Abstract

Contact geometry is the study of odd-dimensional smooth manifolds equipped with contact structures, i.e. hyperplane distributions $\xi = \ker \alpha$ satisfying the contact condition

$$\alpha \wedge (\mathrm{d}\alpha)^n \neq 0.$$

While they originally arise in the study of ODEs and in classical mechanics, the topological study of contact manifolds is a more recent and very active field of research.

A manifold can have multiple different contact structures, which can be either rigid (in which case one speaks of a "tight" manifold) or flexible (in the sense that they satisfy an h-principle). The latter contact manifolds are then called overtwisted. A foundational result of Eliashberg and Borman–Eliashberg–Murphy, roughly speaking, states that overtwisted contact manifolds exist in abundance, namely whenever the manifold admits the topological version of a contact structure (an *almost* contact structure), which is a first obvious obstruction. In dimension three, an almost contact structure is simply an oriented 2-plane field.

To illustrate this dichotomy, consider the sphere S^3 . By a result of Eliashberg, it has precisely one tight contact structure. On the other hand, it has infinitely many overtwisted contact structures, corresponding to the infinitely many homotopy classes of 2-plane fields on the 3-sphere. There are other examples where there are infinitely many or no tight contact structures on a contact manifold.

A further interesting property of contact manifolds comes from the fact that contact geometry is the odd-dimensional counterpart to symplectic geometry. Often, it is possible to view a contact manifold as the boundary of a symplectic manifold. Manifolds that are in this sense "fillable" are always tight. The contrary, however, doesn't need to hold and one can ask the question under which conditions such tight, but non-fillable manifolds exist. The first examples of tight and non-fillable contact manifolds were constructed by Etnyre–Honda in dimension three, and by Massot–Niederkrueger–Wendl in higher dimensions.

More recently, Bowden–Gironella–Moreno–Zhou have shown that there exist homotopically standard, non-fillable but tight contact structures on all spheres S^{2n+1} with n>=2. Starting with a specific open book decomposition of S^{2n-1} , one can construct a contact form on this manifold using a well-known construction by Thurston–Winkelnkemper. Then, according to Bourgeois, this contact structure can be extended to a tight contact structure on $S^{2n-1}\times T^2$. Applying subcritical surgery (preserving the tightness), one can kill the topology of the T^2 -factor and obtain a tight contact structure on S^{2n+1} . Because of the special way of constructing it, one can show that it is non-fillable, but still homotopically standard.

The goal of my master thesis is to give a streamlined explanation of the results of Bowden–Gironella–Moreno–Zhou, including the necessary background needed to understand the main ideas.

Introduction

The Construction: A homotopically standard contact structure on the Sphere

1. Outline

The goal is to find homotopically standard, tight, non-fillable contact structures on the sphere S^{2n+1} for $n \geq 2$. As $(S^{2n+1}, \xi_{\text{std}})$ is defined as the contact boundary of the standard symplectic ball $(D^{2n+2}, \omega_{\text{std}})$, it is fillable by definition. Hence, one needs a different contact structure. This chapter explains how to construct a homotopically standard contact structure on the sphere in a way that at first doesn't seem very intuitive, hoping that in the end it will be different from $\xi_{\rm std}$. Later, in chapters 3 and 4, it will be shown that, as desired, this contact structure really is tight and non-fillable (and hence different from $\xi_{\rm std}$). The starting point for the whole construction is a Milnor open book, i.e. a certain kind of decomposition of S^{2n-1} that comes from Milnors work on hypersurface singularities [Mil69]. By the Giroux correspondence, there is a contact structure on this manifold. In this case, the Giroux correspondence can be realized by an explicit construction due to Thurston-Winkelnkemper. Now, a construction due to Bourgeois ([**Bou02**]) extends the resulting contact structure on S^{2n-1} to $S^{2n-1} \times T^2$. By applying smooth surgery, one can kill the homology in the T^2 -factor and so obtain S^{2n+1} . By an h-principle due to Eliashberg? (citation missing), the surgeries can be realized as contact surgeries. This completes the construction as now one has a contact structure on S^{2n+1} . Finally it will be explained why the contact structure is homotopically standard.

2. Raw material for the construction: The Milnor open book

Define $f \colon \mathbb{C}^n \to \mathbb{C}$ by

$$(z_0,\ldots,z_{n-1})\mapsto z_0^k+z_1^2+\ldots z_{n-1}^2.$$

Consider the sphere $S^{2n-1} \subset \mathbb{C}^n$. The intersection $f^{-1}(0) \cap S^{2n-1}$ is the so called Brieskorn sphere $B = \Sigma_{n-1}(k, 2, ..., 2)$. On the complement $S^{2n-1} \setminus B$, the map

$$\pi_f \colon S^{2n-1} \setminus B \to S^1 \colon (z_0, \dots, z_{n-1}) \mapsto \frac{f(z_0, \dots, z_{n-1})}{|f(z_0, \dots, z_{n-1})|}$$

is a fibration over S^1 , the Milnor fibration. According to ??, this is an open book where the fibers of the Milnor fibration (i.e. the Milnor fibers) form the pages.

3. From open books to contact structures: The Thurston-Winkelnkemper construction

DEFINITION 1 (mapping torus). Let Σ be a smooth manifold with boundary $\partial \Sigma$ and $\phi: \Sigma \to \Sigma$ a diffeomorphism that is equal to the identity close to $\partial \Sigma$. The

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mapping torus $\Sigma(\phi)$ is given by $\Sigma \times [0, 2\pi]/\sim$ where

$$(x, 2\pi) \sim (\phi(x), 0).$$

The generalized mapping torus requires as additional data a smooth function $\overline{\varphi}$: $\Sigma \to \mathbb{R}^+$ that is constant near $\partial \Sigma$. Then,

$$\Sigma_{\overline{\varphi}}(\phi) := \Sigma \times \mathbb{R}/\sim \quad where \quad (x,\theta) \sim (\phi(x), \theta - \overline{\varphi}(x)).$$

Abstract open books. Starting with a mapping torus $\Sigma(\phi)$, we can construct an abstract open book $M(\phi)$ with binding $\partial \Sigma$ (see fig. 1)

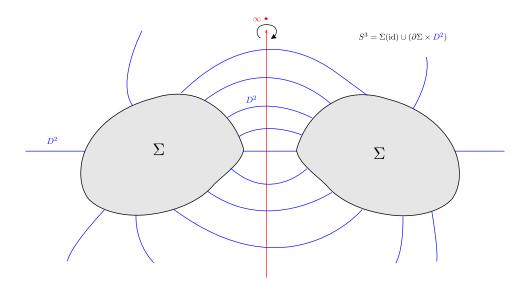


FIGURE 1. abstract open book

We define

$$M(\phi) := (\Sigma(\phi) \cup \partial \Sigma \times D^2) / \sim$$

where we identify

$$[x \in \partial \Sigma, \theta] \sim (x, r = 1, \varphi = \theta)$$

The construction. Let Σ^{2n} be a compact manifold admitting an exact symplectic form $\omega = \mathrm{d}\beta$ s.t. on the boundary $\partial \Sigma$, a contact form β_{∂} is induced (this follows from the conditions requested in Geiges). Let the boundary be connected (i.e. the binding is also connected). Let the monodromy map ϕ be an exact symplectomorphism of (Σ, ω) , equal to the identity near the boundary $\partial \Sigma$ (exactness is not necessary according to Geiges, as it can be obtained via a suitable isotopy of the symplectomorphism). An exact symplectomorphism ϕ of (Σ, ω) is such that

$$\phi^*(\beta) - \beta =: d\overline{\varphi}$$

is exact, i.e. there exists such a function $\overline{\varphi}$ on Σ (of course only defined up to adding a locally constant function. Choose it in such a way that it only takes positive values). The 1-form

$$\alpha\coloneqq\beta+\mathrm{d}\varphi$$

is a contact form on $\Sigma \times \mathbb{R}$:

$$\alpha \wedge (\mathrm{d}\alpha)^n = (\beta + \mathrm{d}\varphi) \wedge \underbrace{(\mathrm{d}\beta)^n}_{=:\Omega} = \beta \wedge \Omega + \mathrm{d}\varphi \wedge \Omega = \mathrm{d}\varphi \wedge \Omega,$$

where Ω is a volume form on Σ (as β is a symplectic form). The $\beta \wedge \Omega$ term vanishes because both are forms on Σ , but Ω is already a top-level form. The resulting form is a wedge product of two volume forms on the product manifolds and therefore a volume form on $\Sigma \times \mathbb{R}$.

Now consider the transformation that induces the generalized mapping torus

$$F := (x, \varphi) \mapsto (\phi(x), \varphi - \overline{\varphi}(x)).$$

Remember that φ only takes positive values, i.e. the mapping torus is well defined. The 1-form α is invariant under this transformation:

$$F^*(\alpha) = F^*(\beta) + F^*(d\varphi) \qquad | \beta \text{ is independent of } \varphi$$

$$= \phi^*(\beta) + dF(\varphi) \qquad | \text{ definition of } \overline{\varphi}, F$$

$$= \beta + d\overline{\varphi} + d\varphi - d\overline{\varphi}$$

$$= \alpha.$$

It follows that α descends to a contact form on $\Sigma_{\overline{\varphi}}(\phi)$.

Now, we describe an adapted gluing construction for the abstract open book coming from a generalized mapping torus. Therefore, we construct a collar neighborhood on the generalized mapping torus s.t. on $[-1,0] \times \partial \Sigma$, the symplectic form is given by $d(e^s\beta_{\partial})$ where s is the collar parameter, d.h. $\beta = e^s\beta_{\partial}$. Why does such a neighborhood exist?

Close to $\partial \Sigma$, ϕ is equal to the identity and therefore $d\overline{\varphi}$ is locally constant (hence constant, as $\partial \Sigma$ is connected). Parametrize the neighborhood so that $\overline{\varphi}$ is constant on $[-1,0] \times \partial \Sigma$.

Now, take a look at

$$M := \left(\Sigma_{\overline{\varphi}}(\phi) \ \dot{\cup} \ \left(\partial \Sigma \times D_2^2 \right) \right) / \sim .$$

A simple linear reparametrization will make the notation a lot easier: As $\overline{\varphi}$ is constant on the neighborhood under consideration, we just pretend $\overline{\varphi} = 2\pi$. Furthermore, we parametrize the boundary $\partial \Sigma$ with $\theta \in S^1$. Then we identify

$$(s, \theta, \varphi) \in [-1, 0] \times \partial \Sigma \times S^1 \subset \Sigma_{\overline{\varphi}}(\phi)$$

with

$$(\theta, s = 1 - r, \varphi) \in \partial \Sigma \times D_2^2 =: \mathcal{N}$$

where (r, φ) are polar coordinates on D_2^2 , i.e. we identify a collar neighborhood of Σ with an annulus in D_2^2 . (See fig. 2)

Now we choose the ansatz

$$\alpha_{\text{ext}} := h_1(r)\beta_{\partial} + h_2(r)d\varphi.$$

for the extension of the contact form over \mathcal{N} . On the gluing area (i.e. $1 \leq r \leq 2$), α_{ext} has to agree with $\alpha = \beta + d\varphi = e^s \beta_{\partial} + d\varphi$, i.e.

$$h_1(r) = e^s = e^{1-r}$$
 $h_2(r) = 1$.

In order to ensure smoothness at r = 0, in a small neighborhood of r = 0 we set $h_1(r) = 2 - r^2$ and $h_2(r) = r^2$, obtaining

$$\alpha_{\rm ext}|_0 = 2\beta_{\partial}$$

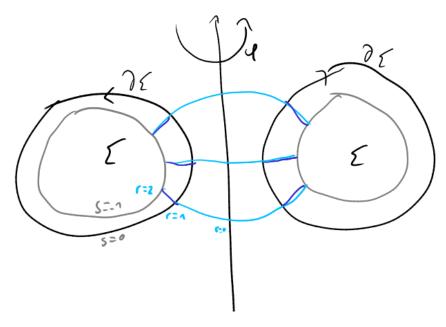


FIGURE 2. Detailled gluing process of the generalized abstract open book

We compute

$$d\alpha_{\rm ext} = h_1'(r)dr \wedge \beta_{\partial} + h_1(r)d\beta_{\partial} + h_2'(r)dr \wedge d\varphi.$$

and

$$(\mathrm{d}\alpha_{\mathrm{ext}})^n = n \cdot \mathrm{d}r \wedge (h_1'(r)\beta_{\partial} + h_2'(r)\mathrm{d}\varphi) \cdot h_1(r)^{n-1}(\mathrm{d}\beta_{\partial})^{n-1} + \underbrace{h_1(r)^n(\mathrm{d}\beta_{\partial})^n}_{0},$$

where the second term vanishes because $(d\beta_{\partial})^n$ is a 2n-form on $\partial \Sigma^{2n-1}$. Finally,

$$\begin{split} \alpha_{\mathrm{ext}} \wedge (\mathrm{d}\alpha_{\mathrm{ext}})^n &= h_1(r)nh_1(r)^{n-1}h_2'(r) \cdot & \beta_{\partial} \wedge \mathrm{d}r \wedge \mathrm{d}\varphi \wedge (\mathrm{d}\beta_{\partial})^{n-1} \\ &+ h_2(r)nh_1(r)^{n-1}h_1'(r) \cdot & \mathrm{d}\varphi \wedge \mathrm{d}r \wedge \beta_{\partial} \wedge (\mathrm{d}\beta_{\partial})^{n-1} \\ &= nh_1(r)^{n-1}(h_1h_2'(r) - h_2h_1'(r)) \cdot & \beta_{\partial} \wedge (\mathrm{d}\beta_{\partial})^{n-1} \wedge \mathrm{d}r \wedge \mathrm{d}\varphi \\ &= nh_1(r)^{n-1}\det\begin{pmatrix} h_1(r) & h_2(r)/r \\ h_1'(r) & h_2'(r)/r \end{pmatrix} \cdot & \beta_{\partial} \wedge (\mathrm{d}\beta_{\partial})^{n-1} \wedge r\mathrm{d}r \wedge \mathrm{d}\varphi \end{split}$$

As β_{∂} is a contact form on $\partial \Sigma$, $\beta_{\partial} \wedge (\mathrm{d}\beta_{\partial})^{n-1}$ is a positive volume form on $\partial \Sigma$. Furthermore, $r\mathrm{d}r \wedge \mathrm{d}\varphi$ is a positive volume form on the disk D_2^2 . As a result, the right term of our result is a volume form on $\mathcal{N} = \partial \Sigma \times D_2^2$. The left term tells us that $h_1(r)$ musn't have any zeros for $r \in [0,2]$ and that $(h_1(r),h_2(r))$ must never be parallel to $(h'_1(r),h'_2(r))$, i.e.

$$h_1(r)^{n-1} \det \begin{pmatrix} h_1(r) & h_2(r)/r \\ h'_1(r) & h'_2(r)/r \end{pmatrix} > 0 \qquad \forall r \in [0, 2].$$

(Close to zero, the determinant is given by $2 \cdot 2 - 0 \cdot 0 = 4 > 0$). Figure 4.7 in [**Gei08**] proves the existence of such a pair of functions h_1 and h_2 .

In total, we obtain that $\alpha_{\rm ext}$ induces the correct orientation on the extension and, as M is connected and orientable, on all of M. In particular, condition (i) of $\ref{eq:condition}$? holds and $\alpha \wedge (\mathrm{d}\alpha)^n = \mathrm{d}\varphi \wedge \Omega$ is a positive volume form on the mapping torus. As $\Omega = (\mathrm{d}\beta)^n = \omega^n$ for ω the symplectic form on Σ , it is a 2n-form and we see that Ω is a positive volume form on Σ . Thus, on Σ , $\mathrm{d}\alpha = \mathrm{d}\beta = \omega$ is a symplectic form that induces the positive orientation of Σ . On $\mathcal N$, we need to check that the form induced by $\mathrm{d}\alpha_{\rm ext}$ on the pages is symplectic with the right orientation. Inside $\mathcal N$, a page is given by the condition $\varphi = {\rm const}$, i.e. $\mathrm{d}\varphi = 0$. We have

$$(d\alpha_{\text{ext}})^n = n \cdot dr \wedge (h'_1(r)\beta_{\partial} + h'_2(r)d\varphi) \cdot h_1(r)^{n-1}(d\beta_{\partial})^{n-1}$$
$$= nh'_1(r)h_1(r)^{n-1}dr \wedge \beta_{\partial} \wedge (d\beta_{\partial})^{n-1}$$

A positive volume form on Σ must be positive on $-\partial_r$, \mathfrak{b} where \mathfrak{b} is a positive basis of a point in $\partial \Sigma$. As $\beta_{\partial} \wedge (\mathrm{d}\beta_{\partial})^{n-1}$ is a positive volume form on $\partial \Sigma$, we obtain

$$(\mathrm{d}\alpha_{\mathrm{ext}})^{n}(-\partial_{r},\mathfrak{b}) = \underbrace{nh_{1}(r)^{n-1}}_{=:A>0} \cdot h'_{1}(r)\mathrm{d}r(-\partial_{r}) \wedge \underbrace{\left[\beta_{\partial} \wedge (\mathrm{d}\beta_{\partial})^{n-1}\right](\mathfrak{b})}_{=:B>0}$$
$$= AB \cdot h'_{1}(r) \cdot -1$$
$$> 0.$$

where in the last line we've used $h_1'(r) = \frac{\mathrm{d}}{\mathrm{d}r}(2-r^2) = -2r < 0$. We have thus verified condition (ii) inside \mathcal{N} and outside \mathcal{N} , on Σ . As a result, is must hold on the whole page. Condition (iii) follows from the fact that on B, $\alpha_{\mathrm{ext}} = 2\beta_{\partial}$ which is a positive contact form on $\partial \Sigma$ and therefore also on B.

4. From M to $M \times T^2$: The Bourgeois construction

It is not very hard to construct a contact structure on the three-torus. When Lutz [Lut79] discovered a contact structure on T^5 , however, it was natural to wonder whether there exists a contact structure on T^{2n+1} for all $n \in \mathbb{N}$. In order to answer this long-standing question, Bourgeois [Bou02] came up with a construction that takes as input a contact structure on M and as output returns a contact structure on $M \times T^2$. To be more precise, it requires a contact structure that is supported by an open book decomposition as an input. However, a result by Giroux and Mohsen [Gei08, Theorem 7.3.5] shows that to every contact structure one can find such an open book decomposition, so that is not an obstruction, but rather part of the construction data.

4.1. General construction. Let dim $M \geq 3$ and (B, π) an open book decomposition of M supporting $(M, \xi = \ker \alpha)$. By definition of an open book, there is a trivial tubular neighborhood $B \times D^2$ around B and there exist a radial coordinate r with r = 0 precisely on B s.t. (r, π) form polar coordinates on this neighborhood. Choose a smooth function ρ of r s.t. $\rho = r$ close to r = 0, $\rho'(r) > 0$ and $\rho = 1$ at ∂D . Extend this function to M by setting $\rho = 1$ on $M \setminus B \times D$ (see fig. 3). As π and ρ are smooth functions on M, one can define the smooth functions $x_1 := \rho \cos \pi$ and $x_2 := \rho \sin \pi$ on all of M. On $B \times D^2$, they coincide with the Cartesian coordinate functions near B. As always with corresponding polar- and cartesian coordinates,

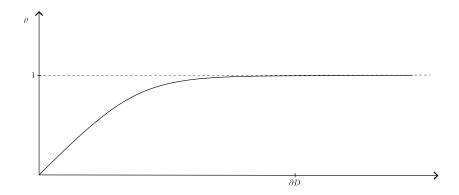


FIGURE 3. Onedimensional sketch of ρ

they satisfy

$$x_1 dx_2 - x_2 dx_1 = \rho^2 \cos^2 \pi d\pi + \rho \cos \pi \sin \pi d\rho + \rho^2 \sin^2 \pi d\pi - \rho \cos \pi \sin \pi d\rho$$
$$= \rho^2 (\cos^2 \pi + \sin^2 \pi) d\pi$$
$$= \rho^2 d\pi$$

and, analogously,

$$\mathrm{d}x_1 \wedge \mathrm{d}x_2 = \rho \mathrm{d}\rho \wedge \mathrm{d}\pi.$$

On $M \times T^2$, choose coordinates (θ_1, θ_2) on the torus part of the manifold. Define

$$\tilde{\alpha} := x_1 d\theta_1 - x_2 d\theta_2 + \alpha.$$

to be a 1-form on M (where α is extended in the obvious way to $M \times T^2$ as the pullback $\pi_1^*\alpha$). This is the candidate for the contact form on $M \times T^2$. Having computed the exterior derivative and its n-th power

$$d\tilde{\alpha} = dx_1 \wedge d\theta_1 - dx_2 \wedge d\theta_2 + d\alpha,$$

$$(d\tilde{\alpha})^n = (n-1)(d\alpha)^{n-1} \wedge (dx_1 \wedge d\theta_1 - dx_2 \wedge d\theta_2)$$

$$- n(n-1)(d\alpha)^{n-2} \wedge dx_1 \wedge d\theta_1 \wedge dx_2 \wedge d\theta_2,$$

one can check the contact condition:

$$\tilde{\alpha} \wedge (\mathrm{d}\tilde{\alpha})^{n} = (x_{1}\mathrm{d}\theta_{1} - x_{2}\mathrm{d}\theta_{2} + \alpha) \wedge (n-1)(\mathrm{d}\alpha)^{n-1} \wedge (\mathrm{d}x_{1} \wedge \mathrm{d}\theta_{1} - \mathrm{d}x_{2} \wedge \mathrm{d}\theta_{2})$$

$$- (x_{1}\mathrm{d}\theta_{1} - x_{2}\mathrm{d}\theta_{2} + \alpha) \wedge n(n-1)(\mathrm{d}\alpha)^{n-2} \wedge \mathrm{d}x_{1} \wedge \mathrm{d}\theta_{1} \wedge \mathrm{d}x_{2} \wedge \mathrm{d}\theta_{2}$$

$$= (n-1)(\mathrm{d}\alpha)^{n-1} \wedge (x_{1}\mathrm{d}x_{2} - x_{2}\mathrm{d}x_{1}) \wedge \mathrm{d}\theta_{1} \wedge \mathrm{d}\theta_{2}$$

$$+ \underbrace{\alpha \wedge (n-1)(\mathrm{d}\alpha)^{n-1} \wedge \mathrm{d}x_{1}}_{2n-\text{form on }M} \wedge \mathrm{d}\theta_{1} - \underbrace{\alpha \wedge (n-1)(\mathrm{d}\alpha)^{n-1} \wedge \mathrm{d}x_{2}}_{2n-\text{form on }M} \wedge \mathrm{d}\theta_{2}$$

$$+ n(n-1)\alpha \wedge (\mathrm{d}\alpha)^{n-2} \wedge \mathrm{d}x_{1} \wedge \mathrm{d}\theta_{1} \wedge \mathrm{d}x_{2} \wedge \mathrm{d}\theta_{2}$$

M has dimension 2n-1, i.e. the middle term is 0

$$= (n-1)(\mathrm{d}\alpha)^{n-1} \wedge \rho^2 \mathrm{d}\phi \wedge \mathrm{d}\theta_1 \wedge \mathrm{d}\theta_2 + n(n-1)\alpha \wedge (\mathrm{d}\alpha)^{n-2} \wedge \rho \mathrm{d}\rho \wedge \mathrm{d}\phi \wedge \mathrm{d}\theta_1 \wedge \mathrm{d}\theta_2.$$

As this expression is a top-dimensional form, it suffices to show that its nowhere zero. For that, one needs to employ the fact that α is supported by (B,π) . By condition (ii) of ??, $(\mathrm{d}\alpha)^{n-1}$ must be a positive volume form on the pages. As explained in that definition, the orientation on M is given by ∂_{ϕ} and the orientation of the page. In particular, $(\mathrm{d}\alpha)^{n-1} \wedge \rho \mathrm{d}\phi$ is a positive volume form on M. Multiplied with a second ρ -factor, it vanishes along B. As $\theta_1 \wedge \theta_2$ is a positive volume form on T^2 , the first term is non-negative everywhere and positive away from

$$\underbrace{B \times 0}_{\subset B \times D^2 \subset M} \times T^2.$$

Let \mathfrak{b} be a basis of the binding B that is positively ordered. Then, $-\partial_r$, \mathfrak{b} and (because the binding is odd-dimensional) \mathfrak{b} , ∂_r are positive bases of the page. Clearly, then,

$$\mathfrak{a} := \mathfrak{b}, \partial_r, \partial_\phi, \partial_{\theta_1}, \partial_{\theta_2}$$

is an ordered basis of $M \times T^2$. Using $\rho'(r) \ge 0$ everywhere, it follows that $d\rho(\partial_r)$ is non-negative. Hence, plugging $\mathfrak a$ into the second term,

$$(n(n-1)\alpha \wedge (d\alpha)^{n-2} \wedge \rho d\rho \wedge d\phi \wedge d\theta_1 \wedge d\theta_2) (\mathfrak{a})$$

$$= n(n-1)\rho \cdot (\alpha \wedge (d\alpha)^{n-2})(\mathfrak{b}) \cdot d\rho(\partial_r) \cdot d\phi(\partial_\phi) \cdot d\theta_1(\partial_{\theta_1}) \cdot d\theta_2(\partial_{\theta_2})$$
>0.

By condition (iii) of $\ref{eq:condition}$, $\alpha \wedge (\mathrm{d}\alpha)^{n-2}$ is positive on B. Therefore, the second term is positive on $B \times 0 \times T^2$ (hence also on a neighborhood) and non-negative everywhere else. This proves the contact condition and $\tilde{\alpha}$ is indeed a contact form on $M \times T^2$.

4.2. Example: Bourgeois and Thurston-Winkelnkemper. Applying the Thurston-Winkelnkemper construction yields a contact form α_{ext} on

$$M = (\Sigma_{\overline{\varphi}}(\phi) \dot{\cup} \mathcal{N}) / \sim$$
.

Now, we apply the Bourgeois construction to it and obtain a contact form

$$\tilde{\alpha}_{\text{ext}} = \alpha_{\text{ext}} + x_1 d\theta_1 - x_2 d\theta_2$$

on $M \times T^2$ where θ_1, θ_2 are coordinates on T^2 and x_1, x_2 are coordinates as described in the section on the Bourgeois construction, i.e. there is a function $\rho: M \to [0, 1]$ that agrees with r near the binding B and we define

$$x_1 \coloneqq \rho \cos(p); \qquad x_2 \coloneqq \rho \sin(p).$$

Inside \mathcal{N} , our projection map p is given by the angular coordinate φ . Therefore,

$$\alpha|_{\mathcal{N}} = \alpha_{\text{ext}} = h_1(r)\beta_{\partial} + h_2(r)d\varphi = h_1(r)\beta_{\partial} + h_2(r)dp.$$

In total, we obtain for the contact form on $\mathcal{U} := \mathcal{N} \times T^2$

$$\tilde{\alpha} = h_1(r)\beta_{\partial} + h_2(r)dp + \rho(r)(\cos(p)d\theta_1 - \sin(p)d\theta_2)$$

Outside \mathcal{U} , the form is given by

$$\tilde{\alpha} = \beta + dp + \cos(p)d\theta_1 - \sin(p)d\theta_2,$$

as $\rho(r) = 1 \ \forall r \geq 1$. Collecting all the conditions on h_1 and h_2 from the last section, we require that

- $h_1(r) = e^s = e^{1-r}, h_2(r) = 1$ in the gluing area $(1 \le r \le 2)$.
- Smoothness around the binding and contact condition on the binding: $h_1(r) = 2 r^2$ and $h_2(r) = r^2$ around r = 0. The $-r^2$ -part is only important to have $h'_1(r) < 0$ around 0 to satisfy the symplectic condition on the pages.
- Contact condition on the tubular neighborhood:

$$h_1(r)^{n-1} \det \begin{pmatrix} h_1(r) & h_2(r)/r \\ h'_1(r) & h'_2(r)/r \end{pmatrix} > 0 \quad \forall r \in [0, 2].$$

4.3. Reeb dynamics. We now define two functions $\mu, \nu: M \to \mathbb{R}$ as follows:

$$\mu = \begin{cases} \frac{\rho'}{\rho' h_1 - \rho h_1'} & \text{inside } \mathcal{U} \\ 0 & \text{outside } \mathcal{U} \end{cases} \quad \text{and} \quad \nu = \begin{cases} \frac{-h_1'}{\rho' h_1 - \rho h_1'} & \text{inside } \mathcal{U} \\ 1 & \text{outside } \mathcal{U} \end{cases}.$$

Both μ and ν are smooth. First of all, we note that both are quotients of smooth functions. We have to make sure that $\rho' h_1 \neq \rho h'_1$. All of the involved functions are non-negative except for h'_1 , which is constantly negative (h_1 needs to be chosen in such a way that h'_1 is nonzero). As a result,

$$\rho' h_1 > 0 > \rho h_1'$$

so the denominator is always nonzero. As ρ is constantly 1 for $r \geq 1$, we see that μ is 0 for $r \geq 1$. In particular, it is smooth everywhere. For ν , setting $\rho \equiv 1$ implies $\rho' = 0$ on $r \geq 1$. Therefore, the denominator simplifies to 1 and ν is smooth, too.

LEMMA 1. The Reeb vector field of the contact form $\tilde{\alpha}$ is given by

$$R = \mu(r)R_B + \nu(r)[\cos(p)\partial_{\theta_1} - \sin(p)\partial_{\theta_2}],$$

where R_B is the Reeb vector field of β_{∂} , the contact form on $\partial \Sigma$.

PROOF. Inside \mathcal{U} , we compute

$$\tilde{\alpha}(R) = h_1(r)\beta_{\partial}(R) + h_2(r)\mathrm{d}p(R) + \rho(r)(\cos(p)\mathrm{d}\theta_1 - \sin(p)\mathrm{d}\theta_2)(R)$$

 β_{∂} and dp are 0 on ∂_{θ_i} , p and θ_i are constant on R_B

$$= h_1 \mu \beta_{\partial}(R_B) + \rho \nu (\cos(p) d\theta_1 - \sin(p) d\theta_2) (\cos(p) \partial_{\theta_1} - \sin(p) \partial_{\theta_2})$$

$$= h_1 \mu + \rho \nu [\cos(p) \cos(p) + \sin(p) \sin(p)]$$

$$= \frac{h_1 \rho'}{\rho' h_1 - \rho h'_1} + \frac{-\rho h'_1}{\rho' h_1 - \rho h'_1}$$

$$= 1$$

Computing $d\tilde{\alpha}$, we immediately drop the dr-terms, as they evaluate to 0 on R.

$$d\tilde{\alpha}(R,\cdot) = h_1(r)d\beta_{\partial}(R,\cdot) - \rho(r)[\sin(p)dp \wedge d\theta_1 + \cos(p)dp \wedge d\theta_2](R,\cdot)$$

 β_{∂} and dp are 0 on ∂_{θ_i} , d $\beta_{\partial}(R_B) = 0$, p is 0 on R_B

$$= \rho(r) \left[\sin(p) d\theta_1(R) dp + \cos(p)\theta_2(R) dp \right]$$

 $d\theta_i$ is 0 on R_B

$$= \rho(r)\nu(r) \left[\sin(p)\cos(p) - \cos(p)\sin(p)\right] dp$$
$$= 0$$

Outside \mathcal{U} , we have

$$R = \cos(p)\partial_{\theta_1} - \sin(p)\partial_{\theta_2}$$
.

Therefore,

$$\tilde{\alpha}(R) = \beta(R) + dp(R) + \cos(p)d\theta_1(R) - \sin(p)d\theta_2(R)$$

 β and dp are 0 on ∂_{θ_i}

$$= [\cos(p)d\theta_1 - \sin(p)d\theta_2](\cos(p)\partial_{\theta_1} - \sin(p)\partial_{\theta_2})$$

= $\cos(p)\cos(p) + \sin(p)\sin(p)$
= 1

Also,

$$d\tilde{\alpha}(R) = d\beta(R, \cdot) - \sin(p)dp \wedge d\theta_1(R, \cdot) - \cos(p)dp \wedge d\theta_2(R, \cdot)$$

 β and dp are 0 on ∂_{θ_i}

$$= \sin(p) d\theta_1(R) dp + \cos(p) d\theta_2(R) dp$$

$$= \sin(p) \cos(p) dp - \cos(p) \sin(p) dp$$

$$= 0$$

From here on, one can compute the Conley-Zehnder-index of the periodic orbits etc. - lots of things that are blackboxed later. Maybe state the results here for transparency.

5. Surgery

Carrying out the constructions of the previous sections, one starts with the Milnor open book decomposition of S^{2n-1} , turns it into a contact manifold via the Thurston-Winkelnkemper construction and then constructs a contact structure on $S^{2n-1} \times T^2$. However, a contact structure on S^{2n+1} is needed!

This is where surgery comes in useful: Surgery is a technique to change manifolds in a controlled way. See [Mil61] The basic idea is that the product manifold $S^p \times S^q$ can be considered either

(1) as the boundary of
$$S^p \times D^{q+1}$$
,

or

(2) as the boundary of
$$D^{p+1} \times D^q$$
.

Hence, given any imbedding of $S^p \times D^{q+1}$ in a manifold M, we can remove the interior and replace it with $D^{p+1} \times S^q$.

6. Homotopy Class of the Contact Structure

DEFINITION 2. An almost contact structure is a cooriented hyperplane field η (with an oriented trivial line bundle complementary to η defining the coorientation) and a complex bundle structure J on η . According to [Gei08, Prop 2.4.5], the space of complex bundle structures compatible with a symplectic form on η is non-empty and contractible. Hence, it suffices to choose a symplectic form ω on η to determine the almost contact structure up to homotopy (as the space of trivial line bundles complementary to η is non-empty and contractible, too).

Tightness

Non-Fillability

Convex decomposition

DEFINITION 3 (ideal Liouville domain). [Gir20, Definition 1] An ideal Liouville domain (F, ω) is a domain F endowed with an ideal Liouville structure ω . This ideal Liouville structure is an exact symplectic form on int F admitting a primitive λ such that: for some (and then any) function $u: F \to \mathbb{R}_{\geq 0}$ with regular level set $\partial F = \{u = 0\}$, the product $u\lambda$ extends to a smooth 1-form on F which induces a contact form on ∂F .

DEFINITION 4 (corresponding Giroux domain). [MNW13, Section 5.3] Given an ideal Liouville domain (F, ω) with primitive λ and function $u: F \to \mathbb{R}_{\geq 0}$ as above, the corresponding Giroux domain is given by

$$F \times S^1_{\theta}$$

endowed with contact structure

$$\ker(u\mathrm{d}\theta + u\lambda)$$

Start with a Bourgeois manifold $BO(\Sigma,...)$. Smoothly, we have

$$BO(\Sigma,...) = OB(\Sigma,...) \times T^2 = [OB(\Sigma,...) \times S^1] \times S^1 =: V \times S^1$$

According to [DG12, Section 6] we obtain a convex decomposition of the first factor

$$V = V_+ \bigcup_{\Gamma} \overline{V}_-,$$

where V_{\pm} are ideal Liouville domains and Γ is the dividing set. In [**DG12**, Section 5.3], the authors explicitly compute Γ and V_{\pm} for the Bourgeois construction and obtain

$$\Gamma = \{y = 0\} = p^{-1}(\{0\}) \cup_B p^{-1}(\{\pi\})$$

and

$$V_{+} = p^{-1}([0,\pi]) \times S^{1}, \qquad V_{-} = p^{-1}([\pi, 2\pi]) \times S^{1},$$

i.e. topologically we get $V_{\pm} = \Sigma \times D^*S_1$. If $\alpha + x d\phi + y d\theta$ is the contact structure on $OB(\Sigma, ...)$, then as explained in [**DG12**, Section 5.3], $\alpha + x d\phi$ is a S_{ϕ}^1 -invariant contact form on Γ ,

$$\omega_{\pm} = \pm \mathrm{d} \left(\frac{\alpha}{y} + \frac{x}{y} \mathrm{d} \phi \right)$$

is an S^1_{ϕ} -invariant symplectic form on V_{\pm} and y is a function with zero level set $\pm \Gamma = \partial V_{\pm}$. Hence, (V_{\pm}, ω_{\pm}) is an ideal Liouville domain with Liouville form

$$\beta_{\pm} = \pm \left(\frac{\alpha}{y} + \frac{x}{y} d\phi\right).$$

According to definition 4, $V_{\pm} \times S_{\theta}^{1}$ endowed with the contact structure

$$\ker(yd\theta + y\beta_{\pm}) = \alpha + xd\phi + yd\theta$$

is the corresponding Giroux domain. Clearly, this is just the restriction of the open book contact structure. Hence, the whole procedure actually yields a splitting into two Giroux domains

$$OB(\Sigma, \dots) = V_{+} \times S_{\theta}^{1} \bigcup_{\Gamma \times S_{\theta}^{1}} V_{-} \times S_{\theta}^{1}$$

Surgery along embedded Giroux domains. Given a embedded Giroux domain, this section describes a procedure to remove its interior and "blow down" the resulting boundary. We will refer to the procedure as "clean cut-out" of the Giroux domain.

These boundary components are always of the form $B = S^1 \times M$. Topologically, blowing down is equivalent to simply gluing in $D^2 \times M$.

This operation can be performed in a way that respects the contact structure, provided that $S^1 \times M$ has a neighborhood of the form $[0, \epsilon)_s \times S^1_t \times M$ where $\alpha_M + sdt$ defines a contact form. In general this holds by [MNW13, Lemma 5.1] if the boundary components are ξ -round hypersurfaces, but it is also possible to show the existence of that neighborhood directly.

Let D be the disk of radius ϵ in \mathbb{R}^2 . The map

$$\Psi \colon (re^{i\theta}, m) \mapsto (s = r^2, t = \theta, m)$$

is a diffeomorphism from $(D \setminus \{0\}) \times M$ to $(0, \epsilon)_s \times S_t^1 \times M$ s.t.

$$\Psi^*(\alpha_M + s dt) = \Psi^*(\alpha_M) + \Psi^*(s dt) = \alpha_M + r^2 d\theta,$$

where the latter contact form can be extended to all of $D^2 \times M$. In summary: If there is such a neighborhood of $M \times S^1$ as described above, we can glue $D \times M$ to $V \setminus B$ to get a new contact manifold in which B has been replaced by M.

Boundary components of Giroux domains are ξ -round hypersurfaces ([MNW13, Section 5.3]), Therefore, after removal of a Giroux domain, we can blow down its boundary components. These two steps together form the clean cut-out.

In [MNW13, Section 6], where the clean cut-out is first introduced, it is shown that it corresponds to a symplectic cobordism. The setting there is actually more general: The authors consider a Giroux domain where already some of the boundary components have been blown down. In our case, the situation is simpler: The Giroux domain $V_- \times S^1 \subset \mathrm{BO}(\Sigma, \ldots)$ is directly obtained from the corresponding ideal Liouville domain V_- by round contactization. Its boundary is given by

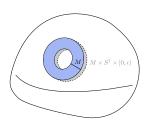
$$\partial V_- \times S^1 = \Gamma \times S^1$$
.

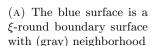
We want a cobordism from $BO(\Sigma,...)$ to the same manifold after clean cut-out of $V_- \times S^1$. Topologically, the cobordism

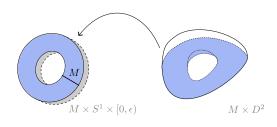
$$W \coloneqq [0,1] \times \mathrm{BO}(\Sigma,\dots) \bigcup_{\{1\} \times (V_- \times S^1)} V_- \times D^2.$$

fulfills the requirements: One boundary is simply $BO(\Sigma,...)$. The other boundary is

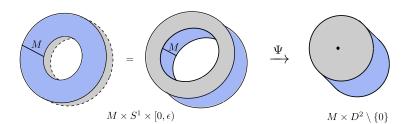
$$[BO(\Sigma, ...) \setminus (V_- \times S^1)] \cup_{\Gamma \times S^1} \Gamma \times D^2,$$







(B) Topologically, the blowdown corresponds to gluing $M \times D^2$ on top of the blue surface.



(c) Visualization of Ψ : The neighborhood of the blue area $M \times S^1 \times \{0\}$ on the left is given by $M \times$ the gray annulus. In the middle, we twist it so that the blue area sits on the inside. Applying Ψ is simply retracting the inner circle of the gray annulus to a point. The blowdown effectively reduces the inner blue surface $M \times S^1$ to $M = M \times \{0\} \subset M \times D^2$.

which is homeomorphic to $BO(\Sigma,...) \setminus (V_- \times S^1)$ with blown down boundary, which is exactly the desired clean cut-out. Using the more general [MNW13, Theorem 6.1], one can realize W as a strong symplectic cobordism.

Why is it strong? What would be the Liouville vector field? Maybe the reasoning is as follows: Start with the symplectization of $BO(\Sigma,...)$. This clearly is a strong cobordism. Maybe the handle attachment can be done in a strong way?

Next, we compute the contact structure on $BO(\Sigma,...)\setminus (V_-\times S^1)$ with blown down boundary $\partial V_-\times S^1=\Gamma\times S^1$.

This should be the contact boundary of $V_+ \times D^2$. Topologically of course that all makes sense,

$$\partial(V_+ \times D^2) = V_+ \times S^1 \cup \partial V_+ \times D^2 = V_+ \times S^1 \cup_{\Gamma \times S^1} \Gamma \times D^2.$$

However, it is unclear to me why $V_+ \times D^2$ should carry a symplectic structure.

It should also be supported by the open book $OB(V_+, id)$, which, again, topologically, makes perfect sense, but I don't understand the contact side of things.

Subcritical surgery and blowdown in one big symplectic cobordism. Now we apply subcritical surgery in the complement of $\overline{V}_- \times S^1$ and remove the Giroux domain $V_- \times S^1$.

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