

# Statistics, Uncertainties and Linear Fitting

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## 1 Uncertainties

Diffraction equation:

$$m\lambda = d\sin\theta$$

Derivation of expression for  $\Delta\lambda$ :

$$\begin{aligned}(\Delta\lambda)^2 &= \left(\frac{1}{m}\sin\theta\right)^2(\Delta d)^2 + \left(\frac{1}{m}d\cos\theta\right)^2(\Delta\theta)^2 \\ &= \left(\frac{1}{m}\right)^2(\sin^2\theta(\Delta d)^2 + d^2\cos^2\theta(\Delta\theta)^2)\end{aligned}$$

$$\Delta\lambda = \frac{1}{m}\sqrt{\sin^2\theta(\Delta d)^2 + d^2\cos^2\theta(\Delta\theta)^2}$$

Given the values:

$$m = 1$$

$$d = (600)^{-1}mm$$

$$\Delta d/d = 1\%$$

$$\theta = 15deg$$

$$\Delta\theta = 0.5deg$$

Evaluation for  $\lambda$ :

$$\lambda = \frac{d\sin\theta}{m} = \frac{[(600)^{-1}mm]\sin(15deg)}{1}$$

$$\lambda = 4.3 \times 10^{-4}m = 430\mu m$$

Evaluation for  $\Delta\lambda$ :

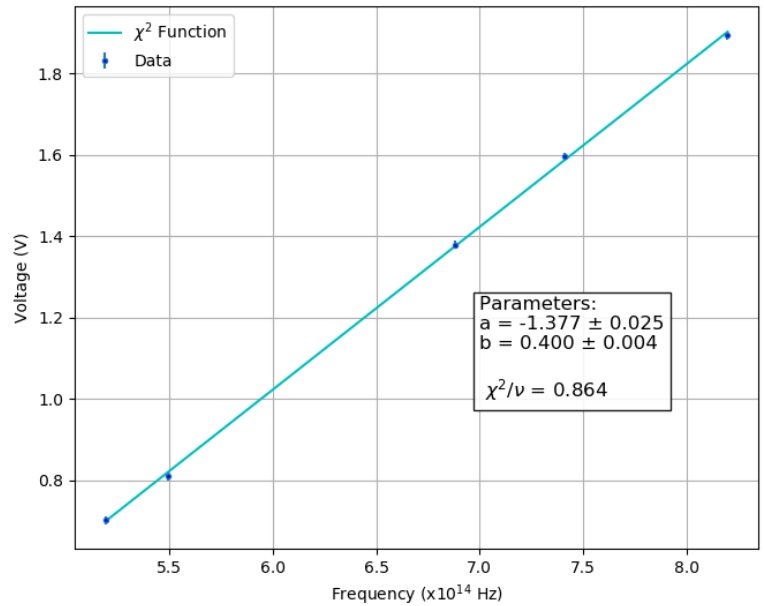
$$\begin{aligned}\Delta\lambda &= \frac{1}{m}\sqrt{\sin^2\theta(\Delta d)^2 + d^2\cos^2\theta(\Delta\theta)^2} \\ &= \frac{1}{1}\sqrt{[\sin^2(15deg)]\left(\frac{1}{6000}m\right)^2 + \left(\frac{1}{600}m\right)^2[\cos^2(15deg)](0.5deg)^2}\end{aligned}$$

$$\Delta\lambda = 8.1 \times 10^{-4}m = 810\mu m$$

## 2 Linear Least Squares Fit

Having programmed an analysis routine on Python and plotting the data and fit, these are the results:

**Least Squares Fitting**



The  $\chi^2/\nu$  is close to the ideal value of 1, so the linear fit is a good approximation to the system's function.

### 3 Geiger Counter - Radiation Decay

#### 3.1 Low Count-Rate Data

The mean number of counts  $\bar{r}$ , and the standard deviation  $s$  without the radioactive sample is:

$$\bar{r} = \frac{1}{N} \sum_{n=1}^N r_n = 1.92$$

$$s = \sqrt{\frac{1}{N-1} \sum_{r=1}^N (r - \bar{r})^2} = 1.37$$

Since we are counting a random event, the data should follow a Poisson distribution. This means that the standard deviation squared should approximate the mean. Given our obtained data:

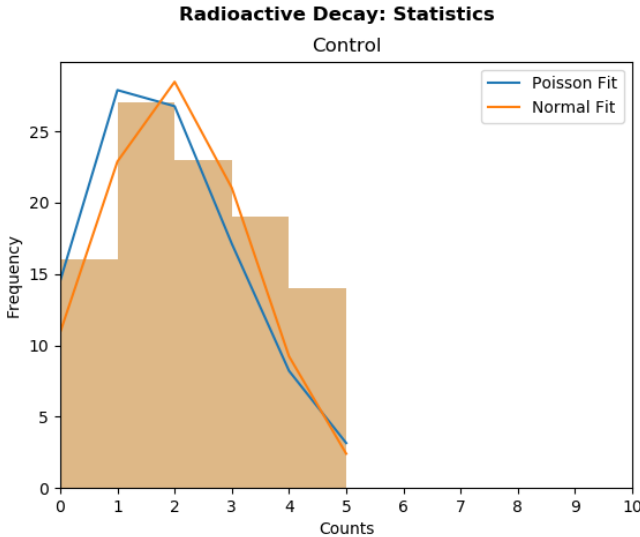
$$s^2 = 1.87 \approx \bar{r} = 1.92$$

The fraction of trials with an absolute deviation from the mean larger than the standard deviation  $s$  is approximately 30% of the sample. If one were to look at the probability of getting a result above that of the mean of the ideal Poisson distribution ( $\approx 2$ ), then:

$$\sum_{r=3}^5 P_{\mu}(r) = \sum_{r=3}^5 \frac{\mu^r e^{-\mu}}{r!} = 29\%$$

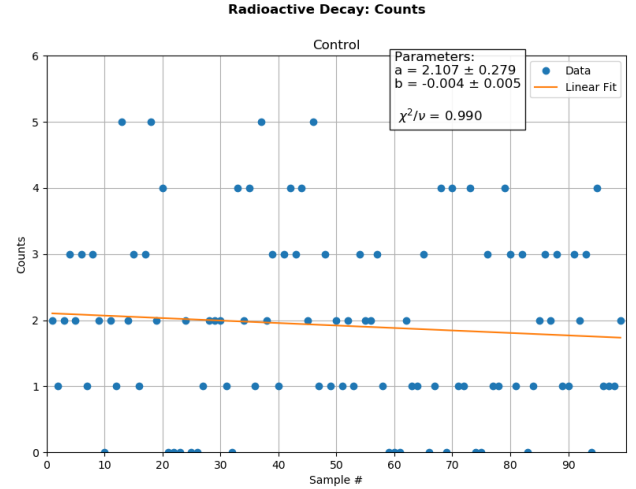
Giving indication that our results are very close to that of a Poisson distribution with the same mean.

The following figure shows a histogram of the data taken, as well as calculated Poisson and Normal distributions from the data's parameters:



As can be seen, the Poisson distribution gives a much better model for the data, showing that the number of counts from the background detected by the Geiger-Müller Counter is entirely random.

The data points of the experiment are shown below, displaying counts against sample number:



A linear Least Squares Fit was used in attempt to recreate the trend of the data. While the experiment assumed that the probability of detecting a radioactive decay product was constant in time ( $a = \bar{r} = 1.92, b = 0$ ), the calculated parameters were very close to the ideal values, within their uncertainties. Therefore, it can be assumed that the detection is constant throughout time.

#### 3.2 High Count-Rate Data

The same process is used to analyze the high count-rate data, which includes the sample.

The mean and the standard deviation of this set:

$$\bar{r} = 54.23$$

$$s = 9.09$$

$$s^2 = 82.54 \neq \bar{r} = 54.23$$

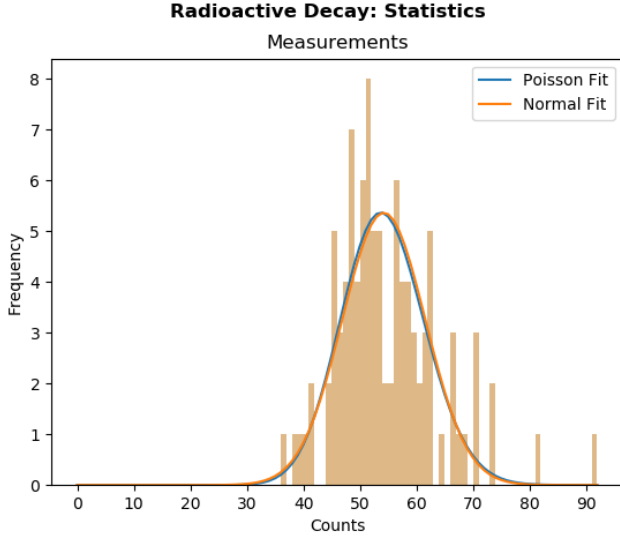
The comparison of the standard deviation and the mean means that this should not follow a Poisson distribution, since now there is a source of decay.

The fraction of trials with an absolute deviation from the mean larger than the standard deviation  $s$  is approximately 26% of the sample. A Poisson distribution would expect to see:

$$\sum_{r=55}^{92} P_{\mu}(r) = 48\%$$

There's a big discrepancy between both of these values as well.

Plotting the histogram of the data and fitting both a Poisson and Normal distribution:



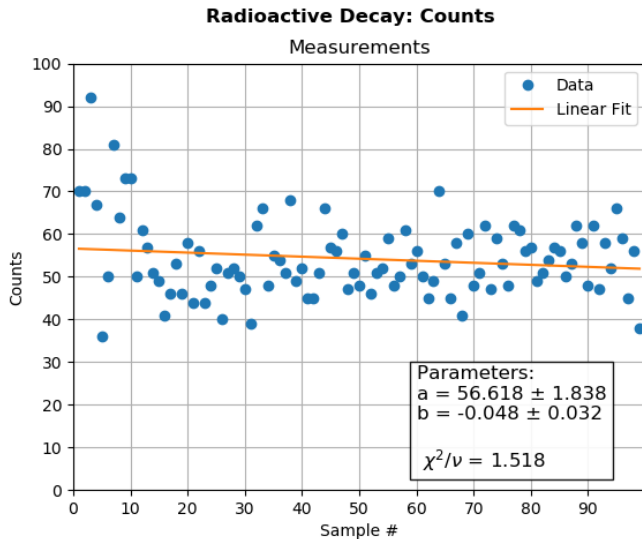
It can be seen that at this point the Poisson and Normal distributions are very similar, which makes sense since the Normal distribution is just a limiting case of the Poisson distribution ( $\mu \gg 1$ ). Their numerical differences on the plot is due to how the Normal distribution is approximated, being:

$$P(r) = \frac{1}{\sigma\sqrt{2\pi}} e^{-(r-\bar{r})^2/2\sigma^2}$$

While that of the Poisson distribution is:

$$P(r) = \frac{\bar{r}^r e^{-\bar{r}}}{r!}$$

The figure below shows the plot of the data points of the experiment and its appropriate Linear Least Squares Fit as well:



A good approximation is achieved, with the parameters being close to their expected values ( $a = \bar{r} = 54.23, b = 0$ ).

### 3.3 Determining the Source Count Rate

With the obtained data, the source's actual count rate can be determined. Knowing that:

$$R_{tot} = 54.23 \pm 9.09 \equiv \text{Source} + \text{Background}$$

$$R_b = 1.92 \pm 1.37 \equiv \text{Background}$$

The total rate of counts from just the source is then:

$$\begin{aligned} R &= R_{tot} - R_b \\ &= 54.23 - 1.92 \\ R &= 52.31 \end{aligned}$$

$$\begin{aligned} \Delta R &= \sqrt{\Delta R_{tot}^2 + \Delta R_b^2} \\ &= \sqrt{(9.09)^2 + (1.36)^2} \\ \Delta R &= 9.19 \end{aligned}$$

$$R = 52.31 \pm 9.19$$