


HW 2.

Laplace's Method, Watson's lemma.

1. Consider the asymptotic behavior of the integral

what if φ 

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$$I(x) = \int_1^3 e^{-x(\frac{9}{t}+t)} dt, \quad x \rightarrow +\infty.$$

(1) Find the leading asymptotic behavior using the Laplace method.

(2) Rewrite the integral using variable $u = 9/t + t - 6$ and find the first two leading terms in the asymptotic expansion (using Watson's lemma).

$$1) I(x) = \int_1^3 e^{x\varphi(t)} dt = \int_1^3 e^{-x(\frac{9}{t}+t)} dt \quad x \rightarrow +\infty$$

$$\varphi(t) = -(\frac{9}{t}+t) \quad t \in [1, 3]$$

$\varphi(t)$ reaches its maximum at $t=3$. $\varphi_{\max} = \varphi(3) = -6$

The leading behavior of $I(x)$ is near $t=3$ as $x \rightarrow +\infty$.

• Step 1:

$$I(x) = \int_1^{3-\varepsilon} e^{-x(\frac{9}{t}+t)} dt + \int_{3-\varepsilon}^3 e^{-x(\frac{9}{t}+t)} dt \quad \varepsilon \rightarrow 0$$

要马过程吗? $\sim \int_{3-\varepsilon}^3 e^{-x(\frac{9}{t}+t)} dt \quad \text{as } x \rightarrow +\infty$

• Step 2:

When $t \in (3-\varepsilon, 3]$, the Taylor expansion of $f(t) = \frac{9}{t} + t$ at

$t=3$ is

$$\begin{aligned} f(t) &= f(3) + f'(3)(t-3) + \sum_{n=2}^{\infty} \frac{(-1)^n n! 9 \cdot 3^{-(n+1)}}{n!} (t-3)^n \\ &= 6 + \sum_{n=2}^{\infty} (-1)^n 3^{-(n-1)} (t-3)^n \end{aligned}$$

$$\text{Thus, } \frac{9}{t} + t \sim 6 - \frac{1}{3}(t-3)^2 + o((t-3)^2)$$

$$\int_{3-\varepsilon}^3 e^{-x(\frac{9}{t}+t)} dt \sim \int_{3-\varepsilon}^3 e^{-x(6 + \frac{1}{3}(t-3)^2)} dt$$

$$= e^{-6x} \int_{3-\varepsilon}^3 e^{-\frac{1}{3}x(t-3)^2} dt$$

$$= e^{-6x} \int_{-\varepsilon}^0 e^{-\frac{1}{3}x t^2} dt$$

$$s = \frac{1}{3}x t^2 \quad = e^{-6x} \cdot \sqrt{\frac{3}{x}} \int_{-\varepsilon}^0 e^{-s^2} ds$$

$$\approx e^{-6x} \cdot \sqrt{\frac{3}{x}} \int_{-\infty}^0 e^{-s^2} ds$$

$$= e^{-6x} \cdot \sqrt{\frac{3\pi}{4x}} \quad x \rightarrow +\infty$$

• Step 3:



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$$\begin{aligned}
 (2) \quad L(x) &= \int_1^3 e^{-x(\frac{9}{t}+t)} dt \\
 &= \int_1^3 e^{-x(\frac{9}{t}+t-6)} e^{-6x} dt \\
 &= e^{-6x} \int_1^3 e^{-ux} dt \\
 u &= \frac{9}{t}+t-6 \Rightarrow t = \frac{u+6 \pm \sqrt{u^2+12u}}{2} \quad u \in [0, 4]
 \end{aligned}$$

$$\text{Since } t \in (1, 3) \quad t = \frac{u+6 - \sqrt{u^2+12u}}{2}$$

$$L(x) = \int_0^4 e^{-ux} d\left(\frac{u+6 - \sqrt{u^2+12u}}{2}\right)$$

$$= \int_0^4 \frac{1}{2} \left(1 - \frac{u+6}{\sqrt{u^2+12u}}\right) e^{-ux} du$$

$$= \frac{1}{2} \int_0^4 e^{-ux} du - \int_0^4 \frac{u+6}{\sqrt{u^2+12u}} e^{-ux} du$$

$$= \frac{1}{2x} (e^{-4x} - 1) - \int_0^4 \frac{u+6}{\sqrt{u^2+12u}} e^{-ux} du$$

$$\begin{aligned}
 \frac{u+6}{\sqrt{u^2+12u}} &= \frac{u+6}{2\sqrt{3} \cdot \sqrt{u} \cdot \sqrt{\frac{u}{12}+1}} \\
 &= \frac{\sqrt{u}}{2\sqrt{3} \sqrt{1+\frac{u}{12}}} + \frac{\sqrt{3}}{\sqrt{u} \sqrt{1+\frac{u}{12}}}
 \end{aligned}$$

$$\begin{aligned}
 \text{The Taylor expansion of } \frac{1}{\sqrt{1+\frac{u}{12}}} &= \sum_{n=0}^{\infty} (-1)^n \frac{(2n-1)!!}{2^{n+1} \cdot 12^n \cdot n!} u^n \\
 &= \sum_{n=0}^{\infty} (-1)^n \frac{1}{12^n} \frac{\Gamma(n+\frac{1}{2})}{n! \Gamma(\frac{1}{2})} u^n
 \end{aligned}$$

Thus,

$$\begin{aligned}
 \frac{u+6}{\sqrt{u^2+12u}} &= \frac{1}{2\sqrt{3}} \cdot u^{\frac{1}{2}} \cdot \sum_{n=0}^{\infty} \frac{(-1)^n}{12^n} \frac{\Gamma(n+\frac{1}{2})}{n! \Gamma(\frac{1}{2})} u^n + \sqrt{3} u^{-\frac{1}{2}} \sum_{n=0}^{\infty} \dots \\
 &= \underbrace{\frac{1}{2\sqrt{3}} \sum_{n=0}^{\infty} \frac{(-1)^n}{12^n} \frac{\Gamma(n+\frac{1}{2})}{n! \Gamma(\frac{1}{2})} u^{\frac{n+1}{2}}}_{\textcircled{1} \quad a_n} + \underbrace{\sqrt{3} \sum_{n=0}^{\infty} \frac{(-1)^n}{12^n} \frac{\Gamma(n+\frac{1}{2})}{n! \Gamma(\frac{1}{2})} u^{n-\frac{1}{2}}}_{\textcircled{2} \quad a_n}
 \end{aligned}$$



By Watson's Lemma.

① $\alpha = \frac{1}{2}, \beta = 1$

$$\begin{aligned} \int_0^4 \frac{u+6}{\sqrt{u^2+12u}} e^{-u\chi} du &\sim \sum_{n=0}^{\infty} \frac{(-1)^n}{12^n} \frac{\Gamma(n+\frac{1}{2})}{n! \Gamma(\frac{1}{2})} \cdot \frac{\Gamma(n+\frac{3}{2})}{\chi^{n+\frac{3}{2}}} \\ &= \sum_{n=0}^{\infty} \frac{(-1)^n}{12^n} \frac{\Gamma(n+\frac{1}{2}) \Gamma(n+\frac{3}{2})}{n! \Gamma(\frac{1}{2}) \chi^{n+\frac{3}{2}}} \\ &= \sum_{n=0}^{\infty} \frac{(-1)^n}{12^n} \frac{(n+\frac{1}{2}) (\Gamma(n+\frac{1}{2}))^2}{n! \Gamma(\frac{1}{2}) \chi^{n+\frac{3}{2}}} \end{aligned}$$

② $\alpha = -\frac{1}{2}, \beta = 1$

$$\begin{aligned} \int_0^4 \frac{u+6}{\sqrt{u^2+12u}} e^{-u\chi} du &\sim \sum_{n=0}^{\infty} \frac{(-1)^n}{12^n} \frac{\Gamma(n+\frac{1}{2})}{n! \Gamma(\frac{1}{2})} \cdot \frac{\Gamma(n+\frac{1}{2})}{\chi^{n+\frac{1}{2}}} \\ &= \sum_{n=0}^{\infty} \frac{(-1)^n}{12^n} \frac{(\Gamma(n+\frac{1}{2}))^2}{n! \Gamma(\frac{1}{2}) \chi^{n+\frac{1}{2}}} \end{aligned}$$

Therefore,

$$\begin{aligned} \mathcal{L}(\chi) &= \frac{1}{2\chi} (e^{-4\chi} - 1) - \int_0^4 \frac{u+6}{\sqrt{u^2+12u}} e^{-u\chi} du \\ &\sim \frac{1}{2\chi} (e^{-4\chi} - 1) - \sum_{n=0}^{\infty} \frac{(-1)^n}{12^n} \frac{(n+\frac{1}{2}) (\Gamma(n+\frac{1}{2}))^2}{n! \Gamma(\frac{1}{2}) \chi^{n+\frac{3}{2}}} \cdot \chi^{-1} \\ &\quad - \sum_{n=0}^{\infty} \frac{(-1)^n}{12^n} \frac{(\Gamma(n+\frac{1}{2}))^2}{n! \Gamma(\frac{1}{2}) \chi^{n+\frac{1}{2}}} \\ &= \frac{1}{2\chi} (e^{-4\chi} - 1) - \left(\sum_{n=0}^{\infty} \left(\frac{n+\frac{1}{2}}{\chi} + 1 \right) \frac{(-1)^n}{12^n} \frac{(n+\frac{1}{2}) (\Gamma(n+\frac{1}{2}))^2}{n! \Gamma(\frac{1}{2}) \chi^{n+\frac{1}{2}}} \right) \end{aligned}$$

The first two leading terms are:

$$\frac{1}{2\chi} (e^{-4\chi} - 1) - \frac{(1+2\chi) \Gamma(\frac{1}{2})}{4\chi^{\frac{3}{2}}} = \frac{1}{2\chi} (e^{-4\chi} - 1) - \frac{\sqrt{\pi} (1+2\chi)}{4\chi^{\frac{3}{2}}}$$



Stationary phase.

2. Use the stationary phase method to find the leading asymptotic behavior of the integral

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$$I(x) = \int_0^\pi \cos(x \sin t - t) dt, \quad x \rightarrow +\infty.$$

Hint: (1) $\cos(x \sin t - t)$ is the real part of $e^{i(x \sin t - t)}$; (2) $\int_0^\infty e^{-iz^2} dz = \frac{\sqrt{\pi}}{2} e^{-i\pi/4}$.

Since $I(x) = \int_0^\pi \cos(x \sin t - t) dt$ is the real part of

$$I_0(x) = \int_0^\pi e^{i(x \sin t - t)} dt = \int_0^\pi e^{-it} e^{ix \sin t} dt$$

We consider $I_0(x)$ first.

$$\varphi(t) = \sin t, \quad f(t) = e^{-it}$$

$$\varphi'(t) = \cos t \quad \text{as } x \rightarrow \infty$$

$$\varphi'(\frac{\pi}{2}) = 0 \quad \varphi''(\frac{\pi}{2}) \neq 0$$

Thus $t = \frac{\pi}{2}$ is a stationary point

$$I_0(x) = \underbrace{\int_0^{\frac{\pi}{2}-\epsilon} e^{i(x \sin t - t)} dt}_{I_1} + \underbrace{\int_{\frac{\pi}{2}-\epsilon}^{\frac{\pi}{2}+\epsilon} e^{i(x \sin t - t)} dt}_{I_2} + \underbrace{\int_{\frac{\pi}{2}+\epsilon}^\pi e^{i(x \sin t - t)} dt}_{I_3}$$

$$\textcircled{1} \quad I_2 = \int_{\frac{\pi}{2}-\epsilon}^{\frac{\pi}{2}+\epsilon} e^{-it} e^{ix \sin t} dt$$

$$\sim \int_{\frac{\pi}{2}-\epsilon}^{\frac{\pi}{2}+\epsilon} e^{-it} e^{ix \left(\varphi(\frac{\pi}{2}) + \frac{\varphi''(\frac{\pi}{2})}{2} (t - \frac{\pi}{2})^2 \right)} dt$$

$$= \int_{\frac{\pi}{2}-\epsilon}^{\frac{\pi}{2}+\epsilon} e^{-it} e^{ix \left(1 - \frac{1}{2} (t - \frac{\pi}{2})^2 \right)} dt$$

$$\stackrel{f(a)}{=} e^{-\frac{\pi}{2}i} \int_{\frac{\pi}{2}-\epsilon}^{\frac{\pi}{2}+\epsilon} e^{ix} e^{-\frac{1}{2}ix(t-\frac{\pi}{2})^2} dt$$

$$= e^{i(x-\frac{\pi}{2})} \int_{\frac{\pi}{2}-\epsilon}^{\frac{\pi}{2}+\epsilon} e^{-\frac{1}{2}ix(t-\frac{\pi}{2})^2} dt$$

$$u = t - \frac{\pi}{2} = e^{i(x-\frac{\pi}{2})} \int_{-\epsilon}^{\epsilon} e^{-\frac{1}{2}ix u^2} du$$



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$$\sim e^{i(\chi - \frac{\pi}{2})} \int_{-\infty}^{+\infty} e^{-\frac{1}{2} i \chi u^2} du$$

$$= 2e^{i(\chi - \frac{\pi}{2})} \int_0^{+\infty} e^{-\frac{1}{2} i \chi u^2} du$$

$$z = \sqrt{\frac{1}{2} \chi} u = \sqrt{\frac{2}{\chi}} e^{i(\chi - \frac{\pi}{2})} \int_0^{+\infty} e^{-i z^2} dz$$

$$= \sqrt{\frac{2\pi}{\chi}} e^{i(\chi - \frac{1}{4}\pi)}$$

② By generalized Riemann - Lebesgue Theorem,

$$I_1 = \int_0^{\frac{\pi}{2}-\varepsilon} e^{-it} e^{i\chi \sin t} dt \sim \frac{e^{-it}}{i\chi \cos t} e^{i\chi \sin t} \Big|_0^{\frac{\pi}{2}-\varepsilon} \quad \chi \rightarrow +\infty$$

$$= O\left(\frac{1}{\chi}\right) \ll O\left(\frac{1}{\sqrt{\chi}}\right) \quad \chi \rightarrow +\infty$$

$$I_3 = \int_{\frac{\pi}{2}+\varepsilon}^{\pi} e^{-it} e^{i\chi \sin t} dt \sim \frac{e^{-it}}{i\chi \cos t} e^{i\chi \sin t} \Big|_{\frac{\pi}{2}+\varepsilon}^{\pi} \quad \chi \rightarrow \infty$$

$$= O\left(\frac{1}{\chi}\right) \ll O\left(\frac{1}{\sqrt{\chi}}\right)$$

No need to prove

Therefore, $I_0(\chi) = \sqrt{\frac{2\pi}{\chi}} e^{i(\chi - \frac{1}{4}\pi)}$

$I(\chi)$ is the real part of $I_0(\chi)$

$$I(\chi) = \sqrt{\frac{2\pi}{\chi}} \cos\left(\chi - \frac{1}{4}\pi\right)$$



14.5/15 Perturbation methods.

3. (1) Compute the first four coefficients in the perturbation series (i.e., $O(1)$, $O(\epsilon)$, $O(\epsilon^2)$ and $O(\epsilon^3)$ terms) to the initial-value problem

$$y' = 2y + 4\epsilon xy, \quad y(0) = 1.$$

(2) Find the exact solution.

(3) Use some software, e.g., matlab, to plot and compare the exact solution and the n -term perturbation expansion for the solution, $n = 1, 2, 3, 4$ (i.e., y_0 , $y_0 + \epsilon y_1$, $y_0 + \epsilon y_1 + \epsilon^2 y_2$ and $y_0 + \epsilon y_1 + \epsilon^2 y_2 + \epsilon^3 y_3$) on $x \in [0, 3]$ when $\epsilon = 0.1$.

1) The problem is

$$\begin{cases} \frac{dy}{dx} = 2y + 4\epsilon xy \\ y(0) = 1 \end{cases} \quad (1)$$

Look for a solution with the form $y(x) = \sum_{n=0}^{\infty} \epsilon^n y_n(x) \quad (2)$

where $y_n(x)$ are to be determined.

Substituting (2) into (1). we have

$$\frac{dy_0}{dx} + \epsilon \frac{dy_1}{dx} + \epsilon^2 \frac{dy_2}{dx} + \dots \sim 2y_0 + (4xy_0 + 2y_1)\epsilon + (4xy_1 + 2y_2)\epsilon^2 + (4xy_2 + 2y_3)\epsilon^3 \quad \epsilon \rightarrow 0$$

Compare coefficient functions:

$$\underline{\epsilon^0 \text{ terms:}} \quad \frac{dy_0}{dx} = 2y_0$$

$$\underline{\epsilon^1 \text{ terms:}} \quad \frac{dy_1}{dx} = 4xy_0 + 2y_1$$

$$\underline{\epsilon^2 \text{ terms:}} \quad \frac{dy_2}{dx} = 4xy_1 + 2y_2$$

...

The initial condition becomes:

$$y_0(0) + \epsilon y_1(0) + \epsilon^2 y_2(0) + \epsilon^3 y_3(0) + \dots \sim 1 \quad \epsilon \rightarrow 0$$

which implies

$$y_0(0) = 1$$

$$y_1(0), y_2(0), y_3(0), \dots = 0$$



- Therefore, the zero order solution satisfies

$$\begin{cases} \frac{dy_0}{dx} = 2y_0 \\ y_0(0) = 1 \end{cases}$$

The solution is $y_0 = e^{2x}$

- The first-order solution satisfies

$$\begin{cases} \frac{dy_1}{dx} = 4x y_0 + 2y_1 = 4x e^{2x} + 2y_1 \\ y_1(0) = 0 \end{cases}$$

The solution is $y_1 = 2x^2 e^{2x}$

- The second-order solution satisfies

$$\begin{cases} \frac{dy_2}{dx} = 4x y_1 + 2y_2 = 8x^3 e^{2x} + 2y_2 \\ y_2(0) = 0 \end{cases}$$

The solution is $y_2 = 2x^4 e^{2x}$

- The third-order solution satisfies

$$\begin{cases} \frac{dy_3}{dx} = 4x y_2 + 2y_3 = 8x^5 e^{2x} + 2y_3 \\ y_3(0) = 0 \end{cases}$$

The solution is $\frac{4}{3} x^6 e^{2x}$ ✓

...

Therefore, the solution is

$$y \sim e^{2x} + 2\varepsilon x^2 e^{2x} + 2\varepsilon^2 x^4 e^{2x} + \frac{4}{3} \varepsilon^3 x^6 e^{2x} \dots \quad \varepsilon \rightarrow 0$$

which means $y^{(n)}(x) = \sum_{n=0}^{\infty} \varepsilon^n y_n(x)$ is an approximation to the exact solution.

$$(2) \begin{cases} y' = 2y + 4\varepsilon xy \\ y(0) = 1 \end{cases}$$

$$\Leftrightarrow \frac{dy}{dx} = (2 + 4\varepsilon x) y$$



$$\Leftrightarrow \frac{1}{y} dy = (2 + 4\varepsilon x) dx$$

$$\Leftrightarrow \ln y = 2x + 2\varepsilon x^2 + C$$

$$\Leftrightarrow y = e^{2x + 2\varepsilon x^2 + C'}$$

Since $y(0) = 1 \Rightarrow C' = 0$

Thus, the exact solution is $y = e^{2x + 2\varepsilon x^2}$, $\varepsilon \rightarrow 0$

(3)

