

# Tricks

$$H_f |x\rangle = f(x) |x\rangle \text{ - eigenvalue equation}$$

$$|\psi\rangle = \sum_x c_x |x\rangle \text{ - eigen vectors are complete linear combination of basis vectors}$$

$$\langle x | \psi \rangle = \langle x | \sum_y c_y |y\rangle = \sum_y \langle x | y \rangle c_y = c_x \langle x | x \rangle = c_x$$

## Expectation Value

$$\begin{aligned} \langle \psi | H_f | \psi \rangle &= \sum_{x,y} c_y^* c_x \langle y | H_f | x \rangle = \sum_{x,y} c_y^* c_x f(x) \langle y | x \rangle \\ &= \sum_x c_x^* c_x f(x) = \sum_x |c_x|^2 f(x) \end{aligned}$$

eigenvalue

$$\text{Probability: } |c_x|^2 = |\langle x | \psi \rangle|^2$$

## Operators

- Hermitian Operators:  $A^\dagger = A$
- Guaranteed real eigenvalues
- Possible measurement are  $\lambda_1, \dots, \lambda_j$
- generic superposition:  $\sum_j c_j^k | \lambda_j^k \rangle$  collapse:  $\frac{\sum_k c_j^k | \lambda_j^k \rangle}{\sqrt{\sum_k |c_j^k|^2}}$

$$A | \lambda_j^k \rangle = \lambda_j^k | \lambda_j^k \rangle$$

multiple eigenvectors per eigenvalue

measure first qubit as zero

$$2 \text{ qubit: } \frac{a_{00} |00\rangle + a_{01} |01\rangle}{\sqrt{|a_{00}|^2 + |a_{01}|^2}}$$

## Computational Basis Operators

$$N = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad \leftarrow 1 \text{ qubit}$$
$$\lambda = 0, 1 \quad \lambda = -1, 1$$

## Expectation Value

$$\langle A \rangle_\psi = \sum_{j,k} |\langle \lambda_j^k | \psi \rangle|^2 \lambda_j \quad \begin{array}{l} \text{probability} \\ \text{value} \end{array} \quad \text{~ typical expected value}$$

$$= \sum_{j,k} \langle \lambda_j^k | \psi \rangle \langle \lambda_j^k | \psi \rangle^* \lambda_j$$

$$= \sum_{j,k} \langle \psi | \lambda_j^k \rangle \langle \lambda_j^k | \psi \rangle \lambda_j$$

$$= \sum_{j,k} \langle \lambda_j^k | \psi \rangle \langle \psi | \lambda_j^k \rangle \lambda_j$$

$$= \sum_{j,k} \langle \lambda_j^k | \psi \rangle \langle \psi | \lambda_j^k \rangle \lambda_j \quad \text{eigenvalue}$$

$$= \sum_{j,k} \langle \lambda_j^k | \psi \rangle \langle \psi | A | \lambda_j^k \rangle$$

$$= \langle \psi | A \sum_{j,k} \langle \lambda_j^k | \psi \rangle | \lambda_j^k \rangle$$

$$= \langle \psi | A \sum_{j,k} \alpha_j^k | \lambda_j^k \rangle$$

$$= \langle \psi | A | \psi \rangle$$

## Variational Principle

(all prob = 1)

$$\langle A \rangle_\psi = \sum_{j,k} |\langle \lambda_j^k | \psi \rangle|^2 \lambda_j \geq \sum_{j,k} |\langle \lambda_j^k | \psi \rangle|^2 \lambda_0 = \lambda_0$$

→ Any observable can be represented as the tensor product of Pauli matrices

Example:  $\langle \psi | \frac{1}{2} Z \otimes I \otimes X - 3 I \otimes Y \otimes Y | \psi \rangle$  (linearity)

$$= \frac{1}{2} \langle \psi | Z \otimes I \otimes X | \psi \rangle - 3 \langle \psi | I \otimes Y \otimes Y | \psi \rangle$$

# VQE General Buildup

1) Eigenvalue Equation

$$H|x\rangle = f(x)|x\rangle$$

2) Wavefunction can be decomposed into linear combination of eigenstates

$$|\psi\rangle = \sum_x c_x |x\rangle$$

3) An expectation value can be calculated as follows:

$$\langle\psi|H|\psi\rangle = \sum_{x,y} c_y^* c_x \langle y|H|x\rangle = \sum_{x,y} c_y^* c_x \overset{\text{eigenvalue eq}}{f(x)} \langle y|x\rangle = \sum_x |c_x|^2 f(x)$$

4) Superscript notation: we are considering multiple eigenvectors for the same eigenvalue

$$A|\lambda_j^k\rangle = \lambda_j |\lambda_j^k\rangle$$

5) Probability for a given state x:

$$\langle x|\psi\rangle = c_x$$

6) Expectation Value for any Hermitian Operator:

$$\langle A \rangle_\psi = \sum_{j,k} |\langle \lambda_j^k | \psi \rangle|^2 \lambda_j = \langle \psi | A | \psi \rangle$$

7) Linearity Trick (A has multiple terms)  $A = \frac{1}{2}Z \otimes I \otimes X - 3I \otimes Y \otimes Y$

$$\langle \psi | A | \psi \rangle = \langle \psi | (\frac{1}{2}Z \otimes I \otimes X - 3I \otimes Y \otimes Y) | \psi \rangle = \frac{1}{2} \langle \psi | Z \otimes I \otimes X | \psi \rangle - 3 \langle \psi | I \otimes Y \otimes Y | \psi \rangle$$

8) Eigenvectors of each component tensored are the states of the tensored operator (we can find eigenvectors and eigenvalues for multiplication)

Example: $Z \otimes I \otimes X$	Eigenvectors	Eigenvalues	Eigenvectors	Values
	$ 0\rangle \otimes  0\rangle \otimes  +\rangle$	$1 \cdot 1 \cdot 1 = 1$	$ 1\rangle \otimes  0\rangle \otimes  +\rangle$	$-1 \cdot 1 \cdot 1 = -1$
	$ 0\rangle \otimes  0\rangle \otimes  -\rangle$	$1 \cdot 1 \cdot -1 = -1$	$ 1\rangle \otimes  0\rangle \otimes  -\rangle$	$-1 \cdot 1 \cdot -1 = 1$
	$ 0\rangle \otimes  1\rangle \otimes  +\rangle$	$1 \cdot -1 \cdot 1 = -1$	$ 1\rangle \otimes  1\rangle \otimes  +\rangle$	$-1 \cdot -1 \cdot 1 = 1$
	$ 0\rangle \otimes  1\rangle \otimes  -\rangle$	$1 \cdot -1 \cdot -1 = 1$	$ 1\rangle \otimes  1\rangle \otimes  -\rangle$	$-1 \cdot -1 \cdot -1 = -1$

9) There is a set of pauli matrices, called the change of basis operator, for which any given probability representation for an eigenstate can be transformed to 000

$$\text{Example: } Z \otimes X \otimes Z \quad \text{eigenstate: } |0\rangle|+\rangle|0\rangle \quad \text{Probability: } |\langle 0| \langle +| \langle 0| \psi \rangle|^2$$

$$\text{Since } |0\rangle|+\rangle|0\rangle = (I \otimes H \otimes I) |0\rangle|0\rangle|0\rangle,$$

$$|\langle 0| \langle +| \langle 0| \psi \rangle|^2 = |\langle 0| \langle 0| \langle 0| (I \otimes H \otimes I)^\dagger |\psi \rangle|^2$$

$$\rightarrow (I \otimes H \otimes I) |\psi \rangle = \overset{\text{(hermitian)}}{(I \otimes H \otimes I)^\dagger} |\psi \rangle \quad \text{For any eigenvector: } |\lambda_A \rangle = (I \otimes H \otimes I) |\lambda_C \rangle$$

→ For any tensor product of pauli matrices, we can create a change of basis operator which maps every eigenstate to an eigenstate of the computational basis

$$\rightarrow \text{Example: prepare } (I \otimes H \otimes I) |\psi \rangle. |000\rangle \text{ represents } |0\rangle|+\rangle|0\rangle$$

10) Goal:  $|\langle \lambda_A | \psi \rangle|^2$

$$\langle \lambda_C | (I \otimes H \otimes I)^\dagger |\psi \rangle = \langle \lambda_A | (I \otimes H \otimes I) (I \otimes H \otimes I)^\dagger |\psi \rangle = \langle \lambda_A | \psi \rangle$$

11) Each eigenvalue doesn't need to be run individually. We prepare the state and calculate relative frequencies for a great enough number of shots

12) Each individual matrix has a change of basis matrix. These matrices can be tensored together to form a deterministic way of switching to the computational basis

$$Z \rightarrow Y: SH \quad Y \rightarrow Z: (SH)^\dagger$$

$$Z \rightarrow X: H \quad X \rightarrow Z: (H)^\dagger$$

$$Z \rightarrow Z: I \quad I \rightarrow Z: I$$

$$\text{Example: } \langle \psi | Z \otimes I \otimes X | \psi \rangle$$

$$\hookrightarrow \text{prepare: } (I \otimes I \otimes H) |\psi \rangle$$

13) Preparation and Measurement are done on a quantum computer. Energy Estimation and Parameter Minimization done on a classical computer

14) Fermionic Hamiltonians --> qubit hamiltonians to be run on quantum computers

# Eigenvalue Equation

$$H|x\rangle = f(x)|x\rangle$$

$|x\rangle$  = computational basis state

## Eigenvectors

Since we are building this matrix from the tensor product of Z's and I's,  $|x\rangle$  ( $|0\rangle, |1\rangle$  for 1 qubit) are guaranteed to be eigenstates

## Wavefunction

→ Can be decomposed into linear combination of basis states

$$|\psi\rangle = \sum_x c_x |x\rangle \quad \sim c_x \text{ is complex coefficient associated w/ each state}$$

## Alternate Probability relation

→ typically, the probability of  $|x\rangle$  is  $|c_x|^2$

$$\langle x|\psi\rangle = \langle x|\sum_y c_y |y\rangle = \sum_y c_y \langle x|y\rangle = 0 \text{ for } y \neq x = c_x$$

1 for  $y=x$

$$\rightarrow \text{Thus, } |c_x|^2 = |\langle x|\psi\rangle|^2$$

## Expectation Value Derivation

$$\langle\psi|H|\psi\rangle =$$

$$\sum_y c_y^* \langle y|H \sum_x c_x |x\rangle =$$

$$\sum_{x,y} c_x c_y^* \langle y|H|x\rangle =$$

$$\sum_{x,y} c_x c_y^* f(x) \langle y|x\rangle =$$

$$\text{Eigenvalue? } \langle y|H = f(x) \langle y|$$

$$\sum_x c_x c_x^* f(x) =$$

same trick as above

$$\sum_x |c_x|^2 f(x) = \sum_x |\langle x|\psi\rangle|^2 f(x)$$

↳ agrees with statistical interpretation (sum of prob. value)

→ For any observable, we are guaranteed to find a set of eigenvectors with attached real eigenvalues. The same arguments from the last page would still apply for any eigenbasis, I chose  $|x\rangle$ , similar to the book, since optimization problems use  $I \otimes Z \otimes I \otimes Z \dots$

Similar, for any observable:  $\langle A \rangle_\psi = \sum_{j,k} |\langle \lambda_j^k | \psi \rangle|^2 \lambda_j$

$$= \langle \psi | A | \psi \rangle$$

### Variational Principle

$$\langle A \rangle_\psi = \sum_{j,k} |\langle \lambda_j^k | \psi \rangle|^2 \lambda_j \geq \sum_{j,k} |\langle \lambda_j^k | \psi \rangle|^2 \lambda_0$$

$= \lambda_0$

Sum of all probabilities  
 $\sum_{j,k} |\langle \lambda_j^k | \psi \rangle|^2 = 1$

### Linearity

Assume a matrix  $A = \frac{1}{2} Z \otimes I \otimes X - 3 I \otimes Y \otimes Y$

### Additive

$$\begin{aligned} \langle \psi | A | \psi \rangle &= \langle \psi | (\frac{1}{2} Z \otimes I \otimes X - 3 I \otimes Y \otimes Y) | \psi \rangle \\ &= \frac{1}{2} \langle \psi | (Z \otimes I \otimes X) | \psi \rangle - 3 \langle \psi | I \otimes Y \otimes Y | \psi \rangle \end{aligned}$$

### Tensor Product

$$(Z \otimes I \otimes X) |\lambda\rangle = \lambda |\lambda\rangle$$

$$|\lambda\rangle = |\lambda_Z\rangle \otimes |\lambda_I\rangle \otimes |\lambda_X\rangle \text{ for all states of } Z, I, X$$

$$\lambda = \lambda_Z \cdot \lambda_I \cdot \lambda_X \text{ for all values of } Z, I, X$$

# Transformations

→ Change of basis operator, for any tensor product of Pauli matrices transforms the eigenstates of our Hamiltonian into computational basis states

Observable:  $Z \otimes X \otimes Z$  eigenstate:  $|0\rangle \otimes |+\rangle \otimes |0\rangle$

$$|0\rangle \otimes |+\rangle \otimes |0\rangle = (I \otimes H \otimes I)(|0\rangle \otimes |0\rangle \otimes |0\rangle)$$

→ this is true for any state

$$\text{inversely, } \langle 0| \otimes \langle +| \otimes \langle 0| = \langle 0| \otimes \langle 0| \otimes \langle 0| (I \otimes H \otimes I)^\dagger$$

$\searrow$   
 $\langle x|$

$|\lambda_A\rangle$  - our basis  $|\lambda_C\rangle$  - computational basis

$$|\lambda_A\rangle = (I \otimes H \otimes I) |\lambda_C\rangle \quad I \otimes H \otimes I \text{ is change of basis operator for } Z \otimes X \otimes Z$$

$$|\langle \lambda_A | \psi \rangle|^2 = |\langle 0 | \langle + | \langle 0 | \psi \rangle|^2 = |\langle 0 | \langle 0 | \langle 0 | (I \otimes H \otimes I)^\dagger | \psi \rangle|^2$$

(for example)

→ Prepare (change of basis)  $|\psi\rangle$  and measure in computational basis  
↳ given the result vary  $|\psi\rangle$

Probability Equivalence

$$\langle \lambda_C | (I \otimes H \otimes I)^\dagger | \psi \rangle =$$

$$\langle \lambda_A | (I \otimes H \otimes I) (I \otimes H \otimes I)^\dagger | \psi \rangle =$$

$$\langle \lambda_A | \psi \rangle$$

↳ probability of any eigenstate of our system

## Qiskit

→ Transforming Max 3SAT Problem

\* Convert to ising: 3 terms: linear, quadratic, constant

$$H = - \sum_{j,k} J_{jk} z_j z_k - \sum_j h_j z_j - \text{constant}$$

Linear:  $z_j$  on  $j^{\text{th}}$  qubit  $I$  on else

Quadratic: 2  $z$ 's on  $i^{\text{th}}, j^{\text{th}}$  qubit  $I$  on else  
 $z_i z_j$

Constant:  $I$  on every qubit

Ansatz:  $(|\psi\rangle)$  Efficient SU2

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### Multiple runs

Stored the final parameters for the ansatz  
Plugged into the SU2 gate and measurement

→ Did that for 500 runs of VQE