7 Systematic uncertainties

The systematic uncertainties considered and propagated in this calibration can be broadly categorised into experimental and modelling systematic uncertainties.

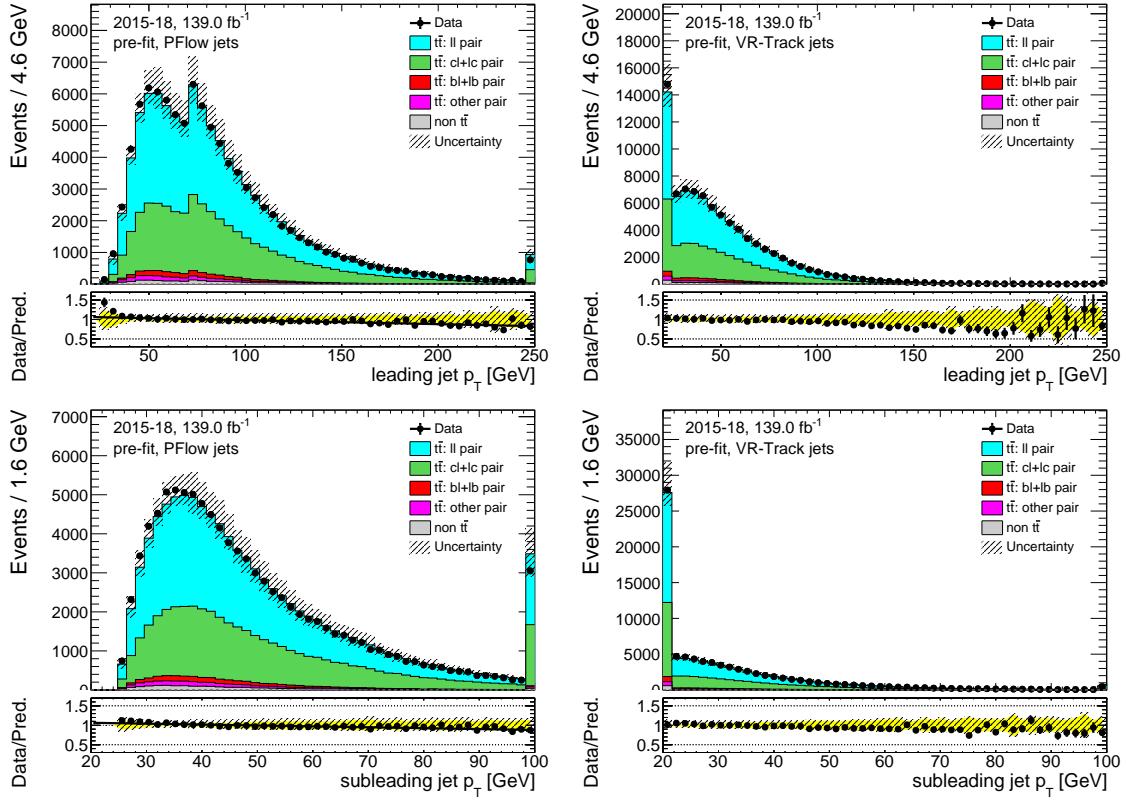


Figure 25: High- p_T selection: data versus simulation of W jets p_T for PFlow jets in the left column and for VR-Track jets in the right column.

6.7.1 Experimental uncertainties

TODO: refer to the analysis Part Experimental uncertainties are related to the detector and estimated using data-driven methods or MC simulations. The lepton energy scale and resolution are corrected to provide better agreement between MC predictions and data, uncertainties due the corrections are considered. Uncertainties are taken into account on the electron and muon trigger, identification and reconstruction efficiencies, and for uncertainties associated with the isolation requirements.

The jet energy scale (JES) uncertainty depends on p_T and η and takes into account uncertainties due to pile-up effects. Uncertainties on the jet energy resolution (JER) are taken into account. Uncertainties on the energy scale and resolution of the electrons, muons, jets and taus are propagated to the calculation of the E_T^{miss} , which also has additional dedicated uncertainties on the scale, resolution, and reconstruction efficiency of tracks not associated to any of the reconstructed objects, along with the modelling of the underlying event. Uncertainties on the b -tagging (mis-tagging) probabilities for b (light) jets are considered both for the tagging jets assigned to the b quark from the top decay and for the jets associated to the hadronically decaying W boson. Supporting material for this section can be found in

	PFlow jets	Track jets
Data	385378	302308
$t\bar{t}$	383520 ± 230	302690 ± 200
Non $t\bar{t}$	12420 ± 120	8570 ± 100
Data/MC	0.973 ± 0.002	0.971 ± 0.002

Table 5: Combined selection: prefit comparison of the number of events in data and in simulation considering the PFlow jets and the VR-Track jets for an inclusive selection.

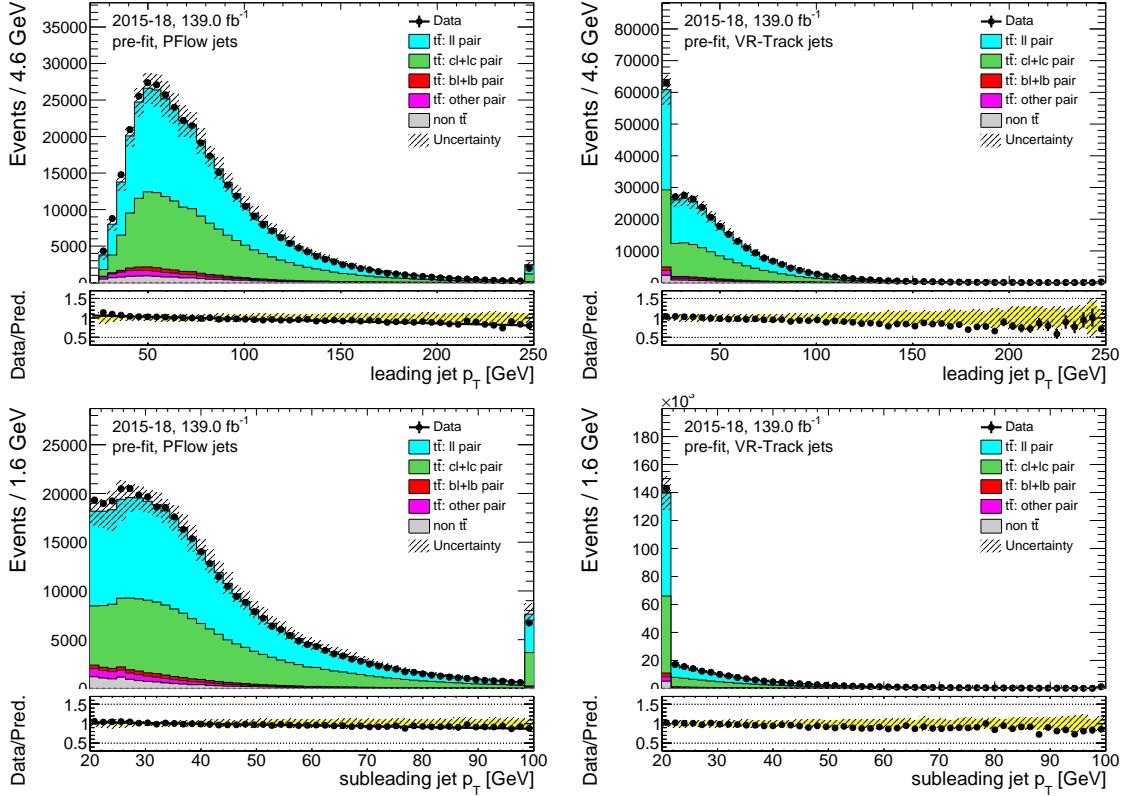


Figure 26: Combined selection: data versus simulation of W jets p_T for PFlow jets in the left column and for VR-Track jets in the right column.

the appendix, Tab.9.

6.7.2 Modelling uncertainties

The uncertainty due to different choices of the parton shower models is estimated by comparing the MC samples generated with nominal parton shower model and with the alternative parton shower model. More specifically, it is derived by comparing the prediction from POWHEG interfaced either to PYTHIA8 or Herwig++. The uncertainty due to additional radiation in the initial state and the final state is estimated by comparing the nominal MC samples with the MC samples with alternative scale of renormalisation and factorisation.

The uncertainty on modelling of initial state radiation (ISR) is assessed with two alternative POWHEG+PYTHIA8 samples. The samples include one with an increase in radiation which has the renormalisation and factorisation scales decreased by a factor of two and the $hdamp$ parameter doubled (which controls the p_T of the first additional emission), while the sample with a decrease in radiation has the scales increased by a factor of two. In all cases, MC-to-MC SFs are taken into account. In addition, the uncertainty due to the variations samples being produced by fast simulation while the nominal samples being produced full simulation is also considered. The comparisons of the nominal $t\bar{t}$ sample and the samples with each systematic uncertainty are shown in Table 6.

	PFlow jets		Track jets	
	Yields	Ratio of difference to nominal sample	Yields	Ratio of difference to nominal sample
$t\bar{t}$ Nominal Data/MC	385378 ± 230 0.973 ± 0.002		302690 ± 200 0.971 ± 0.002	
$t\bar{t}$ AF2 DATA/MC(AF2)	386260 ± 250 0.967 ± 0.002	0.716%	304860 ± 230 0.965 ± 0.002	0.716%
$t\bar{t}$ ISR DATA/MC(ISR)	377130 ± 220 0.989 ± 0.002	-1.665%	297960 ± 200 0.986 ± 0.002	-1.562%
$t\bar{t}$ Herwig DATA/MC(Herwig)	331960 ± 220 1.119 ± 0.002	-13.443%	259940 ± 190 1.126 ± 0.002	-14.123%

Table 6: Comparison of the number of events in data and in simulation considering the PFlow jets and the VR-Track jets for an inclusive selection. The uncertainty due to the variations samples being produced by fast simulation is included in the table as $t\bar{t}$ AF2.

6.8 Under-estimation of $t\bar{t} +$ Heavy flavour background

Despite the fact that the true nature of most of the reconstructed W jets are either c jets or light jets, there is still a very small amount of them are true b jets.

There are two main sources of these true b jets. The first is a W boson decays to a b and a c quark. The second is when the $t\bar{t}$ plus a gluon process (referred to as $t\bar{t} +$ heavy flavour process) is selected, and the gluon splits into a pair a b quarks and one of them is assigned as a W jet. The first source can be excluded by requiring no c jets in the W jets, meaning the true b jet in the W jets can only come from the $t\bar{t} +$ heavy flavour process. This process is underestimated by the MC by about 30% for both the PFlow and VR-Track jets collections, as shown in Table 7 and Figure 27, where an extra cut requiring at least one W jet with $DL1r > 8$ is added to the combined selection to reject most of the true c jets and true light jet. A more thorough study is done in Ref. [77], where the mis-modelling factor is measured to be 1.25 ± 0.25 , which is also consistent with the 30% mismodelling observed

in the previous study. Therefore, events in the simulation in which the top jets and at least one of the W jets are b jets(referred to as 3 true b jets events), are scaled by 1.25 ± 0.25 . All results shown in this chapter have this scale factor implemented, and the full difference between the simulation before applying this scale factor and after is taken as a systematic uncertainty. This uncertainty has been added in quadrature to the systematic uncertainties described in Section 6.7 in all the plots in this chapter.

	PFlow jets	VR-Track jets
Data	1589	1336
$t\bar{t}$	1100 ± 13	940 ± 12
Non $t\bar{t}$	83 ± 6	69 ± 5
Data/MC	1.34 ± 0.04	1.32 ± 0.04

Table 7: Yields of the 2018 data and MC of the combined selection, requiring at least 1 PFlow or track W jet with $\text{DL1r} > 8$ to reject most of the light- and c jets.

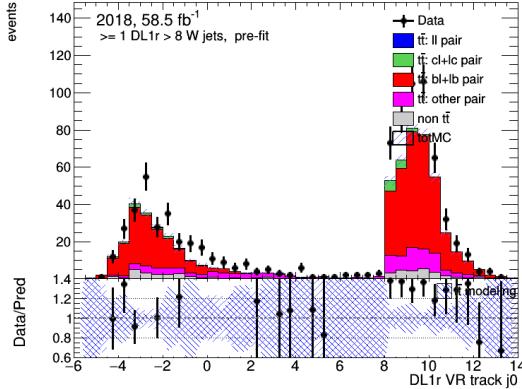


Figure 27: The DL1r score distribution of the leading VR-Track jet, requiring at least 1 VR-Track jets have $\text{DL1r} > 8$ to reject most of the light and the c jets, with $t\bar{t}$ modelling and statistical uncertainties.

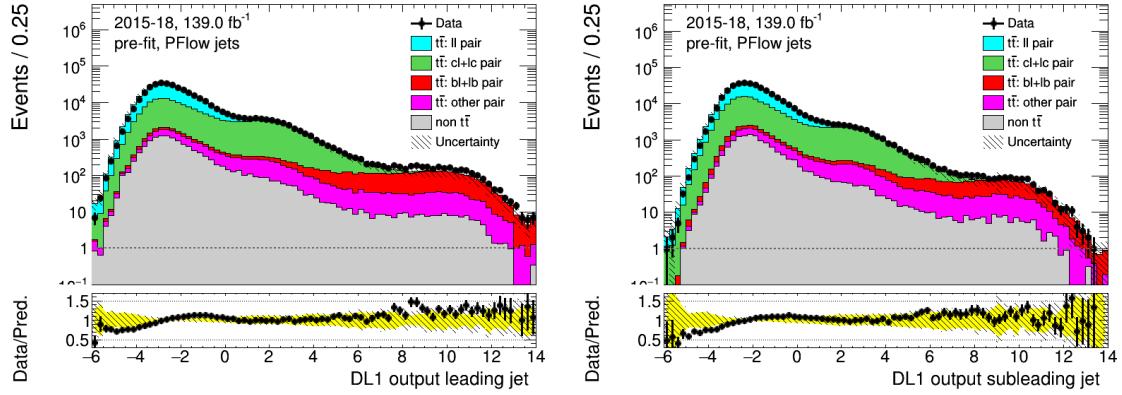
6.9 Results

6.9.1 Overview

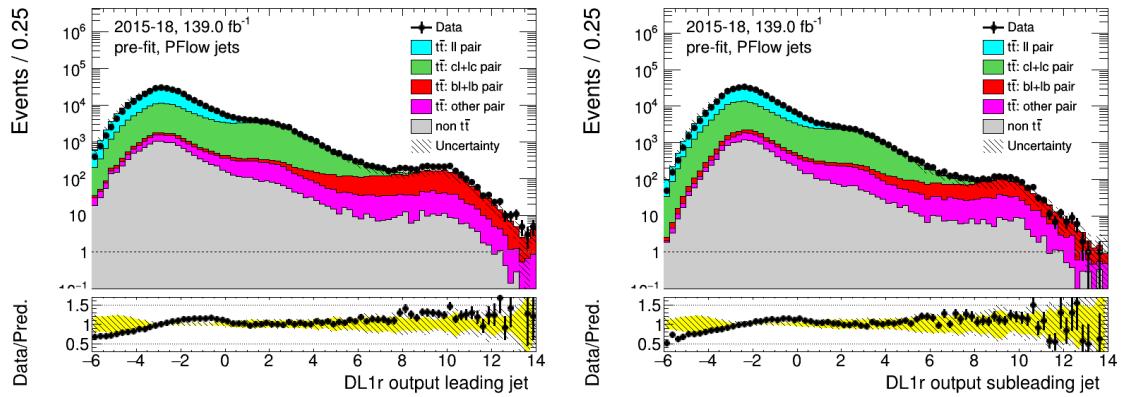
Four rounds of calibrations have been carried out, containing different jet collections, Monte Carlo samples, analysis framework and b jet identification algorithm. In the latest round, the calibration includes the PFlow jet and the VR-Track jet collection, and MV2c10, DL1 and DL1r taggers. The low- p_T selection and the standard selection are carried out for all four calibrations, while the high- p_T selection is only implemented in the latest calibration.

6.9.2 b -tagging algorithms output distribution

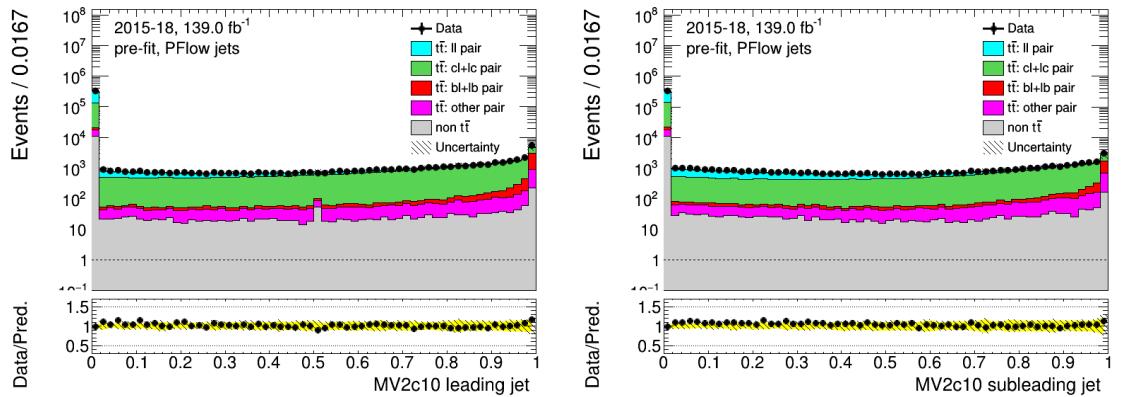
The distributions of the b -tagging algorithm' output of the MC and the data of the latest calibration (December 2020) are shown in Figure 28 for the PFlow jets and Figure 29 for the VR-Track jets, combining the standard selection, low p_T and the high- p_T selection. In these figure, the data events are compared against the simulation. The majority of the events come from $t\bar{t}$ production. There is only a very small fraction of non $t\bar{t}$ events. The W jets pairs are mostly light jets pairs and c jet light jet pairs, and a very small fraction of the pairs are b jet light jet pairs or pairs containing one or more τ hadron(s). The yellow band in the lower pad indicates the overall systematic uncertainties and the black band represents the $t\bar{t}$ modelling systematic uncertainty, which dominates at low b -tagging discriminant (DL1 or $DL1r < 4$). The experimental systematic uncertainty is in general very small. At high b -tagging discriminant (DL1 or $DL1r > 4$), the uncertainty due to the 1.25 ± 0.25 scale factor becomes more important.



(a) DL1 tagger output

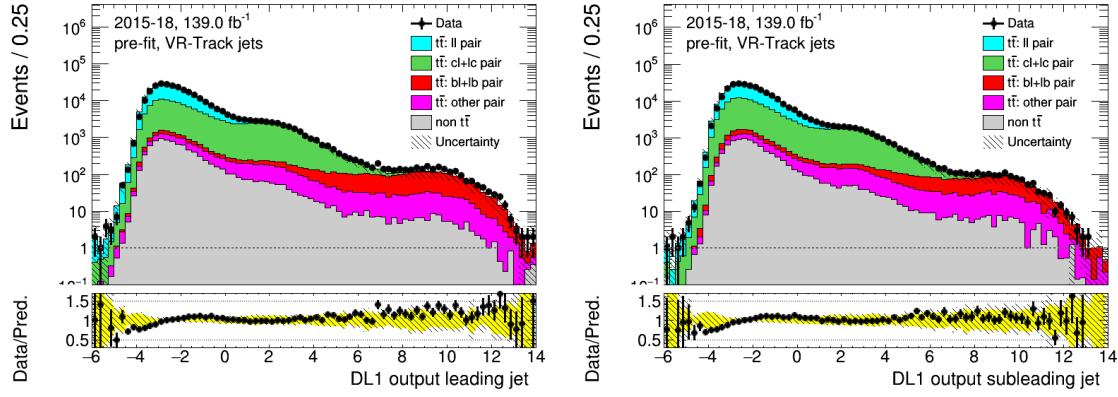


(b) DL1r tagger output

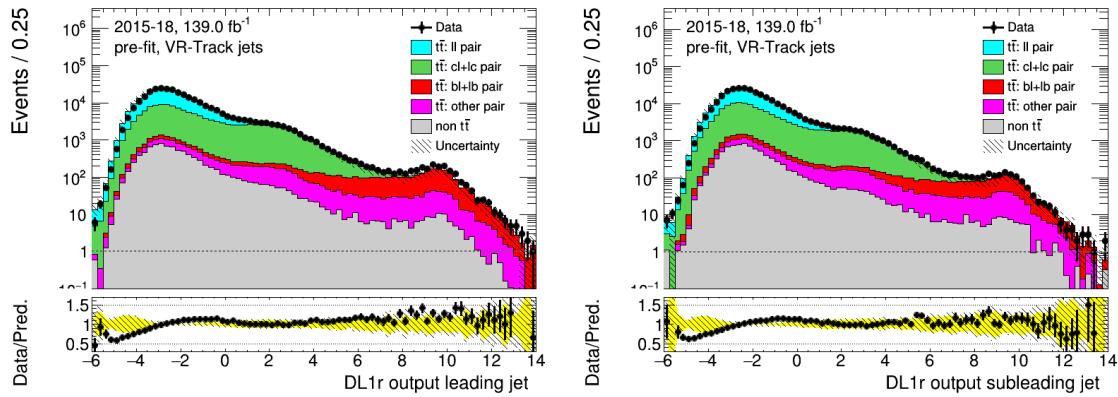


(c) MV2c10 tagger output

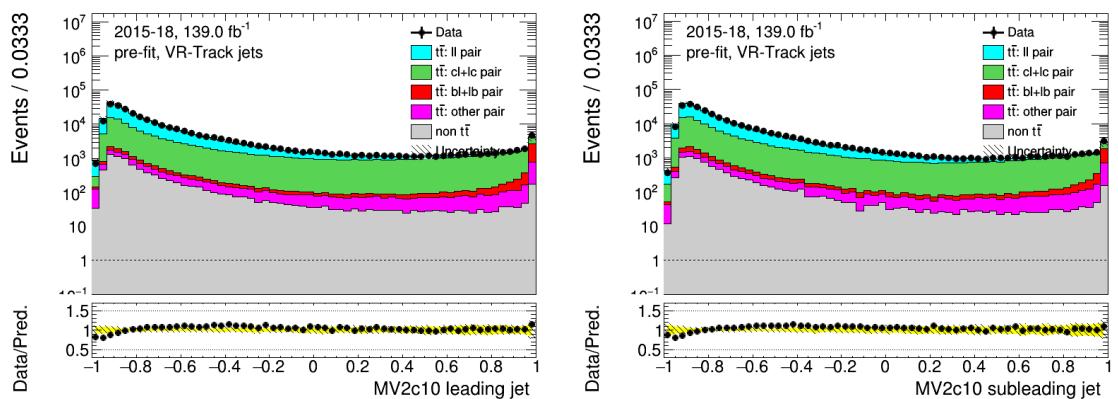
Figure 28: PFlow jets: distributions of the DL1, DL1r and MV2c10 tagger outputs of the combined selection, leading jet in the left column and sub-leading jet in the right column, before fitting or tagging with full uncertainties.



(a) DL1 tagger output



(b) DL1r tagger output



(c) MV2c10 tagger output

Figure 29: VR-Track jets: distributions of the DL1, DL1r and MV2c10 tagger outputs of the combined selection, leading jet in the left column and sub-leading jet in the right column, before fitting or tagging with full uncertainties.

6.9.3 Efficiencies and Scale Factors

The DL1 and DL1r c jet efficiencies and scale factors with systematics uncertainties are calculated with four fixed cut working points for the PFlow and VR VR-Track jets collection in the latest derivation in December 2020.

The c jet mis-tagging efficiencies are shown in Figure 30-33 for the PFlow jet collections and the VR-Track jets with the DL1 and the DL1r tagger. For PFlow jets, these results combine the standard selection, low- p_T selection and the high- p_T selection and for the VR-Track jets they combine the standard selection and the high- p_T selection.

The 1.25 ± 0.25 scale factor is applied on events with 3 true b jets. The overall uncertainties are shown in the red band. The scale factors are shown in Figure 34-37 for the PFlow jets and the VR-Track jets with the DL1 and DL1r tagger. The tighter working points (60%, 70%) show larger uncertainties and bigger deviation from 1, while the looser working points (77%, 85%) have much smaller uncertainty and the simulation is able to recover the data well due to more abundant events statistics. For the PFlow jets, in most of the working points the systematic uncertainties dominate in the low- p_T bins ($p_T < 150$) and the statistical error, represented by the error bars on the markers, become more important in the last bin. For the VR-Track jets the statistical uncertainty is relatively constant for all bins while the systematic uncertainty increases as the p_T increases. To demonstrate the effect on statistics with the high- p_T selection, the fractional statistical uncertainties of 60% working point scale factor are shown in Table 8 for the standard and the combined selection. In some bins the statistical uncertainty can decrease up to 30%, suggesting that the high- p_T selection is successful at increasing events statistics.

	PFlow jets			VR-Track jets		
	Standard selection	High- p_T selection	Fractional decrease	Standard selection	High- p_T selection	Fractional decrease
Bin No.1	3.3%	3.3%	0.0%	5.6%	5.3%	5.7%
Bin No.2	3.1%	2.8%	10.7%	4.2%	3.7%	13.5%
Bin No.3	3.4%	2.6%	30.8%	5.8%	4.9%	18.4%
Bin No.4	12.1%	9.3%	30.1%	7.2%	5.6%	28.6%

Table 8: Comparison of the fractional statistical uncertainty in the DL1r 60% working point scale factor. The p_T range of each bin can be found in section 6.2.

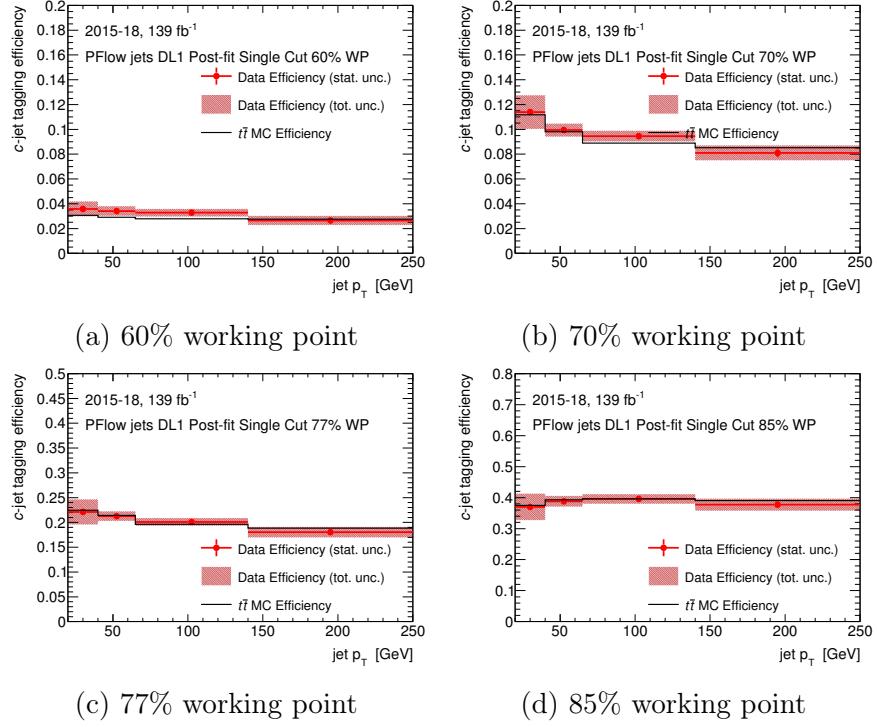


Figure 30: Charm-jet efficiencies for the PFlow jets collection with the DL1 tagger.

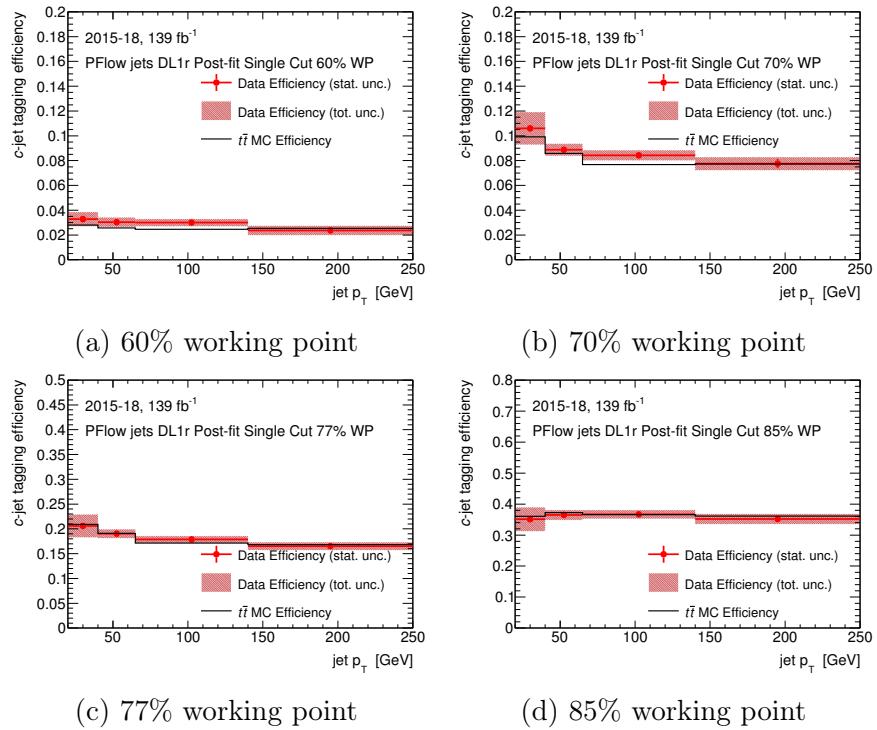


Figure 31: Charm-jet efficiencies for the PFlow jets collection with the DL1r tagger.

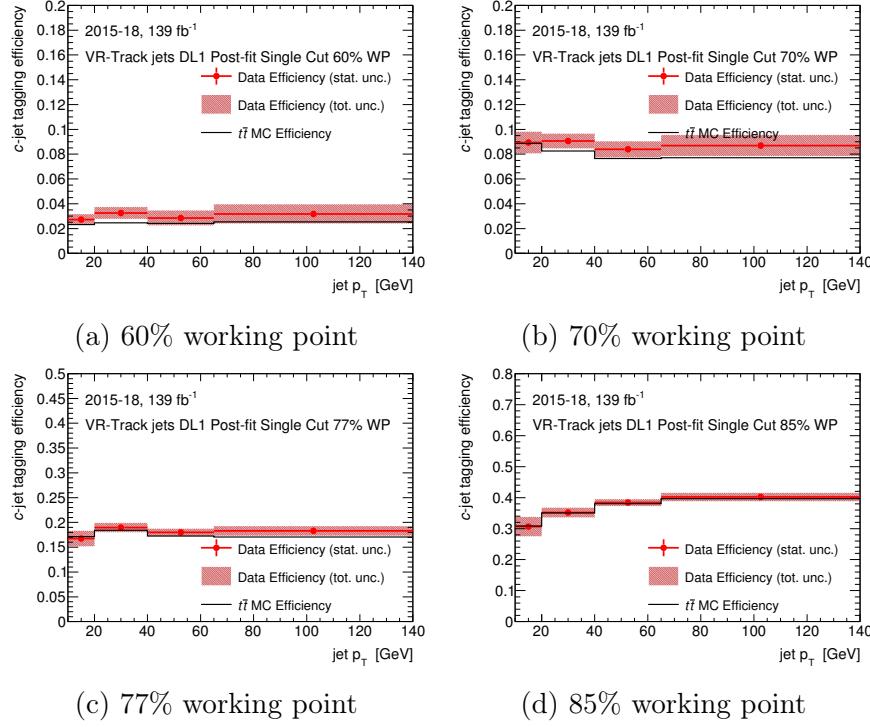


Figure 32: Charm-jet efficiencies for the VR-Track jets collection with the DL1 tagger.

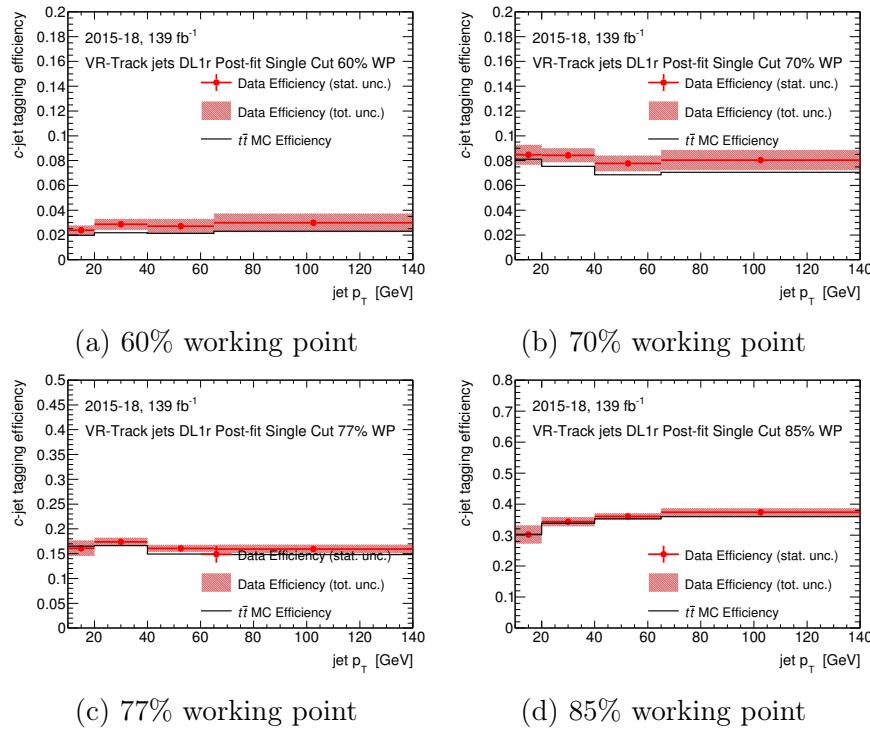


Figure 33: Charm-jet efficiencies for the VR-Track jets collection with the DL1r tagger.

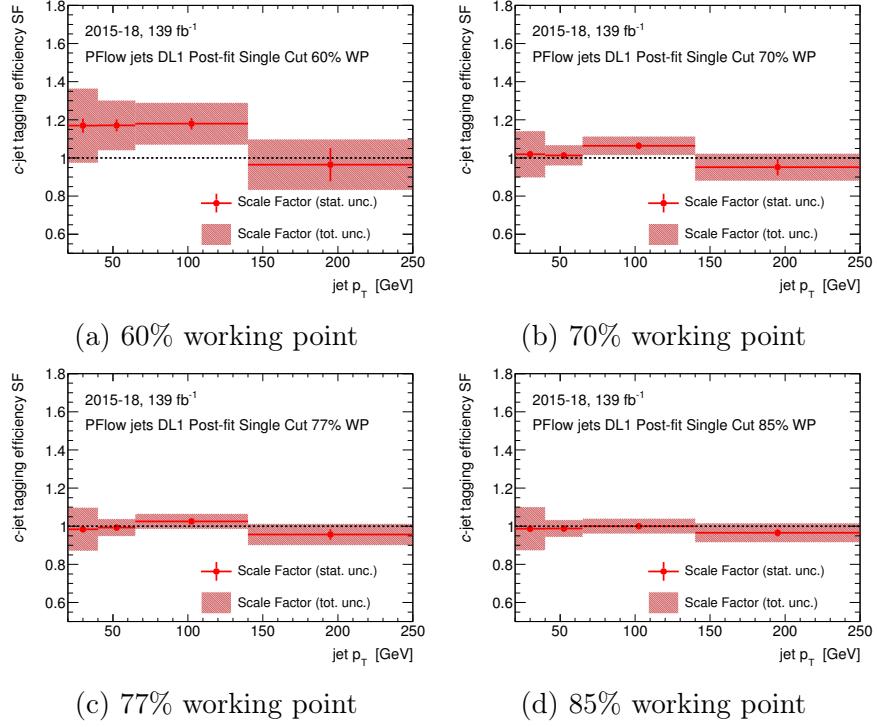


Figure 34: Charm-jet scale factors for the PFlow jets collection with the DL1 tagger.

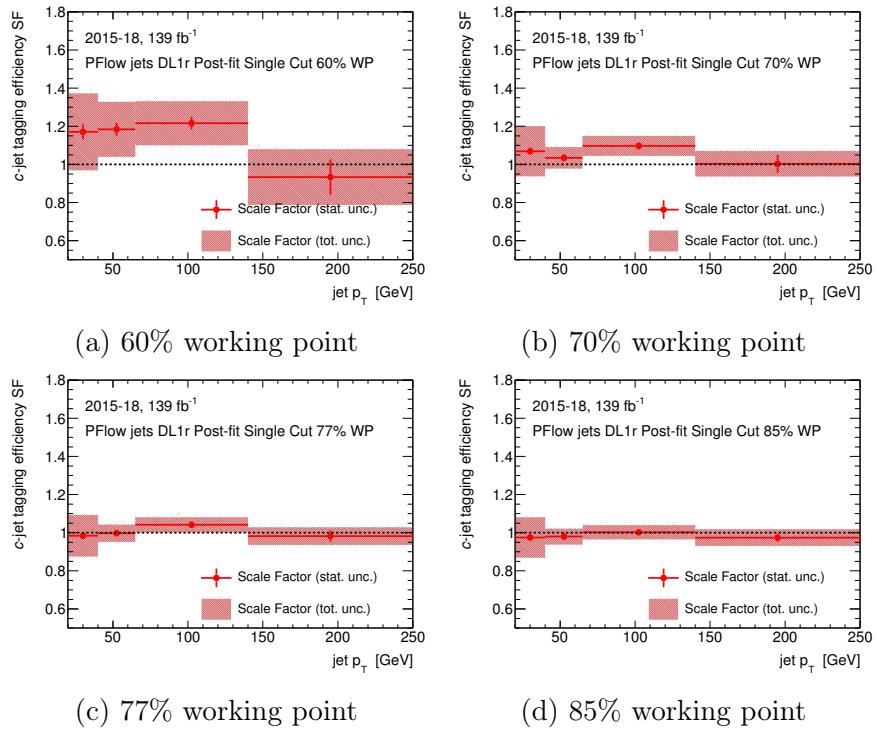


Figure 35: Charm-jet scale factors for the PFlow jets collection with the DL1r tagger.

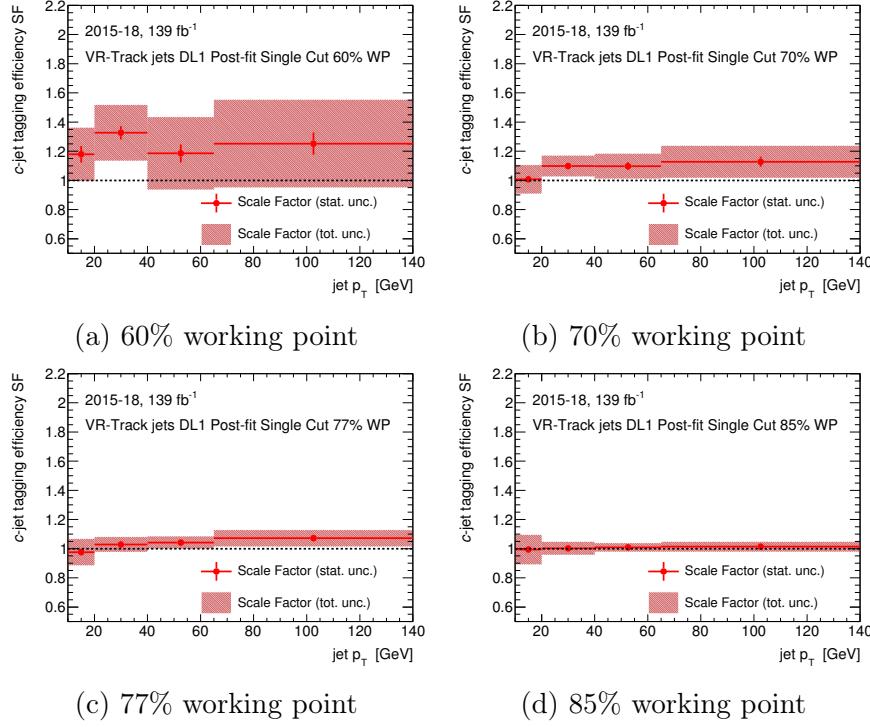


Figure 36: Charm-jet scale factors for the VR-Track jets collection with the DL1 tagger.

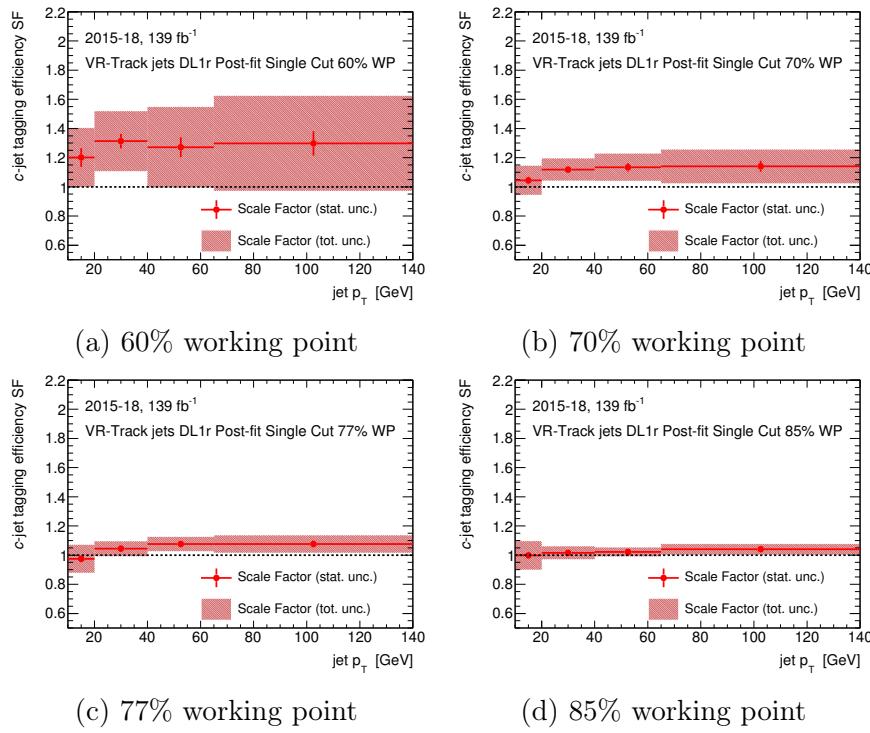


Figure 37: Charm-jet scale factors for the VR-Track jets collection with the DL1r tagger.

7 Search for Higgs boson pair production in the $b\bar{b}\tau\bar{\tau}$ channel

7.1 Data and Monte Carlo samples

7.2 Trigger and event selection

7.3 Background estimation

7.4 Multivariate analysis

7.5 Systematic uncertainties

7.6 Results

8 Summary

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A Supplementary material for c jet calibration

A.1 Additional plots for kinematic variables

A.1.1 Standard selection

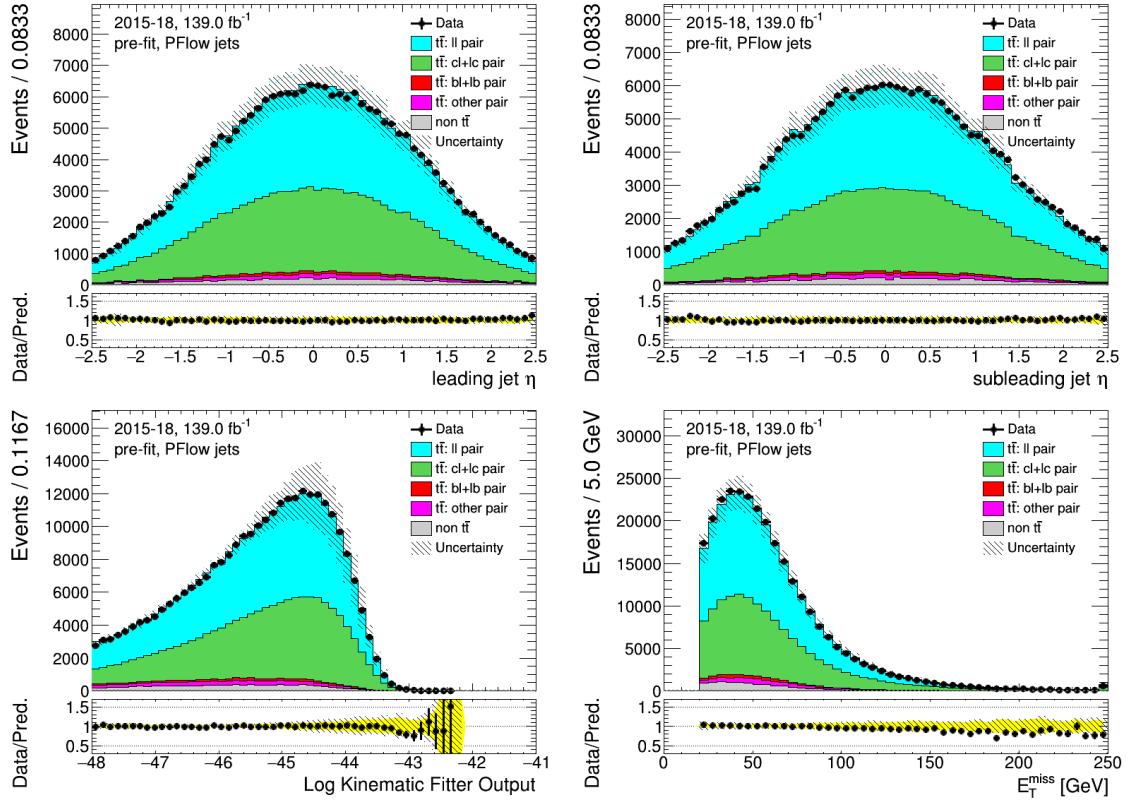


Figure 38: PFlow jets: distributions of the leading and sub-leading jets from W decay, KLFitter output and the transverse missing transverse energy of the standard selection, before fitting or tagging with full uncertainties.

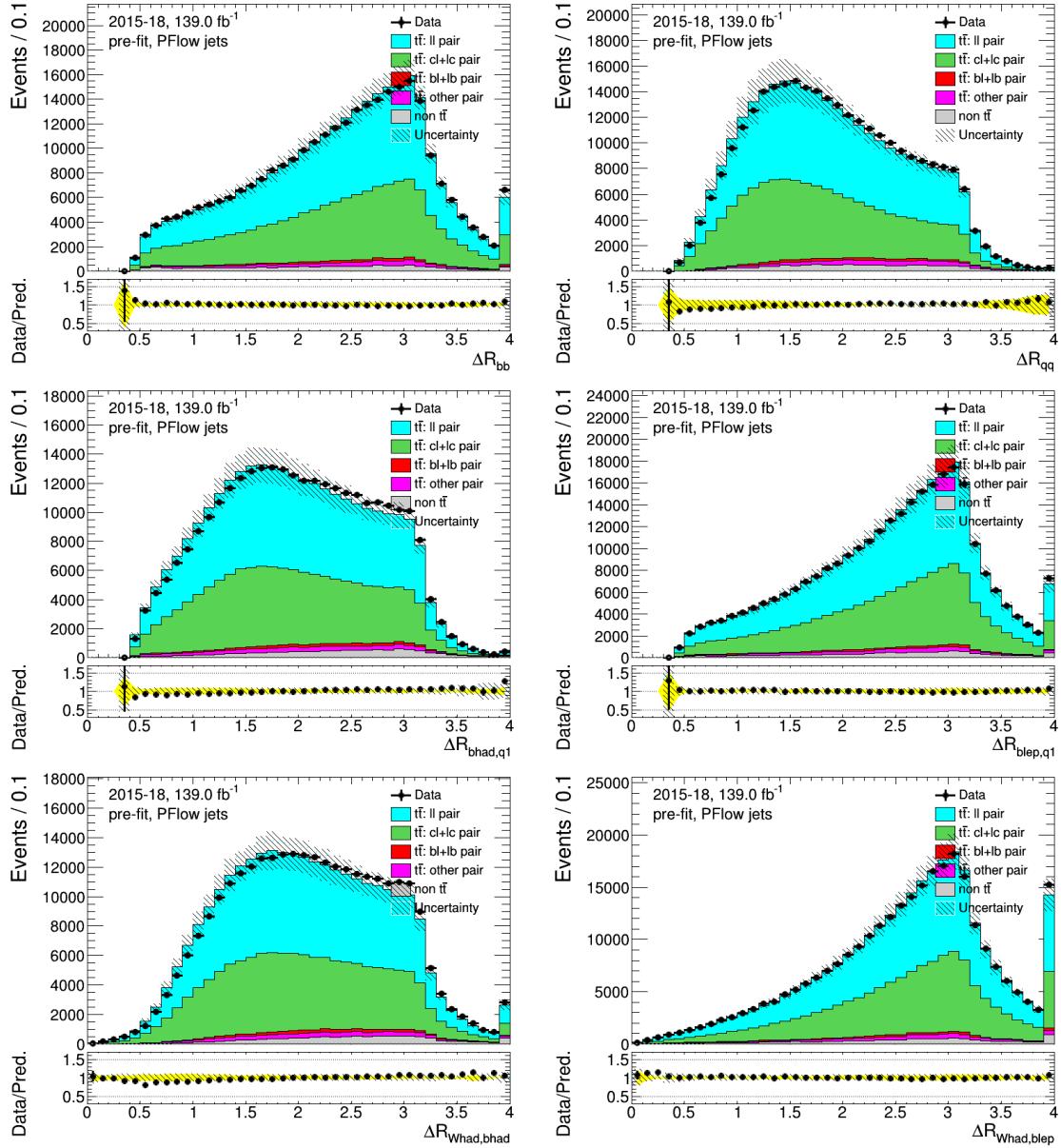


Figure 39: PFlow jets: distributions of angle related variables of the combination of the standard selection, before fitting or tagging with full uncertainties.

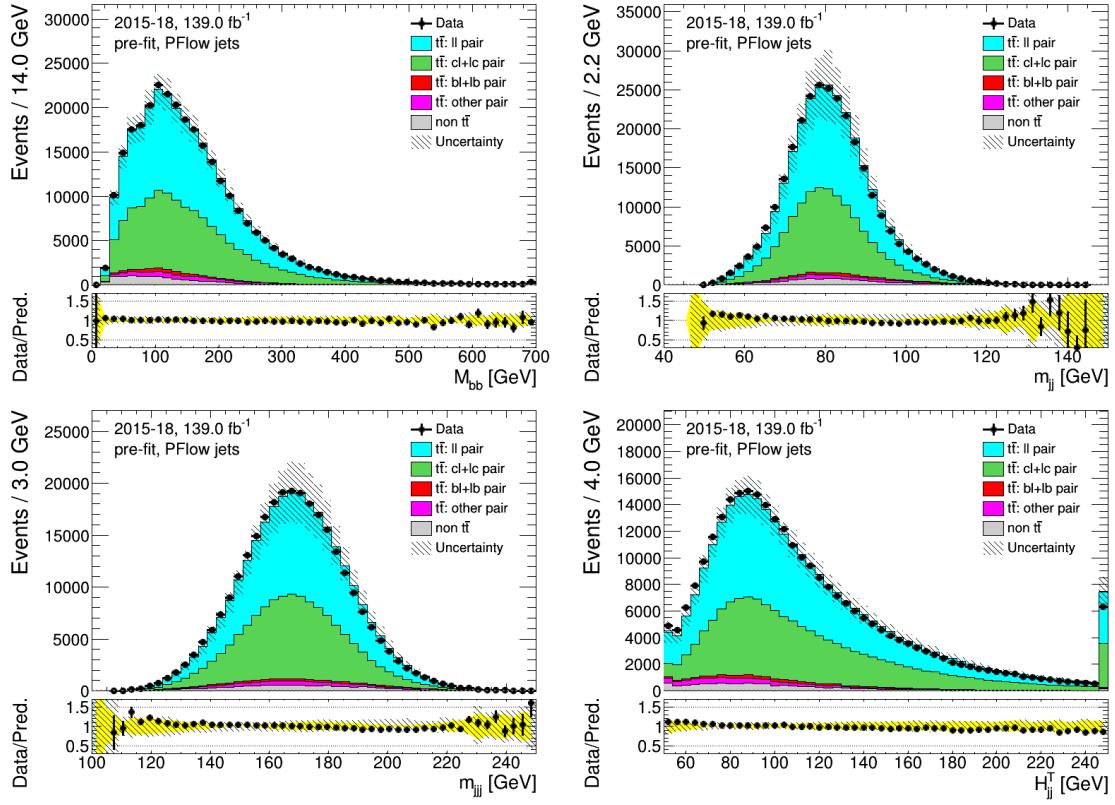


Figure 40: PFflow jets: distributions of mass related variables of the standard selection, before fitting or tagging with stat-only uncertainties.

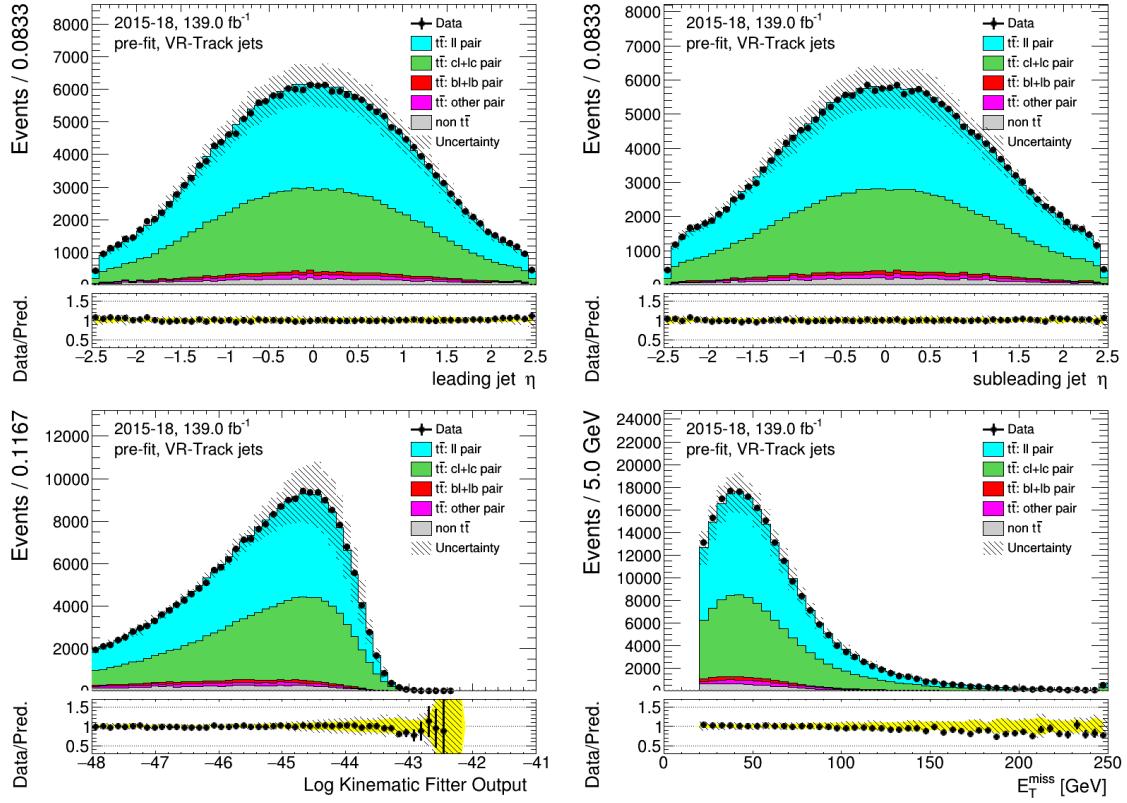


Figure 41: VR-Track jets: distributions of the leading and sub-leading jets from W decay, KLFitter output and the transverse missing transverse energy of the standard selection, before fitting or tagging with full uncertainties.

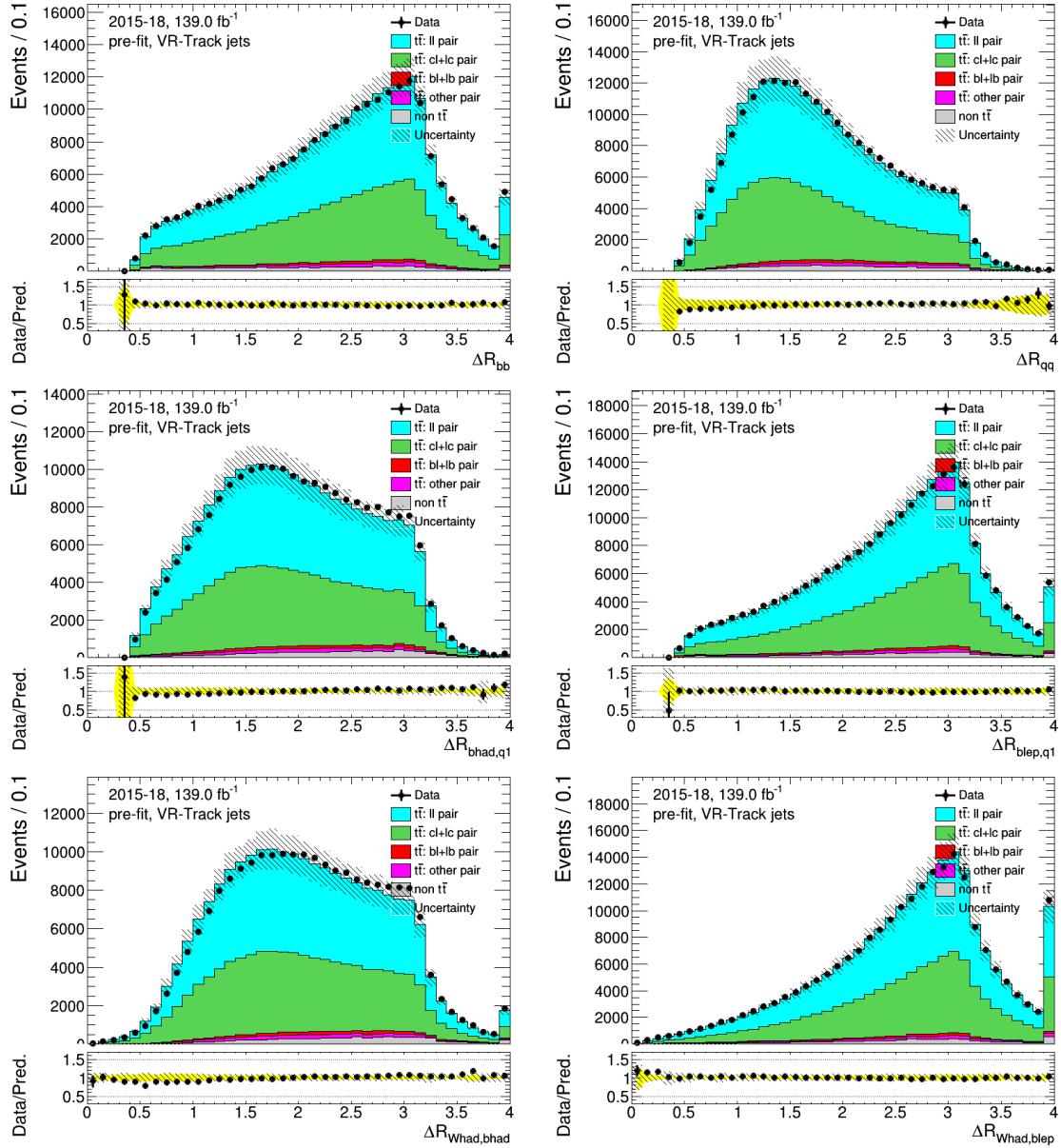


Figure 42: VR-Track jets: distributions of angle related variables of the combination of the standard selection, before fitting or tagging with full uncertainties.

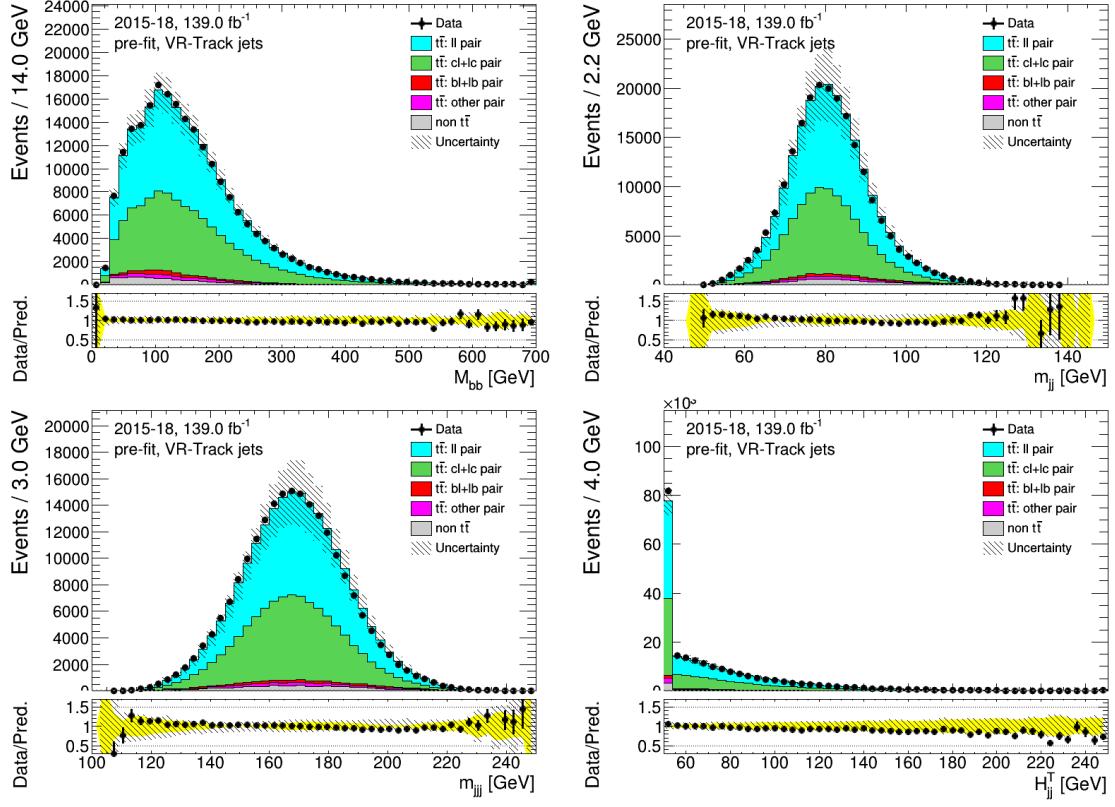


Figure 43: VR-Track jets: distributions of mass related variables of the standard selection, before fitting or tagging with stat-only uncertainties.

A.1.2 Low- p_T selection

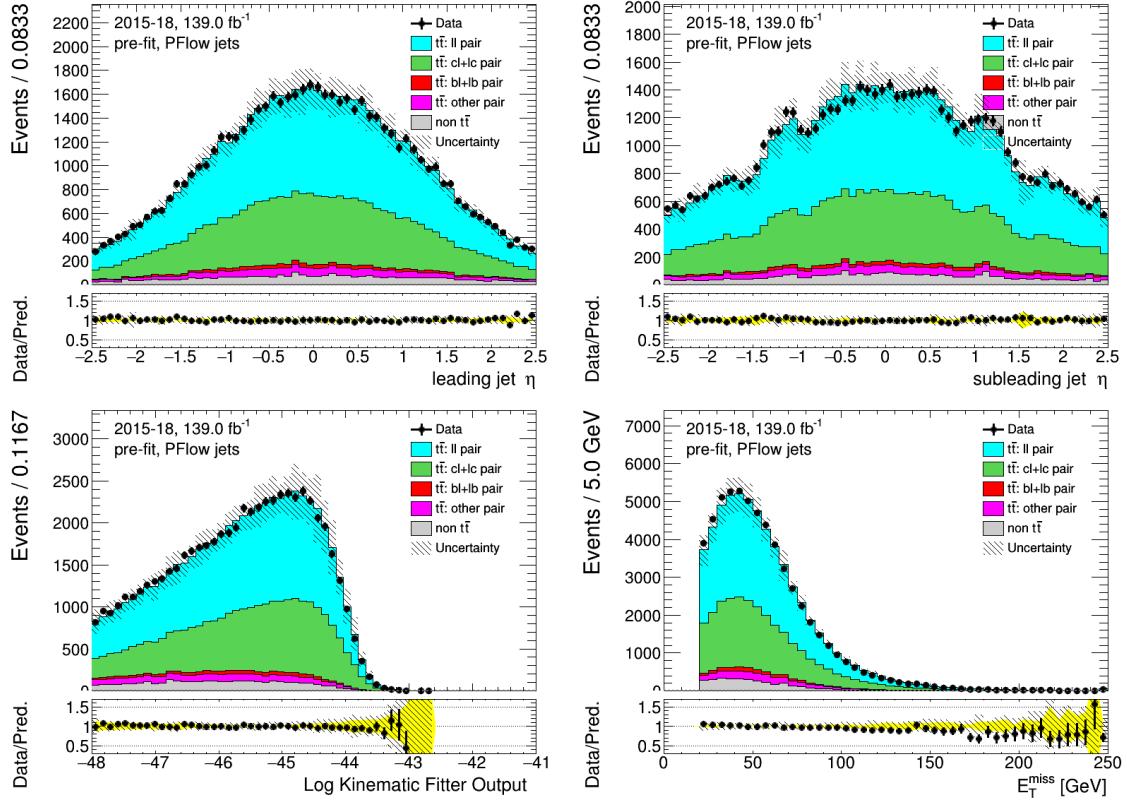


Figure 44: PFlow jets: distributions of the leading and sub-leading jets from W decay, KLFitter output and the transverse missing transverse energy of the low- p_T selection, before fitting or tagging with full uncertainties.

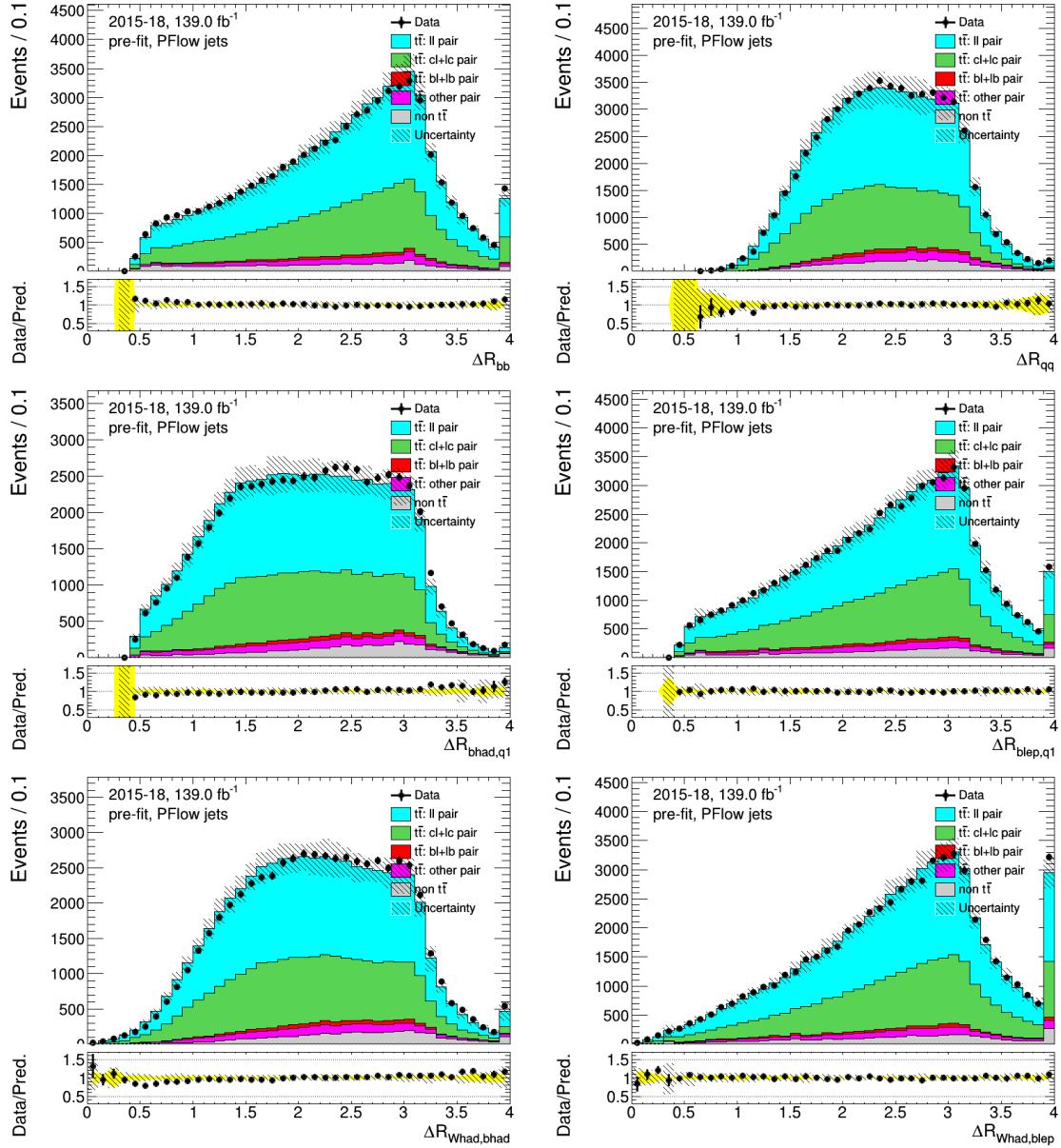


Figure 45: PFlow jets: distributions of angle related variables of the combination of the low- p_T selection, before fitting or tagging with full uncertainties.

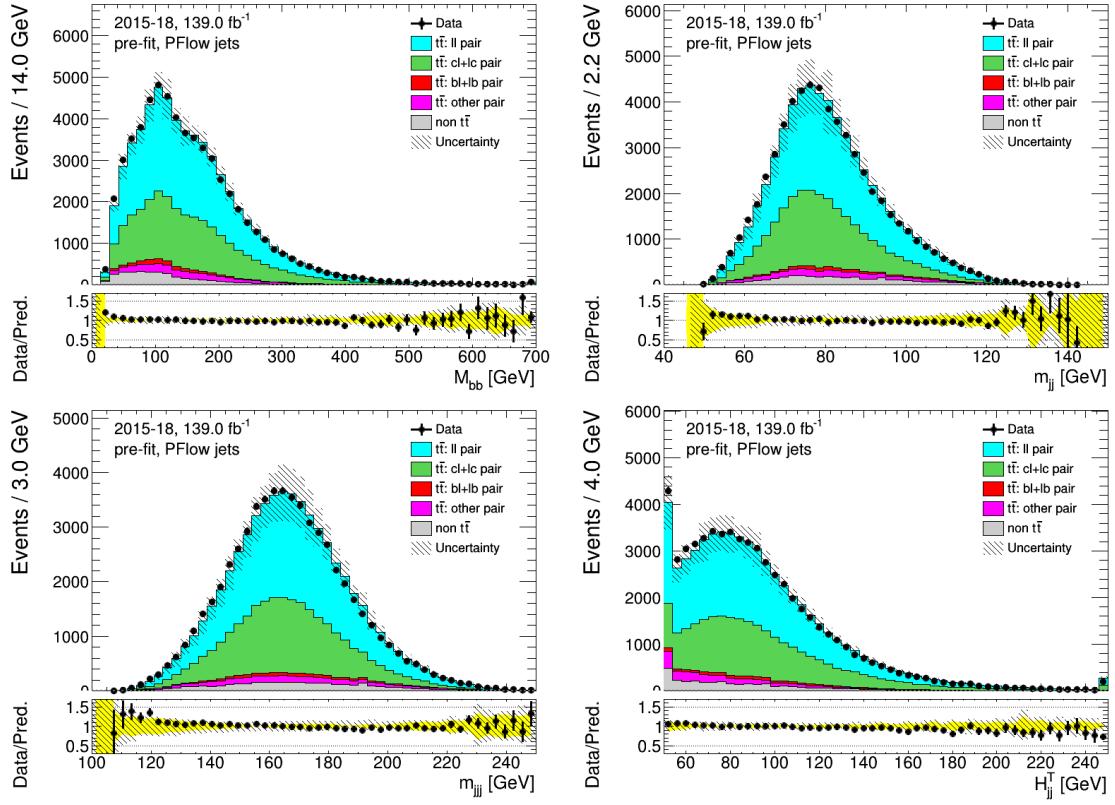


Figure 46: PFflow jets: distributions of mass related variables of the low- p_T selection, before fitting or tagging with stat-only uncertainties.

A.2 High- p_T selection

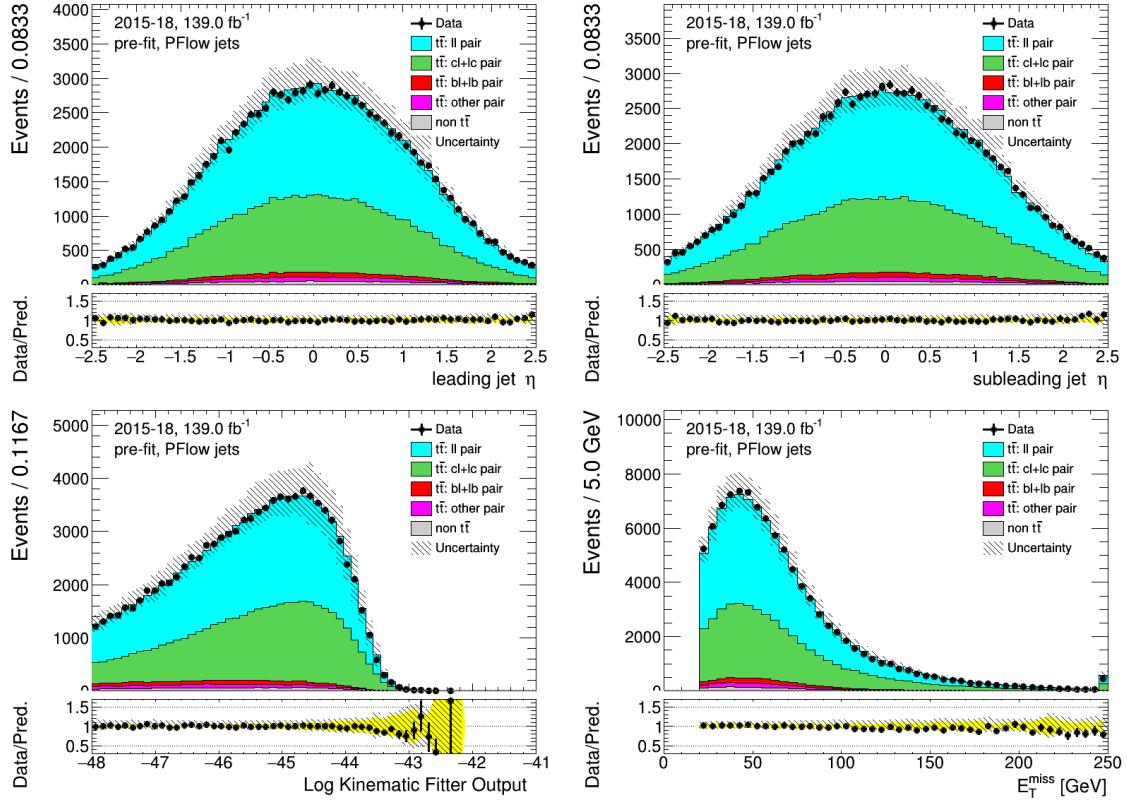


Figure 47: PFflow jets: distributions of the leading and sub-leading jets from W decay, KLFitter output and the transverse missing transverse energy of the high- p_T selection, before fitting or tagging with full uncertainties.

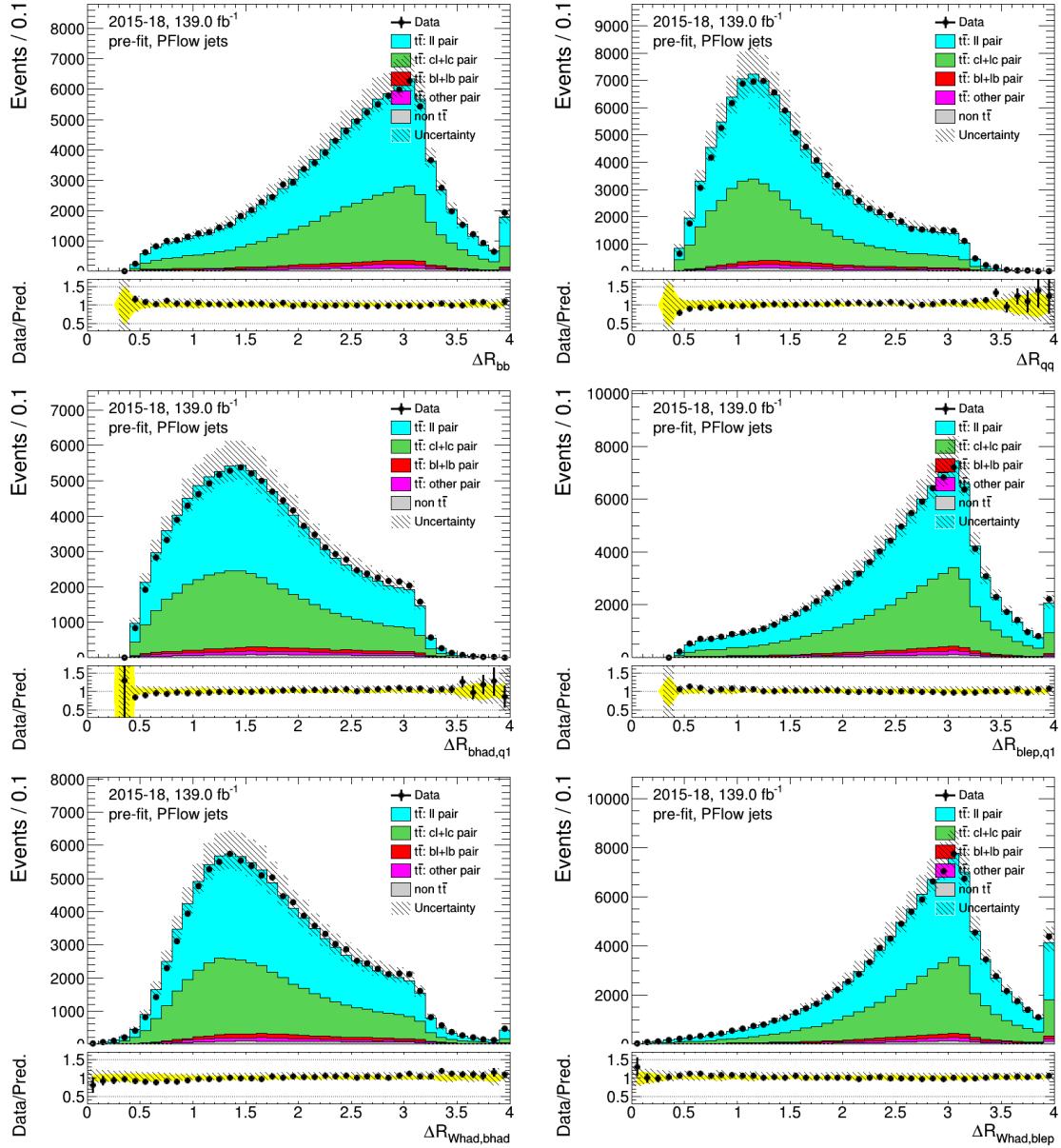


Figure 48: PFlow jets: distributions of angle related variables of the combination of the high- p_T selection, before fitting or tagging with full uncertainties.

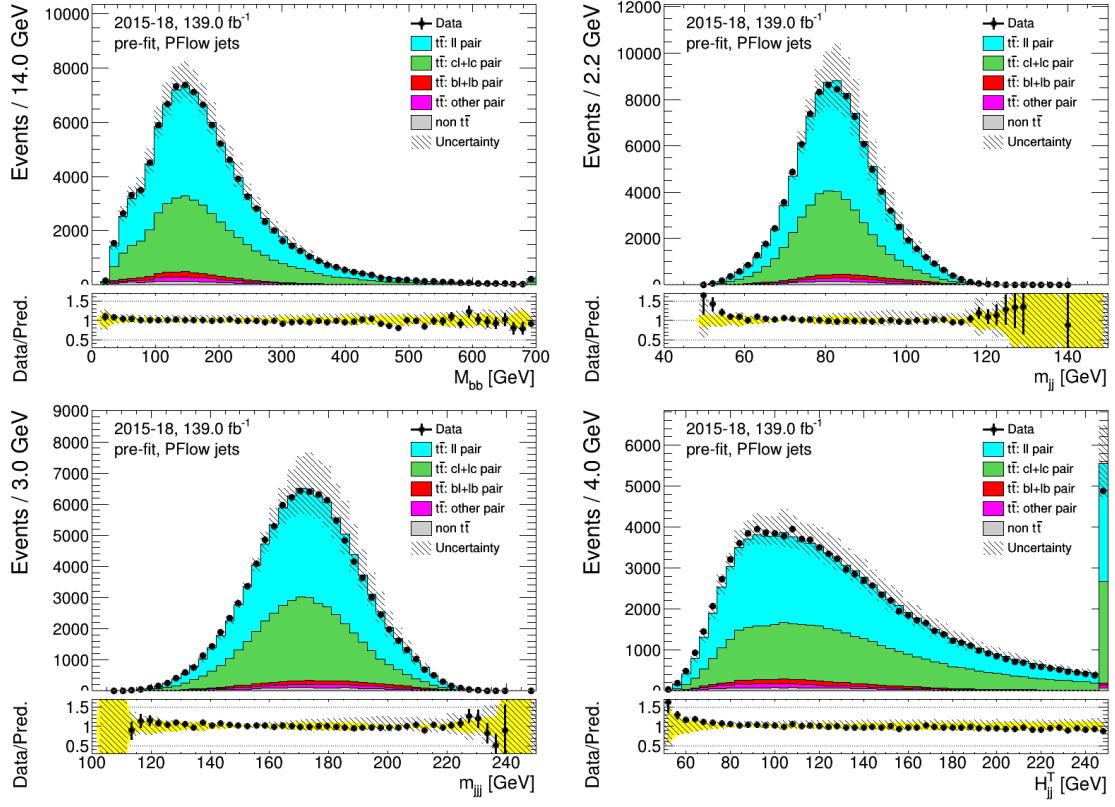


Figure 49: PFlow jets: distributions of mass related variables of the high- p_T selection, before fitting or tagging with stat-only uncertainties.

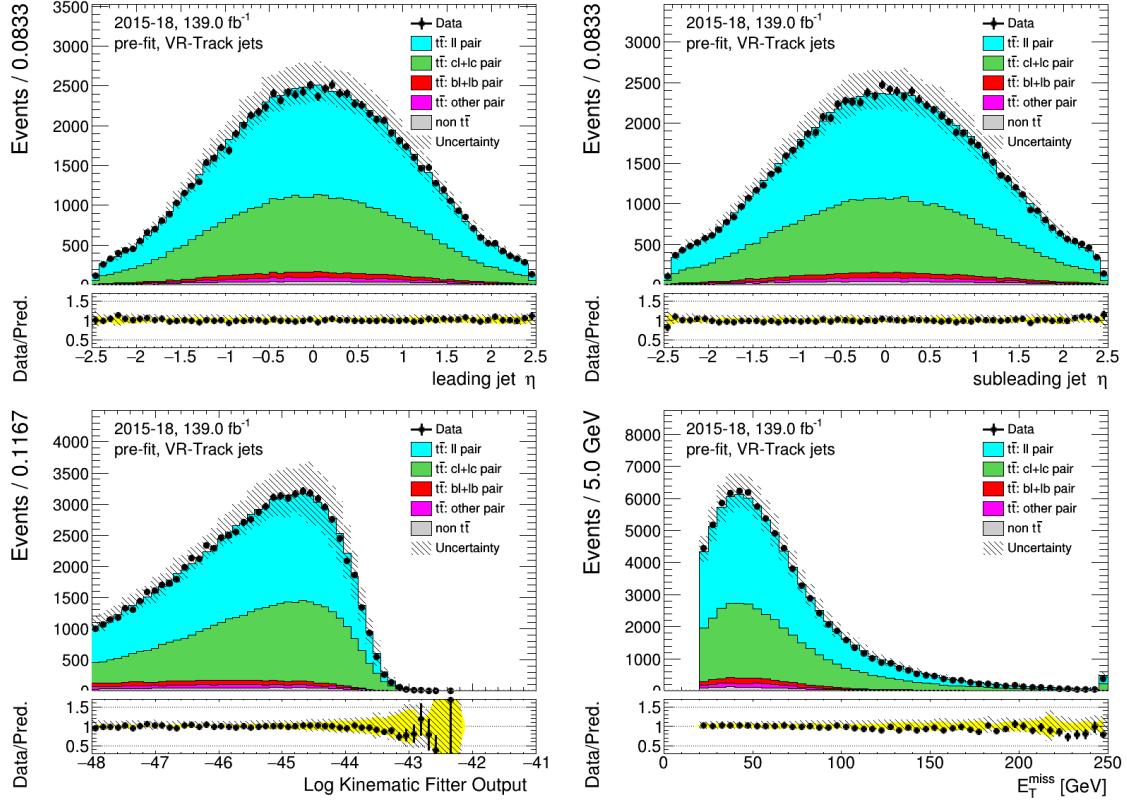


Figure 50: VR-Track jets: distributions of the leading and sub-leading jets from W decay, KLFitter output and the transverse missing transverse energy of the high- p_T selection, before fitting or tagging with full uncertainties.

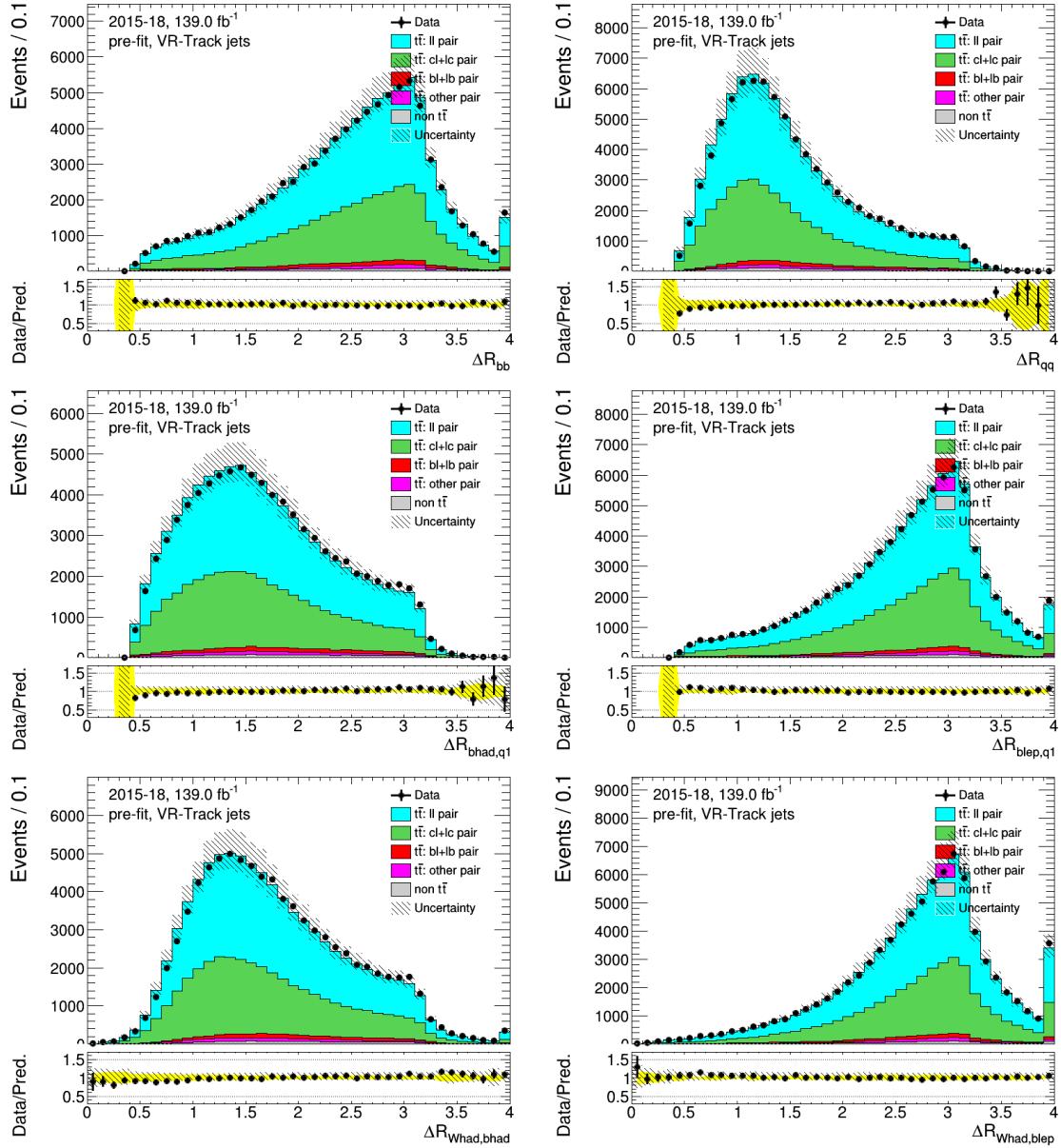


Figure 51: VR-Track jets: distributions of angle related variables of the combination of the high- p_T selection, before fitting or tagging with full uncertainties.

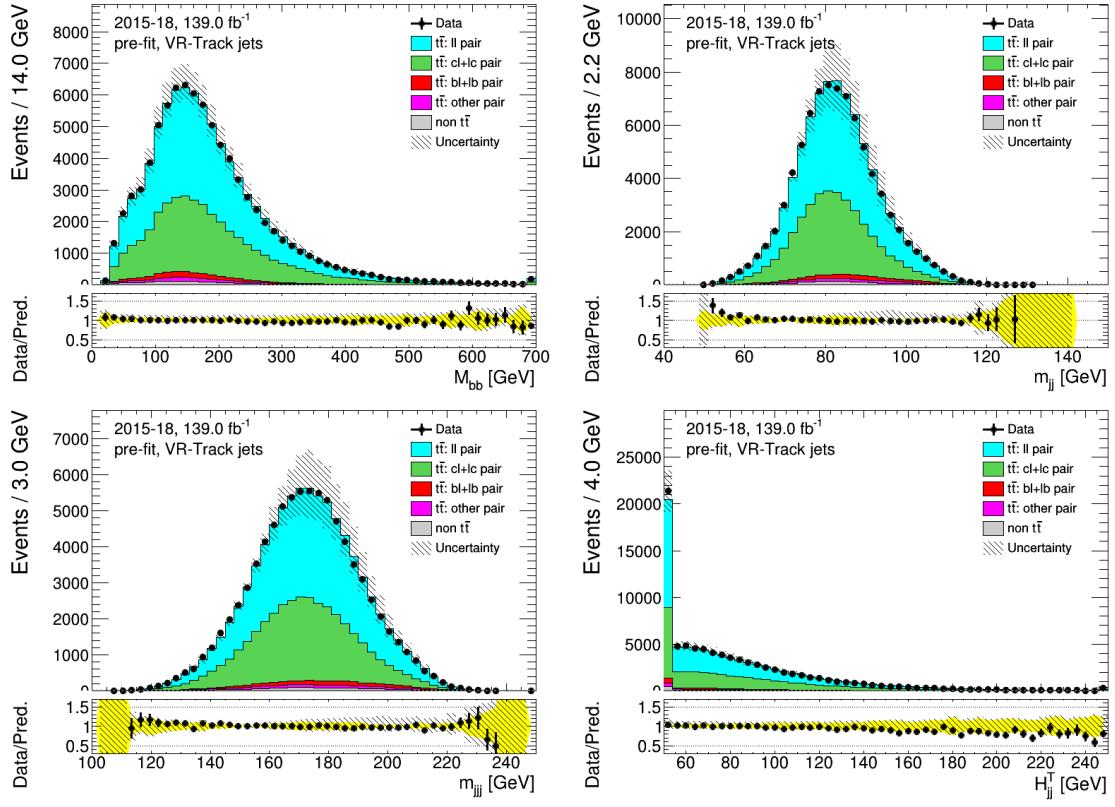


Figure 52: VR-Track jets: distributions of mass related variables of the high- p_T selection, before fitting or tagging with stat-only uncertainties.

A.3 Combined selection

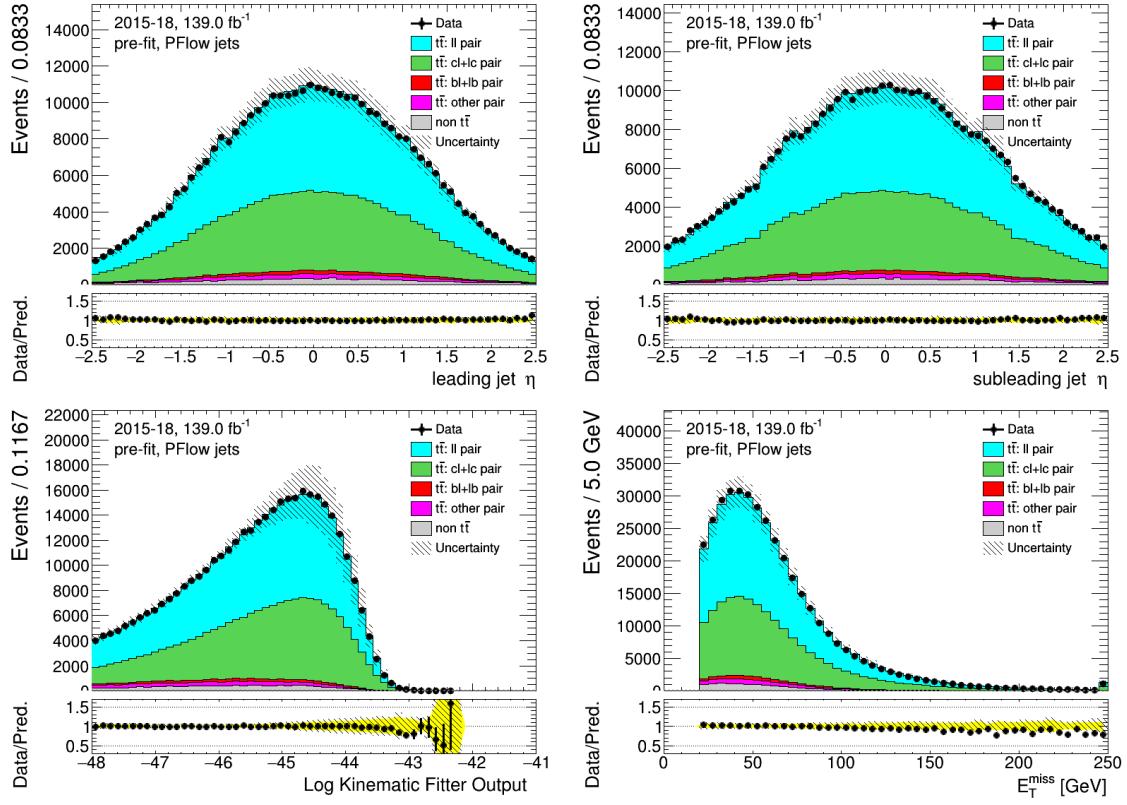


Figure 53: PFlow jets: distributions of the leading and sub-leading jets from W decay, KLFitter output and the transverse missing transverse energy of the combined selection, before fitting or tagging with full uncertainties.

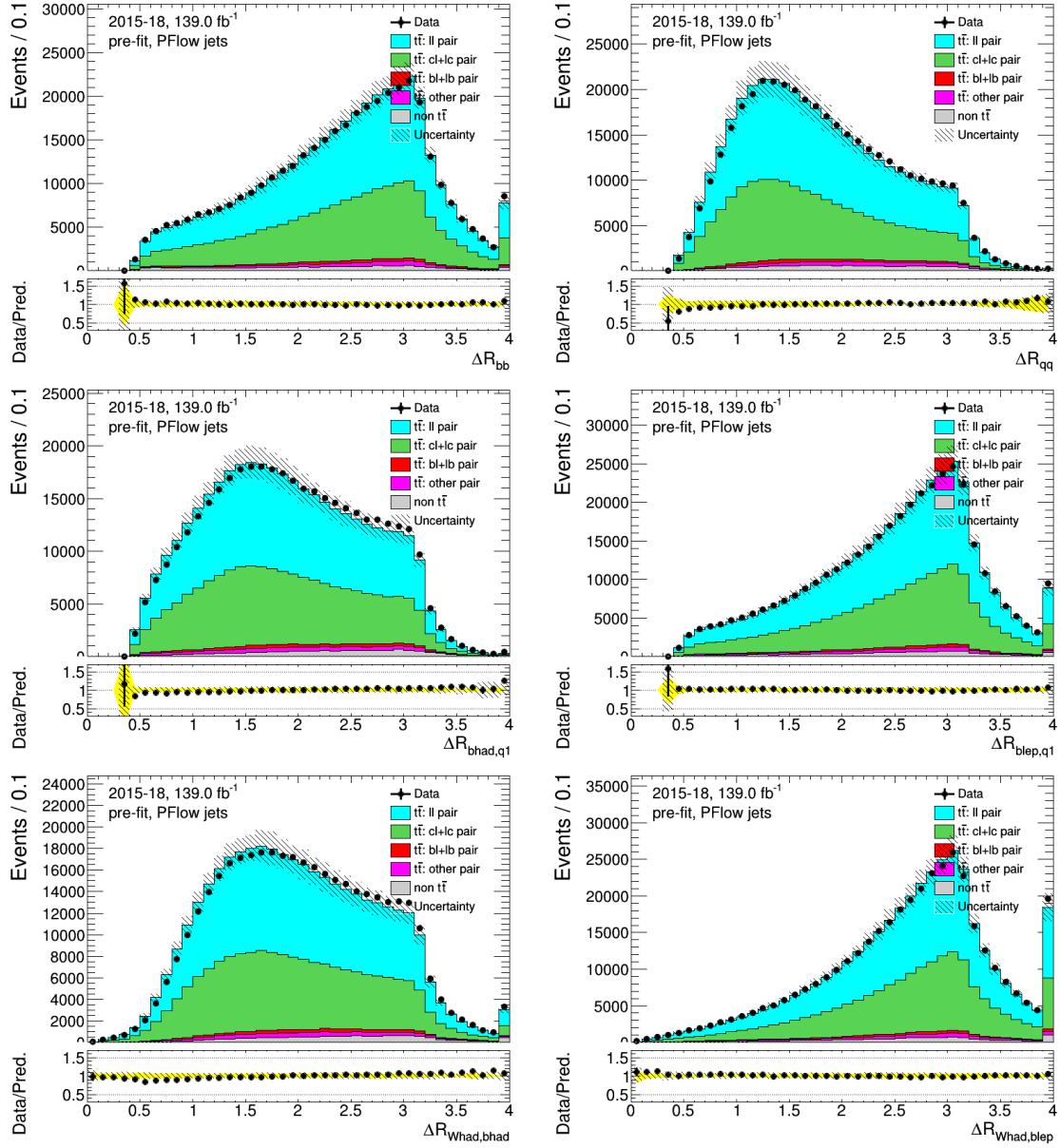


Figure 54: PFlow jets: distributions of angle related variables of the combination of the combined selection, before fitting or tagging with full uncertainties.

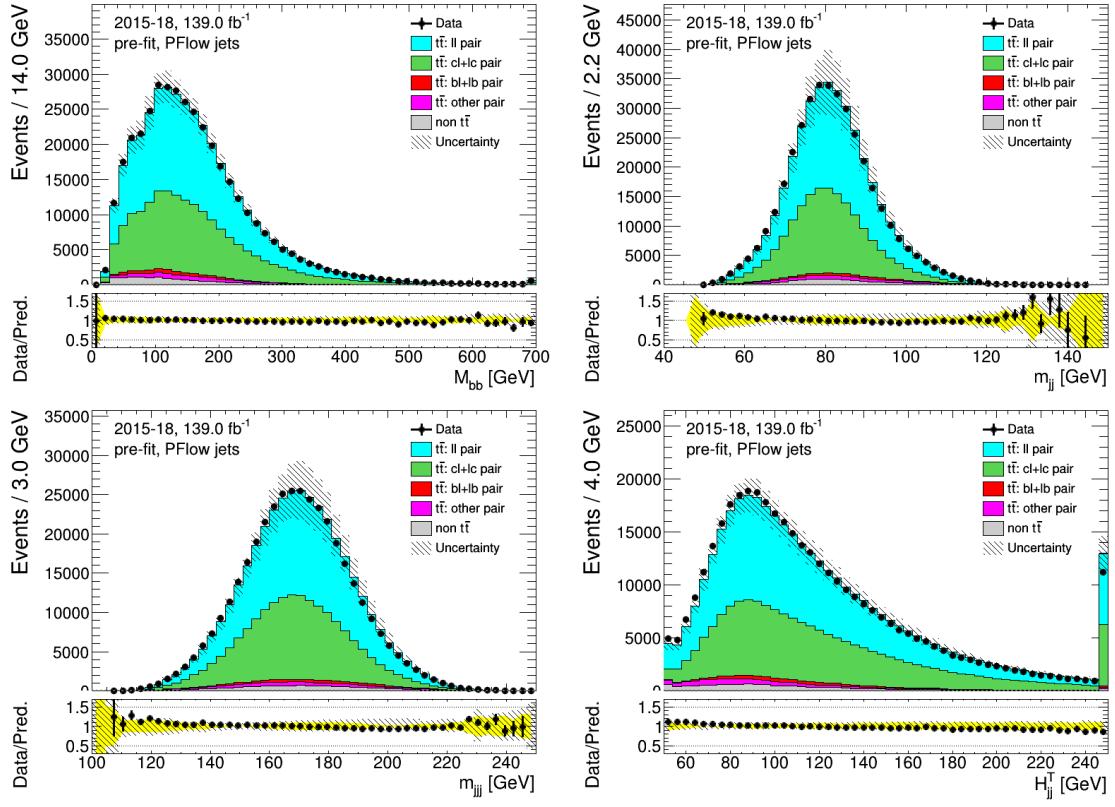


Figure 55: PFflow jets: distributions of mass related variables of the combined selection, before fitting or tagging with stat-only uncertainties.

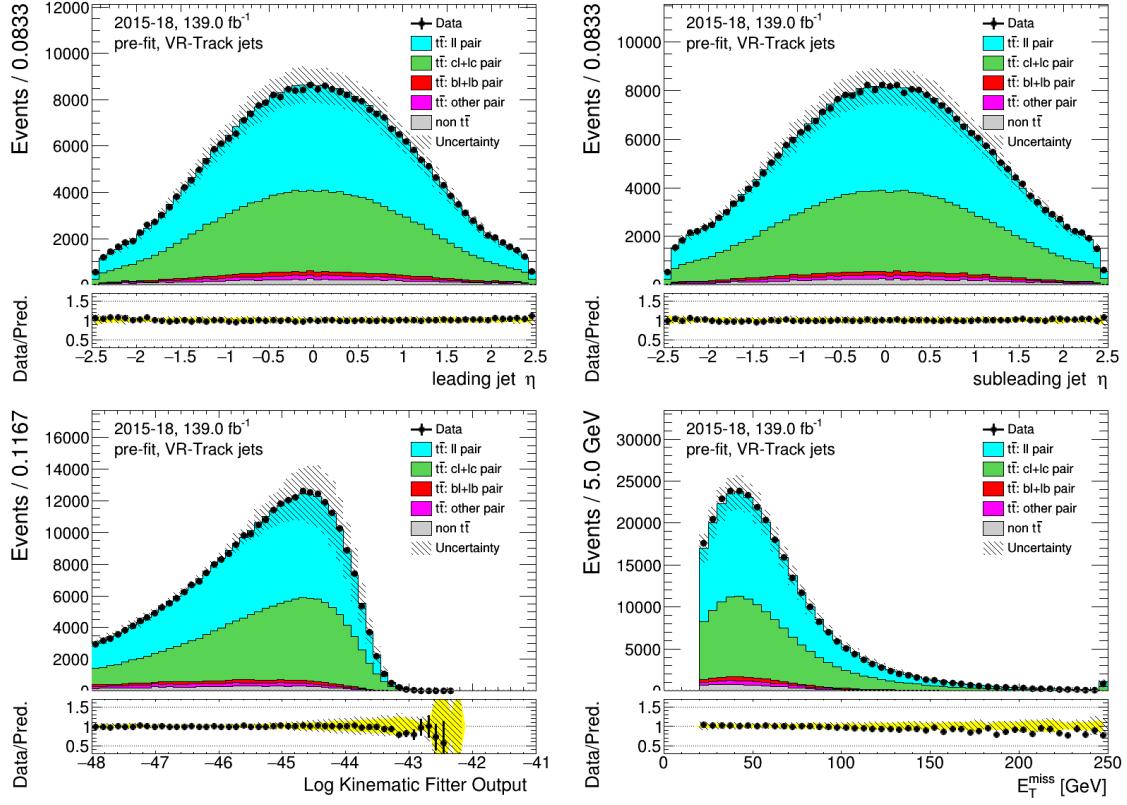


Figure 56: VR-Track jets: distributions of the leading and sub-leading jets from W decay, KLFitter output and the transverse missing transverse energy of the combined selection, before fitting or tagging with full uncertainties.

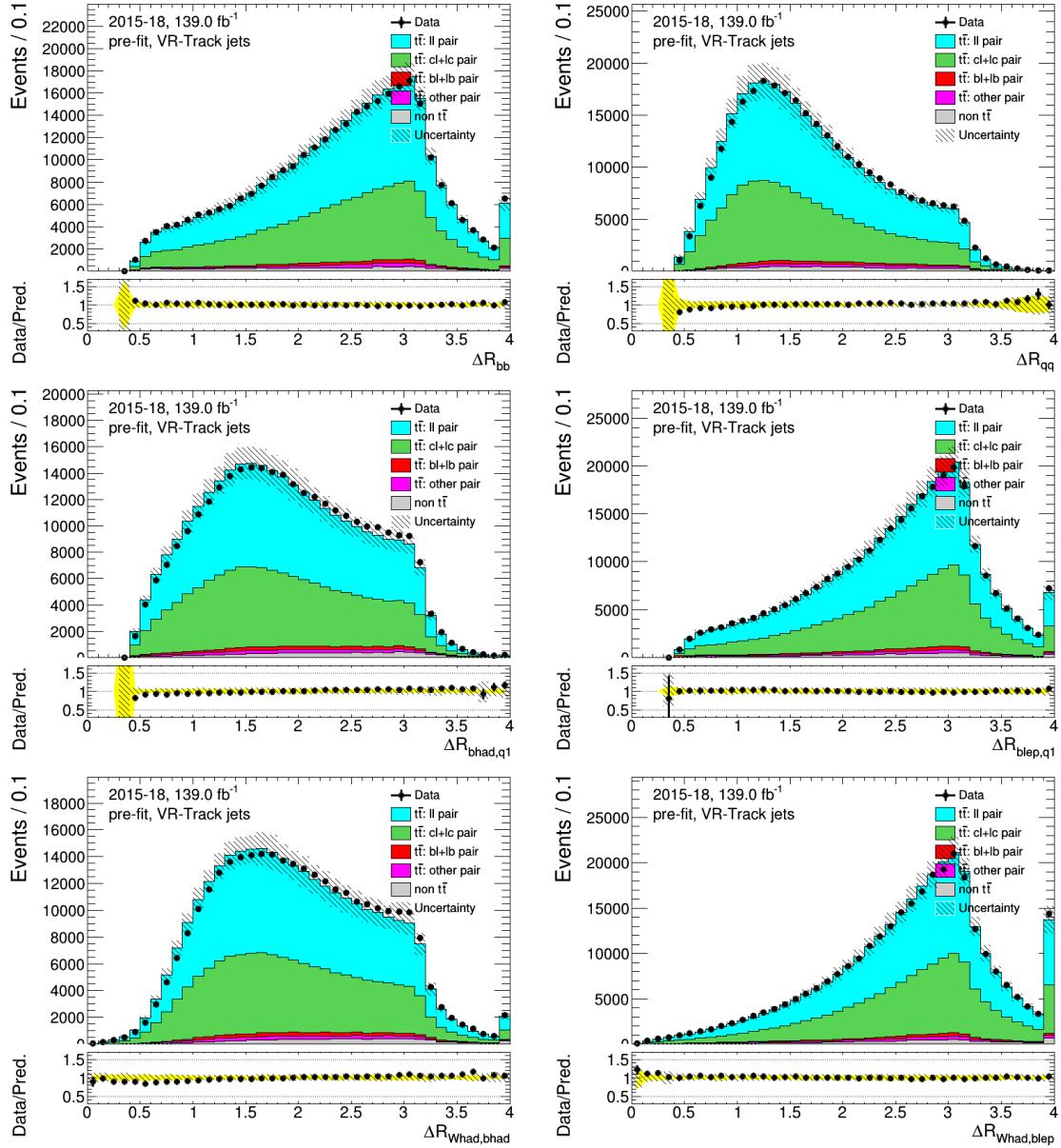


Figure 57: VR-Track jets: distributions of angle related variables of the combination of the combined selection, before fitting or tagging with full uncertainties.

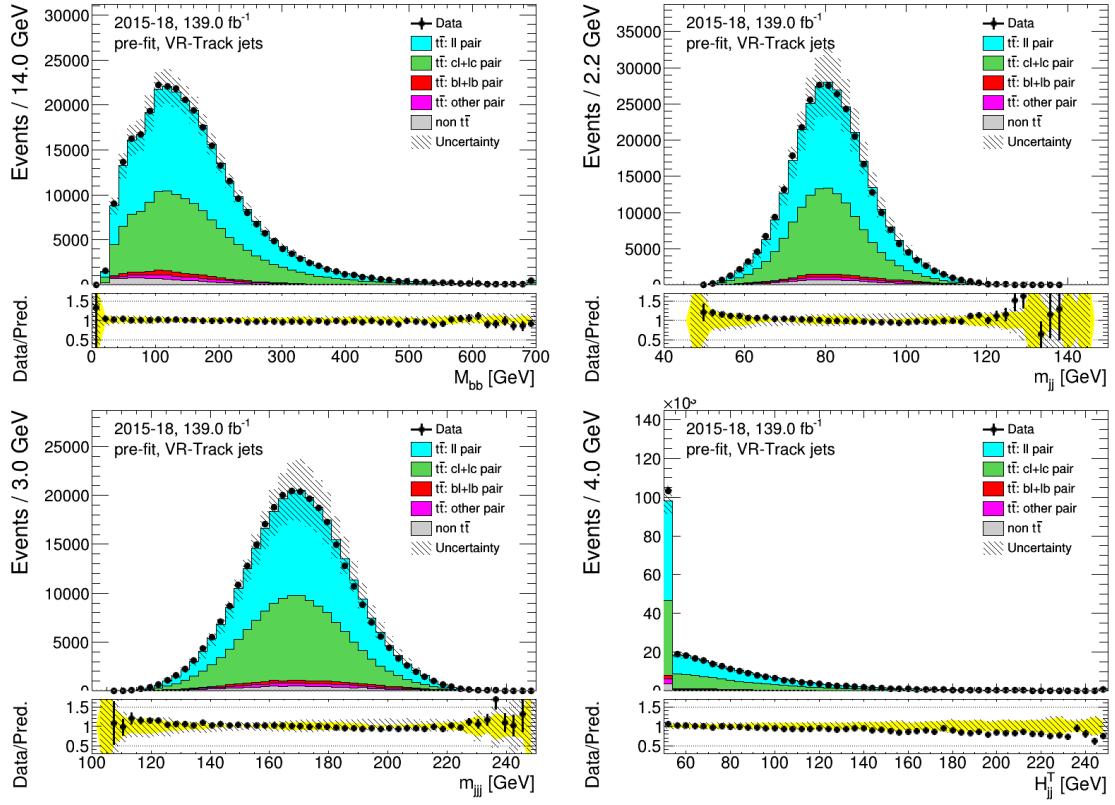


Figure 58: VR-Track jets: distributions of mass related variables of the combined selection, before fitting or tagging with stat-only uncertainties.

A.4 Experimental uncertainties

Systematic uncertainty
EG_RESOLUTION_ALL
MUON_ID
MUON_MS
MET_SoftTrk_ResoPara
MET_SoftTrk_ResoPerp
MET_SoftTrk_ScaleDown
MET_SoftTrk_ScaleUp
JET_Pileup_OffsetNPV
JET_Pileup_RhoTopology
JET_EffectiveNP_Modelling1
JET_EffectiveNP_Modelling2
JET_EffectiveNP_Modelling3
JET_EffectiveNP_Modelling4
JET_EffectiveNP_Statistical4
JET_EffectiveNP_Detector1
JET_JER_EffectiveNP_1
JET_JER_EffectiveNP_2
JET_JER_EffectiveNP_3
JET_JER_EffectiveNP_4
JET_BJES_Response
JET_Flavor_Composition
JET_Flavor_Response

Table 9: List of experimental systematics.