

Kinematic viscosity estimates in reversed-field pinch fusion plasmas

N. Vivenzi^{1,2}, G. Spizzo^{1,3}, M. Veranda¹, D. Bonfiglio^{1,3}, S. Cappello^{1,3} and RFX-mod team.

¹ Consorzio RFX (CNR, ENEA, INFN, Università degli Studi di Padova, Acciaierie Venete SpA), C.so Stati Uniti 4, 35127 Padova, Italy
e-mail: nicholas.vivenzi@studenti.unipd.it

² CRF – University of Padova, Italy

³ CNR - ISTP Padova, Italy



CONSORZIO RFX
Ricerca Formazione Innovazione

ISTP
Consiglio Nazionale delle Ricerche

THEORY OF FUSION PLASMAS
JOINT VARENNA - LAUSANNE INTERNATIONAL WORKSHOP (since 1988)
Villa Monastero, Varenna, Italy

Aim of this work

Framework: Visco-resistive magneto-hydrodynamics (MHD) of reversed-field pinch (RFP) fusion plasmas.

Aim of this poster:

- to study the effects of a non-uniform scalar viscosity in visco-resistive MHD simulations
- to estimate the kinematic viscosity on the basis of RFX-mod experimental data.

Visco-resistive MHD model

SpeCyl [2] solves the visco-resistive MHD model. Hypothesis:

- constant mass density
- negligible pressure ($\beta \ll 1$)
- cylindrical geometry

Input parameters: dimensionless resistivity η and viscosity ν :

SpeCyl equations

$$\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} = \mathbf{j} \times \mathbf{B} + \nu \nabla^2 \mathbf{v}$$

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) - \nabla \times (\eta \mathbf{j})$$

$$\nabla \cdot \mathbf{B} = 0 \quad \nabla \times \mathbf{B} = \mathbf{j}$$

By means of a proper change of coordinates [3] η and ν can be replaced by the magnetic Prandtl number $P = \nu/\eta$ and the Hartmann number $H = 1/\sqrt{\eta\nu}$:

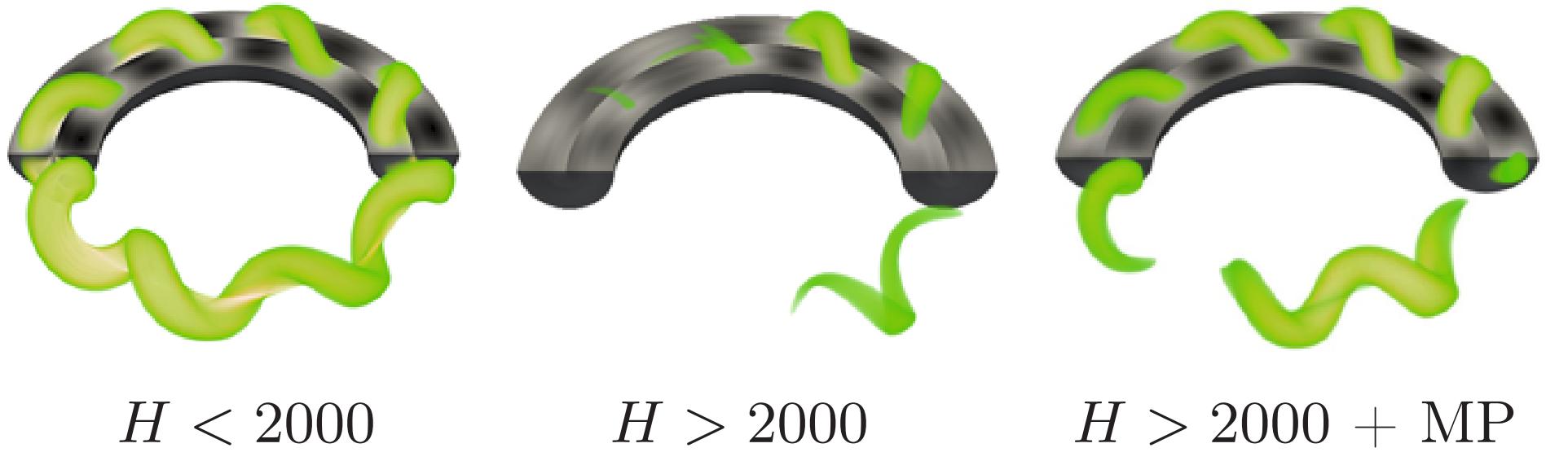
'Rescaled' SpeCyl equations

$$P^{-1} \left(\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} \right) = \mathbf{j} \times \mathbf{B} + H^{-1} \nabla^2 \mathbf{v}$$

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) - \nabla \times (H^{-1} \mathbf{j})$$

⇒ η and ν jointly rule the plasma dynamics through H
⇒ H together with Magnetic Perturbation (finite $b_r(a)$, [4]) determines the plasma helical regime.

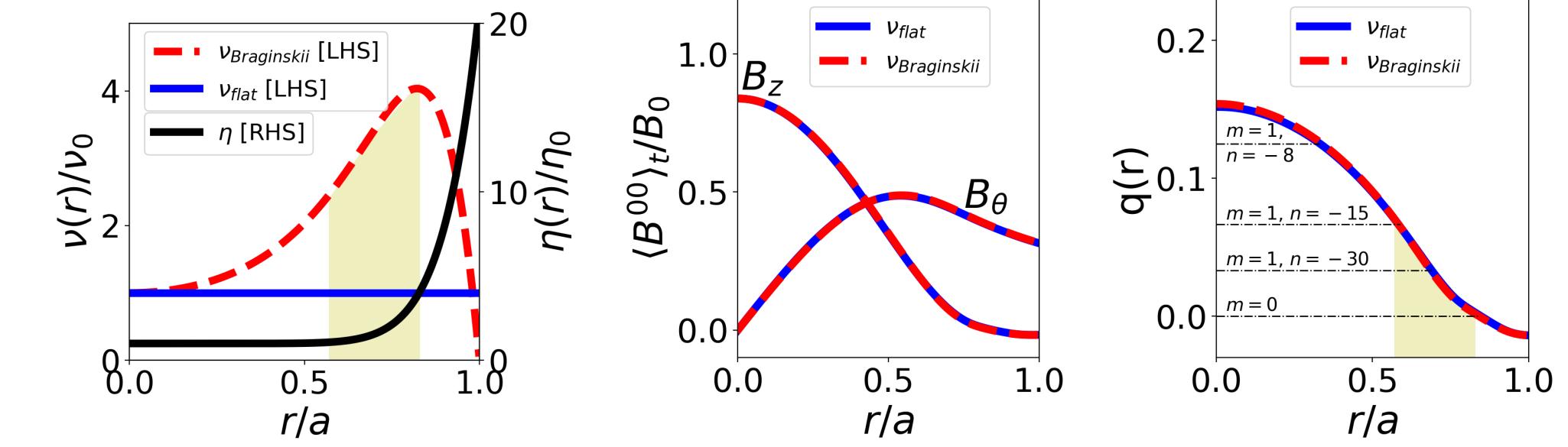
Single Helicity Multiple Helicity Quasi-Single Helicity



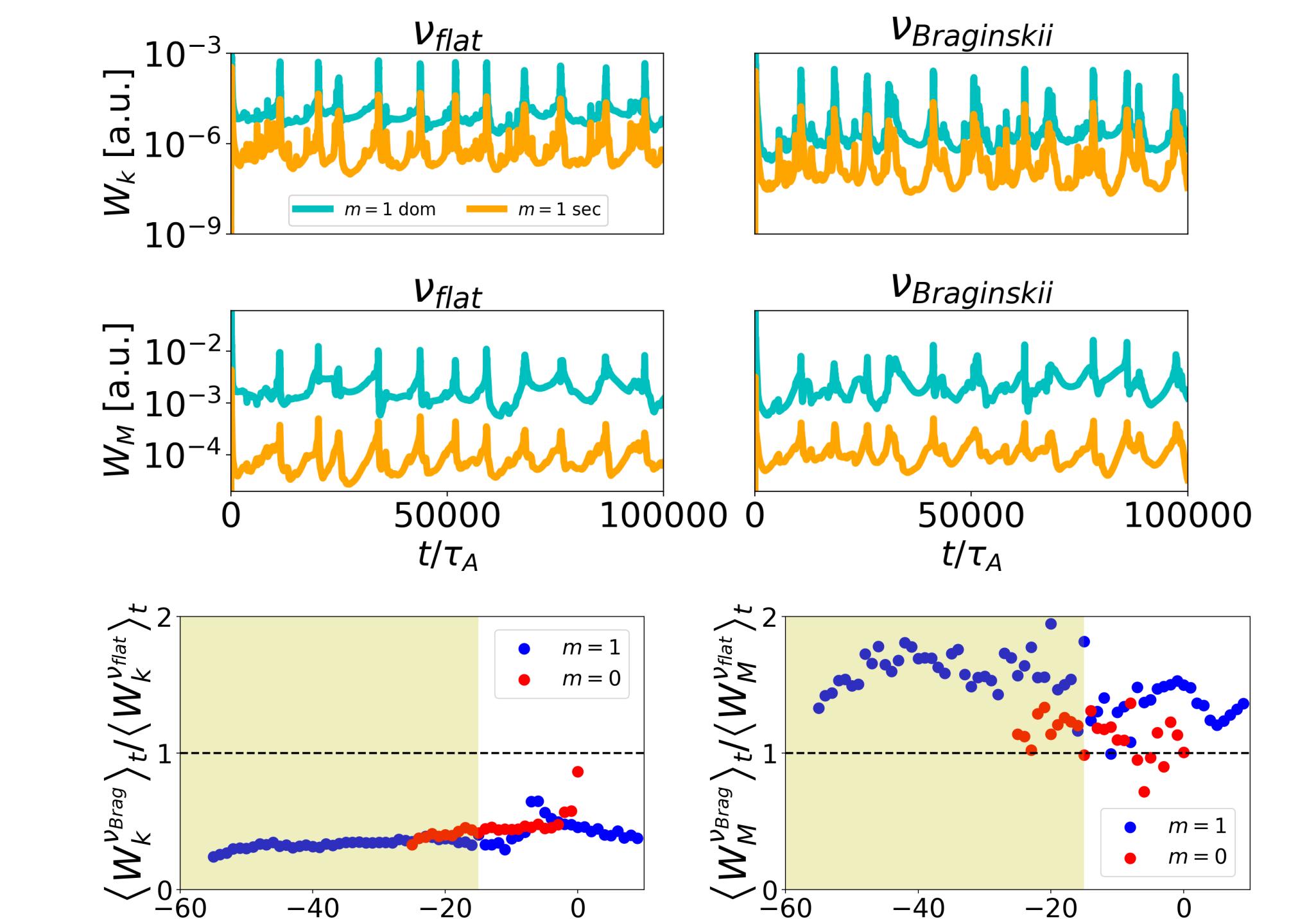
- Resistivity η profile: Spitzer-like $\eta \propto T_e^{-3/2}$, according to typical experimental temperature profiles.
- Viscosity ν profile: uniform in the majority of the simulations. ν represents the major uncertainty parameter among the transport coefficients [5]: different estimates exist according to classical [6] and turbulent theories [7].

MHD simulations: non-uniform viscosity

1) SpeCyl simulations settings: $\eta_0 = 10^{-6}$ $\nu_0 = 10^{-4}$ $b_r(a) = 0$.

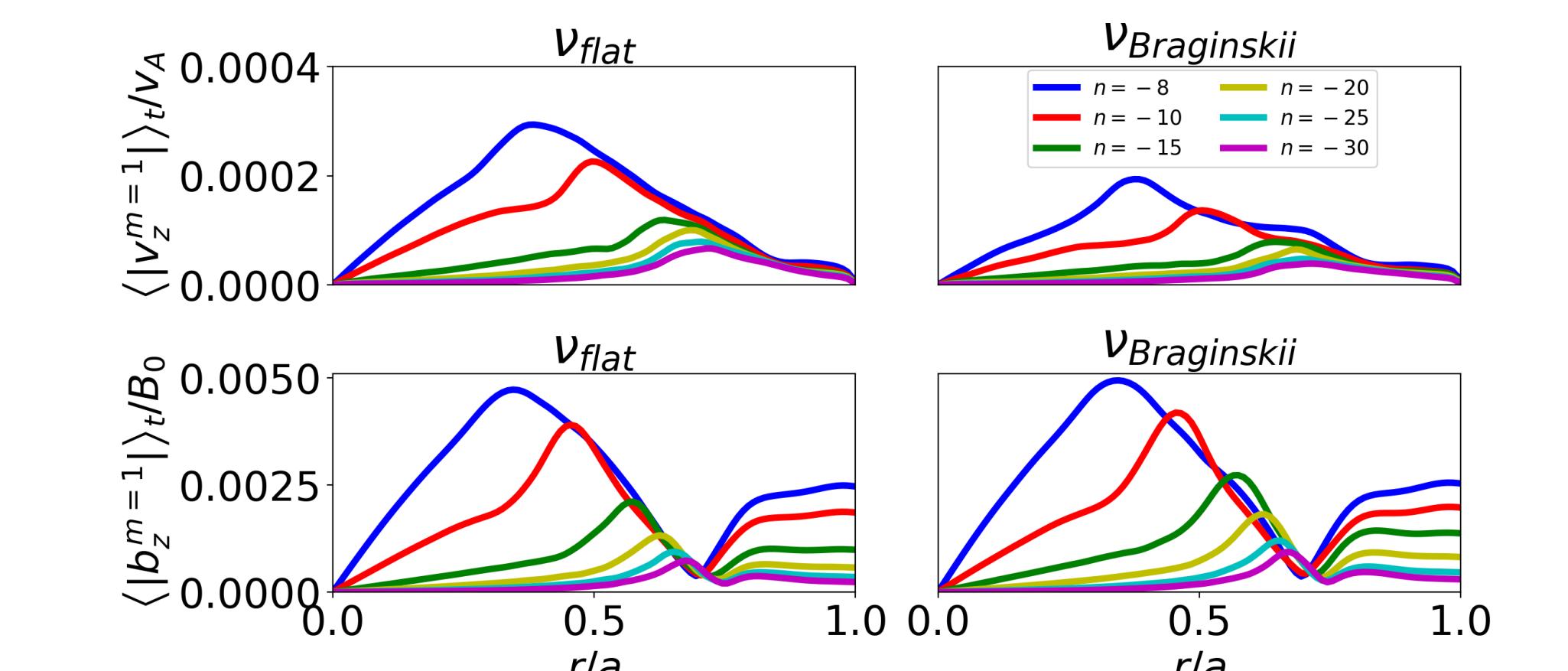


2) Effects on the kinetic and the magnetic energy:



⇒ v_{Brag} profile causes a damp of the kinetic energy W_k and a slight enhancement of the magnetic energy W_M .

3) Effects on the flow and magnetic field eigenfunctions:

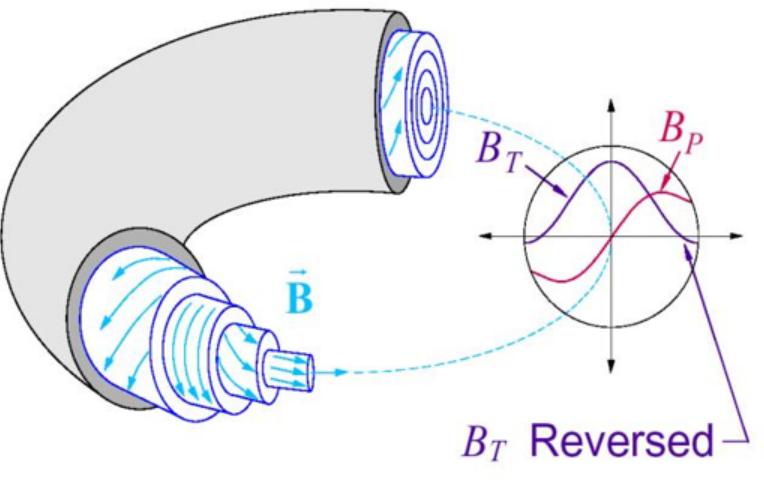


⇒ v_{Brag} profile causes a damp of velocity eigenfunctions and a slight enhancement of magnetic field eigenfunctions.

Conclusions

- Visco-resistive MHD simulations: "Braginskii-like" viscosity profiles produce a moderate damp of the plasma flow and an increased MHD magnetic activity of spectral components resonating in spatial regions where the viscosity profile is higher, confirming a simple picture of the interplay between plasma flow and magnetic field – i.e. $v \downarrow \Rightarrow b^{m,n} \uparrow$.
- RFX-mod experimental data: ν estimates according to the main theories of transport. Braginskii ν_\perp maximizes the SpeCyl simulations - exp data adherence, followed by ν_x . An anomaly in the quantitative comparison with simulations still persists, to which the viscosity profile, the axisymmetric flow effect, a $\nu_\perp - \nu_x$ interaction and a resistivity anomaly can possibly contribute.

Reversed-field pinch configuration



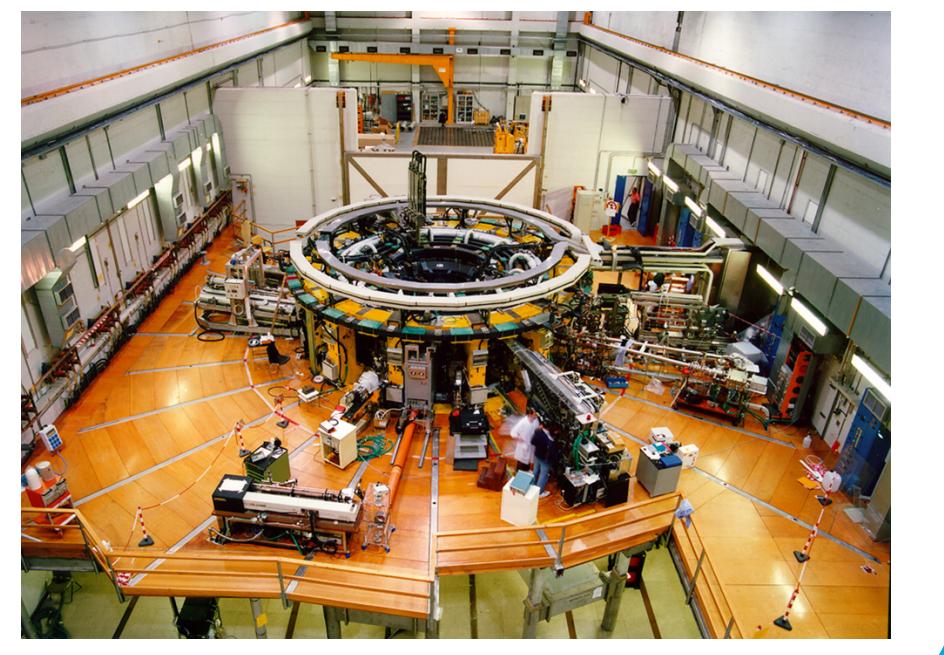
Reversed-Field Pinch (RFP) configuration [1] is characterized by:

- small reversed edge toroidal field ($B_\phi(a) < 0$)
- large plasma current I_p (10 × stronger than tokamak with the same B_ϕ)
- ⇒ high self-organization level
- same order of magnitude for magnetic fields: $B_\phi \approx B_\theta$.

RFX-mod 243 shots database [8]

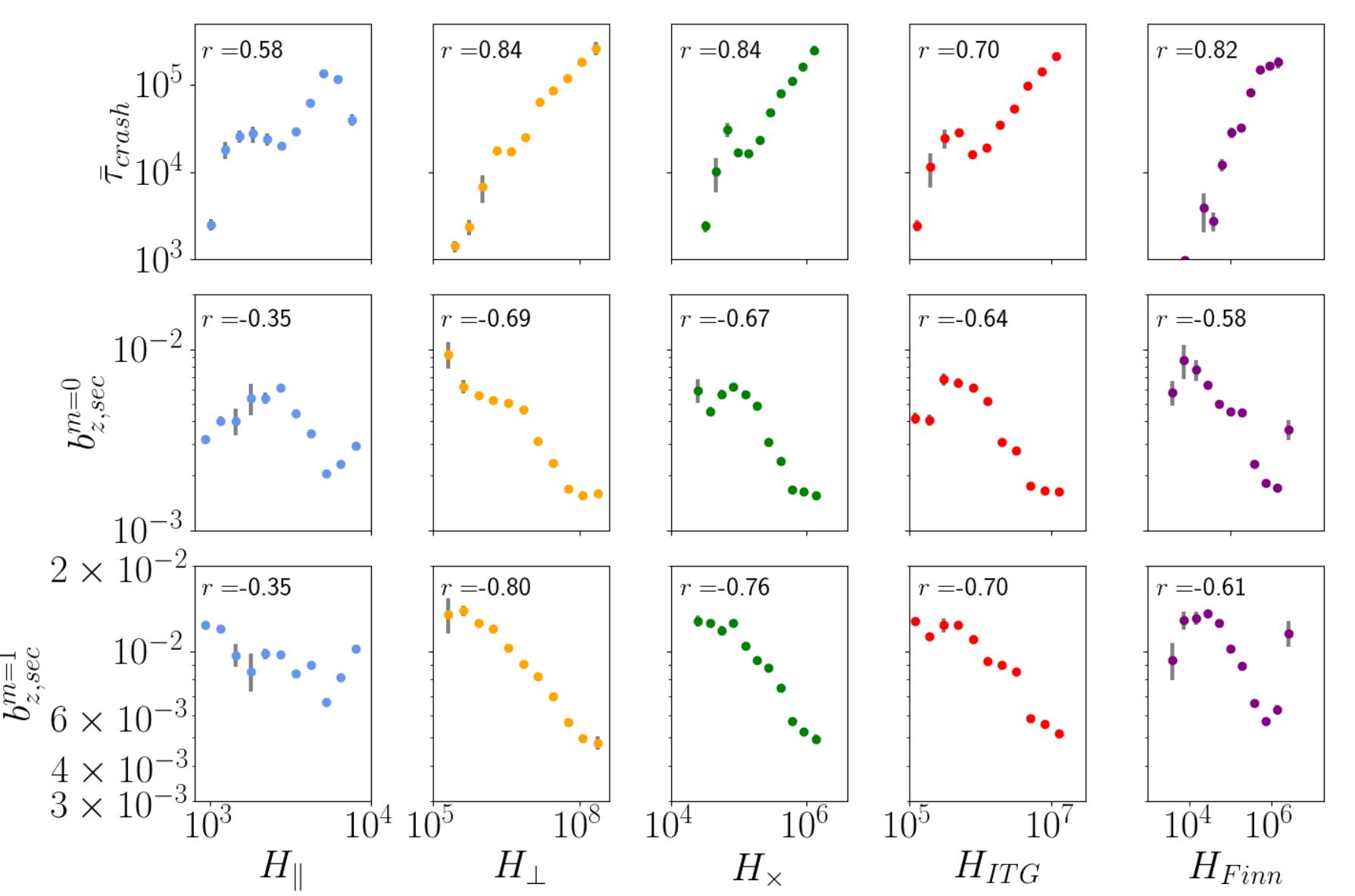
RFP config. parameters

R_0 2 m
 a 0.459 m
 I_p 130 kA – 2.0 MA
 B_0 0.15 T – 2.0 T
 n_e $3.5 \cdot 10^{18} – 2.2 \cdot 10^{20} \text{ m}^{-3}$
 T_e 60 eV – 1.1 keV

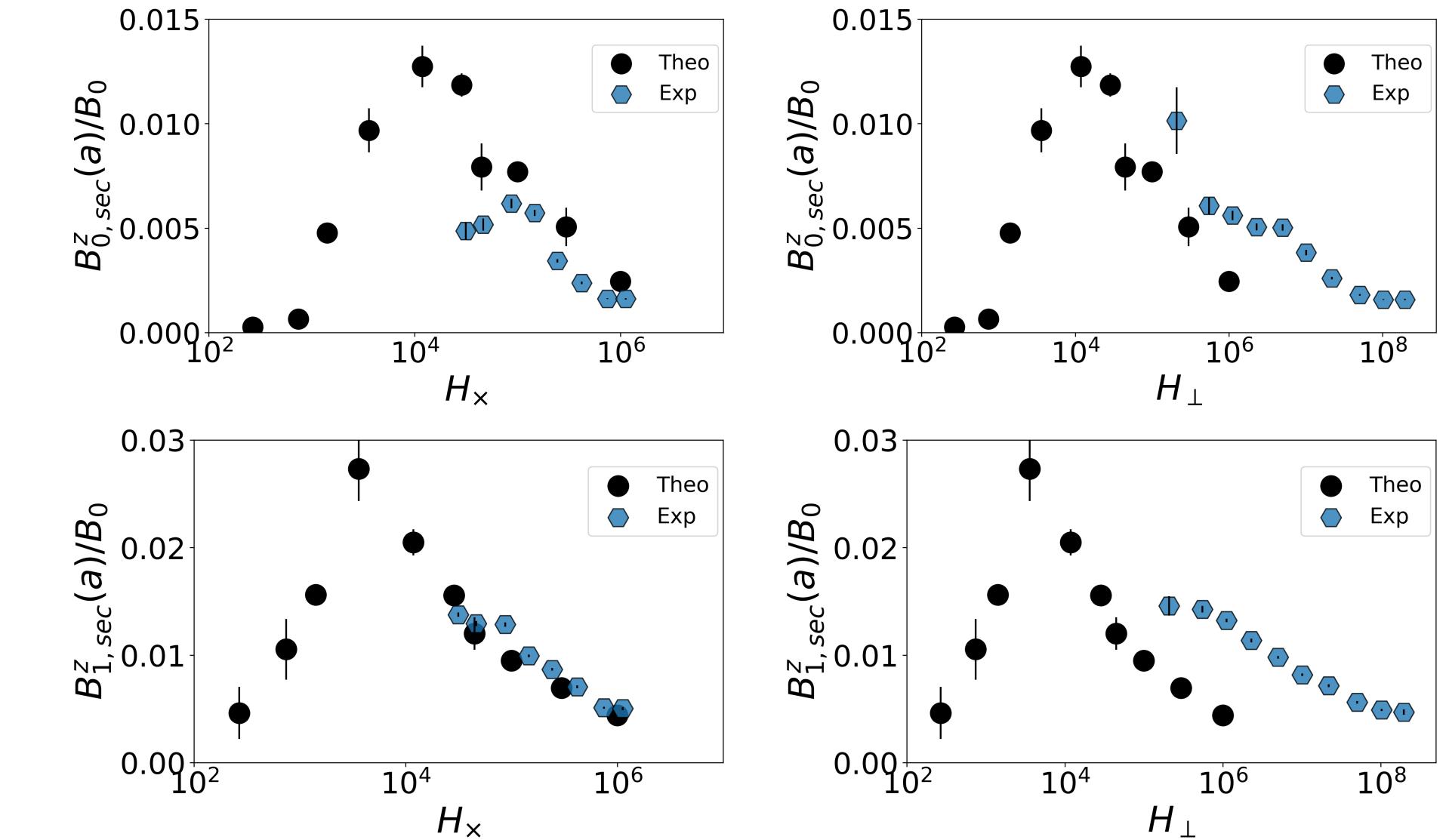


Simulations and data comparison

The normalized time scale of reconnection events $\bar{\tau}_{crash}$, the secondary $m = 0$ and $m = 1$ modes are ruled by the Hartmann number in SpeCyl simulations [12]. For the experimental data:



- ν_\parallel produces small correlation
- ν_{ITG} , ν_{Finn} produce lower correlations than the other estimates
- the ν_\perp and ν_x produce the highest correlation levels.



A non monotonous trend of the secondary $m = 0$ modes for ν_x and an anomaly factor $\delta \sim 250$ for ν_\perp are highlighted by the quantitative comparison.

References

- [1] Marrelli L. et al. *NF*, 61(2), 023001 (2021)
- [2] Cappello S., Biskamp, B. *NF*, 36, 571 (1996)
- [3] Cappello S., Escande, D. F., *PRL*, 85, 3838 (2000)
- [4] Bonfiglio D. et al., *PRL*, 111, 085002 (2013)
- [5] Montgomery D., *PPCF*, 34, 1157 (1992)
- [6] Braginskii S. I., *Rev. of Plasma Phys.*, 1 (1965)
- [7] Balescu R., *Eur. J. Phys.*, 21, 279 (2000)
- [8] Spizzo G. et al., *NF*, 55, 043007 (2015)
- [9] Guo S. C. et al., *PoP*, 1, 2741, (1994)
- [10] Finn J. M. et al., *PoF B*, 5, 1152 (1992)
- [11] Fridström R. et al., *PRL*, 120, 225002 (2018)
- [12] Veranda M. et al., *NF*, 60, 016007 (2020)