

Modeling and simulation of the cavitation phenomenon in space engine turbopumps

Joris Cazé

Directeurs de thèse : Fabien Petitpas, Eric Daniel

Référents CNES : Sébastien Le Martelot, Matthieu Queguineur



Summary

Introduction

Modeling of the cavitation phenomenon

- State-of-the-art
- Two-phase flow approach
- Two-phase flow model
- Blades motion

Numerical method

- Numerical scheme
- MRF fluxes
- Mesh mapping

Results

- Test case setup
- Single-phase flow behavior: pump characteristic
- Two-phase flow: cavitating regime

Conclusions & perspectives

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Introduction: overview

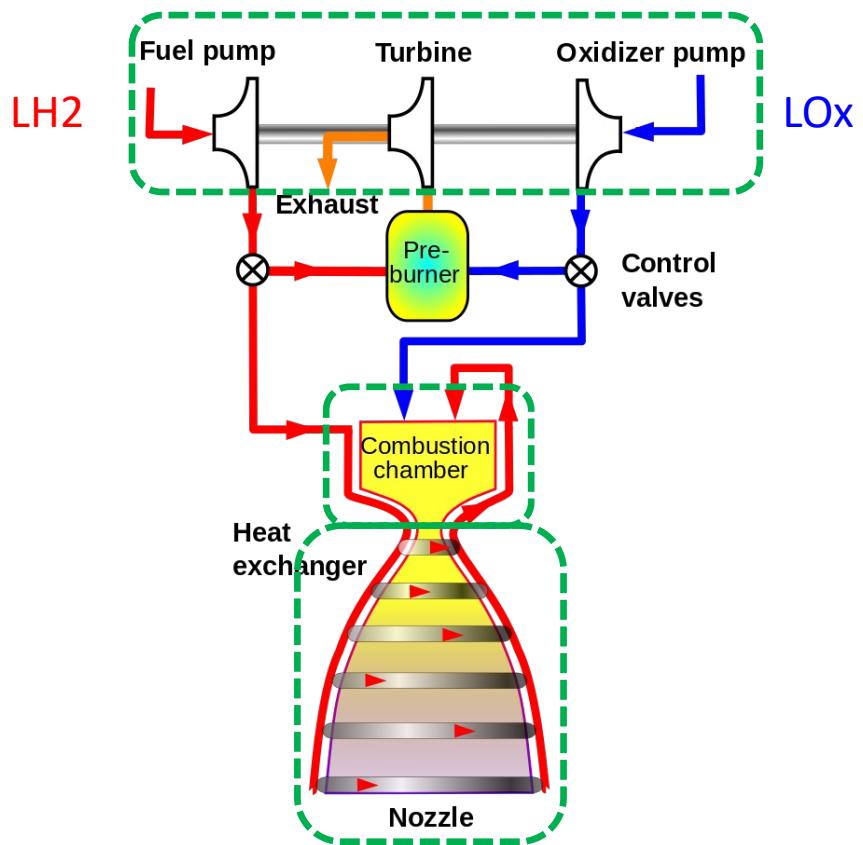
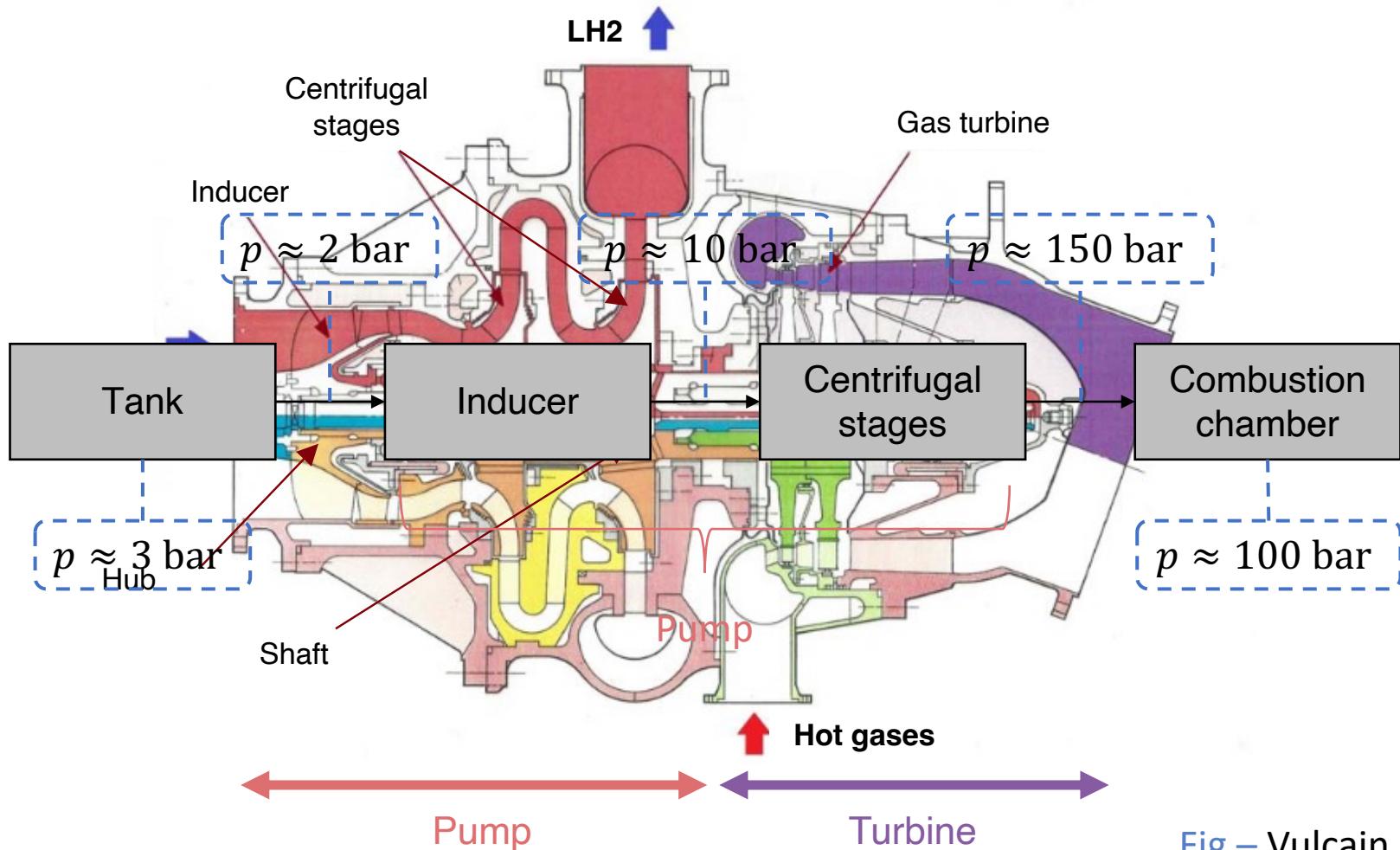


Fig – Vulcain gas generator cycle



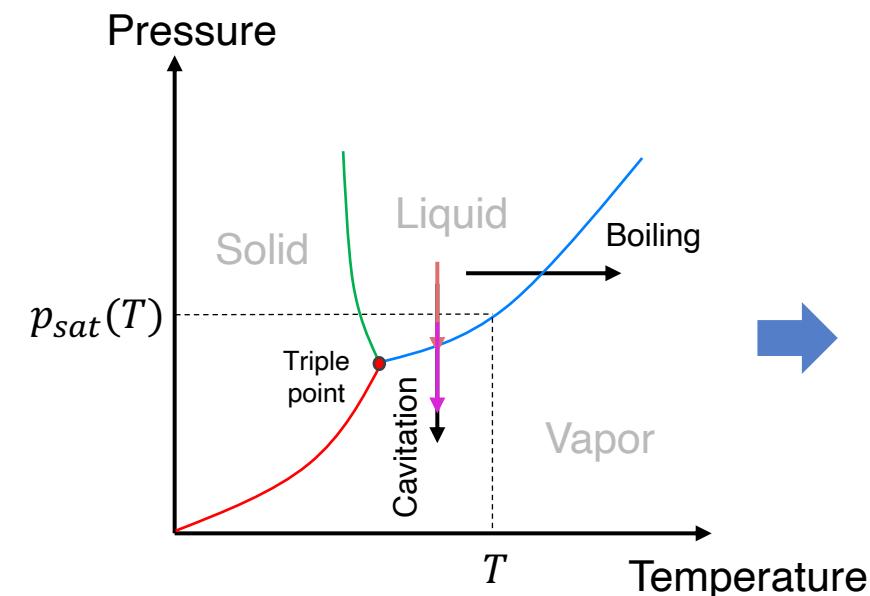
Fig – Vulcain 1 rocket engine

Introduction: turbopump



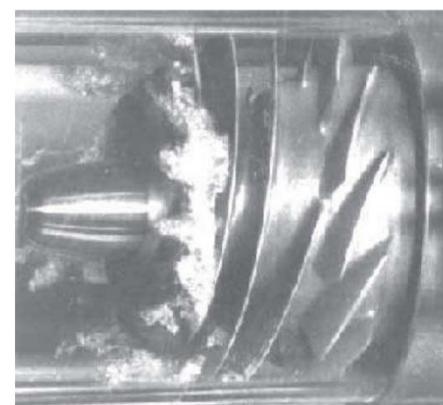
Introduction: cavitation

No cavitation or moderate cavitation



$$p(T) < p_{sat}(T)$$

- Pump suction
- Local overspeed
- Shear stress
- Geometry



Strong cavitation



- Mechanical instabilities
- Performance drop

Fig – Cavitation on SSME LOx inducer [Braisted, 1980]

Introduction: objectives

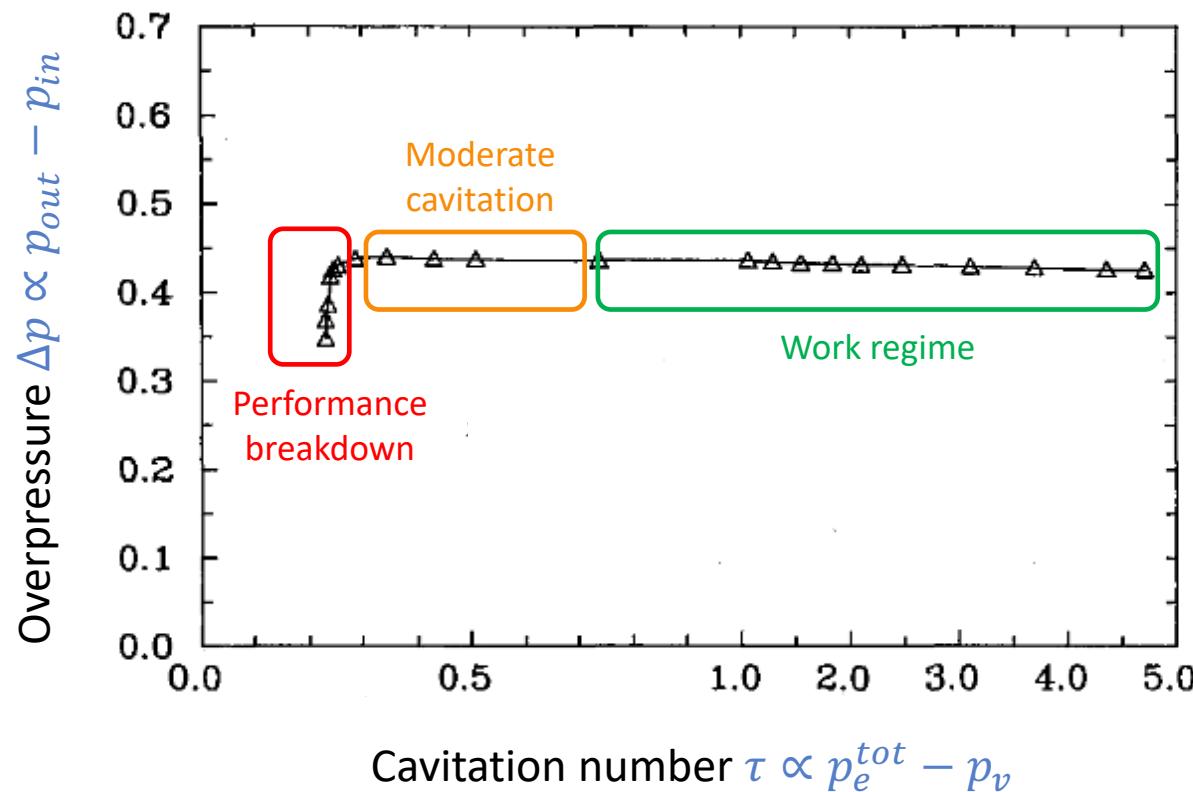
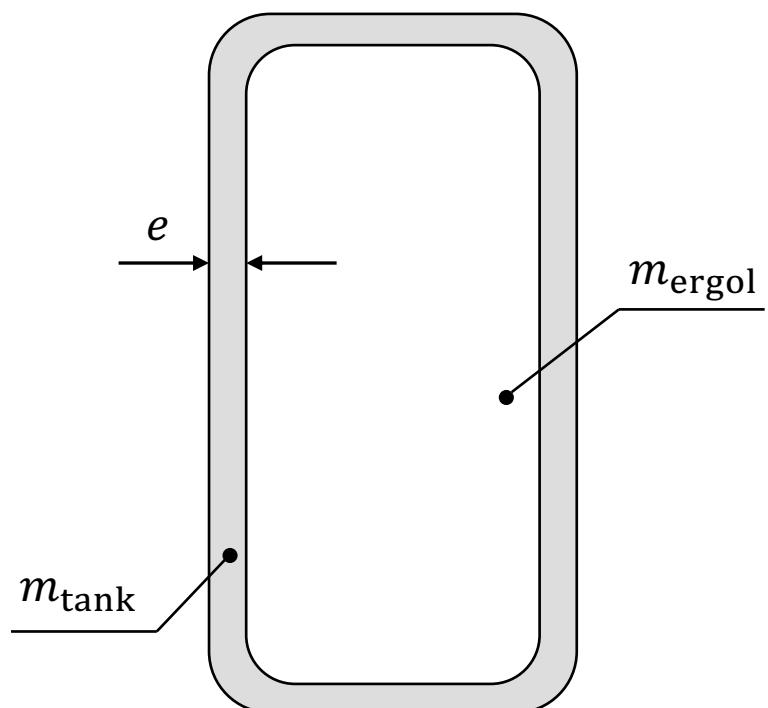


Fig – Experimental performance curve on a centrifugal pump [Franz *et al.*, 1989]

Fig – Illustration of a liquid propellant tank



→ Reduce the rocket dry mass

Introduction: cavitation modeling

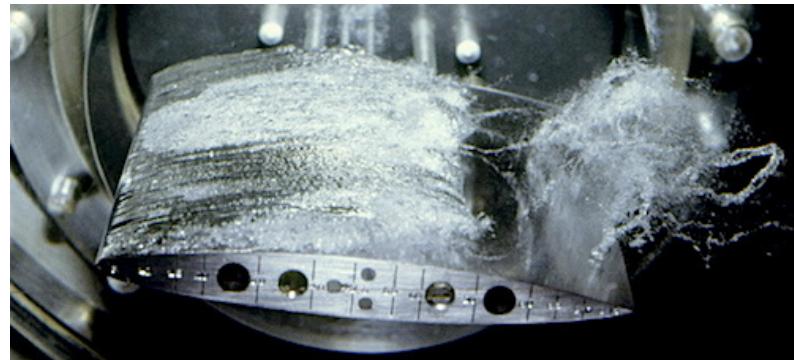


Fig – Cloud cavitation on hydrofoil (left) - Bubble cavitation on propeller (right)
[Franc and Michel, 1995]

Cavitation characteristics:

- Two-phase flow
- Phase change
- (in)compressible regions
- 3D

Pump characteristic:

- Rotor motion

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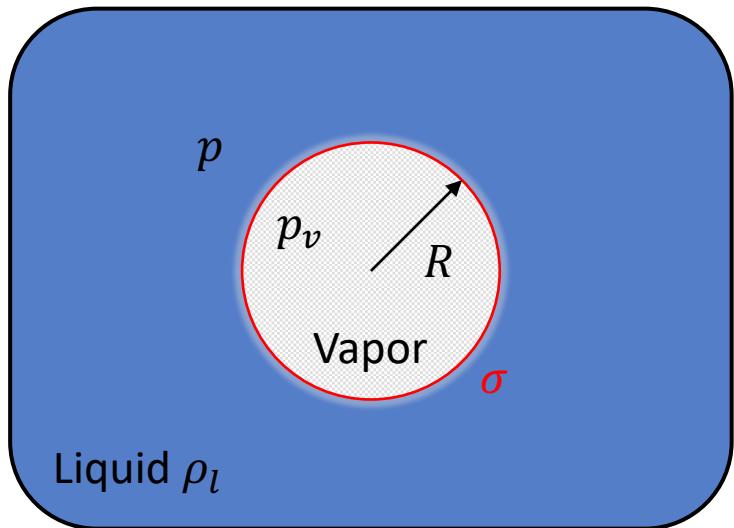
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Modeling: state-of-the-art

Bubble dynamics



Rayleigh-Plesset equation

$$R \frac{d^2 R}{dt^2} + \frac{3}{2} \left(\frac{dR}{dt} \right)^2 + \frac{2\sigma}{R} = \frac{p_v - p}{\rho_l}$$

💡 $\dot{m} = \pm N 4\pi R^2 \rho_v \sqrt{\frac{2}{3} \frac{p_v - p}{\rho_l}}$



Empirical correction
coefficients

[Singhal *et al.*, 2002] [Zhang *et al.*, 2019]

Modeling: state-of-the-art

Barotropic EOS

- Homogeneous model
- 3D RANS → FINE™/Turbo code

$$\rho = \alpha_v \rho_v + (1 - \alpha_v) \rho_l$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = 0$$

$$\frac{\partial \rho \mathbf{u}}{\partial t} + \nabla \cdot (\rho \mathbf{u} \otimes \mathbf{u} + p \mathbf{I}) = \nabla \cdot \boldsymbol{\tau} + \rho \mathbf{F}$$

$$\rho = \rho(p) = \frac{\rho_l + \rho_v}{2} + \frac{\rho_l - \rho_v}{2} \sin \left(\frac{p - p_v}{c_{min}^2} \frac{2}{\rho_l - \rho_v} \right)$$

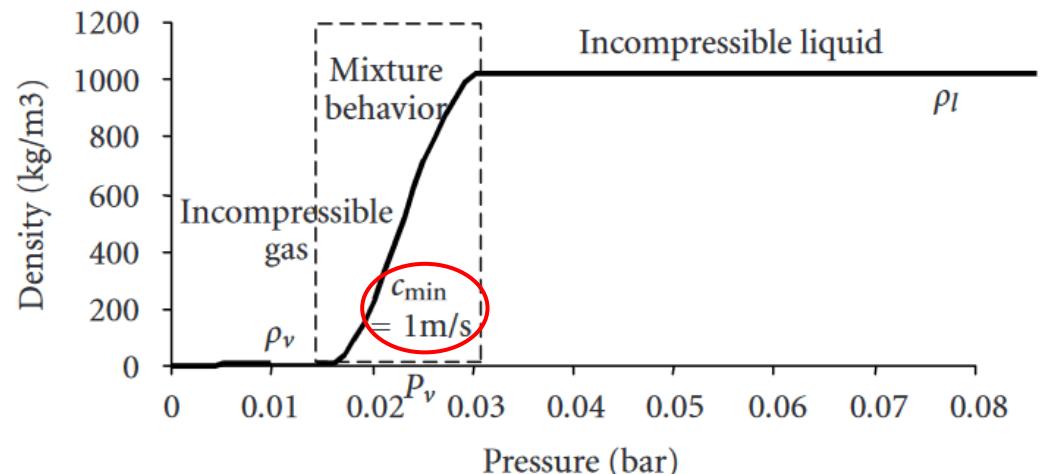


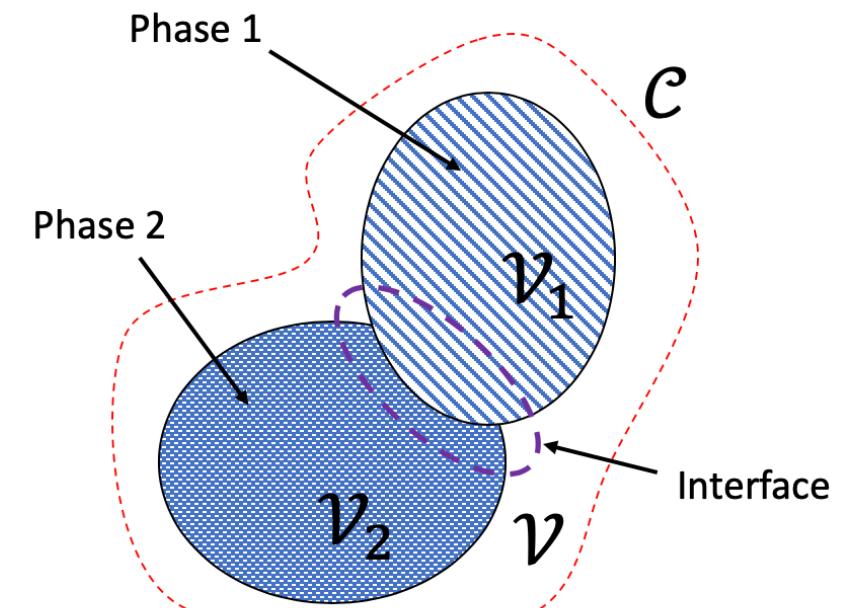
Fig - Barotropic state law $\rho(p)$ for water from [Coutier-Delgosha et al., 2005]

- 👍 No need to define mass transfer
- 👍 Better thermodynamic behavior with the added energy equation [Goncalves et al., 2010]

Modeling: two-phase flow approach

- Hyperbolic model → waves propagation
- Able to handle several interface types L/V or others (multi-species, non-condensable gases)
- Complete thermodynamic description for each phase
- Taking into account the thermodynamic equilibrium

} Phase change



Modeling: two-phase flow model

Velocity equilibrium model ($\mathbf{u}_1 = \mathbf{u}_2$)

Used to study cavitating flow [Petitpas *et al.*, 2011]

For two-phase flow $k = 1, 2$:

$$\partial_t \alpha_1 + \nabla \cdot (\alpha_1 \mathbf{u}) - \alpha_1 \nabla \cdot \mathbf{u} = \mu(p_1 - p_2) + \frac{\dot{m}/\rho_I}{v(g_2 - g_1)/\rho_I}$$

+

EOS
 $e_k = e_k(\rho_k, p_k)$

$$\partial_t(\alpha_k \rho_k) + \nabla \cdot (\alpha_k \rho_k \mathbf{u}) = \pm_k \frac{\dot{m}}{v(g_2 - g_1)}$$

$$\partial_t(\rho \mathbf{u}) + \nabla \cdot (\rho \mathbf{u} \otimes \mathbf{u}) + \nabla(\alpha_1 p_1 + \alpha_2 p_2) = \mathbf{0}$$

$$\partial_t(\alpha_k \rho_k e_k) + \nabla \cdot (\alpha_k \rho_k e_k \mathbf{u}) + \alpha_k p_k \nabla \cdot \mathbf{u} = \mp_k \mu p_I (p_1 - p_2) \pm_k h \frac{\dot{m}}{v(g_2 - g_1)} \pm_k \frac{\dot{Q}}{\theta(T_2 - T_1)}$$



$$\partial_t U + \nabla \cdot F(U) + H(U) \nabla \cdot \mathbf{u} = R(U)$$

[Kapila *et al.*, 2000]
[Saurel *et al.*, 2009]

Modeling: thermodynamical closure

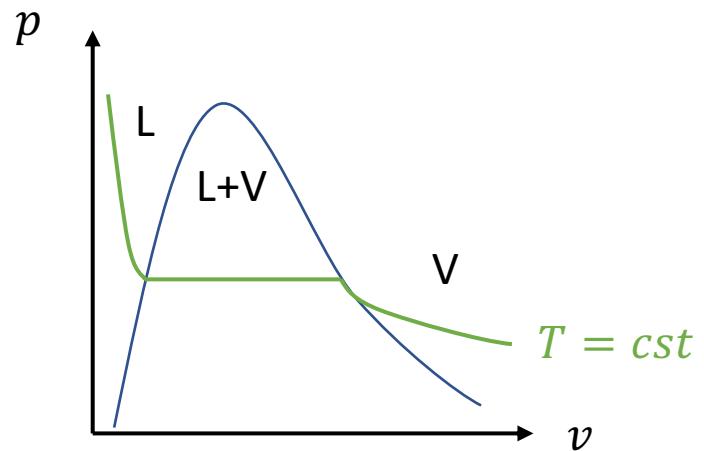


Fig – Typical phase diagram (p, v)

- Each phase is governed by its Equation Of State (EOS)
- EOS \neq Van Der Waals
 - Stiffened Gas
 - Noble-Abel Stiffened Gas
- Calibration of EOS parameters based on experimental saturation curve
[Le Métayer *et al.*, 2004]
[Le Métayer *et al.*, 2016]

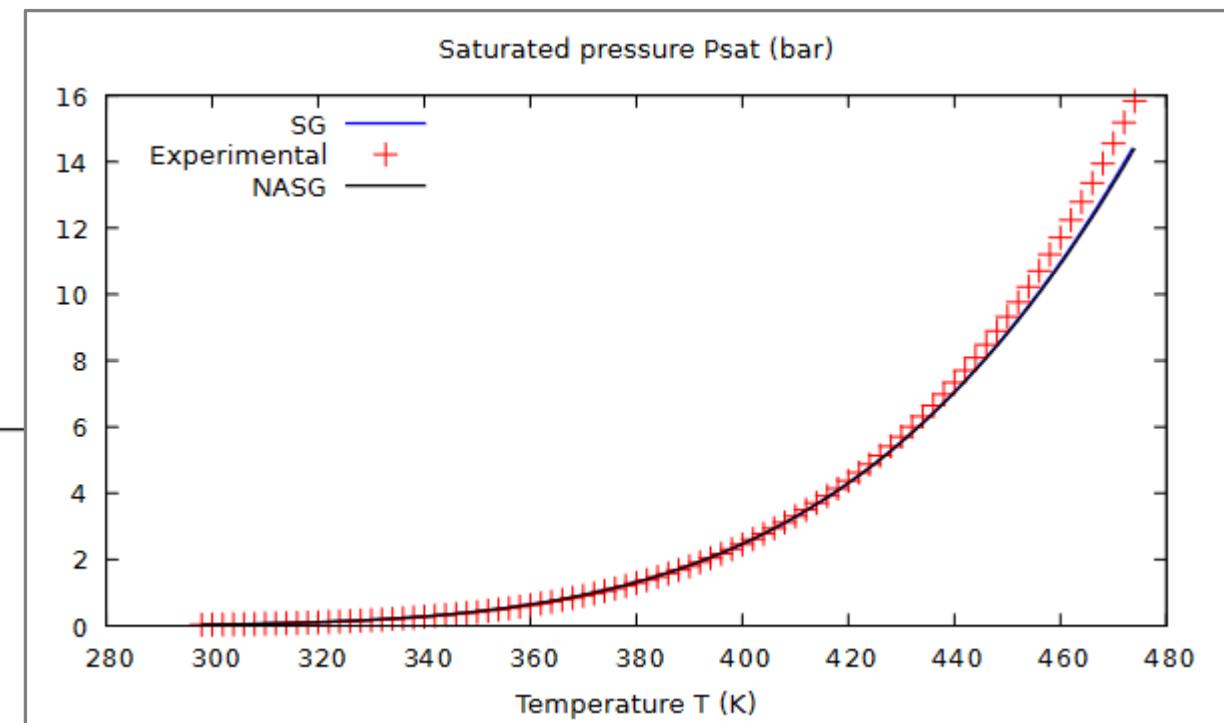
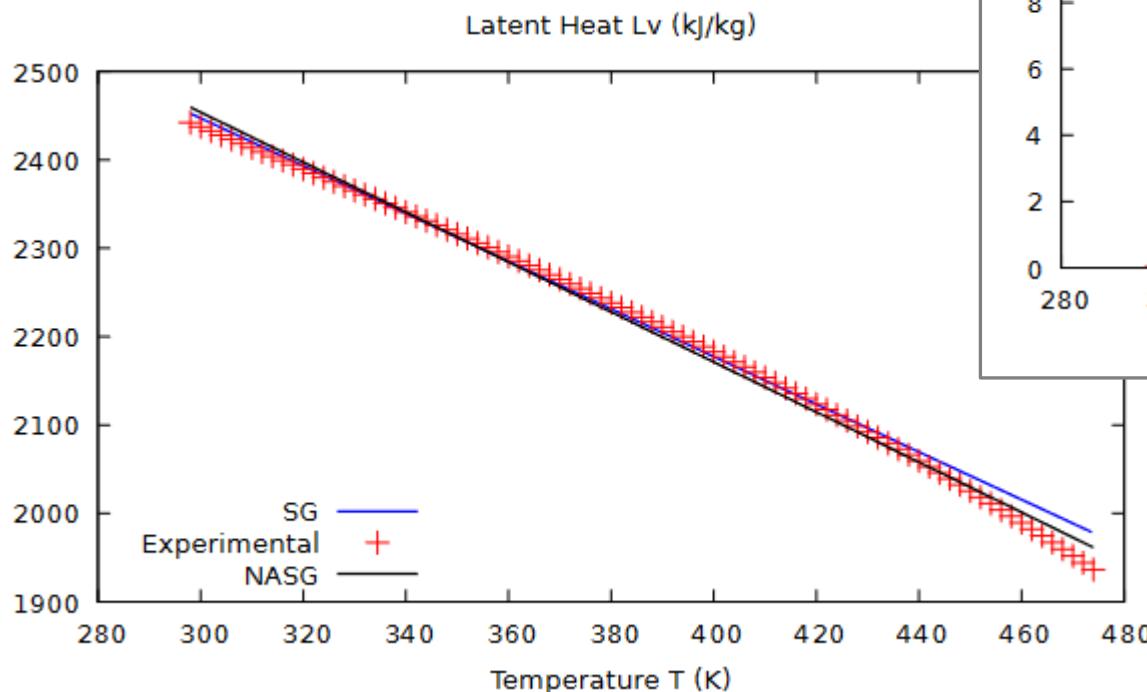
$$p(v, e) = (\gamma - 1) \frac{e - q}{v - b} - \gamma p_\infty$$

Thermal agitation Matter cohesion
Repulsive short distance effects

Modeling: thermodynamical closure

Example Water EOS calibration for

- Stiffened Gas
- Noble-Abel Stiffened Gas



Fluid state is solely based on
thermodynamic considerations

Modeling: blades motion

Moving Reference Frame (MRF) method [Combrinck and Dala, 2014]

Lagrangian derivation based on classical point mechanics

Change of reference frame: **inertial one** → **rotating one**

Rotational effects: **centrifugal force** & **Coriolis force**

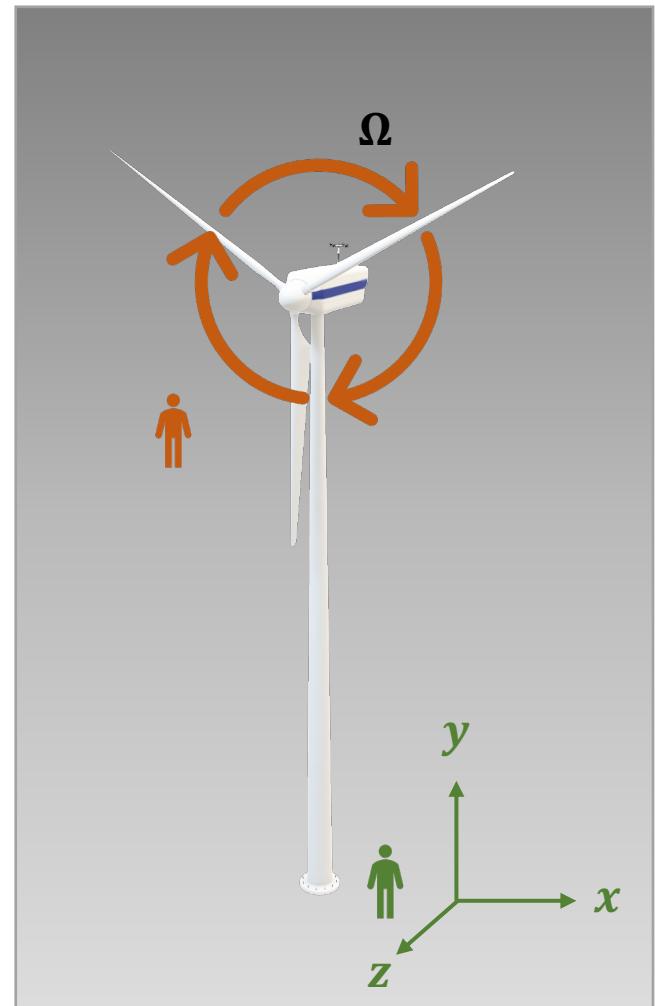
Euler equations:

$$\partial_t \rho + \nabla \cdot \rho \mathbf{u}_r = 0$$

$$\partial_t \rho \mathbf{u}_r + \nabla \cdot (\rho \mathbf{u}_r \otimes \mathbf{u}_r + pI) = -2 \rho \boldsymbol{\Omega} \wedge \mathbf{u}_r - \rho \boldsymbol{\Omega} \wedge (\boldsymbol{\Omega} \wedge \mathbf{x})$$

$$\rightarrow \partial_t \rho E_r + \nabla \cdot ((\rho E_r + p) \mathbf{u}_r) = -\rho \mathbf{u}_r \cdot (\boldsymbol{\Omega} \wedge (\boldsymbol{\Omega} \wedge \mathbf{x}))$$

$$E_r = e + \frac{1}{2} \mathbf{u}_r^2$$



$$\mathbf{u} = \mathbf{u}_r + \boldsymbol{\Omega} \times \mathbf{r}$$

[Cazé et al., 2022]

Modeling: blades motion

Change of reference frame: inertial one → rotating one

Rotational effects: centrifugal force & Coriolis force

Two-phase flow model:

$$\partial_t \alpha_1 + \mathbf{u}_r \cdot \nabla \alpha_1 = \mu(p_1 - p_2) + \frac{\dot{m}}{\rho_I}$$

$$\partial_t(\alpha_k \rho_k) + \nabla \cdot (\alpha_k \rho_k \mathbf{u}) = \pm_k \dot{m}$$

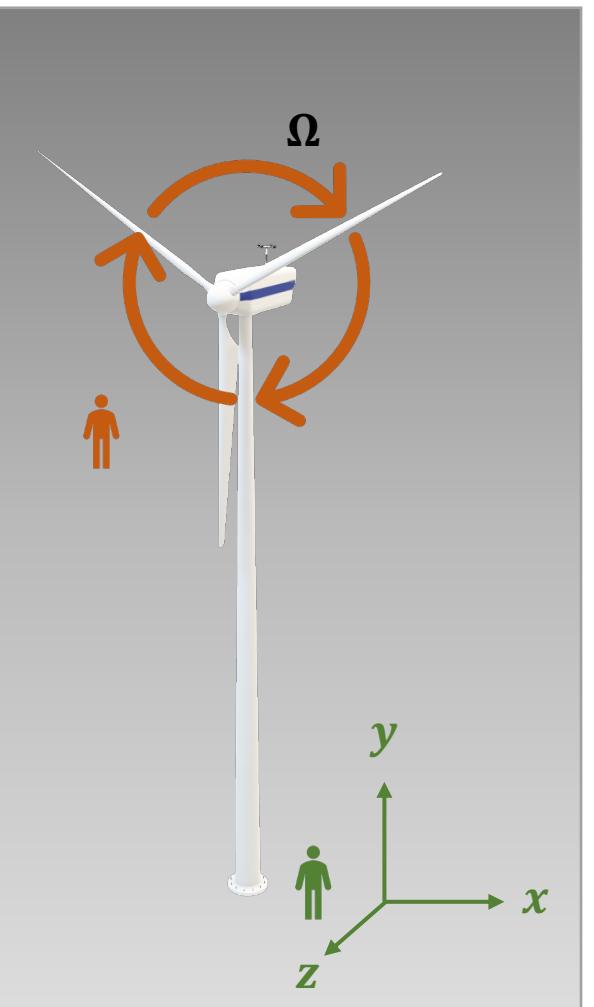
$$\partial_t(\rho \mathbf{u}_r) + \nabla \cdot (\rho \mathbf{u}_r \otimes \mathbf{u}_r) + \nabla(\alpha_1 p_1 + \alpha_2 p_2) = -2 \rho \Omega \wedge \mathbf{u}_r - \rho \Omega \wedge (\Omega \wedge \mathbf{x})$$

$$\partial_t(\alpha_k \rho_k e_k) + \nabla \cdot (\alpha_k \rho_k e_k \mathbf{u}_r) + \alpha_k p_k \nabla \cdot \mathbf{u}_r = \mp_k \mu p_I (p_1 - p_2) \pm_k (\dot{Q} + h_I \dot{m})$$

$$\partial_t \rho E_r + \nabla \cdot ((\rho E_r + p) \mathbf{u}_r) = -\rho \mathbf{u}_r \cdot (\Omega \wedge (\Omega \wedge \mathbf{x}))$$

$$E_r = e + \frac{1}{2} \mathbf{u}_r^2$$

$$e = Y_1 e_1 + Y_2 e_2$$



$$\mathbf{u} = \mathbf{u}_r + \boldsymbol{\Omega} \times \mathbf{r}$$

[Cazé et al., 2022]

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Numerical method: Overview

Solving $\partial_t U + \nabla \cdot F(U) + \Sigma(U) = R(U) + S(U)$ using FV method:

1. Homogeneous system with Godunov's scheme $\partial_t U + \nabla \cdot F(U) + \Sigma(U) = 0$
2. Source step $\partial_t U = S(U)$
3. Phase change $\partial_t U = R_{pTg}(U)$

[Schmidmayer *et al.*, 2020] & [Schmidmayer, Cazé *et al.*, 2022]

Numerical method: numerical scheme

Solving $\partial_t \mathbf{U} + \nabla \cdot \mathbf{F}(\mathbf{U}) + \Sigma(\mathbf{U}) = \mathbf{R}(\mathbf{U}) + \mathbf{S}(\mathbf{U})$ using FV method:

1. Homogeneous system with Godunov's scheme $\partial_t \mathbf{U} + \nabla \cdot \mathbf{F}(\mathbf{U}) + \Sigma(\mathbf{U}) = 0$

$$\mathbf{U}_i^{n+1} = \mathbf{U}_i^n - \frac{\Delta t}{V_i} \left(\sum_{s=1}^{N_s} A_s \mathbf{F}_s^* \cdot \mathbf{n}_s \right) - \frac{\Delta t}{V_i} H(\mathbf{U}_i^n) \left(\sum_{s=1}^{N_s} A_s \mathbf{u}_s^* \cdot \mathbf{n}_s \right)$$

with $\mathbf{F}_s^* = RP(\mathbf{U}_i^n, \mathbf{U}_j^n)$ using approximate Riemann solver

Numerical method: numerical scheme

Solving $\partial_t U + \nabla \cdot F(U) + \Sigma(U) = R(U) + S(U)$ using FV method:

1. Homogeneous system with Godunov's scheme $\partial_t U + \nabla \cdot F(U) + \Sigma(U) = 0$
2. Source step $\partial_t U = S(U)$



Take into account Moving Reference Frame terms

$$U_i^{n+1} = \widetilde{U}_i^n + \Delta t S(U_i^n)$$

or RK2, RK4 scheme

Numerical method: numerical scheme

Solving $\partial_t \mathbf{U} + \nabla \cdot \mathbf{F}(\mathbf{U}) + \Sigma(\mathbf{U}) = \mathbf{R}(\mathbf{U}) + \mathbf{S}(\mathbf{U})$ using FV method:

1. Homogeneous system with Godunov's scheme $\partial_t \mathbf{U} + \nabla \cdot \mathbf{F}(\mathbf{U}) + \Sigma(\mathbf{U}) = 0$
2. Source step $\partial_t \mathbf{U} = \mathbf{S}(\mathbf{U})$
3. Phase change $\partial_t \mathbf{U} = \mathbf{R}_{pTg}(\mathbf{U})$

$$\partial_t \alpha_1 = \mu(p_1 - p_2) + \nu(g_2 - g_1)/\rho_I$$

$$\partial_t (\alpha_k \rho_k) = \pm_k \nu(g_2 - g_1)$$

$$\partial_t (\rho \mathbf{u}) = \mathbf{0}$$

$$\partial_t (\alpha_k \rho_k e_k) = \mp_k \mu p_I (p_1 - p_2) \pm_k h_I \nu (g_2 - g_1) \pm_k \theta (T_2 - T_1)$$

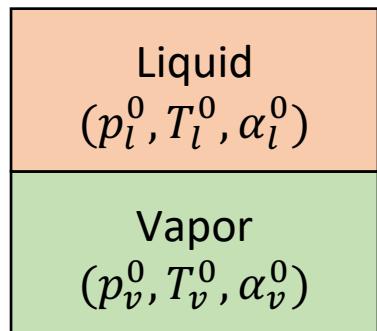
Evaluation of mass transfer terms by
relaxation processes



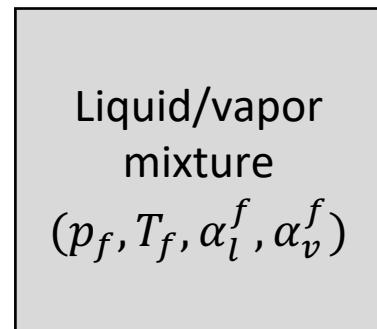
Thermodynamic problem of return to
equilibrium

with $\mu, \nu, \theta \rightarrow \infty$

Numerical method: phase change



Relax. p, T, g



Saturation state

$$\begin{aligned} T_l^f &= T_v^f = T^f \\ p_l^f &= p_v^f = p^f \\ g_l^f &= g_v^f = g^f \end{aligned}$$

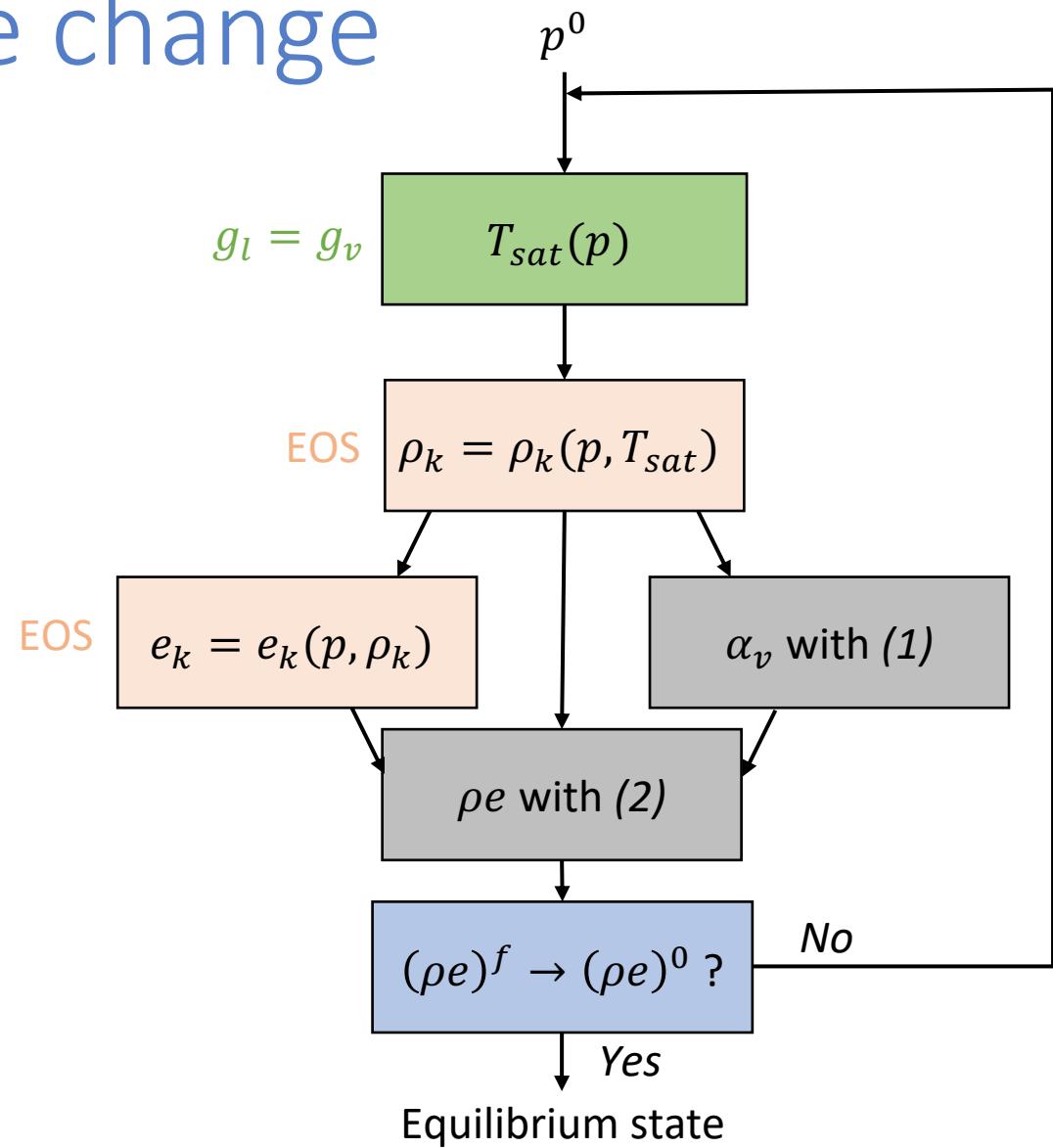
EDO

Mass conservation

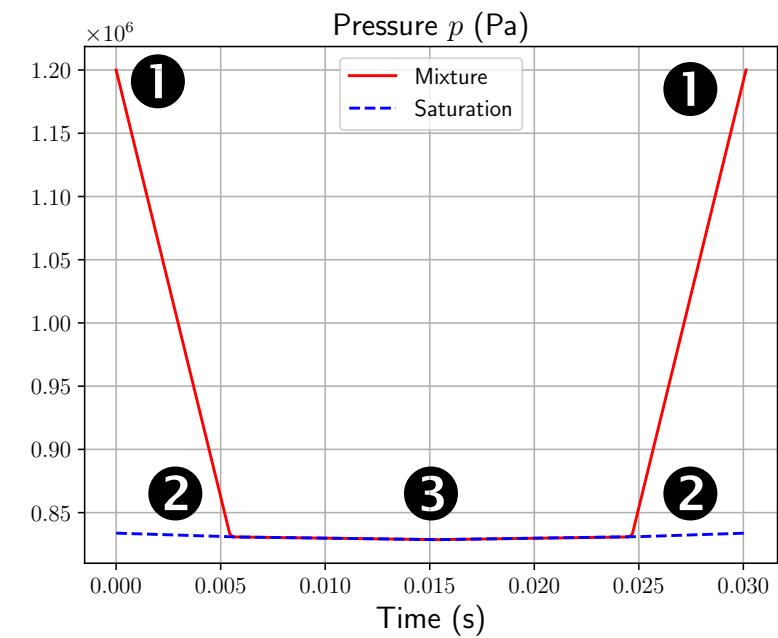
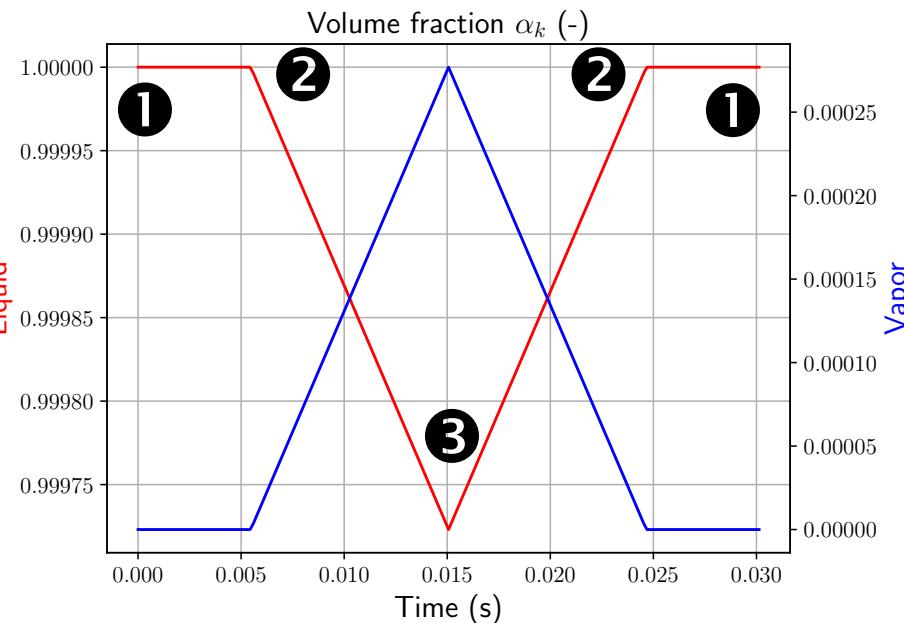
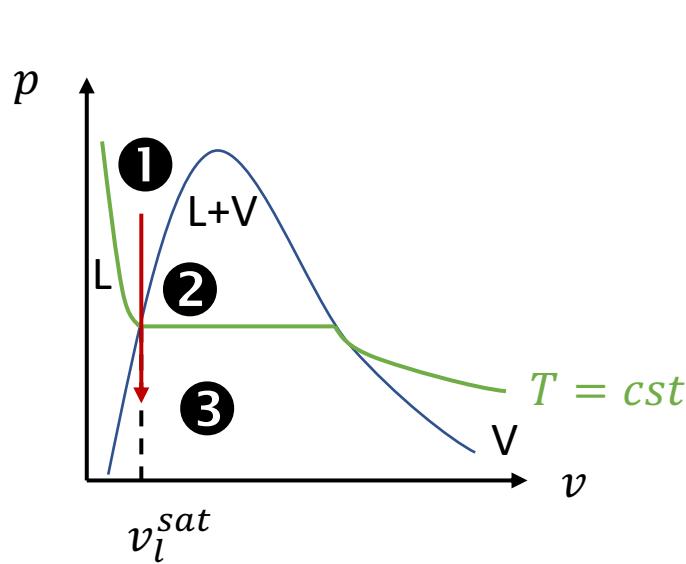
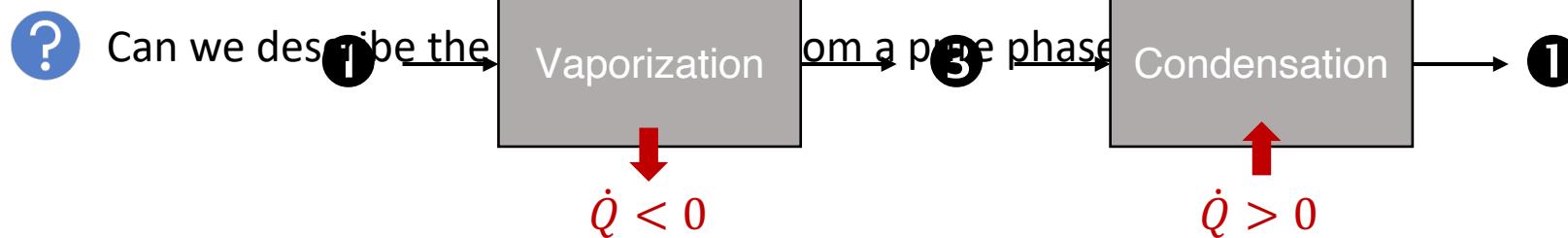
$$\rho^f = \alpha_v^f \rho_v^f + \alpha_l^f \rho_l^f = \rho^0 \quad (1)$$

Energy conservation

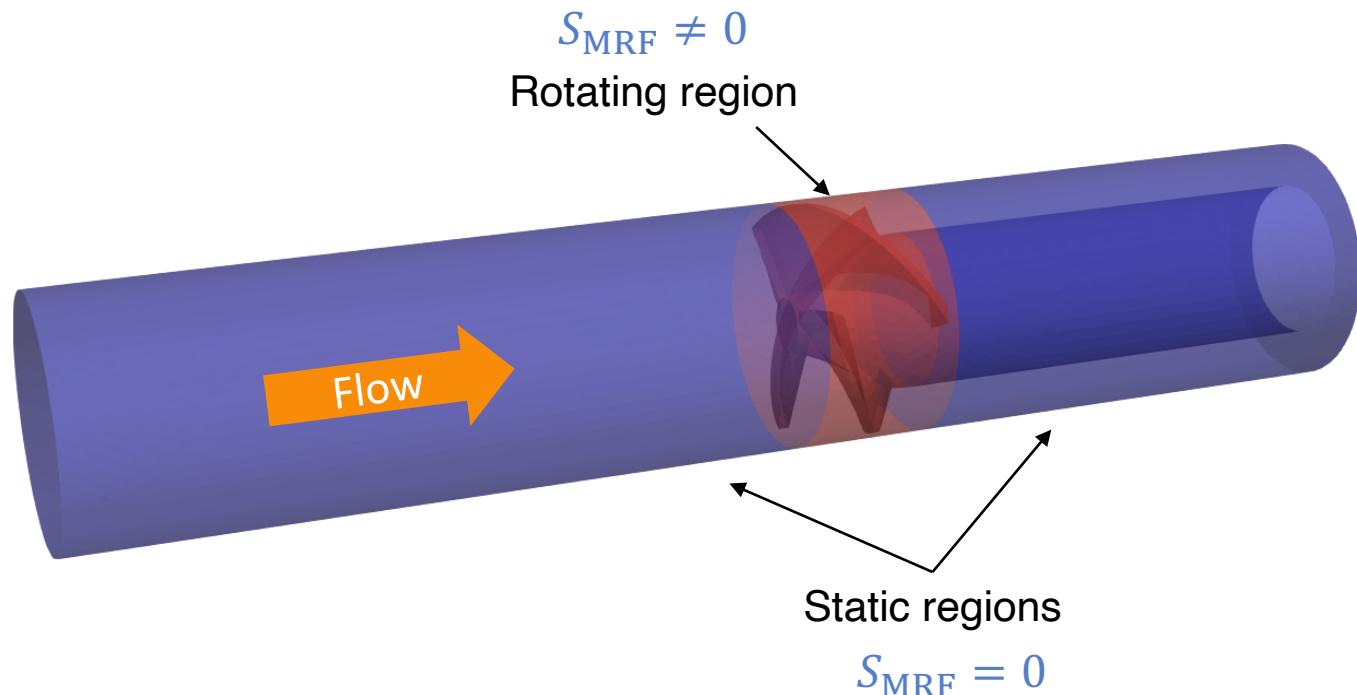
$$(\rho e)^f = (\alpha \rho e)_v^f + (\alpha \rho e)_l^f = (\rho e)^0 \quad (2)$$



Numerical method: 0D phase change validation



Numerical method: MRF fluxes

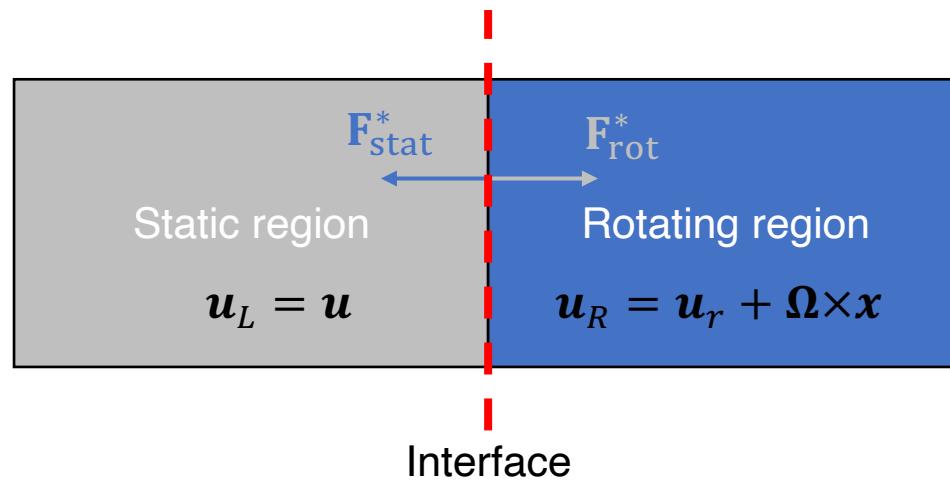


Model in compact form:

$$\partial_t U + \nabla \cdot F(U) + \Sigma(U) = R(U) + S_{\text{MRF}}(U)$$

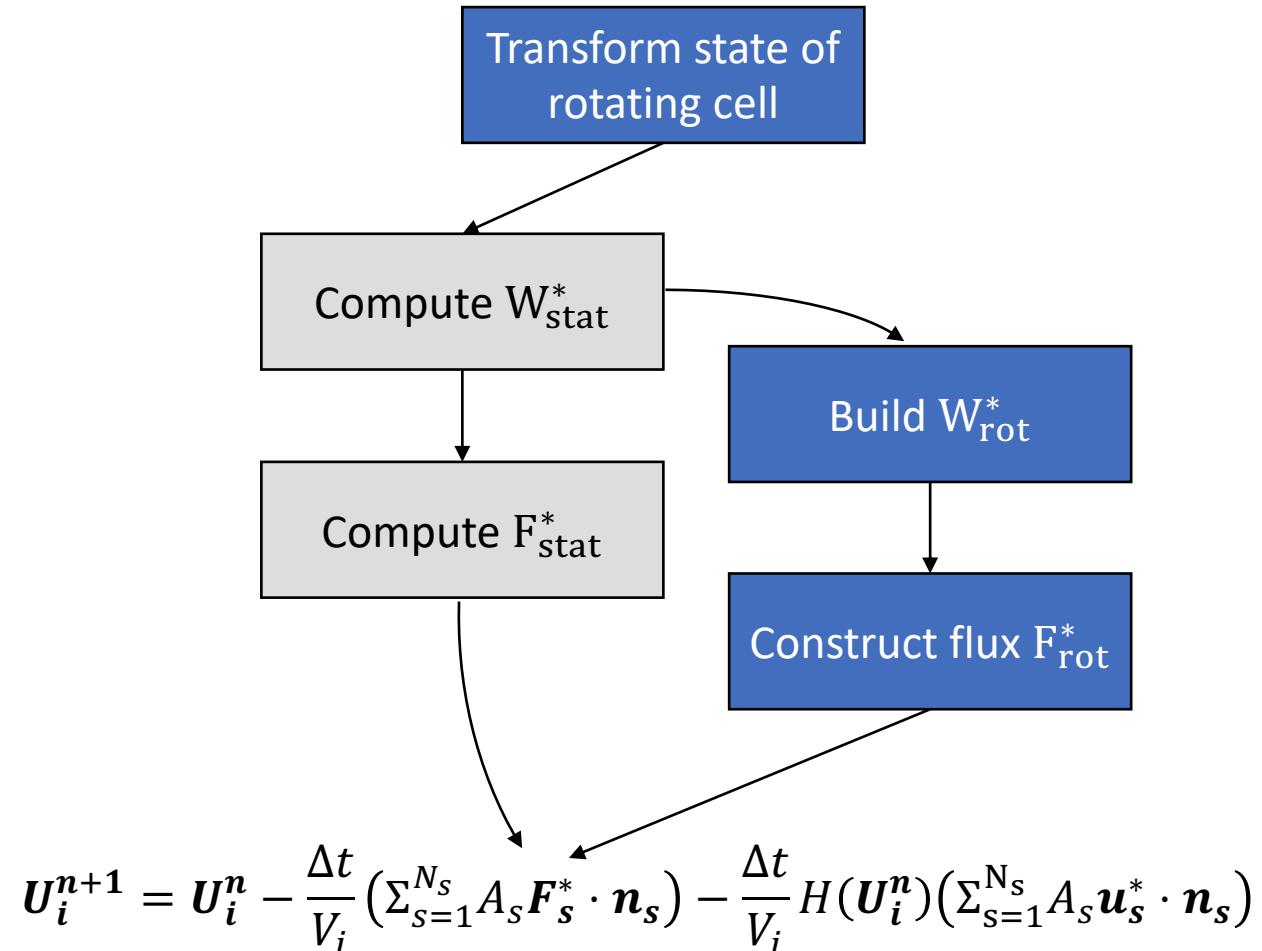
Numerical method: MRF fluxes

💡 Modified Riemann problem



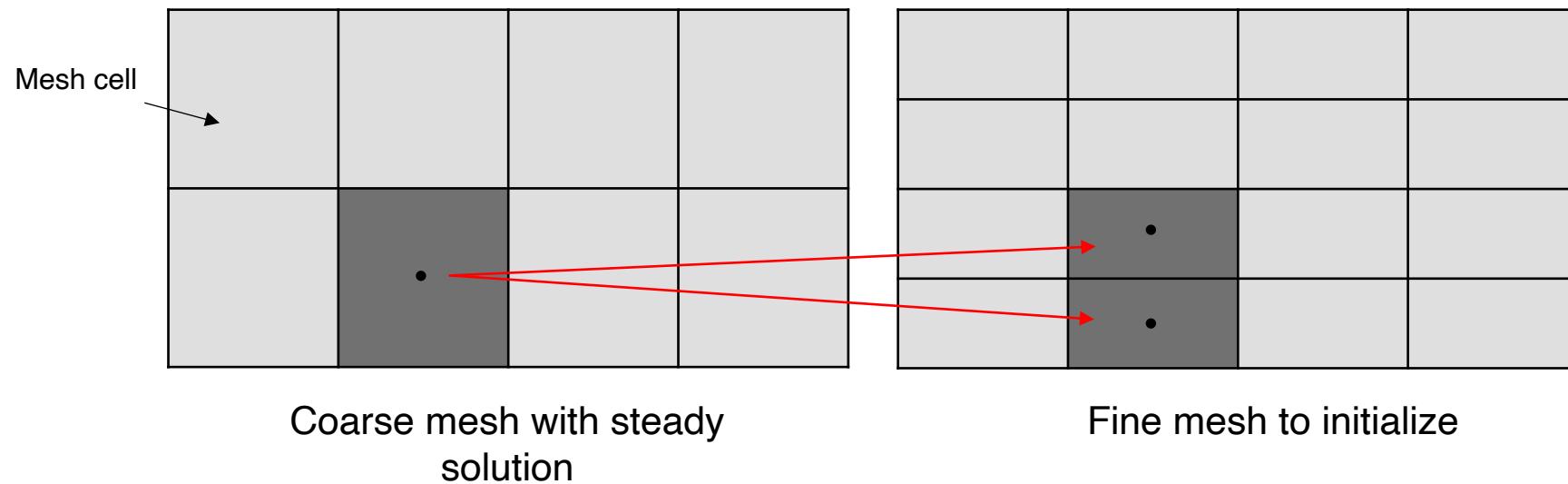
$$\text{Absolute vel.} = \text{Relative vel.} + \text{Rot. vel.}$$

$$\mathbf{u} = \mathbf{u}_r + \boldsymbol{\Omega} \times \mathbf{x}$$

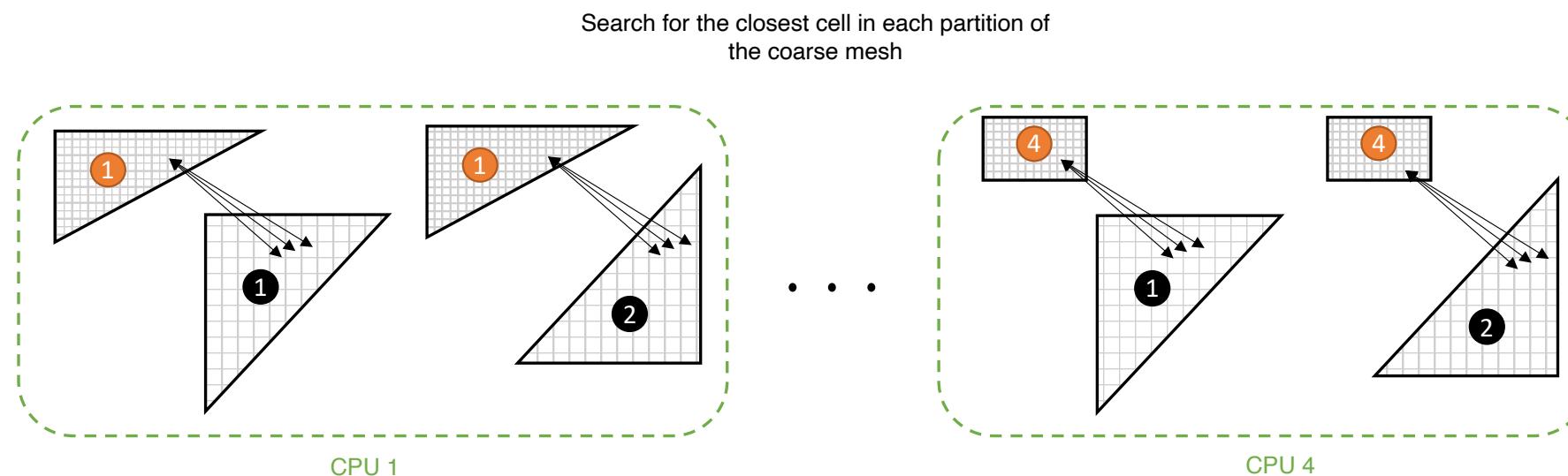
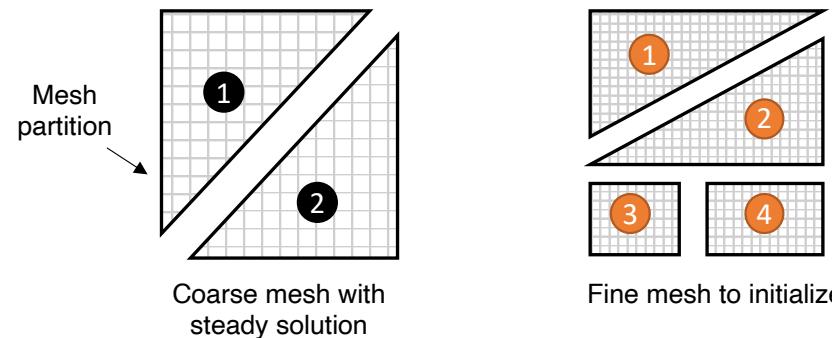


Numerical method: mesh mapping

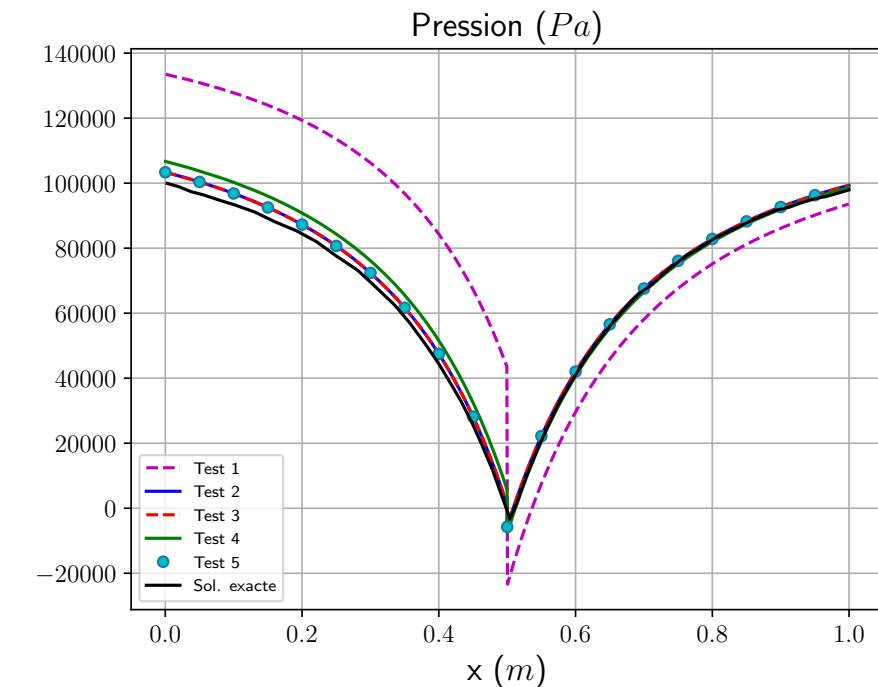
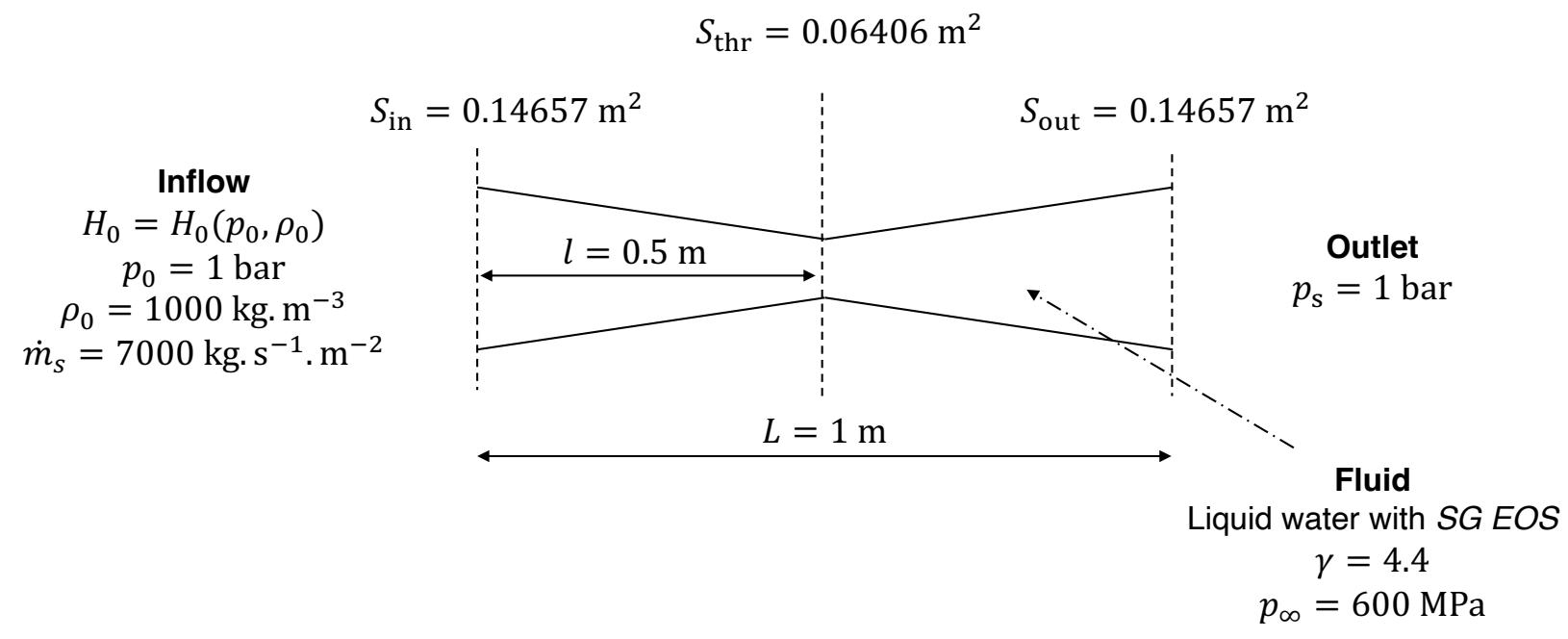
🎯 Reduce computation time



Numerical method: mesh mapping



Numerical method: mesh mapping validation



Test	Mesh	Setup	CPUs	$t_{\text{simu}}/t_{\text{réf}}$	$t_{\text{tot}}/t_{\text{réf}}$
1	1000	Normal	1	1	1
2	10 000	Normal	10	15	15
3	10 000	Mapping on 1	10	8	9

→ Mapping 41% faster

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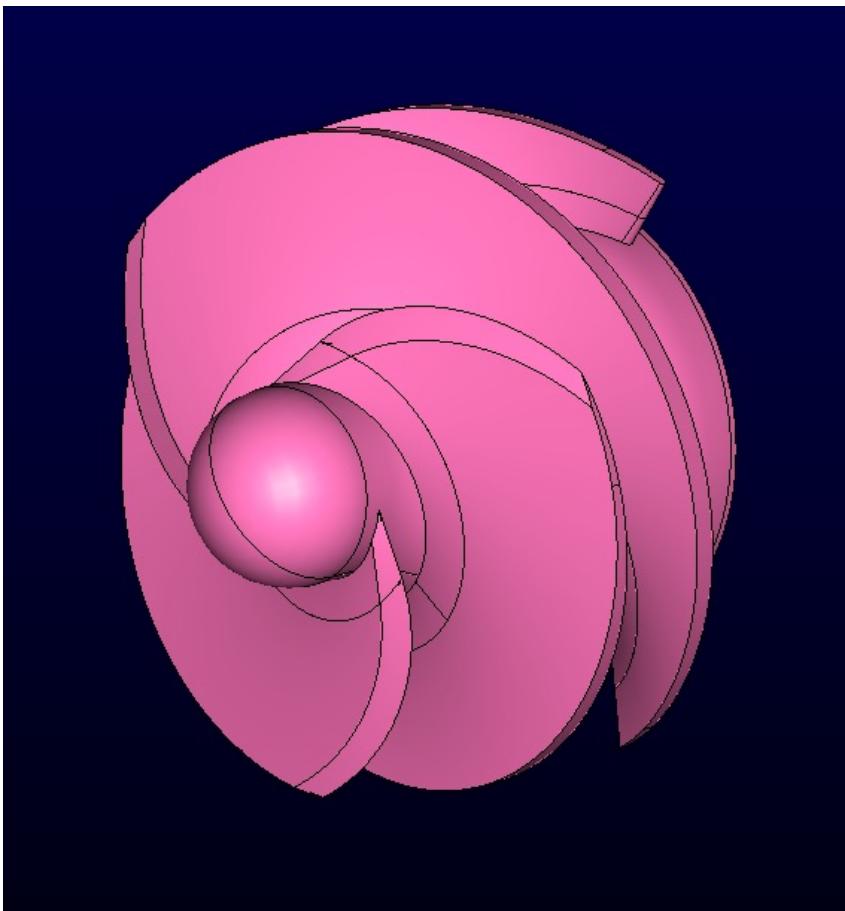
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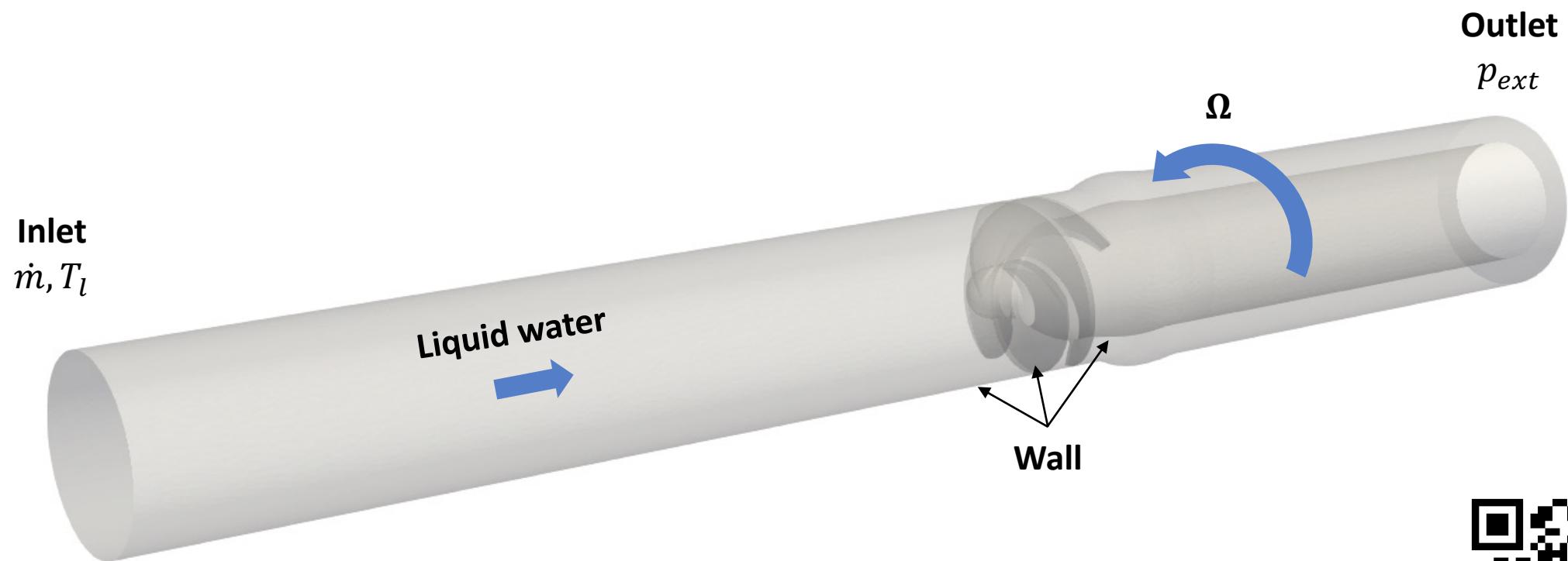
Results: test case setup



- 3 bladed-inducer
- Variable section hub
- Experimental data from



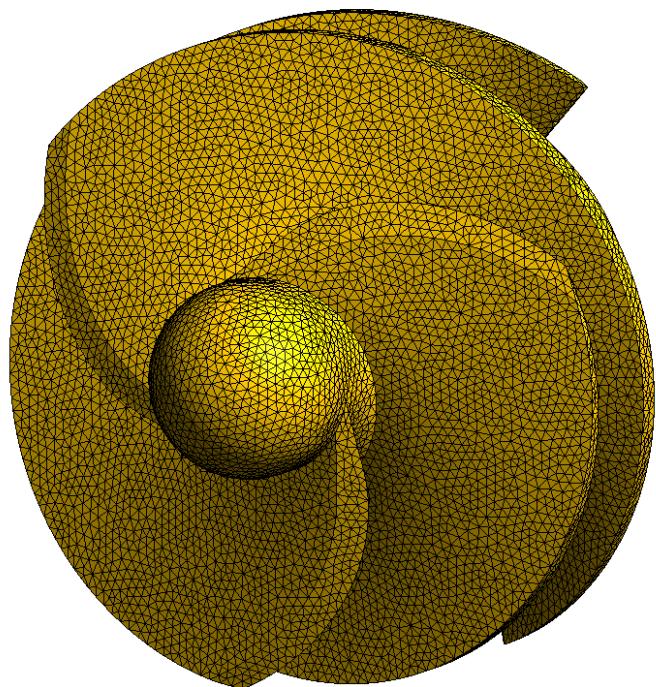
Results: test case setup



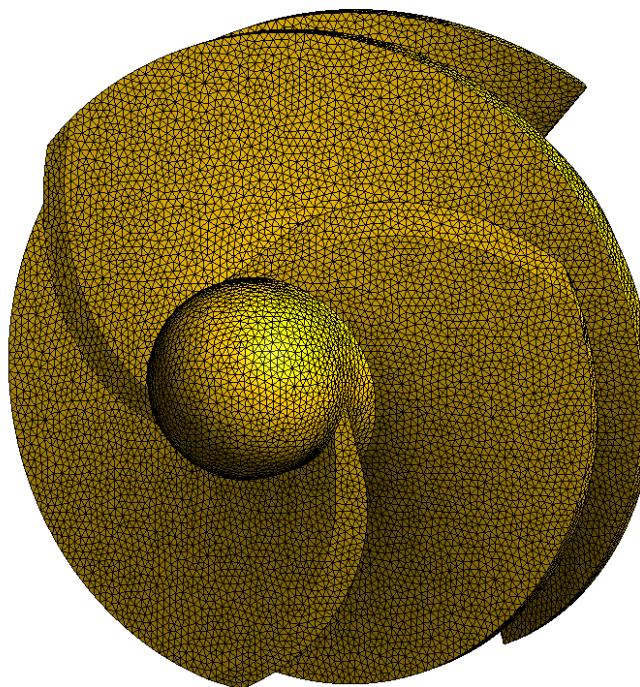
Numerical
implementation in
the open-source
CFD solver **ECOGEN**



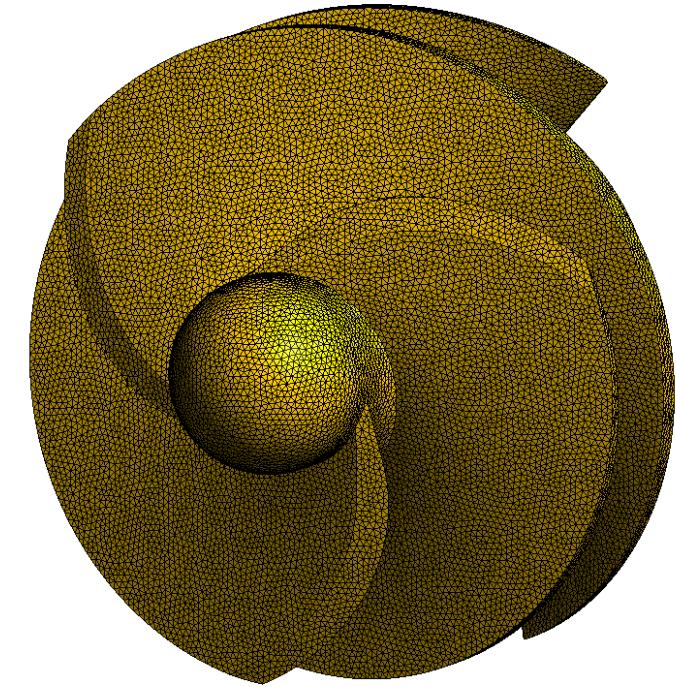
Results: mesh convergence



919k cells



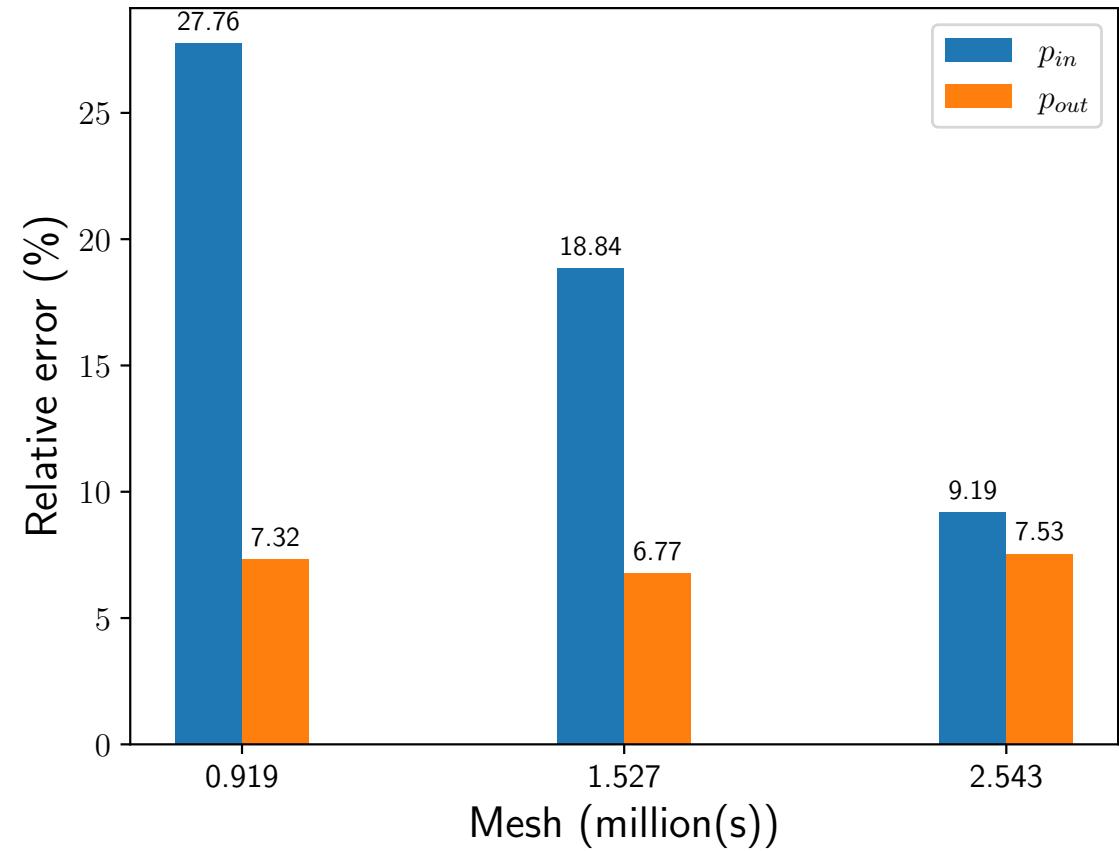
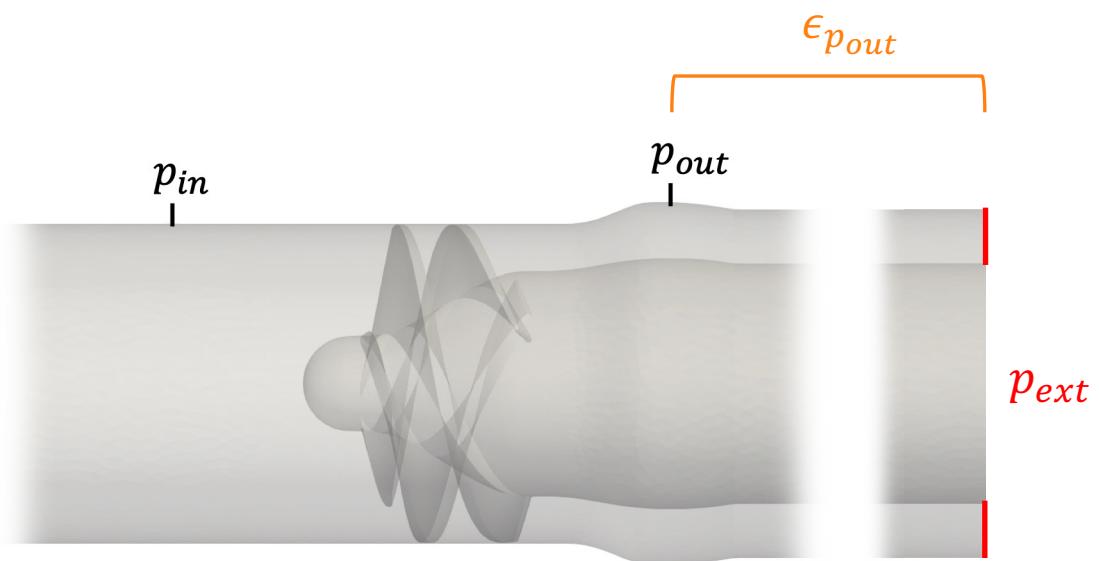
1M527k cells



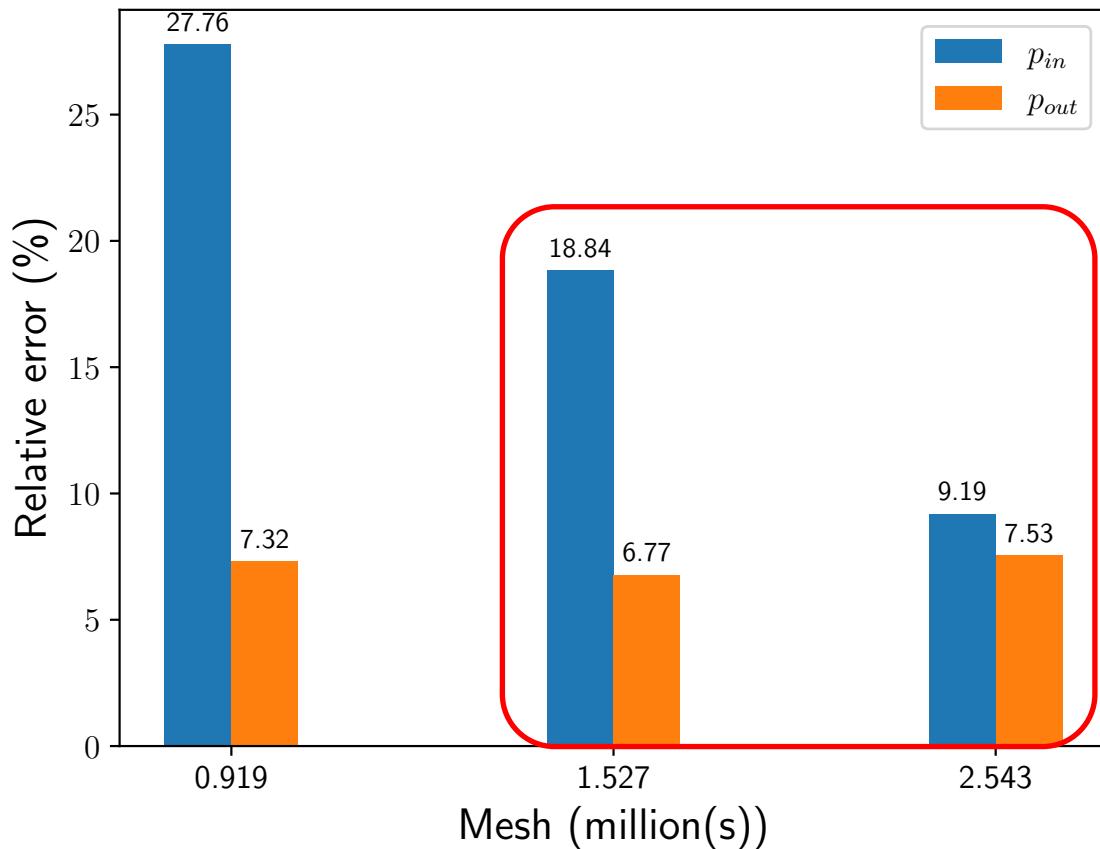
2M543k cells

Results: mesh convergence

- Nominal mass flow rate $\dot{m} = \dot{m}_n$
- Pressure measured at the wall averaged in time/space



Results: mesh convergence



Mesh (millions)	1.527	2.543
Number of CPUs	160	256
Computation time (h)	152	157
CPU time (h)	24440	40192



Trade-off between computational resources / error

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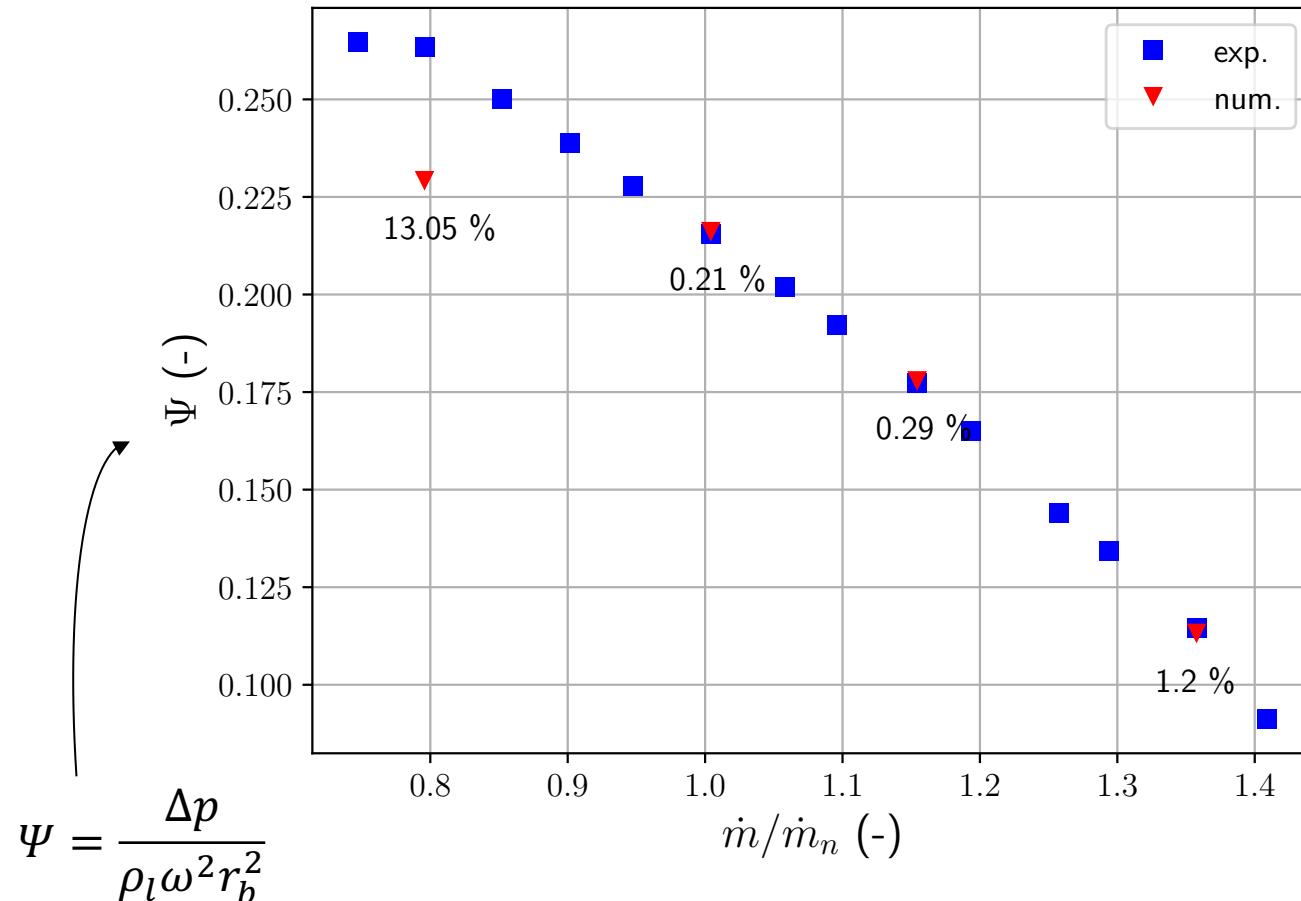
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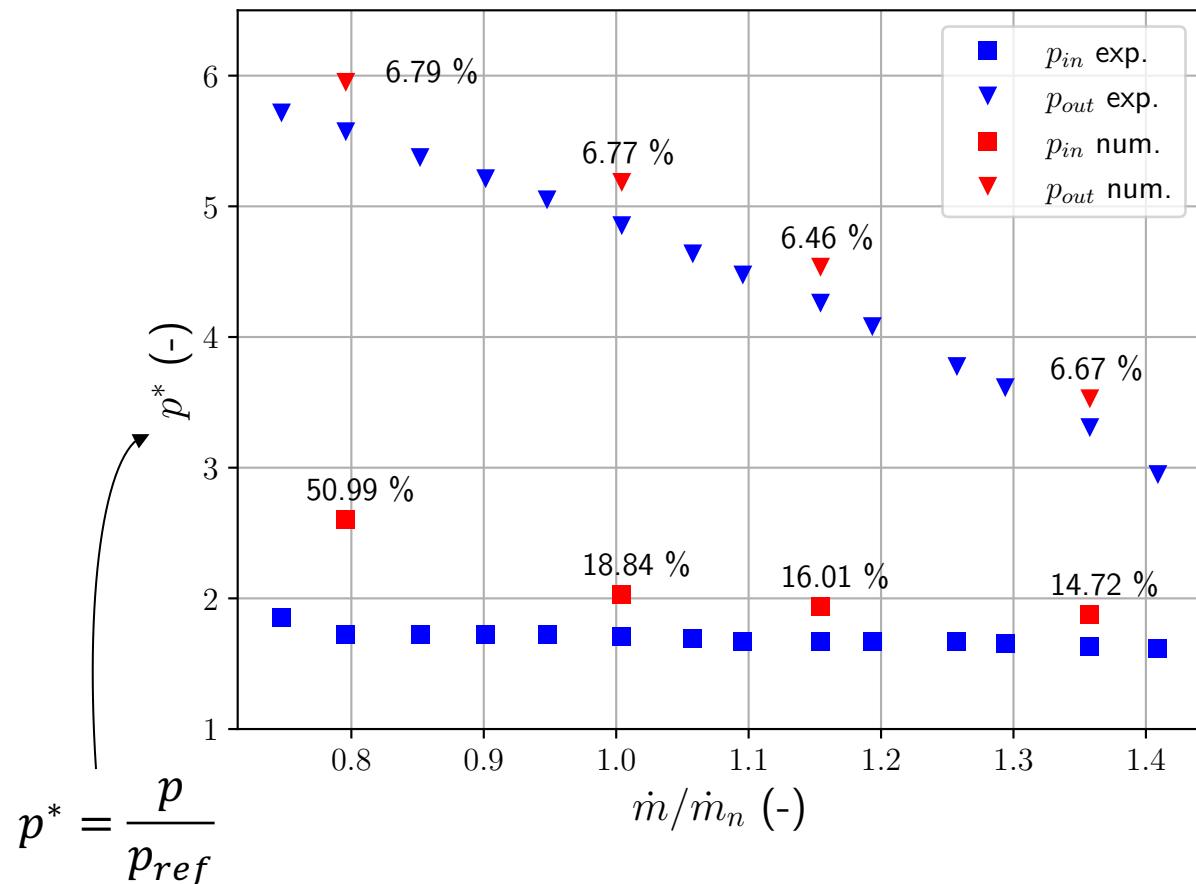
Results: single-phase flow behavior



Flow rate variation \dot{m}/\dot{m}_n

Good match except at low mass flow rate

Results: single-phase flow behavior



Flow rate variation \dot{m}/\dot{m}_n

→ Pressures are overestimated but the inlet/outlet gap is relatively constant except at low massflow rate

Results: single-phase flow behavior

Overpressure error

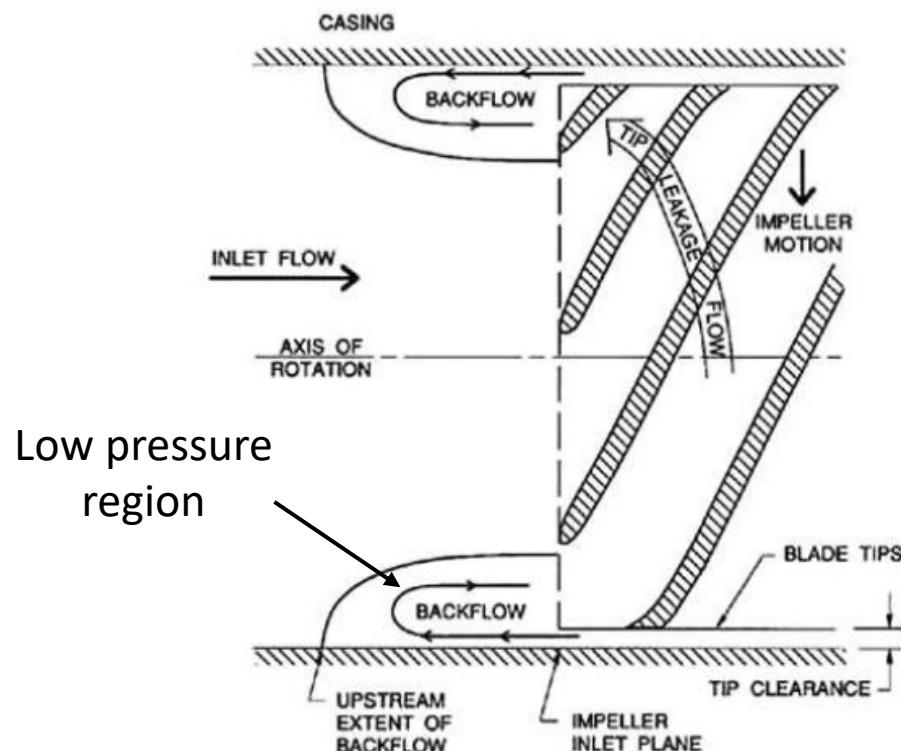
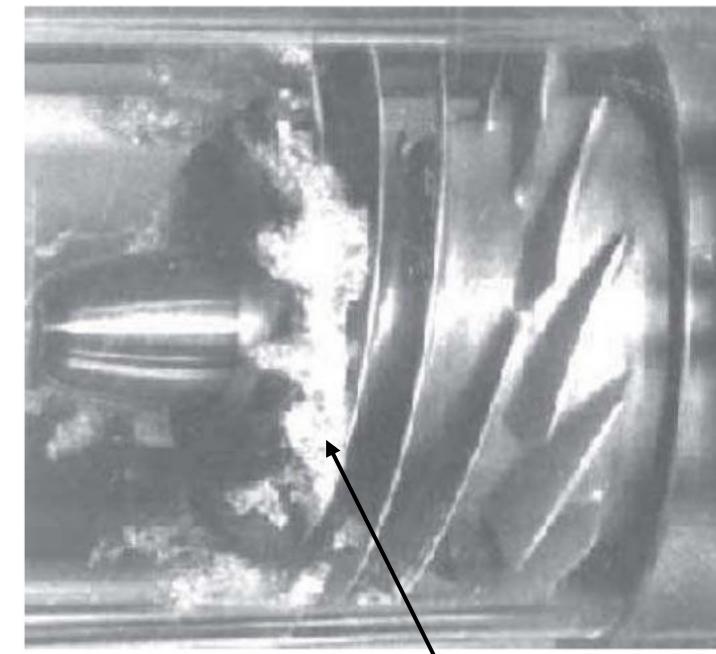


Fig – Illustration of backflow cavitation [Brennen, 2011]

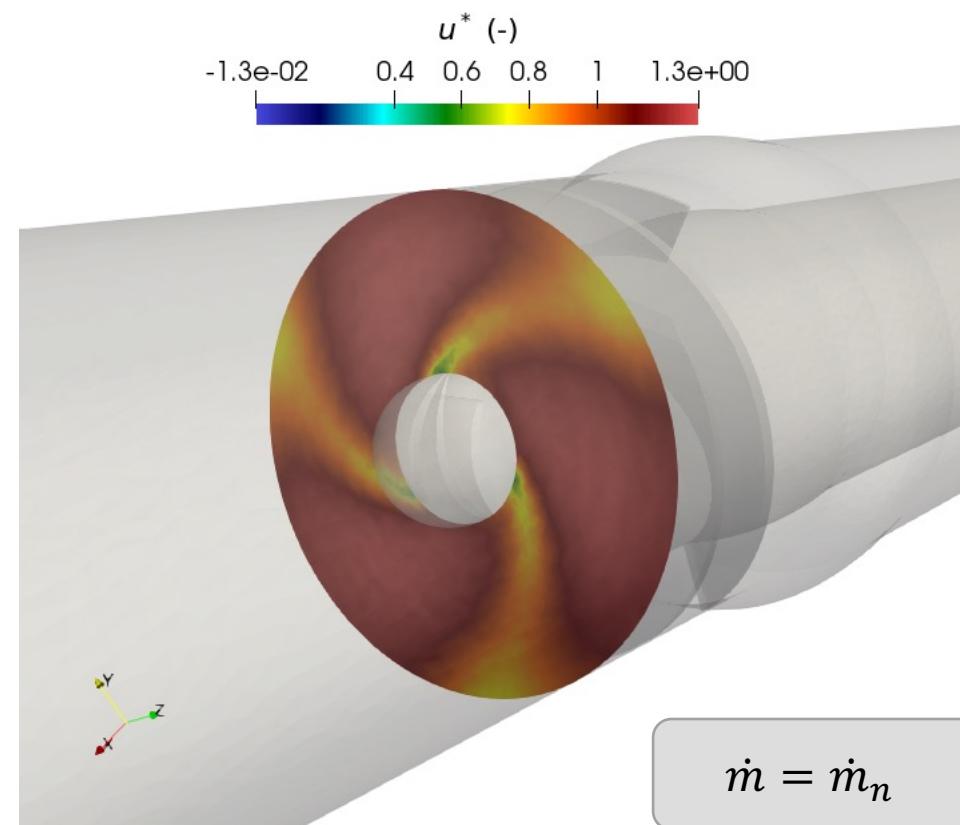
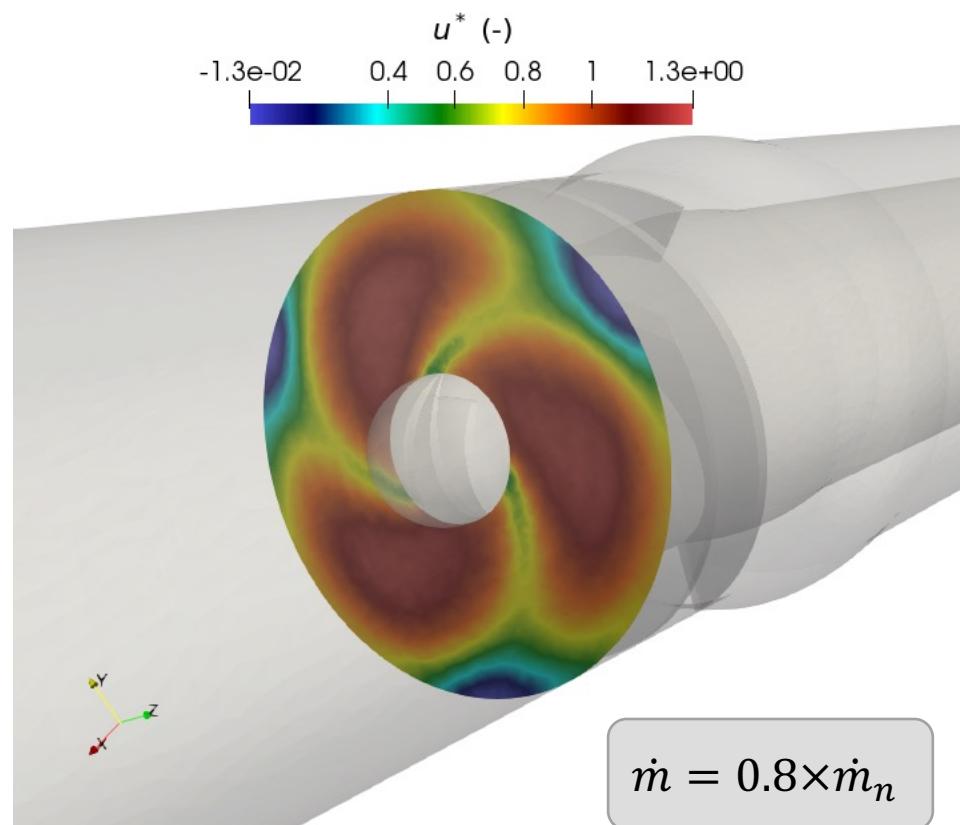
Fig – Backflow cavitation on SSME LOx inducer [Braisted, 1980]



Cavitation at
blade tip

Results: single-phase flow behavior

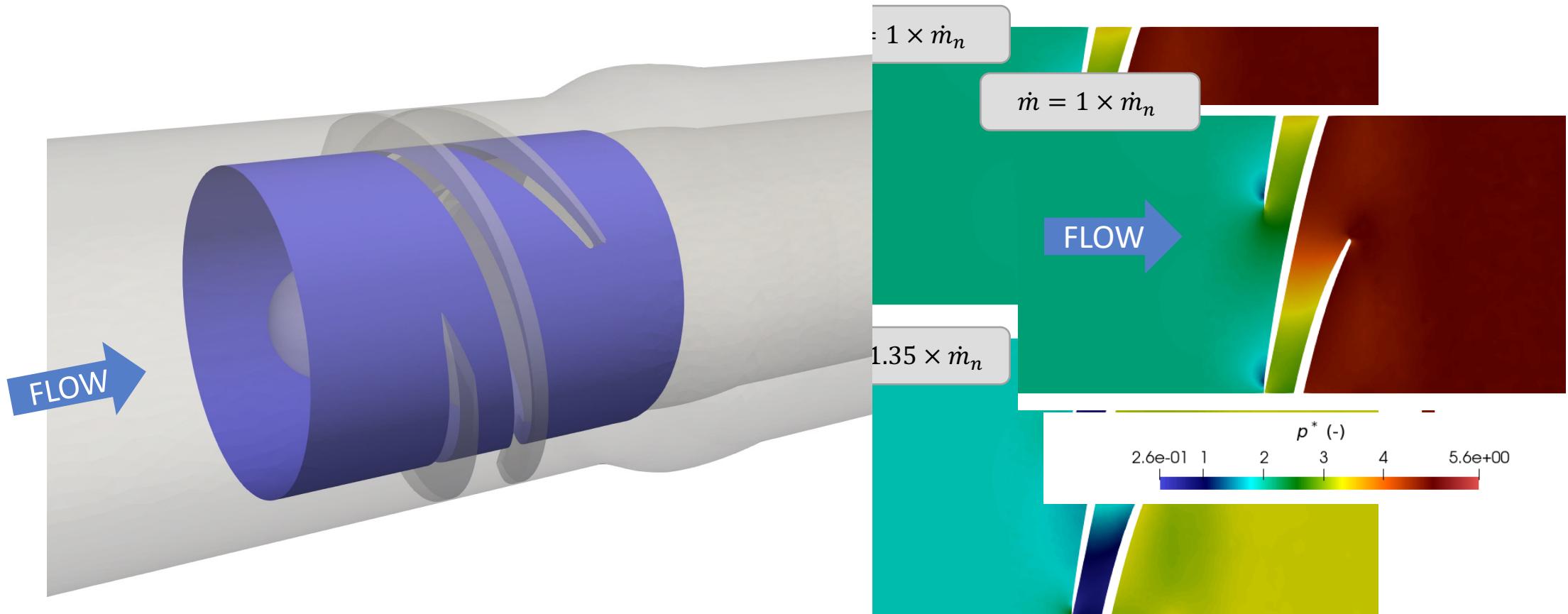
Overpressure error



$$u^* = \frac{u_z}{u_n}$$

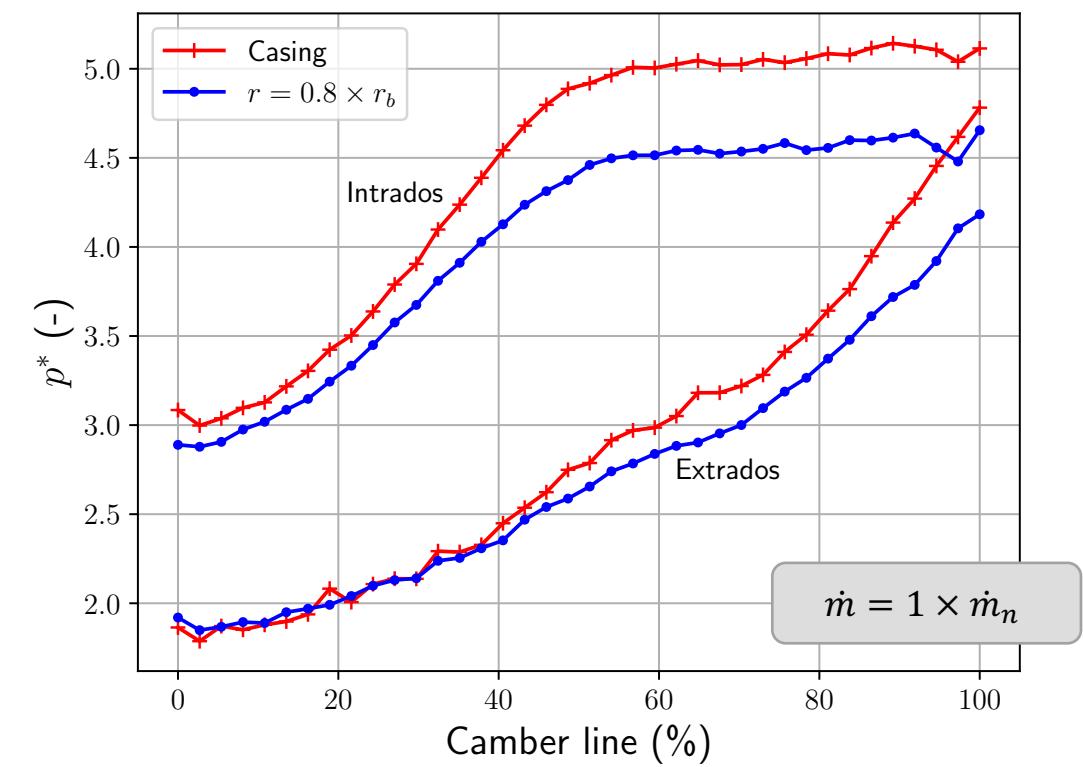
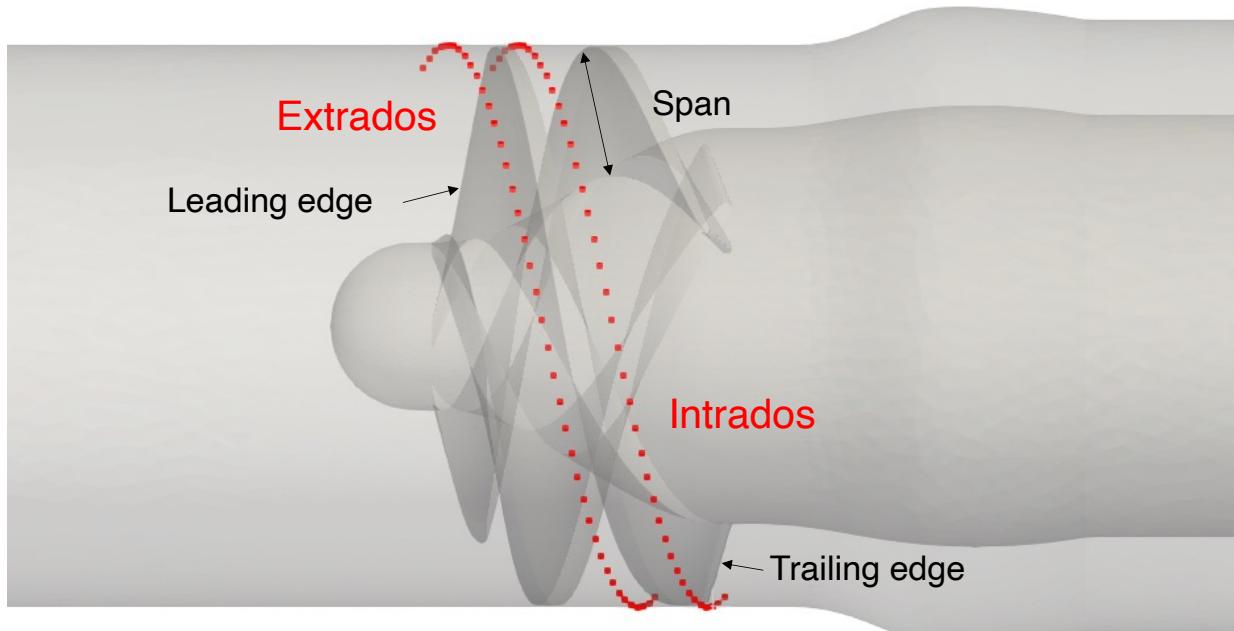
Results: single-phase flow behavior

Study of pressure field



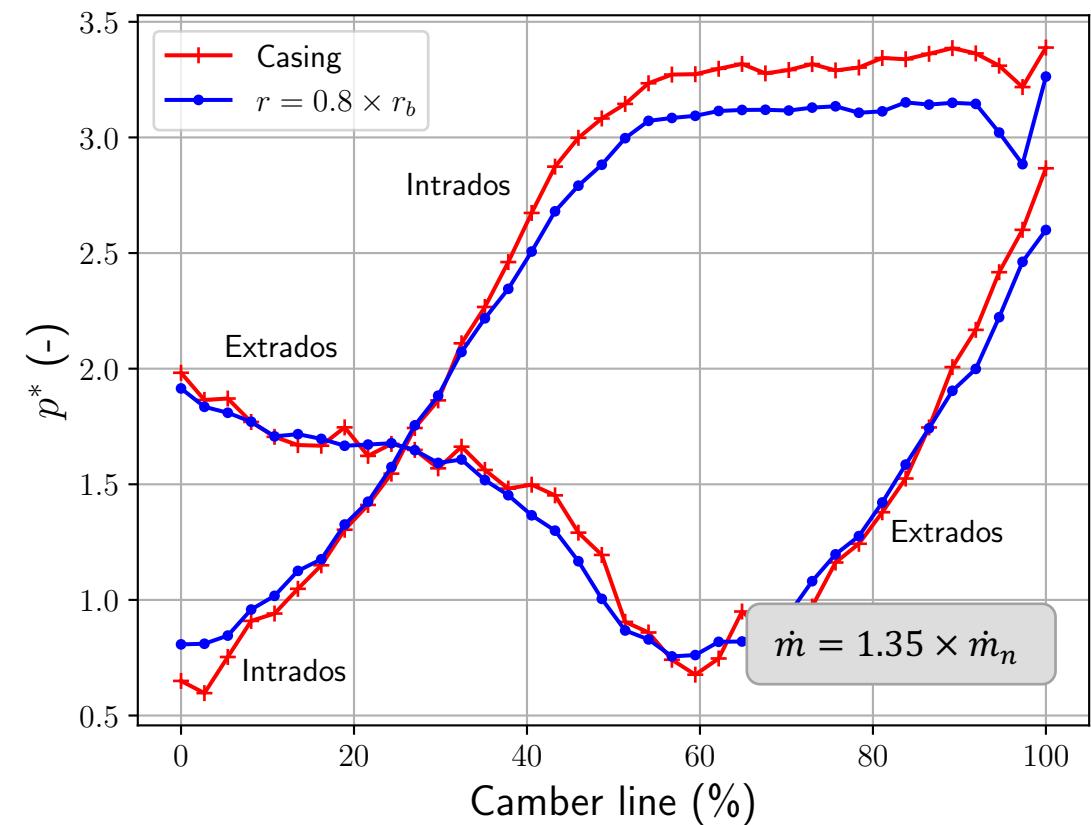
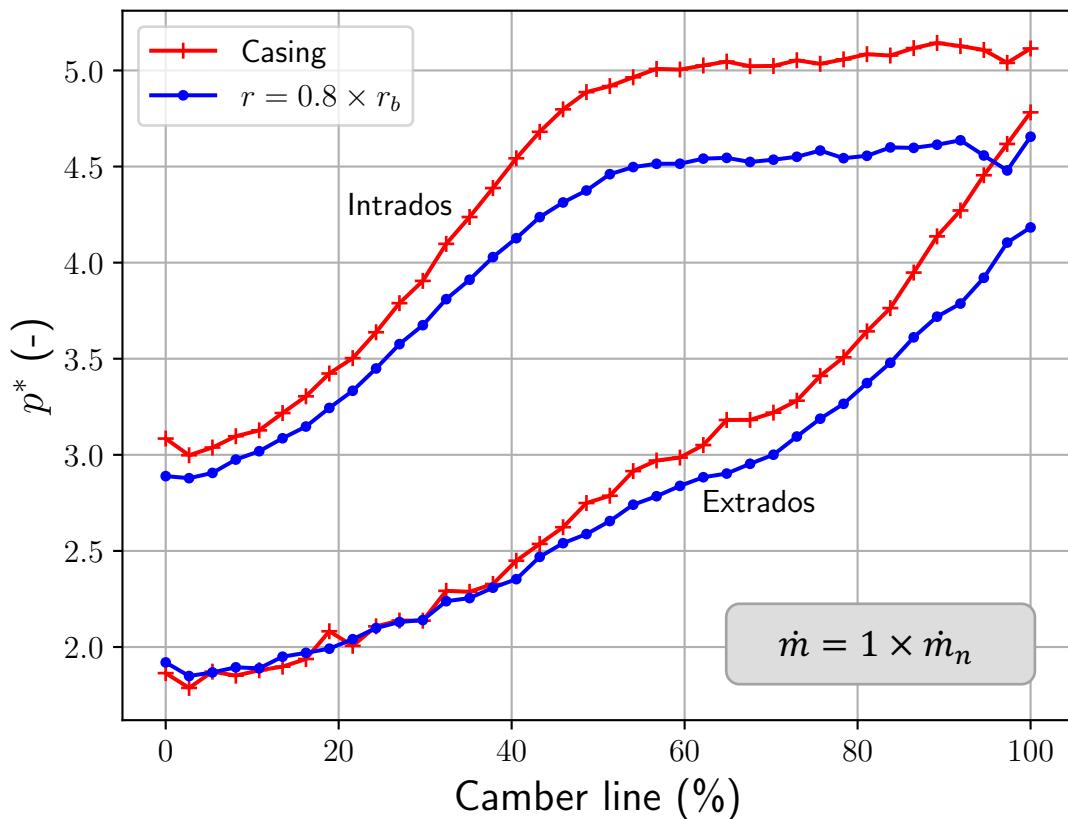
Results: single-phase flow behavior

Blades load



Results: single-phase flow behavior

Blades load



Summary

Introduction

Modeling of the cavitation phenomenon

- State-of-the-art
- Two-phase flow approach
- Two-phase flow model
- Blades motion

Numerical method

- Numerical scheme
- MRF fluxes
- Mesh mapping

Results

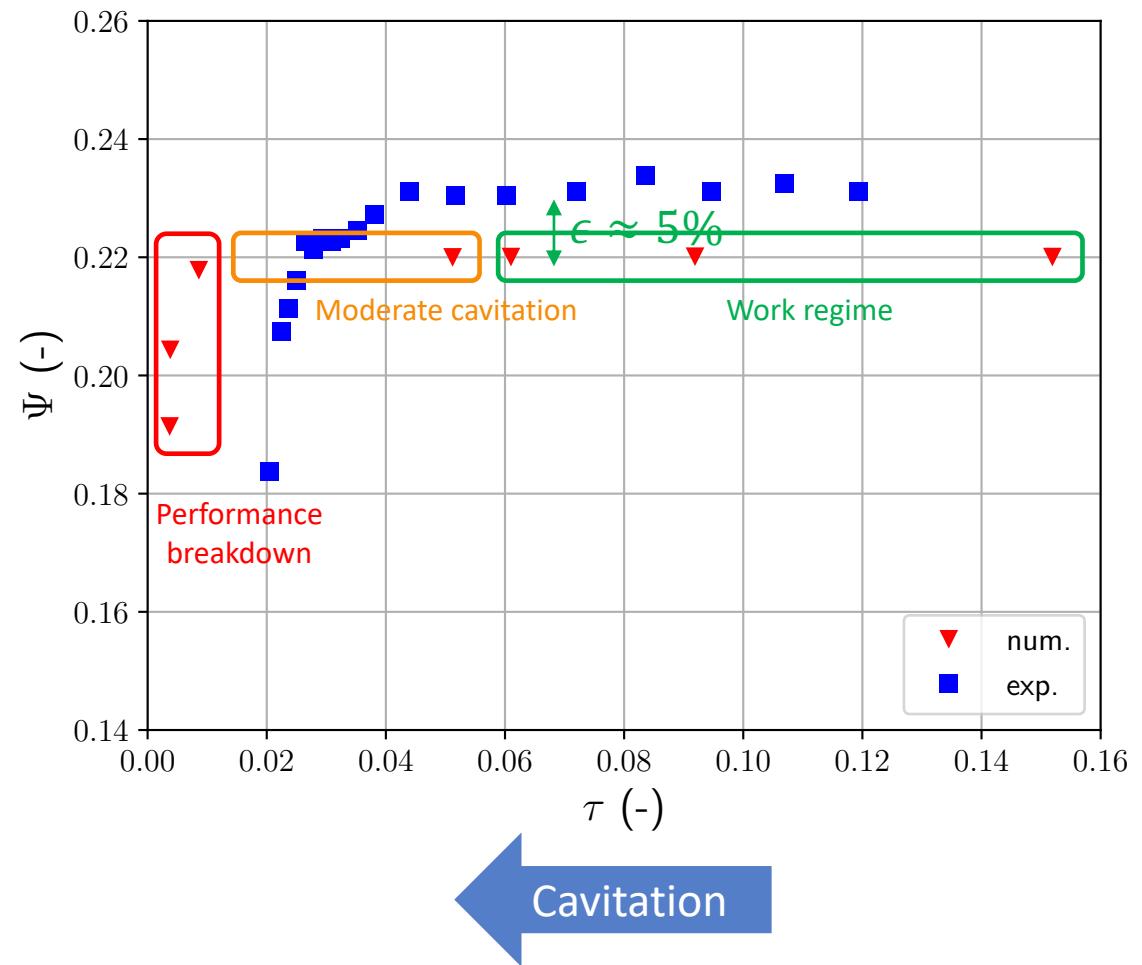
- Test case setup
- Single-phase flow behavior: pump characteristic
- Two-phase flow: cavitating regime

Conclusions & perspectives

Results: two-phase flow in cavitating regime

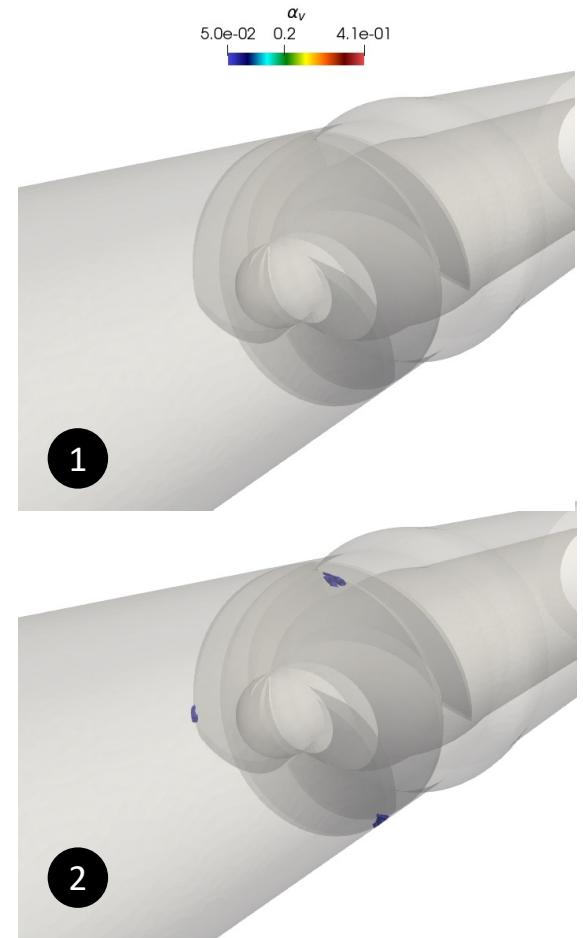
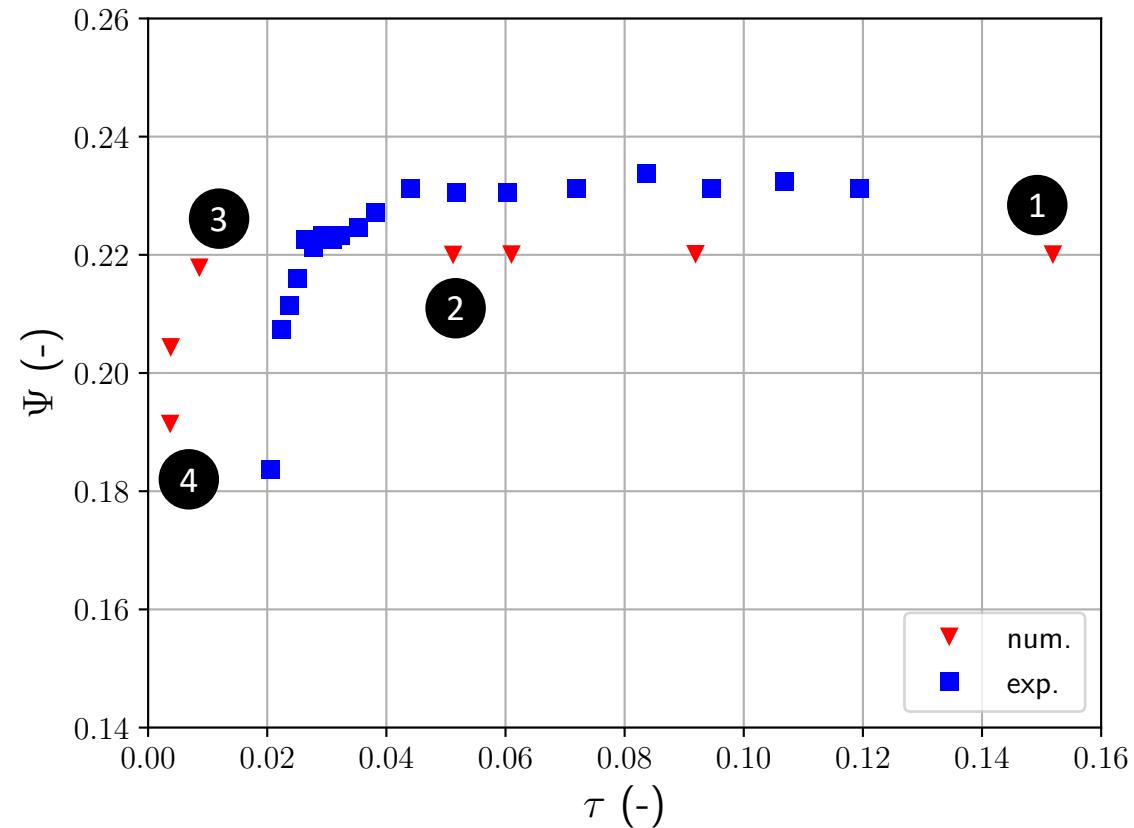
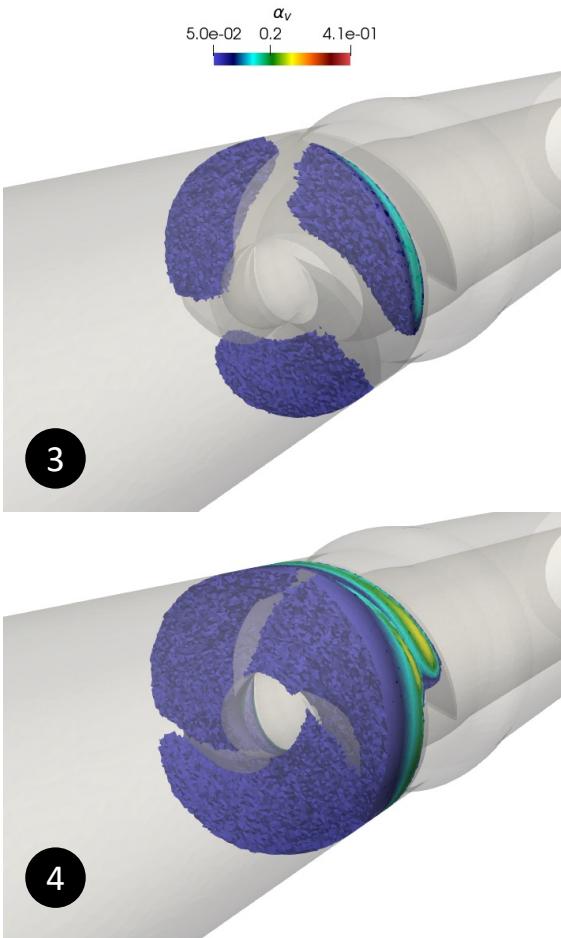
Performance curve

- Variation of the outlet pressure at \dot{m}_n
- $\Psi = \frac{\Delta p}{\rho_l \omega^2 r_b^2}$
- $\tau = \frac{p_e^{tot} - p_v}{\rho_l \omega^2 r_b^2}$



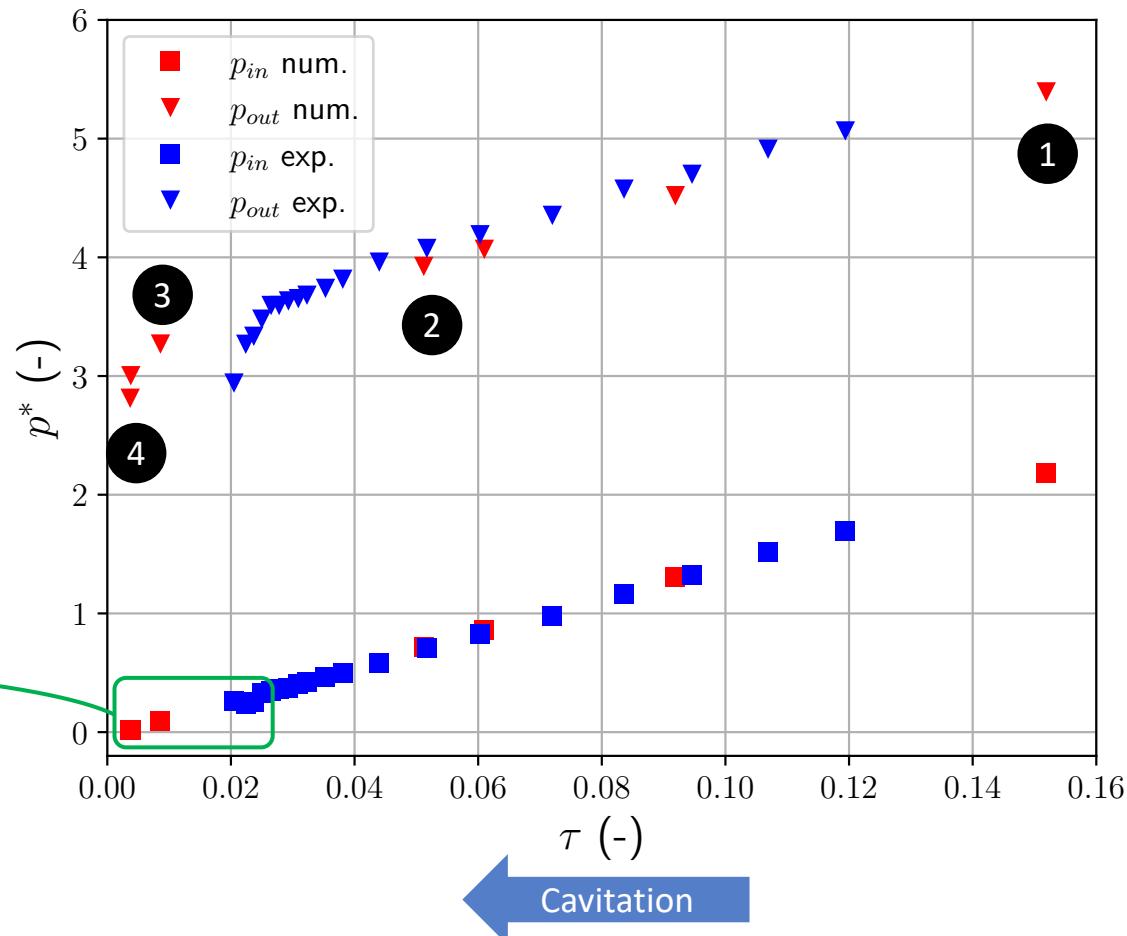
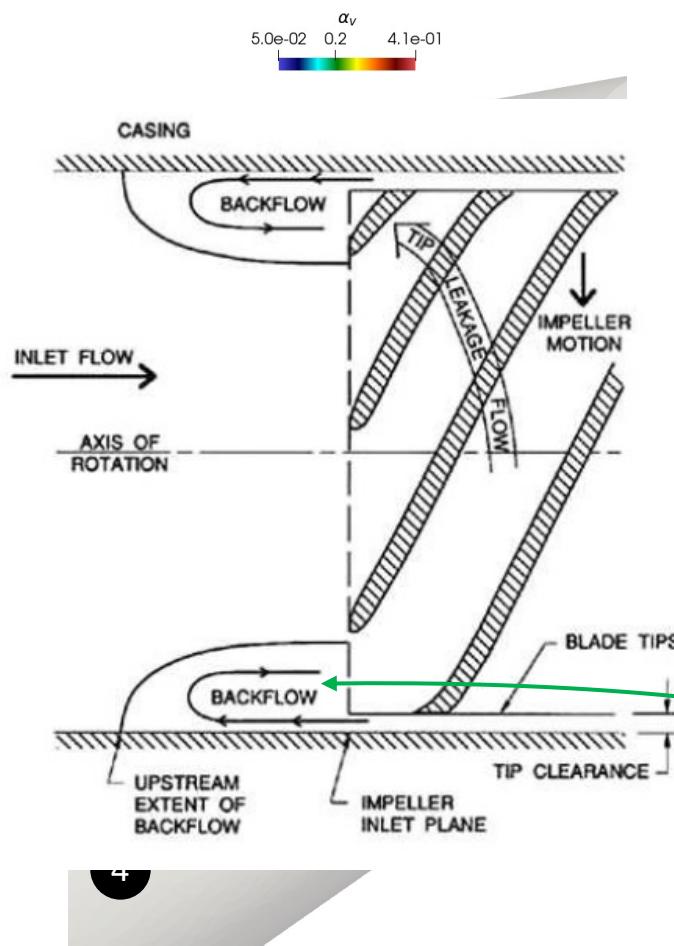
Results: two-phase flow in cavitating regime

Cavitation pockets



Results: two-phase flow in cavitating regime

Cavitation number estimates

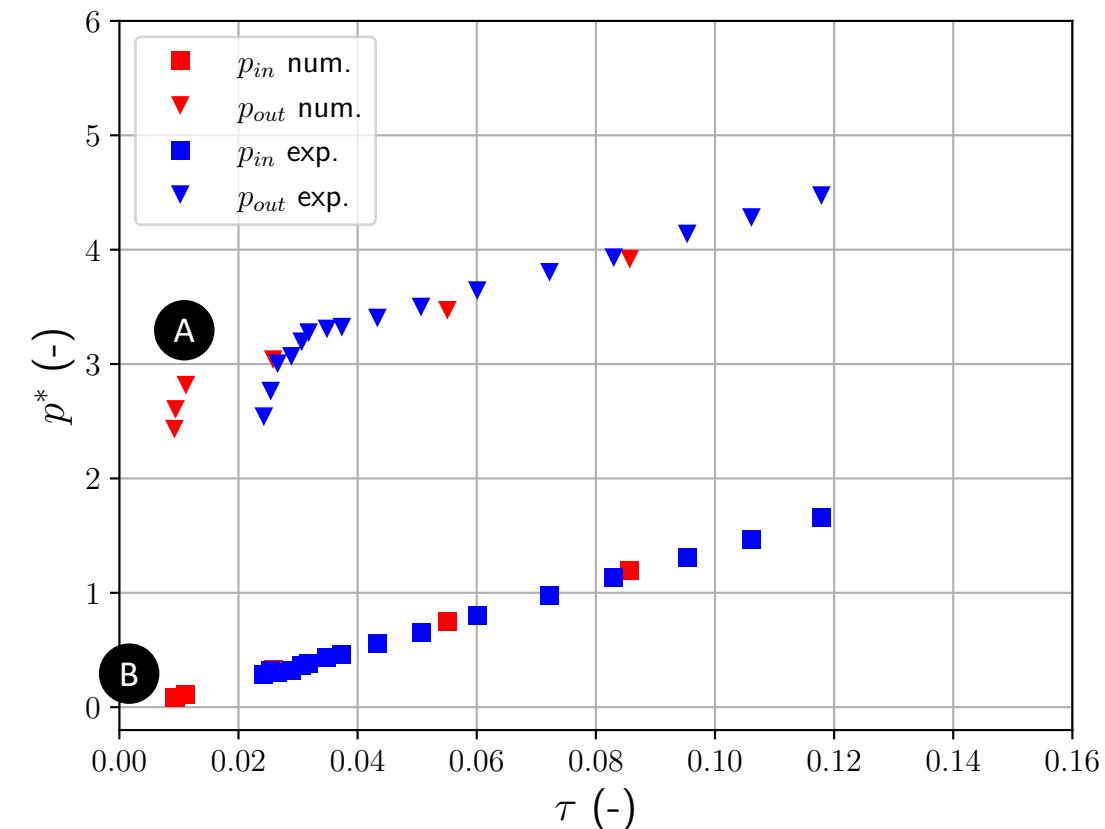
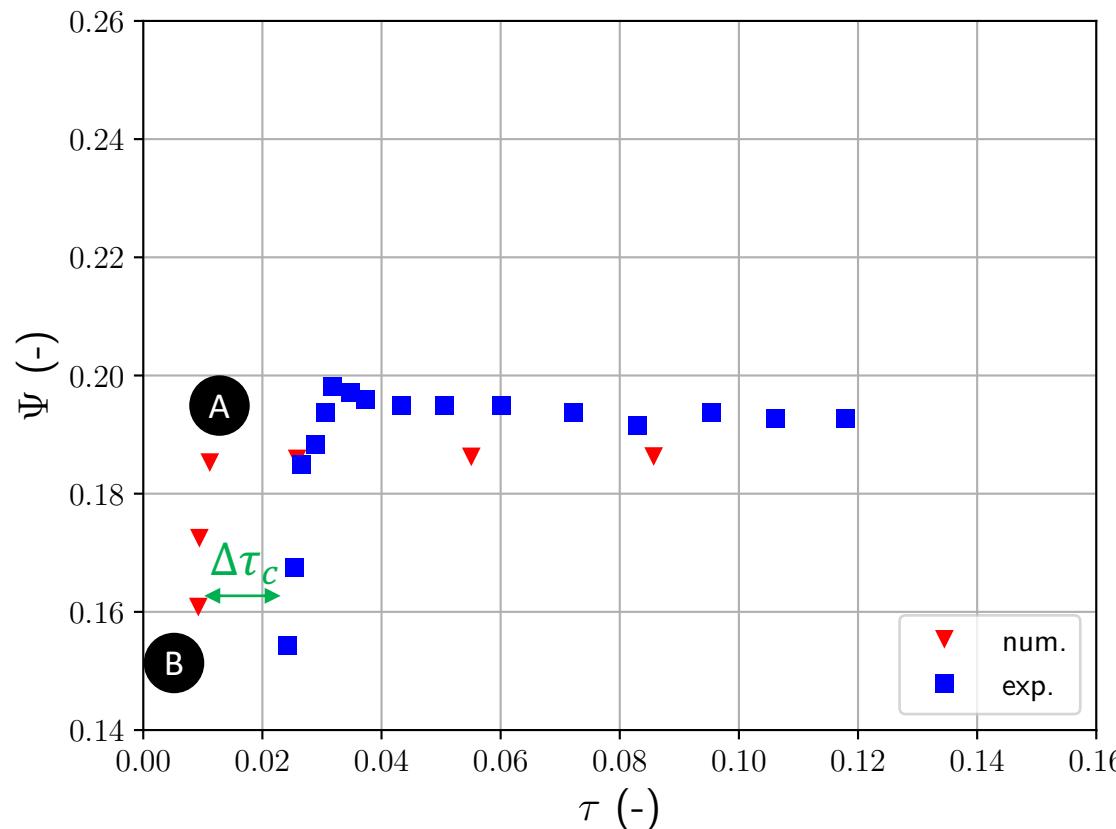


$$\tau = \frac{p_e^{tot} - p_v}{\rho_l \omega^2 r_b^2}$$

$$p^* = \frac{p}{p_{ref}}$$

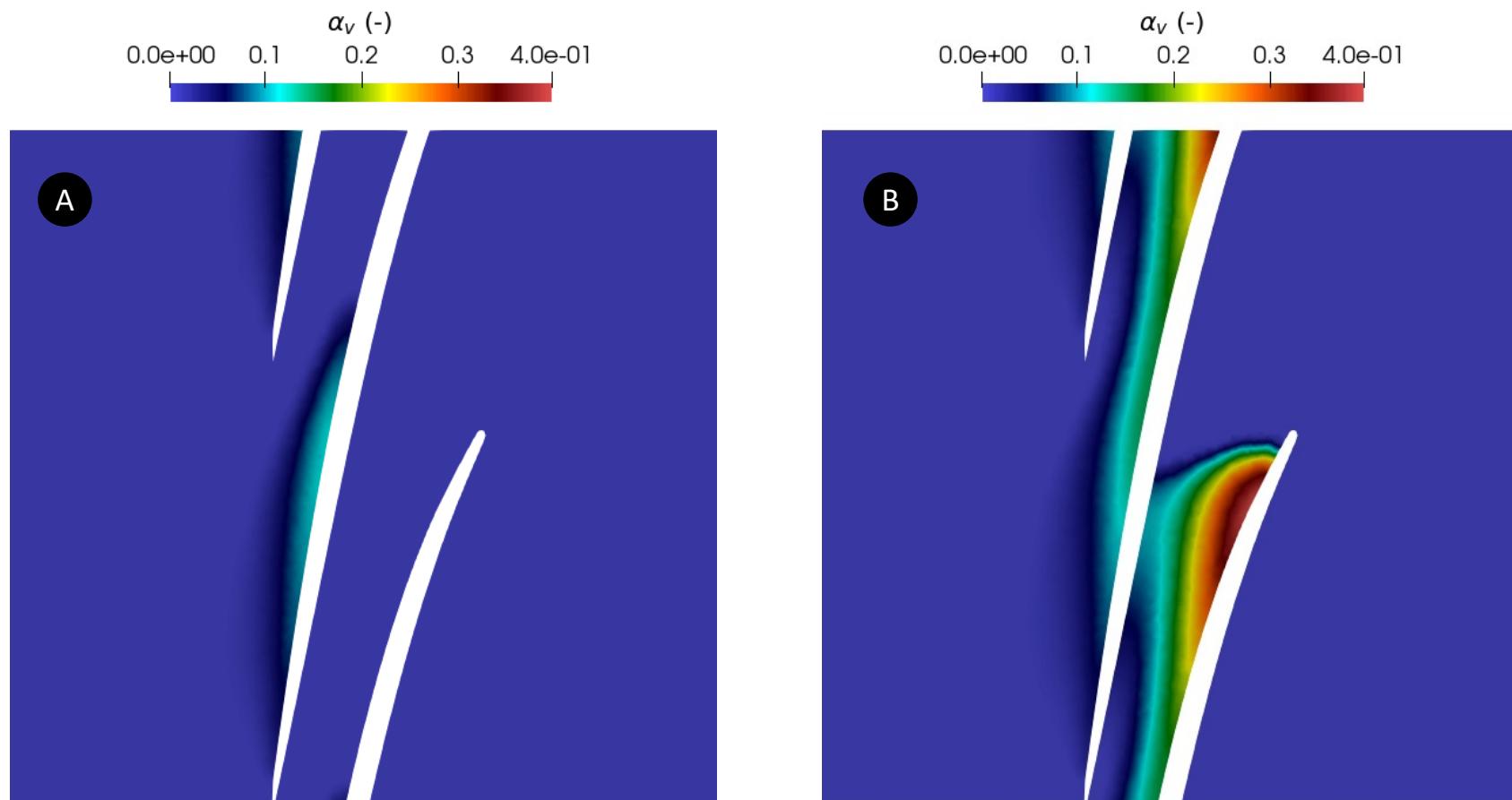
Results: two-phase flow in cavitating regime

Increase massflow rate $\dot{m} = 1.15 \times \dot{m}_n$



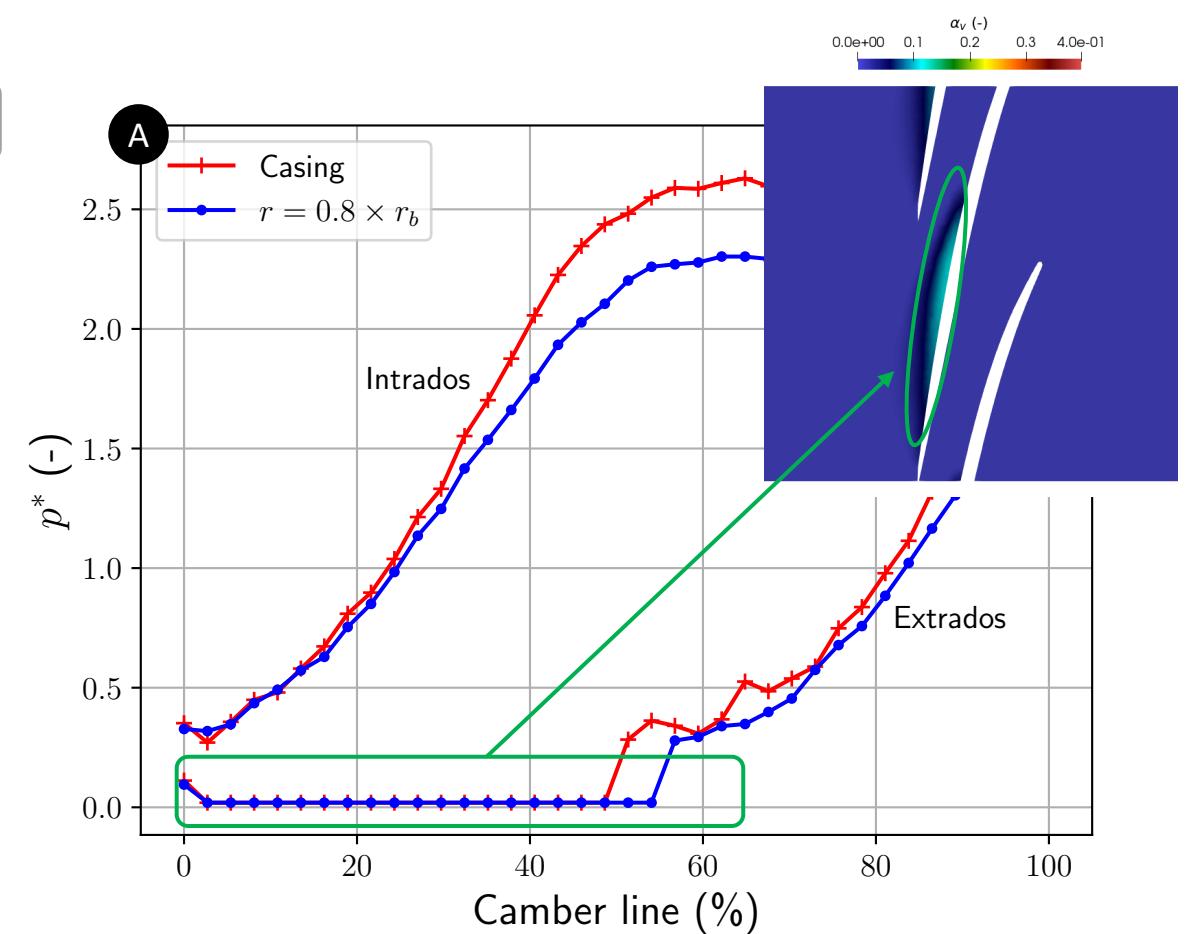
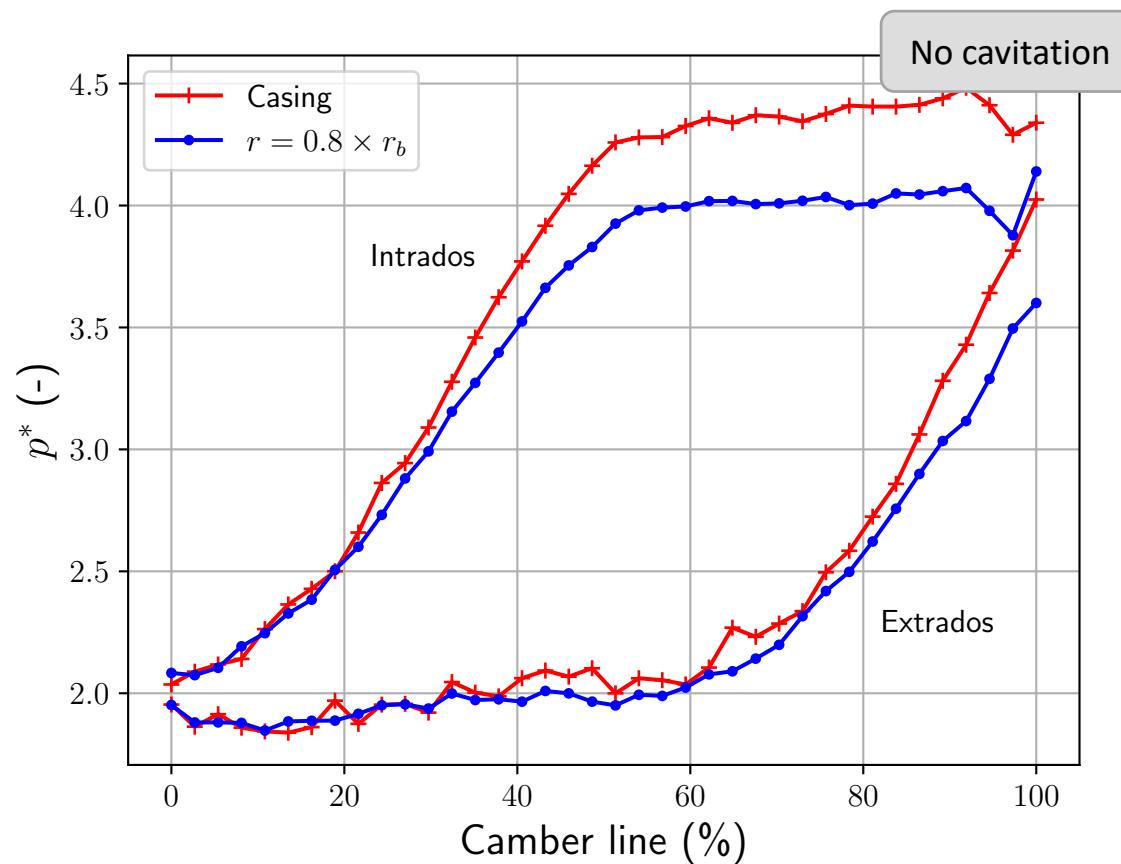
Results: two-phase flow in cavitating regime

Increase massflow rate $\dot{m} = 1.15 \times \dot{m}_n$



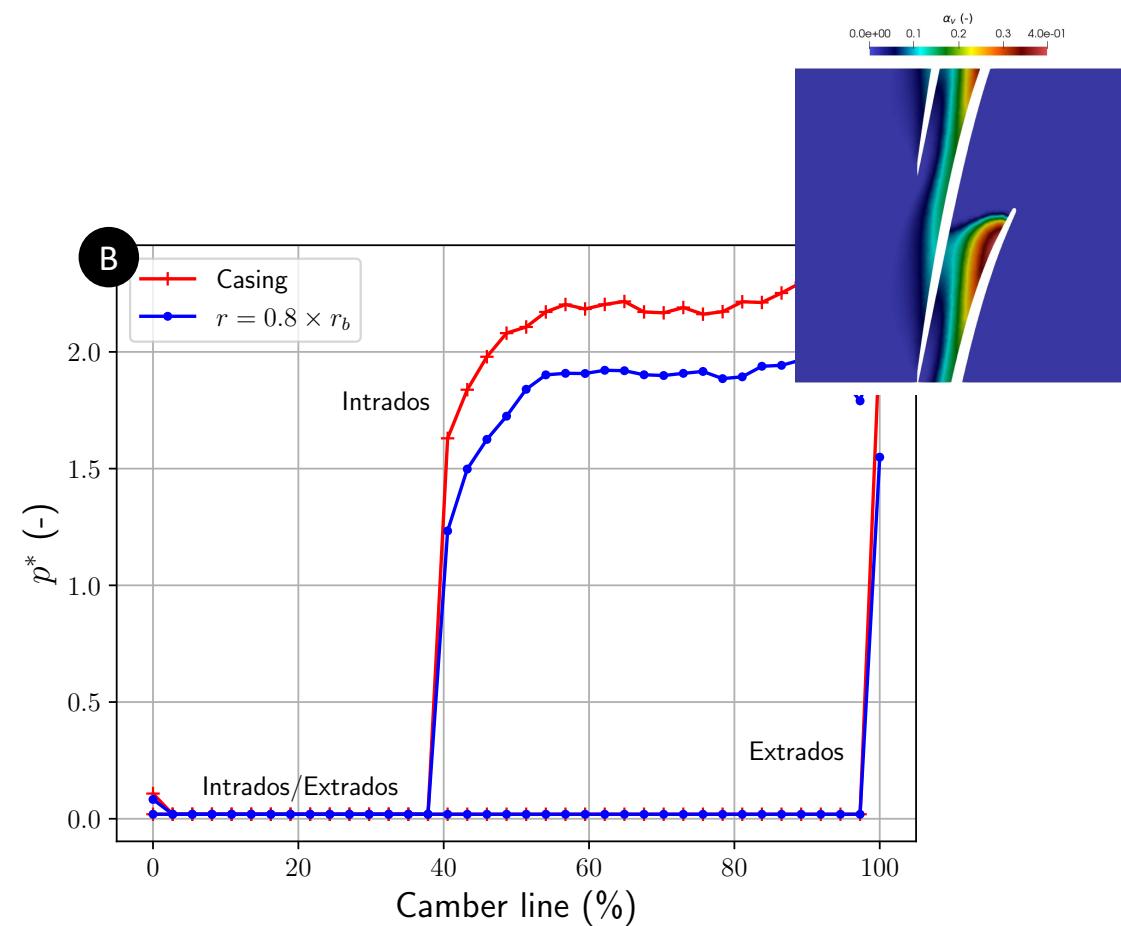
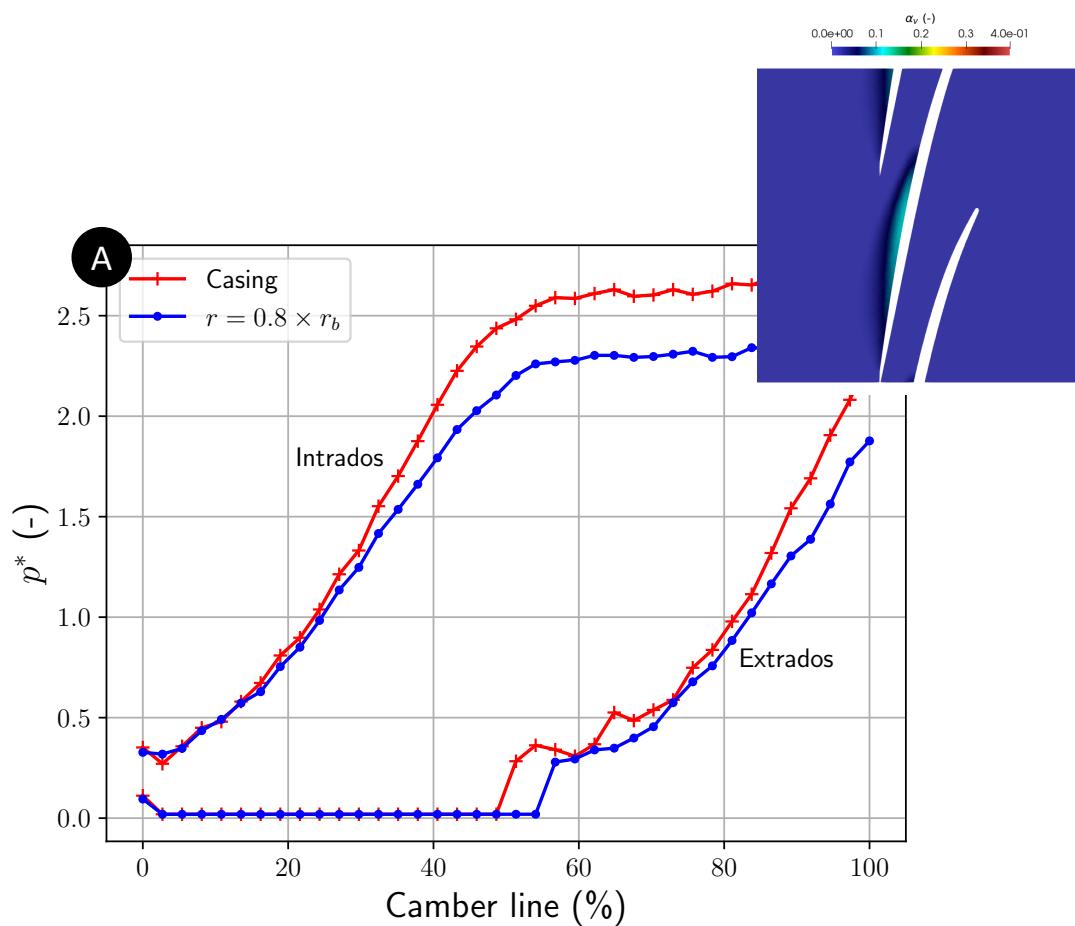
Results: two-phase flow in cavitating regime

Blades load with/without cavitation



Results: two-phase flow in cavitating regime

Blades load on two cavitation regimes



Summary

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Results

- Test case setup
- Single-phase flow behavior: pump characteristic
- Two-phase flow: cavitating regime

Conclusions & perspectives

Conclusions & perspectives

Conclusions

- Two-phase flow model written in a rotating frame
- Phase change based on thermodynamics
- Good estimation of an inducer behavior in non-cavitating regime
- Assessment of the performance breakdown in cavitating regime



Two-phase flow model able to capture cavitation pockets in turbopumps

Perspectives

- Improve performance breakdown
 - Gap between the casing and the blades?
 - Flow temperature?
 - 2nd order numerical scheme?
 - Low-Mach preconditionning?
 - Turbulence?
- Study of the thermodynamic effect responsible of a delay on the cavitation phenomenon

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Thank you for your attention

Thermodynamic effect

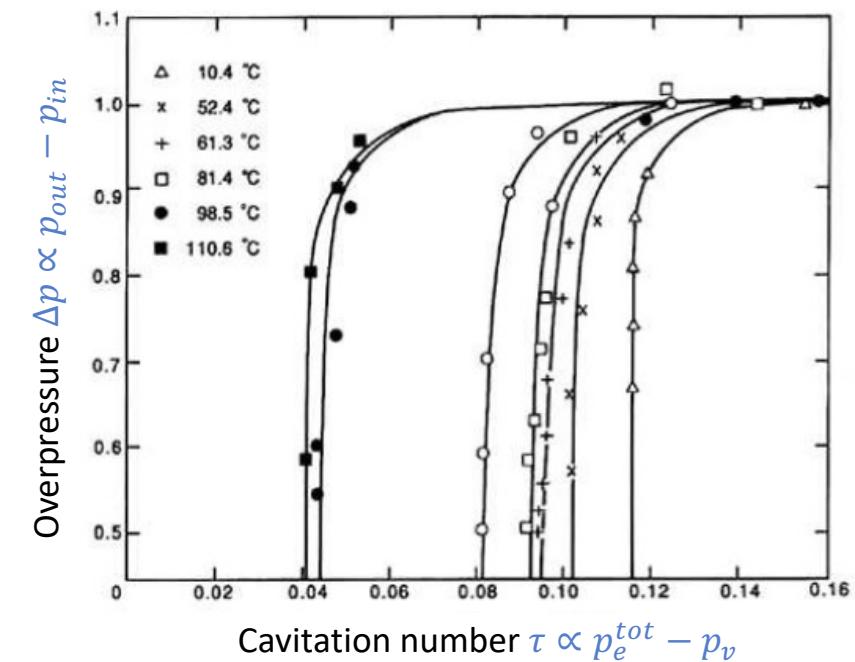
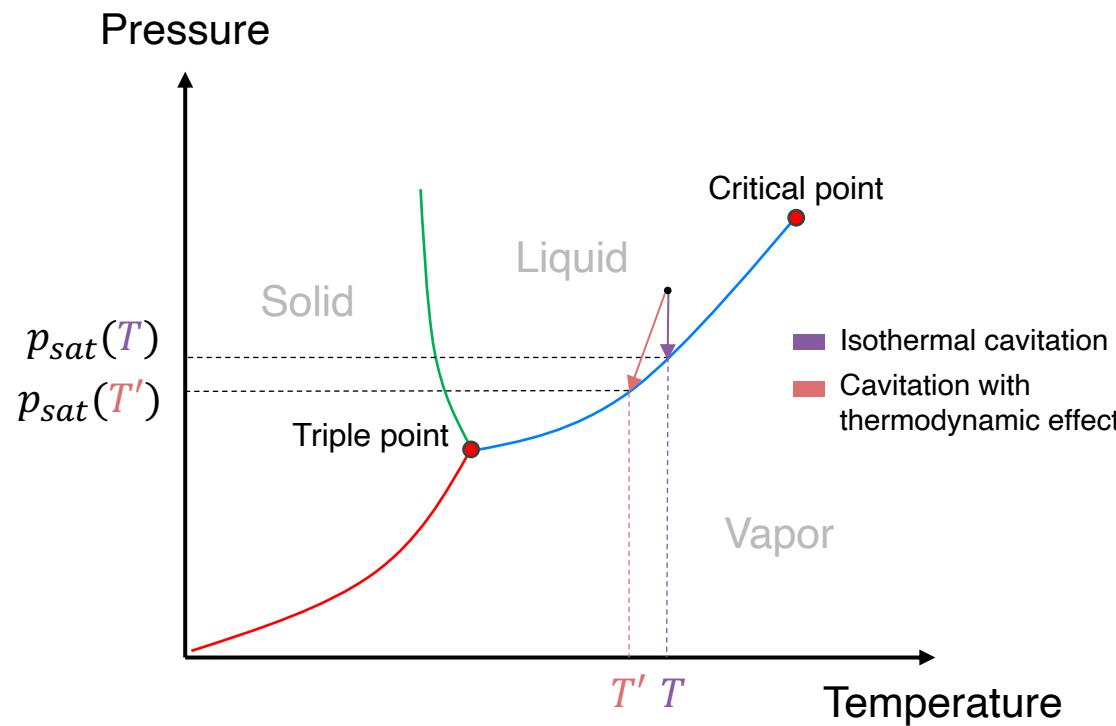


Fig – Performance curve of a water centrifugal pump [Chivers, 1969]

Two-phase flow model based on total energy equations

Internal energy formulation

$$\partial_t \alpha_1 + \nabla \cdot (\alpha_1 \mathbf{u}) - \alpha_1 \nabla \cdot \mathbf{u} = \mu(p_1 - p_2)$$

$$\partial_t(\alpha_k \rho_k) + \nabla \cdot (\alpha_k \rho_k \mathbf{u}) = 0$$

$$\partial_t(\rho \mathbf{u}) + \nabla \cdot (\rho \mathbf{u} \otimes \mathbf{u}) + \nabla(\alpha_1 p_1 + \alpha_2 p_2) = \mathbf{0}$$

$$\begin{aligned}\partial_t(\alpha_1 \rho_1 e_1) + \nabla \cdot (\alpha_1 \rho_1 e_1 \mathbf{u}) + \alpha_1 p_1 \nabla \cdot \mathbf{u} &= -\mu p_I(p_1 - p_2) \\ \partial_t(\alpha_2 \rho_2 e_2) + \nabla \cdot (\alpha_2 \rho_2 e_2 \mathbf{u}) + \alpha_2 p_2 \nabla \cdot \mathbf{u} &= +\mu p_I(p_1 - p_2)\end{aligned}$$

Total energy formulation



$$\begin{aligned}\partial_t(\alpha_1 \rho_1 E_1) + \nabla \cdot (\alpha_1 \rho_1 E_1 \mathbf{u} + \alpha_1 p_1 \mathbf{u}) + \Sigma(\mathbf{U}, \nabla \mathbf{U}) &= -\mu p_I(p_1 - p_2) \\ \partial_t(\alpha_2 \rho_2 E_2) + \nabla \cdot (\alpha_2 \rho_2 E_2 \mathbf{u} + \alpha_2 p_2 \mathbf{u}) - \Sigma(\mathbf{U}, \nabla \mathbf{U}) &= +\mu p_I(p_1 - p_2)\end{aligned}$$

Total mixture energy conservation

[Pelanti and Shyue, 2014]

Moving Reference Frame method

