



# Effects of SOC-dependent electrolyte viscosity on performance of vanadium redox flow batteries



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## HIGHLIGHTS

- The correlations of electrolyte viscosity and SOC are obtained.
- Effect of SOC-dependent electrolyte viscosity is considered in this model.
- This model enables a more realistic simulation of variable distributions.
- It provides accurate estimations of pumping work and system efficiency.

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## ABSTRACT

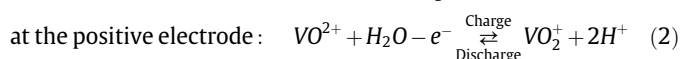
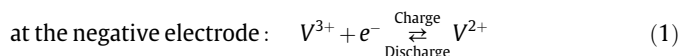
The viscosity of the electrolyte in vanadium redox flow batteries (VRFBs) varies during charge and discharge as the concentrations of acid and vanadium ions in the electrolyte continuously change with the state of charge (SOC). In previous VRFB models, however, the electrolyte has been treated as a constant-viscosity solution. In this work, a mass-transport and electrochemical model taking account of the effect of SOC-dependent electrolyte viscosity is developed. The comparison between the present model and the model with the constant-viscosity simplification indicates that the consideration of the SOC-dependent electrolyte viscosity enables (i) a more realistic simulation of the distributions of overpotential and current density in the electrodes, and (ii) more accurate estimations of pumping work and the system efficiency of VRFBs.

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## 1. Introduction

The all-vanadium redox flow battery (VRFB) was initially proposed by NASA in mid-1970s and developed by Skyllas-Kazacos et al. in the 1980s, using the V(II)/V(III) and V(IV)/V(V) redox couples in sulfuric acid solution as the anolyte and catholyte, respectively [1–3]. This type of battery is particularly suitable for large-scale storage of intermittent power generated from solar cells and wind turbines due to its unique features, including independent scaling of energy and power ratings, a long cycle life, quick response, a low cost, and the tolerance to deep discharge without any risk of damage [4–7].

A typical single vanadium redox flow cell consists of two porous carbon electrodes and two circulating electrolyte solutions separated by an ion-exchange membrane. The reactions that occur in the electrodes can be expressed as:



The energy conversion between electrical energy and chemical potential occurs at the electrodes once the liquid electrolytes flow through the cell. Protons transfer through the membrane, while electrons transfer through an external load to form an electrical circuit during charge or discharge process. As a key component of the battery, the electrolyte needs to be highly concentrated to all the vanadium ions, stable and less viscous, as these properties are closely related to the capacity, power and the parasitic energy loss of the battery [8,9]. The concentration of each vanadium ion with a certain valence in the electrolyte keeps changing during the charge and discharge cycle, while the total vanadium ions concentration is approximately constant in both the positive and negative electrolytes. To describe the relative amount of vanadium ions with different valences, the state of charge (SOC) is defined as:

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$$\text{SOC} = \left( \frac{c_{V(II)}}{c_{V(\text{total})}} \right)_1 = \left( \frac{c_{V(V)}}{c_{V(\text{total})}} \right)_2 \quad (3)$$

where  $c_{V(II)}$  and  $c_{V(V)}$  are the molar concentrations of  $V^{2+}$ ,  $VO_2^+$ ,  $c_{V(\text{total})}$  is the total vanadium ions concentration in a certain electrolyte, and the subscripts 1 and 2 represent the negative and positive electrode, respectively.

Alternatively, SOC can also be defined as:

$$\text{SOC} = \left( 1 - \frac{c_{V(III)}}{c_{V(\text{total})}} \right)_1 = \left( 1 - \frac{c_{V(IV)}}{c_{V(\text{total})}} \right)_2 \quad (4)$$

where  $c_{V(III)}$  and  $c_{V(IV)}$  are the molar concentrations of  $V^{3+}$  and  $VO^{2+}$ .

In addition to vanadium ions at different oxidation states, protons and sulphate ions also appear in the battery. In the normal operation, these spectator ions do not take part in the reactions, but they are essential for maintaining conservation of mass and the charge balance in electrolytes. Taking the charge process as example, when the anolyte and catholyte components change from SOC = 0 to SOC = 1, as shown in Fig. 1, the proton concentration increases by 0.5 M in both the anolyte and catholyte if the secondary ionization of sulfuric acid is ignored.

The viscosity of a solution is related to the components and their respective concentrations [10,11]. For a VRFB, the variation in the electrolyte component concentrations will result in a continuous change in the electrolyte viscosity during the charge and discharge cycle [12–16]. Such a SOC-dependent viscosity is expected to cause different effects, as opposed to those resulting from the constant viscosity assumption [17–23], on the velocity field, the concentration and overpotential distributions, and the battery performance. The objective of this work is to study the effects of SOC-dependent electrolyte viscosity by developing a new model as presented below.

## 2. Mathematical model

Fig. 2 shows the schematic illustration of a single vanadium redox flow cell, which consists of a Nafion 115 membrane in the middle, two graphite felts as electrodes on two sides of the membrane and two current collectors. Sulfuric acid solution containing vanadium ions is used as the positive and negative electrolytes, which are stored in respective reservoirs and pumped to the porous electrodes of the cell. The geometric parameters of the VRFB are listed in Table 1.

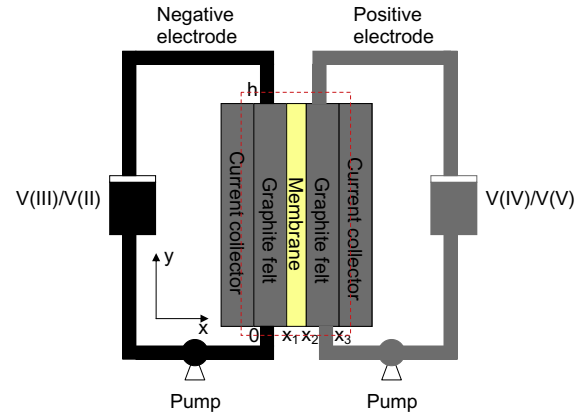


Fig. 2. Schematic of the modeling domain of a VRFB.

### 2.1. Simplifications and assumptions

We consider a two-dimensional system with the x-axis located in the lower boundary of the battery and the y-axis located at the interface between the negative current collector and graphite felt. Other simplifications and assumptions used in the present work are as follows:

- (1) As indicated in Eqs. (3) and (4), in general, SOC varies with time. However, in the case when the reservoir is sufficiently big, the change in SOC is relatively small. As a result, the transient charge and discharge processes can be simplified as a steady-state process.
- (2) An isothermal condition is assumed in the entire domain.
- (3) The membrane is impermeable to all ions and species, except for protons.
- (4) Possible side reactions, such as hydrogen and oxygen evolutions, are not considered.

### 2.2. Governing equations

#### 2.2.1. Transport through the porous electrode

To model the mass transport in the porous electrode, we express the molar flux of each species,  $\tilde{N}_i$  (with  $i$  representing  $V^{2+}$ ,  $V^{3+}$ ,  $VO^{2+}$ ,  $VO_2^+$ ,  $H^+$ ,  $HSO_4^-$ ), in terms of the modified Nernst-Planck equation as:

$$\tilde{N}_i = -D_i^{\text{eff}} \nabla c_i - F z_i c_i \nabla \phi_i + \tilde{u} c_i \quad (5)$$

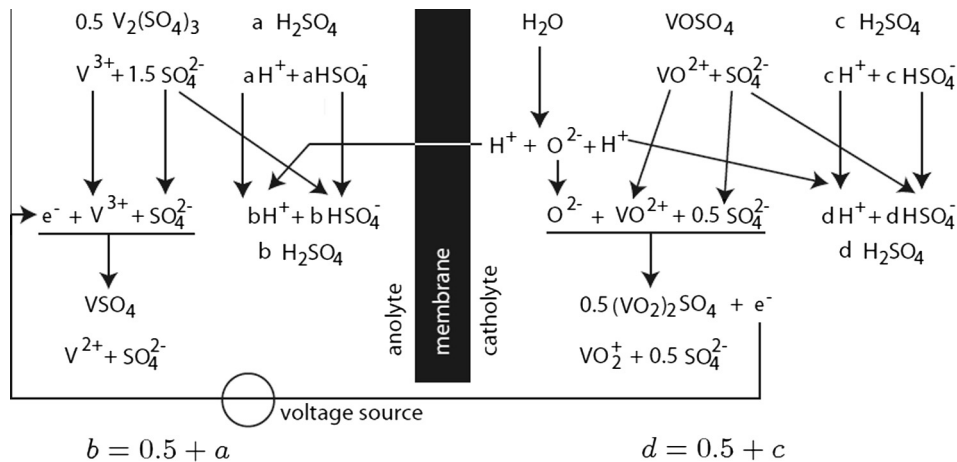


Fig. 1. Illustration of the anolyte and catholyte component variations from SOC = 0 to SOC = 1.

**Table 1**  
Geometric and operating parameters.

Parameters	Symbols	Value	Unit	Reference
Specific surface area	$A_r$	$2 \times 10^5$	$\text{m}^{-1}$	[29]
Electrode thickness	$\delta_e$	$3 \times 10^{-3}$	m	
Polymer electrolyte membrane	$\delta_{\text{mem}}$	$1.8 \times 10^{-4}$	m	[29]
Electrode height	$h$	$2 \times 10^{-2}$	m	
Electrode width	$w$	$2 \times 10^{-2}$	m	
Operating temperature	$T$	300	K	
Electrode porosity	$\varepsilon_0$	0.7	1	
Electrode fiber diameter	$d_f$	17.6	$\mu\text{m}$	
Initial $\text{V}^{2+}$ ion concentration	$c_{\text{V(II)}}^0$	1050	$\text{mol m}^{-3}$	
Initial $\text{V}^{3+}$ ion concentration	$c_{\text{V(III)}}^0$	450	$\text{mol m}^{-3}$	
Initial $\text{VO}^{2+}$ ion concentration	$c_{\text{V(IV)}}^0$	450	$\text{mol m}^{-3}$	
Initial $\text{VO}_2^+$ ion concentration	$c_{\text{V(V)}}^0$	1050	$\text{mol m}^{-3}$	
$\text{HSO}_4^-$ ion concentration	$c_{\text{HSO}_4}^0$	2000	$\text{mol m}^{-3}$	
Initial $\text{H}^+$ ion concentration	$c_{\text{H}^+}^0$	2000	$\text{mol m}^{-3}$	
Current density	$I$	400	$\text{A m}^{-2}$	
Electrolyte volumetric flow rate	$Q$	0.3	$\text{ml s}^{-1}$	
Pump efficiency	$\psi_{\text{pump}}$	0.9	1	

where  $c_i$  represents the concentration of each species,  $F$  is the Faradaic constant,  $z_i$  and  $\xi_i$  are the valence and ionic mobility of each species,  $\phi_l$  is the ionic potential in the electrolyte, and

$$D_i^{\text{eff}} = D_i \varepsilon^{1.1} (1 + 1.46 \times 10^{-3} \text{Pe}^2) \quad (6)$$

is the effective diffusion coefficient of each species which includes the effect of flow dispersion, with  $\varepsilon$  and  $D_i$  representing, respectively, the porosity of the electrode and the free-space diffusivity of each species [24]. The superficial electrolyte velocity  $\vec{u}$  can be expressed using Darcy's law as:

$$\rho \frac{v}{K} \vec{u} = -\nabla p \quad (7)$$

where  $K$  is the permeability of the porous electrode,  $v$  is the kinematic viscosity of the electrolyte and  $p$  is the liquid pressure. The liquid pressure is determined from an overall mass balance as:

$$\nabla \cdot \vec{u} = 0 \quad (8)$$

The flow of charged species gives the current in electrolyte solution:

$$\vec{i}_l = z_i F \vec{N}_i \quad (9)$$

The electrolyte is considered to be electrically neutral:

$$\sum_i z_i c_i = 0 \quad (10)$$

Combining Eqs. 5, 9 and 10, we can express the total current density in the electrolyte as:

$$\vec{i} = \sum_i \vec{i}_i = -F \sum_i z_i D_i^{\text{eff}} \nabla c_i - F^2 \sum_i z_i^2 c_i \xi_i \nabla \phi_l \quad (11)$$

The conservation of species can be written as:

$$\nabla \cdot \vec{N}_i = \dot{R}_i \quad (12)$$

where  $\dot{R}_i$  is the generation rate for species  $i$  [17].

Combining Eqs. (5) and (12) gives:

$$\nabla \cdot (\vec{u} c_i) = \nabla \cdot (D_i^{\text{eff}} \nabla c_i) + \nabla \cdot (F z_i c_i \xi_i \nabla \phi_l) + \dot{R}_i \quad (13)$$

Due to the conservation of charge, the charge entering the electrolyte solution is balanced by the charge leaving the solid phase:

$$\nabla \cdot \vec{i} = \nabla \cdot \vec{i}_l = -\nabla \cdot \vec{i}_s \quad (14)$$

where  $\vec{i}_l$  is the current density in the electrolyte solution and  $\vec{i}_s$  is the current density in the solid matrix.

The electronic potential in the solid matrix is given by Ohm's law:

$$-\sigma_s^{\text{eff}} \nabla^2 \phi_s = \nabla \cdot \vec{i}_s \quad (15)$$

where  $\sigma_s^{\text{eff}}$  and  $\phi_s$  are the effective conductivity and electronic potential of the solid matrix.

## 2.2.2. Transport through the membrane

As proton is the only mobile ion, the current conservation equation is:

$$\vec{N}_{\text{H}^+} = -\frac{\sigma_{\text{mem}}}{F} \nabla \phi_{\text{mem}} \quad (16)$$

where  $\sigma_{\text{mem}}$  and  $\phi_{\text{mem}}$  are the conductivity and electric potential of the ion-exchange membrane. The values of the constants related to mass and charge transport can be found in Table 2. It is worth noting that the constant value of electrolyte viscosity used for comparison represents the electrolyte viscosity of negative half-cell at the SOC of 1.0.

## 2.2.3. Reaction kinetics

The electrochemical reactions taking place at the solid surfaces of the porous graphite electrode can be expressed using the Butler-Volmer Equation as [22]:

$$j_1 = \varepsilon A_v F k_1 (c_{\text{V(III)}}^s)^{\alpha_{-1}} (c_{\text{V(II)}}^s)^{\alpha_{+1}} \left\{ \exp\left(\frac{\alpha_{+1} F \eta_1}{RT}\right) - \exp\left(-\frac{\alpha_{-1} F \eta_1}{RT}\right) \right\} \quad (17)$$

and

$$j_2 = \varepsilon A_v F k_2 (c_{\text{V(V)}}^s)^{\alpha_{-2}} (c_{\text{V(IV)}}^s)^{\alpha_{+2}} \left\{ \exp\left(\frac{\alpha_{+2} F \eta_2}{RT}\right) - \exp\left(-\frac{\alpha_{-2} F \eta_2}{RT}\right) \right\} \quad (18)$$

Eqs. (17) and (18) are for the negative and positive electrodes, respectively, where  $A_v$  is specific surface area of the porous electrode,  $k_1$  and  $k_2$  are the standard rate constants for negative and positive electrochemical reactions,  $c_i^s, i \in \{\text{V}^{2+}, \text{V}^{3+}, \text{VO}^{2+}, \text{VO}_2^+\}$  are the vanadium-species concentrations at the liquid–solid interfaces of the porous region,  $\alpha_+$  and  $\alpha_-$  are the anodic and cathodic transfer coefficients,  $\eta_1$  and  $\eta_2$  are the overpotentials in the negative and positive electrodes, respectively. The electrochemical constants and coefficients are listed in Table 3.

The interface concentrations  $c_i^s$  can be related to the bulk concentrations  $c_i$  by balancing the electrochemical reaction rate with the rate of mass transfer of the reactants to (or from) the electrode surface. For the negative electrode during discharge, the balance is:

**Table 2**  
Transport properties.

Parameters	Symbols	Value	Unit	Reference
V <sup>2+</sup> ion diffusivity in electrolyte	$D_{V(II)}$	$2.4 \times 10^{-10}$	$m^2 s^{-1}$	[30]
V <sup>3+</sup> ion diffusivity in electrolyte	$D_{V(III)}$	$2.4 \times 10^{-10}$	$m^2 s^{-1}$	[30]
VO <sup>2+</sup> ion diffusivity in electrolyte	$D_{V(IV)}$	$3.9 \times 10^{-10}$	$m^2 s^{-1}$	[30]
VO <sub>2</sub> <sup>+</sup> ion diffusivity in electrolyte	$D_{V(V)}$	$3.9 \times 10^{-10}$	$m^2 s^{-1}$	[30]
Water diffusivity in membrane	$D_{H_2O}^{eff}$	$5.75 \times 10^{-10}$	$m^2 s^{-1}$	[17]
Proton diffusivity in membrane	$D_{H^+}^{eff}$	$3.5 \times 10^{-10}$	$m^2 s^{-1}$	[31]
HSO <sub>4</sub> <sup>-</sup> ion diffusivity in membrane	$D_{HSO_4^-}^{eff}$	$2.2 \times 10^{-10}$	$m^2 s^{-1}$	[17]
Carman-Kozeny constant	$K_{CK}$	5.55		[17]
Constant electrolyte kinematic viscosity	$\nu$	$1.2 \times 10^{-6}$	$m^2 s^{-1}$	
Electrolyte density	$\rho$	1400	$kg m^{-3}$	
Conductivity of the solid phase	$\sigma_m$	1000	$S m^{-1}$	
Conductivity of the membrane	$\sigma_{mem}$	10	$S m^{-1}$	

**Table 3**  
Electrochemistry properties.

Parameters	Symbols	Value	Unit	Reference
Cathodic transfer coefficient for reaction (1)	$\alpha_{-,1}$	0.5		Assumed
Anodic transfer coefficient for reaction (1)	$\alpha_{+,1}$	0.5		Assumed
Cathodic transfer coefficient for reaction (2)	$\alpha_{-,2}$	0.5		Assumed
Anodic transfer coefficient for reaction (2)	$\alpha_{+,2}$	0.5		Assumed
Standard rate constant for reaction (1)	$k_1$	$1.75 \times 10^{-7}$	$m s^{-1}$	[17]
Standard rate constant for reaction (2)	$k_2$	$3.0 \times 10^{-9}$	$m s^{-1}$	[31]
Equilibrium potential for reaction (1)	$E_1^0$	-0.26	V	[22]
Equilibrium potential for reaction (2)	$E_2^0$	1.004	V	[22]

$$k_m(c_{V(II)} - c_{V(II)}^s) = \varepsilon k_1 \left[ c_{V(II)}^s \exp\left(\frac{\alpha_{+,1} F \eta_1}{RT}\right) - c_{V(II)}^s \exp\left(-\frac{\alpha_{-,1} F \eta_1}{RT}\right) \right] \quad (19)$$

and

$$k_m(c_{V(III)} - c_{V(III)}^s) = \varepsilon k_1 \left[ c_{V(III)}^s \exp\left(\frac{\alpha_{-,1} F \eta_1}{RT}\right) - c_{V(III)}^s \exp\left(-\frac{\alpha_{+,1} F \eta_1}{RT}\right) \right] \quad (20)$$

where  $k_m$  is the mass transfer coefficient and can be expressed as [24]:

$$k_m = D_i(2 + 1.534 Re^{0.912})/d_p \quad (21)$$

Combining Eqs. (19) and (20), we can express the concentrations of V(II) and V(III) ions at the liquid–solid interface as:

$$c_{V(II)}^s = \frac{k_m c_{V(II)} + \varepsilon k_1 k_m c_{V(III)} \exp\left(-\frac{\alpha_{-,1} F \eta_1}{RT}\right)/A_1}{B_1 - C_1/A_1} \quad (22)$$

and

$$c_{V(III)}^s = \frac{k_m c_{V(III)} + \varepsilon k_1 k_m c_{V(II)} \exp\left(-\frac{\alpha_{+,1} F \eta_1}{RT}\right)/B_1}{A_1 - C_1/B_1} \quad (23)$$

where

$$A_1 = k_m + \varepsilon k_1 \exp\left(\frac{\alpha_{-,1} F \eta_1}{RT}\right) \quad (24)$$

$$B_1 = k_m + \varepsilon k_1 \exp\left(\frac{\alpha_{+,1} F \eta_1}{RT}\right) \quad (25)$$

and

$$C_1 = (\varepsilon k_1)^2 \exp\left(\frac{-\alpha_{+,1} - \alpha_{-,1} F \eta_1}{RT}\right) \quad (26)$$

A similar approach can be used to solve the concentrations of V(IV) and V(V) at the liquid–solid interface of positive electrode.

The overpotentials in Eqs. (17) and (18) are defined as:

$$\eta_1 = \phi_{m,1} - \phi_{s,1} - E_1 \quad (27)$$

and

$$\eta_2 = \phi_{m,2} - \phi_{s,2} - E_2 \quad (28)$$

where  $E_1$  and  $E_2$  are thermodynamic equilibrium potentials for reactions (1) and (2), respectively, and can be approximated using the relevant Nernst equation as:

$$E_1 = E_1^0 + \frac{RT}{F} \ln \frac{c_{V(III)}}{c_{V(II)}} \quad (29)$$

and

$$E_2 = E_2^0 + \frac{RT}{F} \ln \frac{c_{V(V)} c_{H^+}^2}{c_{V(IV)}} \quad (30)$$

where the equilibrium potentials  $E_1^0$  and  $E_2^0$  are given in Table 3. Then the cell voltage can be written as:

$$E_{cell} = E_2^0 - E_1^0 + \frac{RT}{F} \ln \frac{c_{V(V)} c_{V(II)} c_{H^+}^2}{c_{V(IV)} c_{V(III)}} - |\eta_1| - |\eta_2| - IR_{cell} \text{ (discharge)} \quad (31)$$

and

$$E_{cell} = E_2^0 - E_1^0 - \frac{RT}{F} \ln \frac{c_{V(V)} c_{V(II)} c_{H^+}^2}{c_{V(IV)} c_{V(III)}} + |\eta_1| + |\eta_2| + IR_{cell} \text{ (charge)} \quad (32)$$

where  $I$  is the current density and  $R_{cell}$  is the cell electrical resistance.

### 2.3. Boundary conditions

The boundary conditions can be specified by referring to Fig. 2 as follows:

At  $x = 0$ ,

$$\frac{\partial \phi_{s,1}}{\partial x} = 0, \quad \phi_{m,1} = 0 \quad (33)$$

$$\frac{\partial p}{\partial x} = 0, \quad \frac{\partial p}{\partial y} = 0 \quad (34)$$

$$\frac{\partial c_i}{\partial x} = 0, \quad i = V^{2+}, V^{3+}, H^+, HSO_4^- \quad (35)$$

At  $x = x_1$ ,

$$\phi_{s,1} = \phi_{mem} \quad (36)$$

$$\vec{N}_i \cdot \vec{n} = 0, \quad i = V^{2+}, V^{3+}, HSO_4^- \quad (37)$$

$$\vec{i}_s \cdot \vec{n} = \vec{i}_{mem} \cdot \vec{n} \text{ (for proton)} \quad (38)$$

At  $x = x_2$ , the interfacial conditions are similar to those at  $x = x_1$ ,

$$\phi_{s,2} = \phi_{mem} \quad (39)$$

$$\vec{N}_i \cdot \vec{n} = 0, \quad i = VO^{2+}, VO_2^+, HSO_4^- \quad (40)$$

$$\vec{i}_{mem} \cdot \vec{n} = \vec{i}_s \cdot \vec{n} \text{ (for proton)} \quad (41)$$

At  $x = x_3$ ,

$$\frac{\partial \phi_{s,2}}{\partial x} = 0, \quad -\sigma_m^{eff} \frac{\partial \phi_{m,2}}{\partial x} = I \quad (42)$$

$$\frac{\partial p}{\partial x} = 0, \quad \frac{\partial p}{\partial y} = 0 \quad (43)$$

$$\frac{\partial c_i}{\partial x} = 0, \quad i = VO^{2+}, VO_2^+, H^+, HSO_4^- \quad (44)$$

At  $y = 0$ ,

$$\frac{\partial \phi_s}{\partial y} = 0, \quad \frac{\partial \phi_m}{\partial y} = 0 \quad (45)$$

$$c_i = c_i^0, \quad i = V^{2+}, V^{3+}, H^+ \text{ and } HSO_4^- \text{ when } 0 < x < x_1$$

$$i = VO^{2+}, VO_2^+, H^+ \text{ and } HSO_4^- \text{ when } x_2 < x < x_3 \quad (46)$$

$$\vec{u} \cdot \vec{n} = u_{in} \quad (47)$$

where

$$c_i^0 = c_{total} \cdot SOC \text{ (for } V^{2+} \text{ and } VO_2^+) \quad (48)$$

$$c_i^0 = c_{total} \cdot (1 - SOC) \text{ (for } V^{3+} \text{ and } VO^{2+}) \quad (49)$$

At  $y = h$ ,

$$\frac{\partial \phi_s}{\partial y} = 0, \quad \frac{\partial \phi_m}{\partial y} = 0 \quad (50)$$

$$\frac{\partial c_i}{\partial y} = 0, \quad i = V^{2+}, V^{3+}, H^+ \text{ and } HSO_4^- \text{ when } 0 < x < x_1;$$

$$i = VO^{2+}, VO_2^+, H^+ \text{ and } HSO_4^- \text{ when } x_2 < x < x_3 \quad (51)$$

$$p = p_{out} \quad (52)$$

At all the boundaries except the inlet/outlet, the pressure satisfies the Neumann condition:

$$\nabla p \cdot \vec{n} = 0 \text{ (except inlet/outlet)} \quad (53)$$

The interfacial boundary conditions (Eqs. (38) and (41)) are based on the current balance to associate the variables in both electrodes for iterative solution.

### 3. The correlations of the SOC-dependent viscosity of anolyte and catholyte

To consider the effect of the SOC-dependent viscosity of both anolyte and catholyte, the correlations of the respective electrolyte viscosity in terms of SOC are required. In this work, the correlations are obtained through a theoretical analysis. The basic idea for predicting the electrolyte viscosity is to utilize the viscosities of binary solutions (*i*-H<sub>2</sub>O) which compose the VRFB electrolyte. From Eyring's absolute rate theory, for an electrolyte solution the viscosity can be expressed as [25]:

$$v_{mix} = v_{id} \exp \left( \frac{\Delta G^\neq}{RT} \right) \quad (54)$$

with

$$v_{id} = \exp \left( \sum_i x_i \ln v_i^{pure} \right) \quad (55)$$

where  $v_{mix}$  and  $v_{id}$  are the kinematic viscosities of a mixed solution and an ideal solution, respectively. Superscript *pure* together with subscript *i* represent the property of pure substance.  $\Delta G^\neq$  is the molar excess activation free energy,

$$\Delta G^\neq = \beta RT \sum_i z_i x_i \ln a_i \quad (56)$$

where  $\beta$  is a system-specific constant for non-ideal solutions,  $z_i a_i$  and  $x_i$  are the stoichiometric coefficient, the activity and molar fraction of the component.

If an electrolyte is made up of *n* electrolyte solutes and one solvent (generally water), Eq. (54) can be rewritten as:

$$\ln v_{mix} = \sum_i x_i \ln v_i^{pure} + x_{H_2O} \ln v_{H_2O}^{pure} + \frac{1}{RT} \Delta G^\neq \quad (57)$$

with

$$x_i = m_i / \left( (1000/M_{H_2O}) + \sum_{i=1}^n m_i \right) \quad (58)$$

where  $m_i$  is the molality (mol/kg).

Similarly, for a binary solution, the viscosity can be expressed as:

$$\ln v_i^{(i,0)} = x_i^{(i,0)} \ln v_i^{pure} + x_{H_2O}^{(i,0)} \ln v_{H_2O}^{pure} + \frac{1}{RT} \Delta G_i^{\neq(i,0)} \quad (59)$$

where the variables with the superscript (*i,0*) together with the subscript *i* to denote the quantities of *i* in the binary solution (*i*-H<sub>2</sub>O) having the same water activity as that of the mixed solution.

Combining Eqs. (56) and (59), the viscosity of the electrolyte can be expressed as:

$$\ln v_{mix} = \sum_i \frac{x_i}{x_i^{(i,0)}} \ln v_i^{(i,0)} + \beta \sum_i z_i x_i \ln \left( \frac{z_i m_i}{\sum_i z_i m_i} \right) \quad (60)$$

or

$$\ln \mu_{mix} = \sum_i \frac{x_i}{x_i^{(i,0)}} \ln \mu_i^{(i,0)} + \beta \sum_i z_i x_i \ln \left( \frac{z_i m_i}{\sum_i z_i m_i} \right) + \left( \ln \rho_{mix} - \sum_i \frac{x_i}{x_i^{(i,0)}} \ln \rho_i^{(i,0)} \right) \quad (61)$$

where  $\rho$  is the density. It is worth noting that in a dilute solution system, the last two terms in the right-hand side of Eq. (61) can be eliminated (Note the value of  $\rho_i^{(i,0)}$  is close to that of  $\rho_{mix}$  and  $\sum_i \frac{x_i}{x_i^{(i,0)}} = 1$ ) as an approximation. Based on the experimental data in [15] and the data we measured with a viscometer (Brookfield®), the relationships between the dynamic viscosity (Pa · s) and concentration (mol/L) for VSO<sub>4</sub>, V<sub>2</sub>(SO<sub>4</sub>)<sub>3</sub>, VOSO<sub>4</sub> and (VO<sub>2</sub>)<sub>2</sub>SO<sub>4</sub> solutions (with water as solvent) can be fitted as follows:

$$\mu_{b,V(II)} = (1.362 + 0.251c_{V(II)} + 0.15c_{V(II)}^2)/1000 \quad (62)$$

$$\mu_{b,V(III)} = (1.913 - 0.855c_{V(III)} + 1.05c_{V(III)}^2)/1000 \quad (63)$$

$$\mu_{b,V(IV)} = (2.751 + 1.61c_{V(IV)} + 2.05c_{V(IV)}^2)/1000 \quad (64)$$

and

$$\mu_{b,V(V)} = (2.875 - 0.45c_{V(V)} + 1.5c_{V(V)}^2)/1000 \quad (65)$$

where  $\mu_{b,i}$  is the dynamic viscosity of  $i$  in the binary solution ( $i$ -H<sub>2</sub>O).

According to Eq. (61), we can obtain the expressions for the viscosities of anolyte and catholyte in a VRFB:

$$\begin{aligned} \ln \mu_1 = & \frac{X_{V(II)}}{X_{V(II)}^{(i,0)}} \cdot \ln((1.362 + 0.215 \times c_{V,total} \times SOC + 0.15 \\ & \times (c_{V,total} \times SOC)^2)/1000) + \frac{X_{V(III)}}{X_{V(III)}^{(i,0)}} \cdot \ln((1.913 - 0.855 \\ & \times c_{V,total} \times (1 - SOC) + 1.05 \\ & \times (c_{V,total} \times (1 - SOC))^2)/1000) + \frac{X_{H_2SO_4}}{X_{H_2SO_4}^{(i,0)}} \cdot \ln((0.941 \\ & + 0.144 \times (c_{H_2SO_4} + c_{V,total} \times 0.5 \times SOC) + 0.025 \\ & \times (c_{H_2SO_4} + c_{V,total} \times 0.5 \times SOC)^2)/1000) + \frac{X_{H_2O}}{X_{H_2O}^{(i,0)}} \\ & \cdot \ln \mu_{H_2O} \end{aligned} \quad (66)$$

$$\begin{aligned} \ln \mu_2 = & \frac{X_{V(V)}}{X_{V(V)}^{(i,0)}} \cdot \ln((2.875 - 0.45 \times c_{V,total} \times SOC + 1.51 \\ & \times (c_{V,total} \times SOC)^2)/1000) + \frac{X_{V(IV)}}{X_{V(IV)}^{(i,0)}} \cdot \ln((2.751 + 1.61 \\ & \times c_{V,total} \times (1 - SOC) + 2.05 \\ & \times (c_{V,total} \times (1 - SOC))^2)/1000) + \frac{X_{H_2SO_4}}{X_{H_2SO_4}^{(i,0)}} \cdot \ln((0.941 \\ & + 0.144 \times (c_{H_2SO_4} + c_{V,total} \times 0.5 \times SOC) + 0.025 \\ & \times (c_{H_2SO_4} + c_{V,total} \times 0.5 \times SOC)^2)/1000) + \frac{X_{H_2O}}{X_{H_2O}^{(i,0)}} \\ & \cdot \ln \mu_{H_2O} \end{aligned} \quad (67)$$

To validate Eqs. (66) and (67), we consider a VRFB electrolyte consisting of 1.5 M VOSO<sub>4</sub> and 3.0 M H<sub>2</sub>SO<sub>4</sub> (SOC = 0) as reported elsewhere [13]. The dynamic viscosity according to Eq. (67) is 1.69 mPa·s which is close to the experimental data of 1.97 mPa·s. We also measured the catholyte viscosities at SOC = 0.5 (1.58 mPa·s) and SOC = 1 (1.47 mPa·s). The calculated values from Eq. (67) are 1.43 and 1.34 mPa·s respectively. For the anolyte, the measured values at SOC = 0 and SOC = 1 are 1.85 and 1.36 mPa·s respectively. The comparison indicates that the equations can accurately predict the electrolyte viscosity.

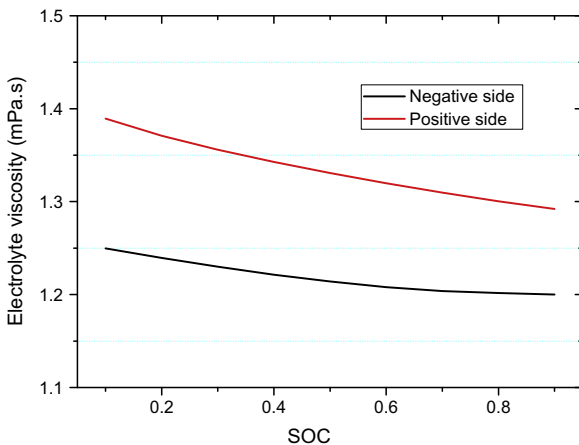


Fig. 3. Variations in the electrolyte viscosity with SOC in the positive and negative half-cells.

Eqs. (66) and (67) are plotted in Fig. 3 (The electrolytes consist of 1.5 M vanadium ions and 2.0 M H<sub>2</sub>SO<sub>4</sub>). The equations are integrated into the numerical model.

#### 4. Results and discussion

The variables in the conservation equations were iteratively solved using the self-written code based on the finite-volume method. The relative error tolerance was set to 1.0E–6.

##### 4.1. Model validation

The numerical solutions are validated against the experimental data in the literature [26]. Since the cell voltages were measured as a function of time, the time-dependant experimental data need to be transformed to the SOC-dependant data for the purpose of comparison. The transformation is made according to [19]:

$$SOC = SOC^0 + \frac{t}{t_0}(SOC^e - SOC^0) \quad (68)$$

where  $t_0$  is the total charge (or discharge) time,  $SOC^0$  and  $SOC^e$  represent, respectively, the value of SOC at the beginning and end of the discharge (or charge). The comparison between the numerical and experimental data is shown in Fig. 4. It can be seen that the model captures the trend well. The small discrepancies in the cell voltage between the numerical results and the experimental data at high SOC may be caused by the presence of side reactions, which are not accounted in the present model. In addition, the cell performance predicted by the present model is compared with that predicted from the model with a constant electrolyte viscosity. Under the same operating conditions, the cell voltage solved from the latter is lower during the charge process and higher during the discharge process. Hence, ignoring the variation of electrolyte viscosity will overestimate the cell performance in the whole range of SOC.

##### 4.2. Pressure drop

Fig. 5 shows the pressure drop distributions in the positive and negative electrodes at SOC = 0.7. As can be seen, the pressure drop through the positive electrode reaches as high as 2978 Pa while that in the negative side is 2676 Pa. The difference is induced by the divergence in viscosity, as the viscosity of the positive electrolyte is higher than that of the negative electrolyte. By comparison, applying the commonly used assumption that the electrolyte viscosity is a constant (1.2 mPa·s) the pressure drop in either the

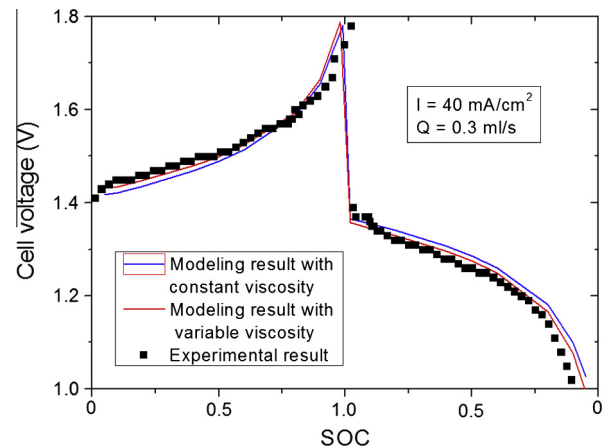
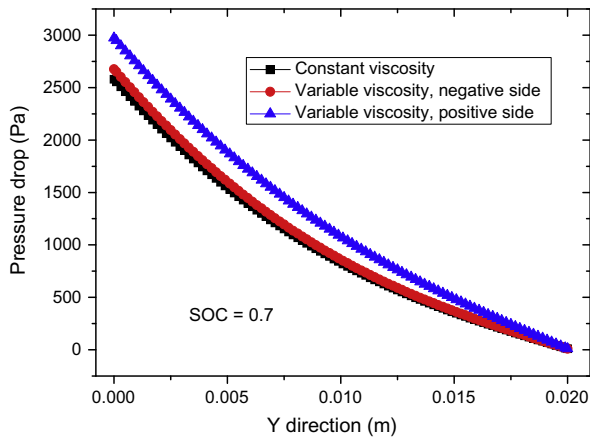


Fig. 4. Cell performance comparison between numerical solutions and experimental data.





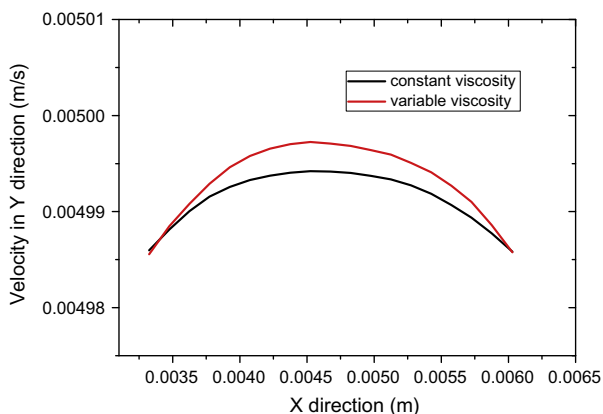
**Fig. 5.** Pressure drops across the electrode resulting from the present model and the constant-viscosity model.

positive or the negative side is 2590 Pa. Moreover, the velocities in the positive and negative electrodes can be obtained. Fig. 6 shows the distributions of the velocity at the middle of the positive electrode. It can be seen that a larger viscosity causes sharper velocity profile, but the difference is rather small. The velocity distribution is used to solve other variables such as the ions concentrations, electrode overpotential and local current density.

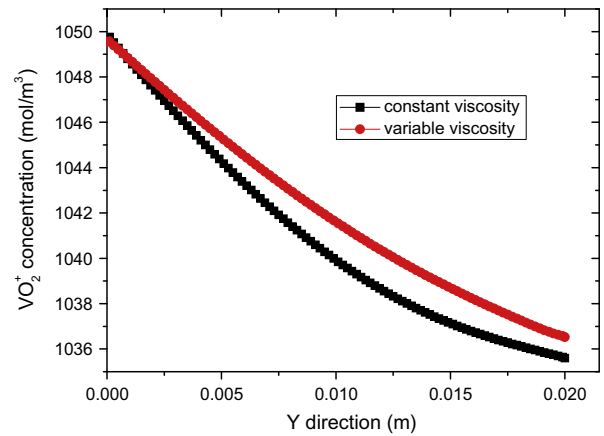
#### 4.3. Distributions of concentration, overpotential and local current density

The distributions of  $\text{VO}_2^+$  concentration at the middle line of the positive electrode at the discharge current density of  $40 \text{ mA/cm}^2$  are presented in Fig. 7 (at SOC = 0.7). During discharge,  $\text{VO}_2^+$  is consumed along the flow direction such that the concentration decreases from the inlet to the outlet. At any position within the electrode, the  $\text{VO}_2^+$  concentration resulting from the present model is higher than that from the constant-viscosity model. This difference is caused by the difference in the velocity distribution inside the porous electrode.

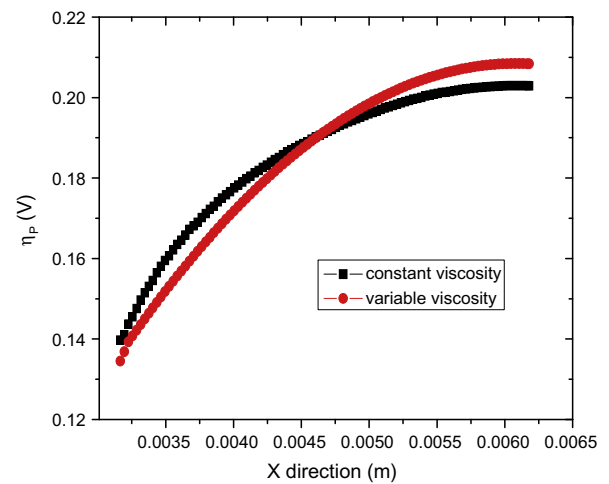
Fig. 8 shows the overpotential distributions obtained through two models in the cross-section of  $y = 0.01 \text{ m}$  under the discharge current density of  $40 \text{ mA/cm}^2$ . It is worth mentioning that according to the definition given by Eq. (28), the overpotential of the positive electrode is less than zero during discharge. For convenience, here we show the absolute value of the overpotential. As can be seen, the overpotential rises from the region near the membrane



**Fig. 6.** Distributions of the velocity at the middle of the positive electrode.



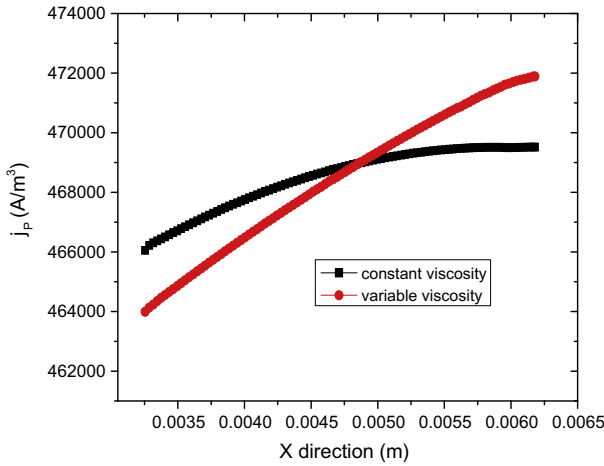
**Fig. 7.** Distributions of the  $\text{VO}_2^+$  concentration resulting from the present model and the constant-viscosity model.



**Fig. 8.** Distributions of the overpotential in the positive electrode resulting from the present model and the constant-viscosity model.

to the region close to the current collector in order to maintain the current. The difference in the overpotential along the electrolyte flow direction is extremely small, since the concentration difference from the inlet to the outlet is small. The overpotential distribution resulting from the present model is steeper than that from the constant-viscosity model. In detail, the overpotential obtained through the present model is higher in the region near the membrane, whereas in the region close to the current collector, the distribution of overpotential is reverse. At the interface of current collector and positive electrode ( $x = 0.00618 \text{ m}$ ), the absolute value of overpotential obtained through the present model is  $0.209 \text{ V}$ , while the value is  $0.20 \text{ V}$  when applying a constant-viscosity model.

The distribution of the local current density in the cross-section of  $y = 0.01 \text{ m}$  is shown in Fig. 9. The local current density changes remarkably in the electrode thickness direction. Less current is generated in the region near the membrane and the maximum current density occurs in the region close to the current collector. The current density resulting from the present model is steeper than that from the constant-viscosity model. The reason leading to this divergence can be explained as follows. The electrochemical reaction rate is determined by both the local overpotential and local reactant concentrations. As the concentration difference is rather small (see Fig. 7), the quantity of local current density is primarily



**Fig. 9.** Distributions of the local current density resulting from the present model and the constant-viscosity model.

affected by the local overpotential distribution. As shown in Fig. 8, the overpotential obtained from the present model distributes more steeply than does that from the constant-viscosity model, causing the corresponding distribution of current density.

#### 4.4. System efficiency

Under practical operating conditions, the pumps continuously circulate the electrolyte between tanks and cell, consuming a part of energy in the VRFB system, which lowers the system efficiency. Including this effect, the energy-based system efficiency can be defined as [27,28]:

$$\psi_{\text{energy}} = \frac{W_{\text{disch}} - W_{\text{pump,disch}}}{W_{\text{char}} + W_{\text{pump,char}}} \quad (69)$$

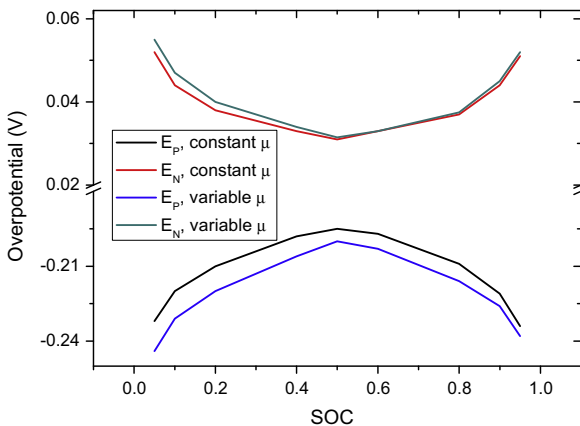
where  $W_{\text{disch}}$  is the VRFB output work during discharge and  $W_{\text{char}}$  is the input work during the charge process. The pumping works can be expressed as:

$$W_{\text{pump,disch}} = Qt_{\text{disch}} \Delta p / \psi_{\text{pump}} \quad (70)$$

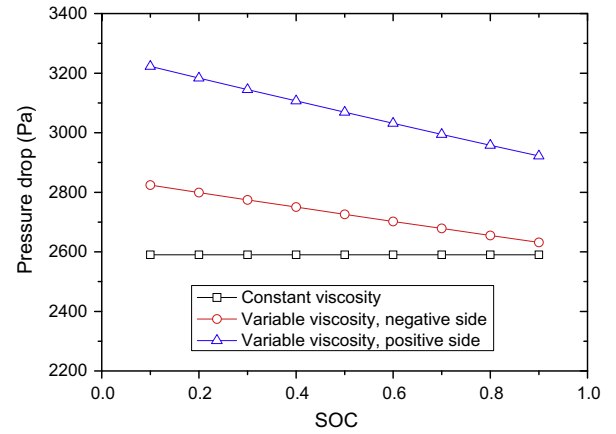
and

$$W_{\text{pump,char}} = Qt_{\text{char}} \Delta p / \psi_{\text{pump}} \quad (71)$$

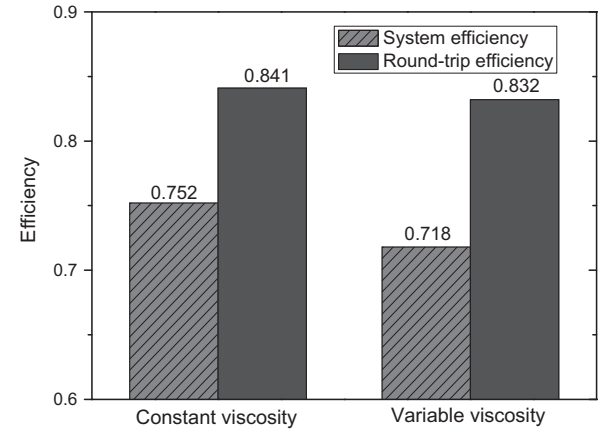
where  $Q$  is the flow rate,  $\psi_{\text{pump}}$  is the pump efficiency and  $\Delta p$  is the pressure drop through the battery.



**Fig. 10.** Variations in the overpotential with SOC.



**Fig. 11.** Variations in the pressure drop with SOC.



**Fig. 12.** Energy-based efficiency and the round-trip efficiency resulting from the present model and the constant-viscosity model.

To solve  $W_{\text{disch}}$  and  $W_{\text{char}}$ , the cell voltages in the whole range of SOC is needed. Fig. 10 shows the variation in the positive and negative half-cell overpotentials in the VRFB obtained through the present model and the constant-viscosity model, respectively. Since the local electrochemical reaction rate is related with the concentrations of two species (See Eqs. (17)), the minimum value of overpotential appears around SOC of 0.5, while the maximum value is at the lowest or highest SOC. The discrepancy in overpotentials resulting from two models is larger in positive half-cell, where the sluggish kinetics results in a larger overpotential. After obtaining the values of overpotential, the corresponding cell voltages can be solved, as already shown in Fig. 4. At all SOC, the cell voltage obtained through the present model is higher during the charge process and is lower during the discharge process. However, the difference is less than 20 mV.

In order to obtain the values of  $W_{\text{pump}}$ , the pressure drops at various SOC are solved and shown in Fig. 11. Both the pressure drops for the negative and positive sides keep constant in the full range of SOC for the constant-viscosity model. In contrast, for the case of variable viscosity, the pressure drops keep decreasing as the SOC rises. Moreover, at each SOC the pressure drop is larger in the positive half-cell since the electrolyte viscosity in positive side is larger than that in negative side.

With the solved pumping work and battery work for a charge–discharge cycle, the energy-based system efficiencies obtained through the present model and the constant- electrolyte-viscosity



model are shown in Fig. 12. The efficiency is 0.752 for the former and is 0.718 for the latter. For comparison, the round-trip efficiencies (*RTE*, the ratio of the energy output during discharge to the energy input during charge process within a charge–discharge cycle) for the two cases are also illustrated. For the former, *RTE* is 0.841; for the latter it becomes 0.832. The results show that the system efficiency including the pumping energy loss can reflect the energy generation and consumption in a VRFB more veritably, since the pumping energy is not neglectable in practical operations. Moreover, the difference between the energy-based system efficiency and *RTE* for that resulting from the present model is larger than that from the constant-viscosity model, implying that the previous studies have underestimated the effect of electrolyte viscosity in VRFBs.

## 5. Conclusions

In this work, a 2-D mass-transport and electrochemical model for a VRFB is developed, which takes the effect of SOC-dependent electrolyte viscosity into account. The model is used to investigate the distributions of vanadium ions concentration, overpotential and local current density of a single vanadium redox cell. Compared with the results resulting from a constant-electrolyte-viscosity model, the results from this model exhibit higher pressure drop (especially in the positive half-cell), steeper distributions of overpotential and local current density in the electrodes. The comparison of modeling results indicates that the consideration of the SOC-dependent electrolyte viscosity enables a more realistic simulation and more accurate estimations of pumping work and the system efficiency of VRFBs.

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