#### Supporting information: Shoreface dynamics formulation and sensitivity analysis

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#### S1. Shoreface dynamics formulation

In this supplementary material, we use a cross-shore energetics sediment transport formulation for wave-driven suspended sediment transport presented by Ortiz and Ashton [2013] to develop an analytical estimation for the proposed depth-average shoreface response rate  $K[L^2/T]$ . Following the energetics-based approach of Bowen [1980], Ortiz and Ashton [2013] show that the wave-driven cross-shore suspended sediment transport formulation for shallowwater waves can be formulated as:

$$q_s(z) = -\phi \frac{H(z)^5 g^2}{32w_s} \left( \frac{15}{4z^3} + \frac{9T^2 g}{16\pi^2 z^4} + \alpha_L(z) \frac{g^{1/2}}{w_s z^{5/2}} \right).$$
 (S1)

Under linear Airy wave theory the formulation becomes:

$$q_s(z) = -\phi \frac{\pi^5 H(z)^5}{w_s T^4 \sinh^5(2\pi z/L(z))} \left[ \frac{15}{4L(z)} + \frac{9}{4L(z)\sinh^2(2\pi z/L(z))} + \frac{\alpha_L(z)}{w_s T} \right], \tag{S2}$$

where z[L] is the local water depth , T[T] is the wave period, H[L] the local wave height, L[L] the local wave length,  $w_s[L/T]$  the settling velocity,  $\alpha_L(-)$  the local shoreface slope,  $g[L/T^2]$  the acceleration by gravity, and the coefficient  $\phi[T^2/L]$  is:

$$\phi = \frac{16e_s C_s \rho}{15\pi (\rho_s - \rho)g},\tag{S3}$$

where  $e_s$  is the suspended sediment transport efficiency factor,  $C_s(-)$  the friction factor,  $\rho\left[M/L^3\right]$  the seawater density, and  $\rho_s\left[M/L^3\right]$  the sediment density. After rearranging (S1) and

(S2) in the form  $q_s(z) = K_L(\alpha_{e,L} - \alpha_L)$ , the local, depth-dependent equilibrium slope  $\alpha_{e,L}(-)$ , and local shoreface response rate  $K_L\left[L^2/T\right]$  under the shallow-water waves assumption are:

$$\alpha_{e,L} = \frac{3w_s}{4z^{1/2}g^{1/2}} \left(5 + \frac{3T^2g}{4\pi^2z}\right),\tag{S4}$$

$$K_L = \phi \frac{H^5 g^{5/2}}{32 w_s^2 z^{5/2}} \,. \tag{S5}$$

Under linear Airy wave theory the expressions become:

$$\alpha_{e,L} = \frac{3w_s T}{4L} \left( 5 + \frac{3}{\sinh^2(2\pi z/L)} \right),\tag{S6}$$

$$K_{L} = \frac{\phi}{w_{s}^{2}} \frac{\pi^{5} H^{5}}{T^{5} \sinh^{5}(2\pi z/L)}.$$
 (S7)

The local wave height H in (S5) and (S7) can be calculated in terms of the deep wave height  $H_0[L]$  using the standard dispersion relationship.

Using (S3) and (S4) or (S6) and (S7), and assuming the parameter values included in Table S1, the shoreface response rate  $K_L$  and the slope  $\alpha_L$  as a function of depth in Figure S1. In this paper, however, we conceptualize the shoreface response rate by depth-averaging using the estimation:

$$K = \frac{\int_{z^{*}}^{D_{T}} K_{L} dz}{D_{T} - z^{*}},$$
 (S9)

where the limits of integration are: the shoreface depth  $D_T$ , which for our time scale of interest [Stive et al., 1991; Wright et al., 1991; Stive and de Vriend, 1995; Wright, 1995] is  $D_T$ , and the upper limit depth  $z^*[L]$ . Note that in the region  $0 < z < z^*$ , surf-zone processes dominate and the

processes are strongly non-linear and occur on time scales below those of interest here [*Carruthers et al.*, 2013], particularly we are interested in interactions with the deeper shoreface. For the parameter values in Table S1, we obtain  $K = 4,464 \,\mathrm{m}^2/\mathrm{y}$  under shallow-water waves, and  $K = 1,250 \,\mathrm{m}^2/\mathrm{y}$  under linear Airy wave theory. This motivates our choice to explore barrier dynamics for K spanning the range  $0-10,000 \,\mathrm{m}^2/\mathrm{y}$ .

### **Table**

Table S1. State variables and their dimensions.

Symbol	Meaning	Value
T	wave period	8 s
$H_{0}$	deep wave height during fair weather	1 m
$W_s$	settling velocity	0.033 m/s
g	gravity	$9.8 \text{ m/s}^2$
$e_s$	suspended sediment transport efficiency factor	0.01
$C_{s}$	friction factor	0.01
$\rho$	seawater density	$1,040 \text{ kg/m}^3$
$ ho_{s}$	sediment density	$2,650 \text{ kg/m}^3$
$z^*$	Upper limit depth	5 m
$D_{\!\scriptscriptstyle T}$	shoreface depth	15 m

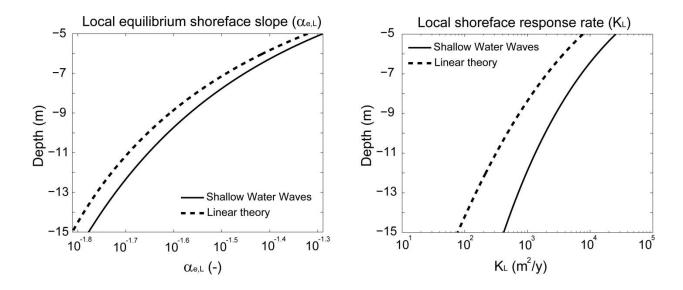


Figure S1. Plots of the local equilibrium slope  $\alpha_{e,L}$ , and the local shoreface response rate  $K_L$  as a function of water depth z.

## S2. Sensitivity to initial conditions:

All the results presented in the main manuscript assume that the initial barrier geometry is at static equilibrium (equation (29) in the main text). In this section, we investigate different initial conditions, finding that this assumption plays a minor role on the type of barrier response. In particular, we find that independently of the initial conditions the system eventually attains a constant amplitude and frequency of oscillation (Figure S2).

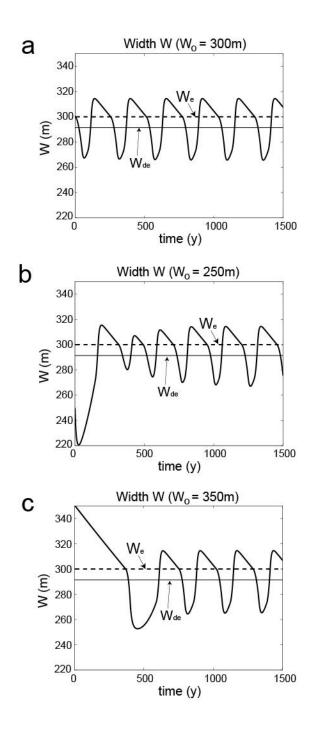


Figure S2. Time evolution of the barrier width for a barrier system undergoing 'periodic retreat' for different initial barrier widths. The input parameter values are included in Table S2, and the initial conditions in Table S3.

### S3. Sensitivity analysis to input parameters:

Changes in the sea-level rise rate and the bathymetric profile landward of the subaerial barrier back-barrier slope play a major role on the type of barrier response, as discussed in the main text (Figures 11 and 12). In this supplementary material, for completeness we present results from sensitivity analyses focusing on parameters found to play a minor role in the type of barrier response. These are the maximum deficit volume  $V_{d,\max}$ , shoreface toe depth  $D_T$ , equilibrium (or critical) barrier width  $W_e$  and height  $H_e$ , and equilibrium (or critical) shoreface slope  $\alpha_e$ .

#### 3.1 Maximum deficit volume

A change in the maximum deficit volume  $V_{d,\max}$  has a minor effect on the height and width drowning regions, but it enhances dynamic equilibrium in lieu of discontinuous retreat (Figure S3).

#### 3.2 Shoreface Toe Depth

The depth of the shoreface toe  $D_T$  is directly linked to wave, climate, sediment size, and the time scale under consideration [Stive and de Vriend, 1995; Cowell et al., 1999]. Given the simple geometry of a constant slope continental shelf, as we increase  $D_T$ , we concomitantly increase the back-barrier depth. By increasing the accommodation depth, this results in a reduction of the efficiency of the overwash flux and a reduction of the shoreface response rate. Thus, if we doubled  $D_T$ , maintaining the same barrier response would require doubling the amount of overwash and shoreface response rate (Figure S4).

# 2.3 Critical barrier width

As we might expect, narrow barriers are more prone to width drowning (Figure S5). Beyond a threshold barrier width, however, changes in  $W_e$  have little effect on barrier response.

## 2.4 Critical height and shoreface slope

Changes in  $H_{\scriptscriptstyle e}$  (Figure S6) or  $\alpha_{\scriptscriptstyle e}$  (Figure S7) have only a modest effect on barrier response.

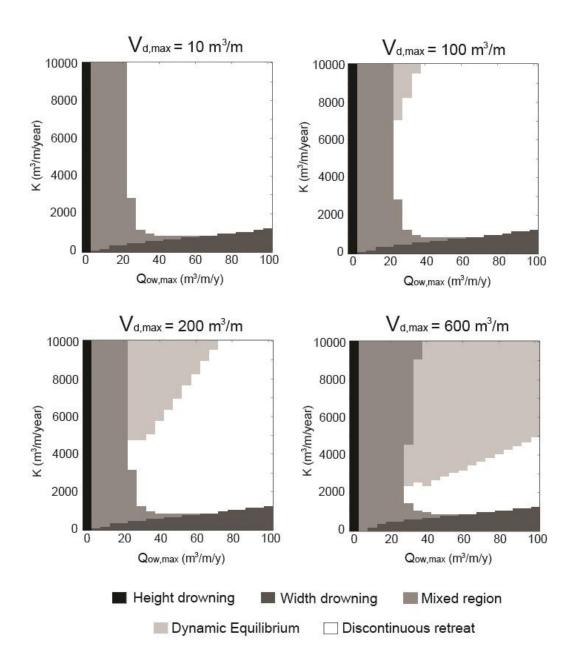


Figure S3. Phase diagrams of barrier evolution type for different values of the maximum deficit  $\text{volume } V_{d,\max} \text{ . Parameter values are included in Table S2.}$ 

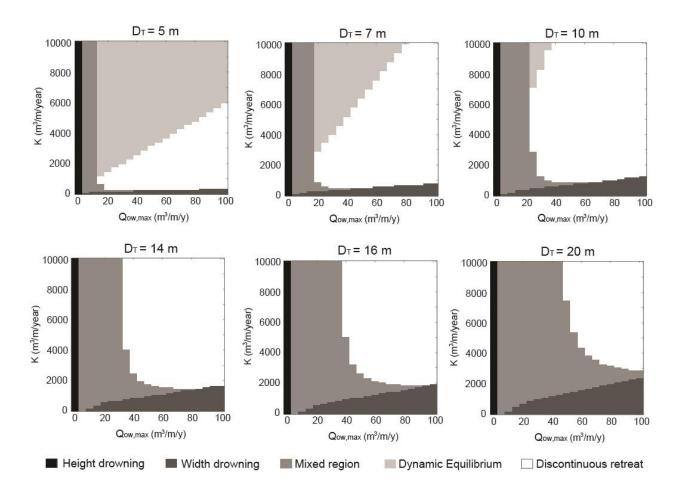


Figure S4. Phase diagrams of barrier evolution type for different values of the shoreface toe depth  $D_t$ . Note that an increase in  $D_t$  reduces the efficiency of  $Q_{OW, max}$  and K. For instance, the  $D_t = 20$  m regime diagram 'zooms' into region of the  $D_t = 10$  m regime diagram, within the range  $Q_{OW, max} = (0-50 \text{ m}^3/\text{m/y})$  and  $K = (0-5,000 \text{ m}^3/\text{m/y})$ . Parameter values are included in Table S2.

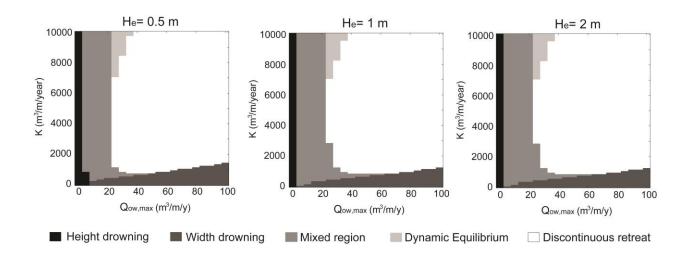


Figure S5. Phase diagrams of barrier evolution type for different values of the critical (or equilibrium) height  $H_e$ . Parameter values are included in Table S2.

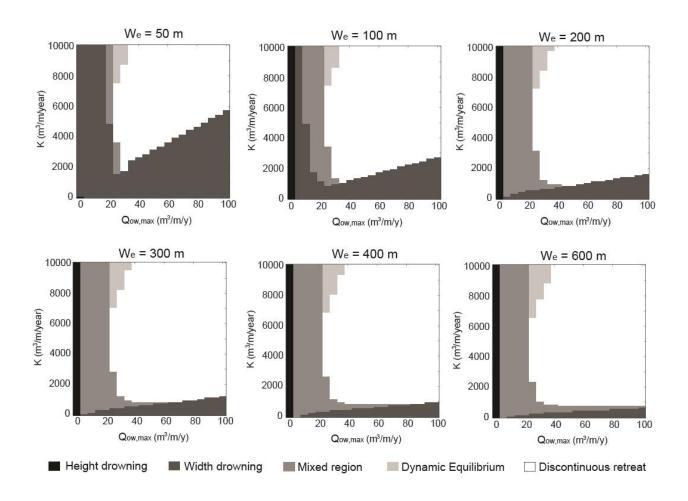


Figure S6. Phase diagrams of barrier evolution type for different values of the critical (or equilibrium) width  $W_e$ . Parameter values are included in Table S2.

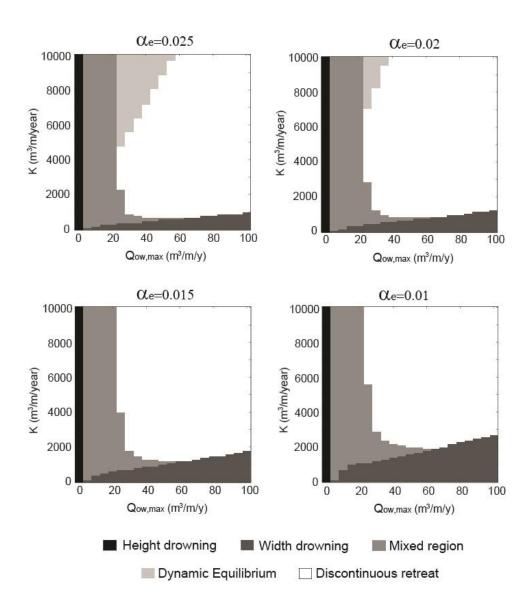


Figure S7. Phase diagrams of barrier evolution type for different values of the equilibrium shoreface slope  $\alpha_e$ . Parameter values are included in Table S2.

Table S2. Input parameters used in Figures S2 to S7.

Figure	K	$Q_{\scriptscriptstyle OW, ext{max}}$	$V_{d,\max}$	$lpha_{_e}$	β	$W_{e}$	$H_{_{\it e}}$	$D_{\scriptscriptstyle T}$	$\dot{z}$
	$(m^3/m/y)$	$(m^3/m/y)$	$(m^3/m)$	(-)	(-)	( <i>m</i> )	<i>(m)</i>	( <i>m</i> )	(mm/y)
S2	5,000	70	300	0.02	0.001	300	2	10	2
<b>S</b> 3	varies	varies	varies	0.02	0.001	300	2	10	2
S4	varies	varies	100	0.02	0.001	300	2	varies	2
<b>S</b> 5	varies	varies	100	0.02	0.001	300	varies	10	2
<b>S</b> 6	varies	varies	100	0.02	0.001	varies	2	10	2
S7	varies	varies	100	varies	0.001	300	2	10	2

Table A3. Initial Conditions.

Figure	$\alpha_0$ (-)	$W_0$ $(m)$	$H_0$ $(m)$
S2a	$\alpha_{\scriptscriptstyle e}$	$W_e$	$H_{e}$
S2b	$lpha_{_e}$	250	$H_{\scriptscriptstyle e}$
S2c	$\alpha_{\scriptscriptstyle e}$	350	$H_{_{\it e}}$

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