# Clamped-Free Bar Model for ARCH-COMP 2021 AFF category

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## 1 Model description

We consider a uni-dimensional wave propagation problem from [1]. The problem is governed by the partial differential equation

$$EA\frac{\partial^2 u}{\partial x^2}(x,t) - \rho A\frac{\partial^2 u}{\partial t^2}(x,t) = 0, \tag{1}$$

where u(x,t) is the displacement of the point in position x at time t, considering the axis shown in Figure 1. The model consists of a bar of length L=200 and cross-section area A=1. The bar is formed by a linear elastic material with Young modulus  $E=30\times 10^6$  and density  $\rho=7.3\times 10^{-4}$ .

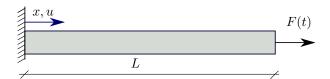


Figure 1: Example 2: diagram of the clamped-free bar excited by end load

The bar is considered to be initially at rest, with u(x,0)=0 and  $\frac{\partial u}{\partial t}(x,0)=0$  for all  $x\in[0,L]$ . The boundary conditions are u(0,t)=0, corresponding to the fixed end, and  $\sigma(L)A=F(t)$ , for the free end, where  $\sigma(x,t)=E\frac{\partial u}{\partial x}(x,t)$ . The free end is submitted to a step force F(t)=10000H(t), where H(t) is the heaviside function.

The analytical solution of this problem in the continuum can be obtained using mode superposition [2] and it is given by:

$$u(x,t) = \frac{8FL}{\pi^2 EA} \sum_{s=1}^{\infty} \left\{ \frac{(-1)^{s-1}}{(2s-1)^2} \sin \frac{(2s-1)\pi x}{2L} \left( 1 - \cos \frac{(2s-1)\pi \mu t}{2L} \right) \right\}, \quad (2)$$

where  $\mu = \sqrt{E/\rho}$ .

Considering N two-node finite elements with linear interpolation, the following system of linear ODEs is obtained:

$$M\frac{\partial^2 u(t)}{\partial t^2} + Ku(t) = F(t), \qquad t \in [0, T], \tag{3}$$

where the  $N \times N$  stiffness and mass matrices are respectively:

where  $\ell = L/N$  is the length of each element.

#### 2 Transformation to first order

Since M is invertible, Eq. (3) can be rewritten as a system of first order ODEs. We introduce the auxiliary variables x = [u, u'], then multiply the equation Mu''(t) + Ku(t) = F(t) by  $M^{-1}$ :

$$u'' = -M^{-1}Ku(t) + M^{-1}F(t)$$
(4)

The resulting system is:

$$x'(t) = \frac{d}{dt} \begin{bmatrix} u \\ u' \end{bmatrix} = \begin{bmatrix} 0 & I \\ -M^{-1}K & 0 \end{bmatrix} \begin{bmatrix} u \\ u' \end{bmatrix} + \begin{bmatrix} 0 \\ M^{-1}F(t) \end{bmatrix} = Ax(t) + f(t)$$
 (5)

The matrices A and b are provided in MAT format for problem sizes N=100, N=500 and N=1000.

# 3 Reachability settings

There are two versions of this benchmark:

- CB21C: constant input F(t).
- CB21F: the inputs F(t) can change arbitrarily over time.

The bar is considered to start at rest:  $x_i = 0$  for all i = 1, ..., 2N. In both CB21C and CB21F, the forcing term is  $F(t) \in [9900, 10100]$ . Three scenarios are considered with an increasing number of dimensions: N = 100, N = 500, N = 1000.

### 4 Requirements

The tools should report the computation time for the time horizon T=0.01s. Discrete-time analysis should use a step size  $9.88\times 10^{-7}$ . The accuracy of the results is measured by reporting the maximum value of the position (resp. velocity) at nodes 70, 350 and 700 for the scenarios with N=100, N=500 and N=1000 respectively.

Tools should plot the velocity at node 700 as a function of time for the N=1000 scenario (or the corresponding node for a smaller instance) for the time interval  $[8.15\times 10^{-3}, 8.40\times 10^{-3}]$ .

### 5 Damped case

To model a realistic beam, we consider adding a damping matrix to the system.

$$C = aK + bM, (6)$$

with  $a = b = 10^{-6}$ . The system matrix then becomes

$$x'(t) = \frac{d}{dt} \begin{bmatrix} u \\ u' \end{bmatrix} = \begin{bmatrix} 0 & I \\ -M^{-1}K & -M^{-1}C \end{bmatrix} \begin{bmatrix} u \\ u' \end{bmatrix} + \begin{bmatrix} 0 \\ M^{-1}F(t) \end{bmatrix} = Ax(t) + f(t). \tag{7}$$

### References

- [1] Mohammad Mahdi Malakiyeh, Saeed Shojaee, and Klaus-Jürgen Bathe. The bathe time integration method revisited for prescribing desired numerical dissipation. *Computers & Structures*, 212:289–298, 2019.
- [2] Michel Géradin and Daniel J Rixen. *Mechanical vibrations: theory and application to structural dynamics*. John Wiley & Sons, 2014.