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# Thermodynamic assessment of the Al–Cu–Zn system, part II: Al–Cu binary system



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### ABSTRACT

The thermodynamic description of the Al–Cu system is reassessed, and for the first time all available experimental data have been comprehensively compared with calculated results. Key distinctions from previous Calphad descriptions are worked out. The new model for the  $\gamma_1$  and  $\gamma_2$  phases partially reflects the crystal structure while being sufficiently simple for a realistic joining with the  $\gamma$ -CuZn phase in the Al–Cu–Zn ternary system. The  $\eta_1/\eta_2$  phase transition and the  $\alpha_2$  phase have been introduced in the description for the first time. An improved agreement of the calculated phase diagram with experimental data, very good agreement with the extensive thermodynamic properties of the liquid phase and reasonable agreement for solid phase thermodynamic properties is demonstrated.

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## 1. Introduction

The thermodynamic description of the Al-Cu system is essential for focused development of Al-based and Cu-based alloys. It is also a key binary edge system for Zn-Al-Cu ternary alloys and, as such, important for Zn-based alloys and thermodynamic databases. In the first thermodynamic description of Al-Cu system [1], four intermediate  $\theta$  (Al<sub>2</sub>Cu),  $\eta$  (AlCu),  $\gamma$  (Al<sub>4</sub>Cu<sub>9</sub>) and  $\beta$  (Bcc) phases were considered, and the  $\theta$ ,  $\eta$  and  $\gamma$  phases are treated as stoichiometric phases. Later, Murray [2] improved the thermodynamic description of [1] based on a review of the experimental work on Al-Cu system by 1985. Besides these four intermetallic phases, Chen et al. [3] introduced two more stoichiometric phases, the  $\delta_1$ and  $\alpha_2$  phases, into their thermodynamic modeling. Currently, it is the thermodynamic description of the Al-Cu system [4] offered in the public COST507 database that has been most widely used and accepted in multicomponent databases. Unfortunately, only the optimized parameters were published in the COST507 book [4]. Nothing is found on the experimental data evaluation, the thermodynamic model selection and, most importantly, the comparison of calculated results with experimental data for the Al-Cu system in the COST507 publication [4] or any later publication from that author. Based on that dataset [4], Liang and Chang [5] simplified the  $\gamma_1$ -Al<sub>4</sub>Cu<sub>9</sub> (cP52, P-43m) phase from a three sublattice model to a single sublattice model together with the  $\gamma\text{-Cu}_8\text{Zn}_5$  (cl52, *I*-43m) phase in the Cu–Zn system, thus obtaining a complete solid solution range of the  $\gamma$  phase in the Al–Cu–Zn ternary system. Miettinen [6] accepted all the models used as [5] and revised some parameters for their Cu–Al–Zn system. However, their thermodynamic description [6] only focused on the Cu-rich corner, and did not provide parameters for Al-rich phases. But if introducing the missing parameters from [4], then the calculated Al–Cu phase diagram in full composition range is completely off. Witusiewicz et al. [7] revised some parameters of the  $\gamma_1\text{-Al}_4\text{Cu}_9$  phase and the liquid phase of [4] in order to avoid an artifact below room temperature and the temperature dependence of the liquid enthalpy of mixing in [4].

Although several thermodynamic descriptions of the Al–Cu binary system are available in the literature [1–7], none of them provided a comprehensive comparison between results of thermodynamic calculations and available experimental data. In addition, some new experimental results after 1991, the actual year of the COST507 assessment [4] also need to be taken into consideration. It is the purpose of this study to provide this thorough thermodynamic reassessment of the Al–Cu system, which will also serve as a key binary edge system of the Al–Cu–Zn ternary.

## 2. Critical review of experimental literature data

## 2.1. Crystal structure of solid phases

Table 1 [8–11] compiles the crystal structure information of all solid phases. The phase names used as Greek symbols in the text and in the TDB file of this work are provided in the first column.

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**Table 1**Solid phases in the Al–Cu binary system.

Phase notation/Name in TDB [This work]	Sublattice model [This work]	Sublattice Model in [4,7]	Phase name/ Temperature range (°C) <sup>a</sup>	Pearson symbol/Space group/Prototype <sup>a</sup>	Maximum reported composition range $(at\%)^a$
(Al) Fcc	(Al,Cu)1	(Al,Cu)1	(Al) ≤ 660.45	cF4	0-2.48% Cu
				Fm-3m	
(Cu) Fcc	(Al,Cu)1	(Al,Cu)1	(Cu) ≤ 1084.62	Cu cF4	80.3-100% Cu
(cu) rec	(Al,Cu) i	(Ai,Cu)i	(Cu) ≤ 1084.02	Fm-3m	80.3-100% Cu
				Cu	
θ Al2Cu	(Al)2 (Al,Cu)1	(Al)2 (Al,Cu)1	$\theta$ , CuAl <sub>2</sub> $\leq$ 591	tI12	31.9-33.0% Cu
				I4/mcm	
				CuAl <sub>2</sub>	
$η_1$ Eta1	(Al,Cu)1 (Cu)1	(Al,Cu)1 (Cu)1	η <sub>1</sub> , CuAl(h) 624-560	oS32 <sup>b</sup>	49.8-52.4% Cu
				Cmmm	
				Cu <sub>0.94</sub>	
Eta2	(Al)48 (Cu)52	(Al,Cu)1 (Cu)1	$\eta_2$ , CuAl(r) $\leq$ 561	Al <sub>0.94</sub> mC20	49.8-52.3% Cu
η <sub>2</sub> Eta2	(AI)46 (CII)52	(AI,Cu)I (Cu)I	$\eta_2$ , CuAi(1) $\leq 361$	C2/m	49.6-32.3% Cu
				CuAl(r)	
ζ Al9Cu11	(Al,Cu)9 (Cu)11	(Al)9 (Cu)11	ζ <sub>1</sub> , CuAl(h) <sup>c</sup>	oF88-4.7°	55.2-59.8% Cu <sup>a</sup>
\$	(,,- (,	()- ()	590–530 <sup>a</sup>	Fmm2	Single phase
			Stable at 550 and 400	Cu <sub>47.8</sub>	57.5% Cu
			(10 days) <sup>c</sup>	Al <sub>35.5</sub>	(550 °C) <sup>c</sup>
ζ Al9Cu11	(Al,Cu)9 (Cu)11	(Al)9 (Cu)11	$\zeta_2$ , CuAl(r) <sup>d</sup>	oI24-3.5 <sup>d</sup>	55.2–56.3% Cu
			≤ 570 <sup>a</sup>	Imm2	55–58% Cu <sup>d</sup>
			Decomposes at 400 forming	Cu <sub>11.5</sub> Al <sub>9.0</sub>	single phase
P.o.	(A1 C+1)1	(A1 C)1	ζ <sub>1</sub> d	Cubic?	56.8% Cu <sup>d</sup> 59.4–62.1% Cu
$\varepsilon_1$ Bcc $\varepsilon_2$ Eps2	(Al,Cu)1 (Al,Cu)1 (Cu)1	(Al,Cu)1 (Al,Cu)1 (Cu)1	$\varepsilon_1$ , $Cu_{1-x}Al_x$ (958-848) $\varepsilon_2$ , $Cu_{2-x}Al$ (850-560)	hP6	55.0-61.1% Cu,
ε <sub>2</sub> Eps2	(Al,Cu)1 (Cu)1	(Ai,Cu)i (Cu)i	ε <sub>2</sub> , Cu <sub>2-x</sub> Ai (830-300)	P6 <sub>3</sub> /mmc	55.0-01.1% Cu,
				Ni <sub>2</sub> In	
				(NiAs type)	
$\delta_1$ Al5Cu8	(Al)5 (Cu)8	(Al)2 (Cu)3	$\delta$ , $Cu_{1-x}Al_x$	hR*	59.3-61.9% Cu
				R-3m	
$\gamma_2$ Gam2	(Cu)2 (Cu)3 (Al,Cu)8	(Al)4 (Al,Cu)1 (Cu)8	$\gamma_0$ , $Cu_{1-x}Al_x$ 1037-800	cI52	59.8-69% Cu
				I-43m	
Com 1	(C-)2 (C-)2 (ALC-)2	(41)4 (41 0-)4 (6.30	C. Al. 900	Cu <sub>5</sub> Zn <sub>8</sub>	C2 C0% C::
γ <sub>1</sub> Gam1	(Cu)2 (Cu)3 (Al,Cu)8	(Al)4 (Al,Cu)1 (Cu)8	$\gamma_1$ , $Cu_9Al_4 < 890$	cP52 P-43m	62–68% Cu
				P-43m Cu <sub>9</sub> Al <sub>4</sub>	
в Всс	(Al,Cu)1	(Al,Cu)1	β, Cu <sub>3</sub> Al (h) 1049-559	cl2	70.6-82.0% Cu
p Dec	(11,Cu)1	(1 ii,Cu) i	p, ca3m (m) 1043-333	Im-3mW	70.0 02.0% Cu
α <sub>2</sub> Alpha2	(Al)23 (Cu)77	_	$\alpha_2$ , $Cu_{1-x}Al_x < 363$	(TiAl <sub>3</sub> type)	76.5-78% Cu

a Data in the last three columns are generally from MSIT report [8] and accepted as assessment of experimental information.

For convenience, the corresponding phase names used in [2,8] are shown in column 4. The crystal structures of the solid phases are also generally accepted from the Materials Science International Team (MSIT) report [8]. Here we only present the experimental results published later than 2004 or not included in [8].

In the earlier review [2], the  $\eta_1$  phase was reported as orthorhombic structure with Pearson symbol of oP16 and space group of *Pban* [12] or with Pearson symbol of oC16 and space group *Cmmm* from another unknown source. In [8], the  $\eta_1$  phase is reported just as o\*32. Recently, Ponweiser and Richter [9] determined the crystal structure of  $\eta_1$ –CuAl $_{1-\delta}$  phase as orthorhombic (space group *Cmmm*), their data are accepted here.

In [2], the high temperature stable  $\zeta_1$  phase was reported as hP42 with space group P6/mmm cited from [12], and the room temperature  $\zeta_2$  phase was possibly monoclinic cited from [13]. In the review [8], it is stated that the  $\zeta_1$  phase (oF88, Fmm2) is stable at temperature range from 530 to 590 °C, and the  $\zeta_2$  phase (oI24, Imm2) is stable at temperature below 570 °C, referring to a conference abstract [14]. However, in the work published later by the same group, Gulay and Harbrecht [10,11] reported that the  $\zeta_1$ -Al<sub>3</sub>Cu<sub>4</sub> (oF88, Fmm2) phase is stable at 400 °C and 550 °C [11], while the  $\zeta_2$ -Al<sub>3</sub>Cu<sub>4- $\delta$ </sub> (oI24, Imm2) phase is stable at 530 °C but

metastable at 400 °C [10]. At 450 °C, the mixture of  $\zeta_1$  and  $\zeta_2$  phases is seen, thus they proposed that the eutectoid decomposition temperature of  $\zeta_1$  into  $\zeta_2$  and  $\eta_1$  is between 400 °C and 450 °C, and the  $\zeta_2$  phase is metastable at ambient temperature [10]. Since it is controversial to distinguish the  $\zeta_1$  and  $\zeta_2$  phase, we treat them as a single  $\zeta$  phase for simplicity.

The  $\alpha_2$  phase is reported to form below 363 °C at about 77.5 at% Cu with an ordered Fcc long-period superlattice structure [2,8]. The  $\alpha_2$  phase is considered in this work as opposed to any previous Calphad assessment. Additional metastable phases are discussed in detail in the review [2].

## 2.2. Thermodynamic data

Table 2 [9,15–38] collects all the original experiment work on thermodynamic properties of Al–Cu system phase. The enthalpy of mixing of the liquid phase was firstly measured by Kawakami [15], which is accepted in [39]. Hultgren [16] published a set of data provided by Yasava and Itagaki as private communication, followed by their publication in [17]. Yasava's data [16,17] are much less exothermic than Kawakami's data [15,39]. The later published liquid enthalpy of mixing data [18–21,35] measured by high

<sup>&</sup>lt;sup>b</sup> Crystal structure of  $\eta_1$  is from [9].

 $<sup>^{</sup>c}$   $\zeta_{1}$  is perceived as high-temperature phase [2, 8], marked data are of from [11], notation  $(\zeta_{2}/\zeta_{1})$  differs in [9].

<sup>&</sup>lt;sup>d</sup>  $\zeta_2$  is perceived as low-temperature phase [2, 8], marked data are of from [10], notation  $(\zeta_2/\zeta_1)$  differs in [9].

**Table 2**Summary of the experimental investigations for the Al–Cu binary system.

Quantity measured	Method	Temperature (°C)	Composition (at% Cu)	Refs.
$\Delta_m^{Liquid}H$	Calorimetry Calorimetry Adiabatic calorimetry	1130 1100 1100	10–90 10–90 11–92	[15] [16] [17]
		1084	4–44	[18] cited
	Calorimetry	1194	5.2-80.1	in [19] [19]
	Calorimetry	1177	9.8-39.9	[20]
	Drop calorimetry	700	0-40	[21]
$\Delta_f^{solid}H$	Calorimetry	20	15.7-92	[22]
$\Delta_f^{-}$ 11	Calorimetry	325	50-66.7	[23]
	Calorimetry	30	63-86	[24]
	Calorimetry	500	0.4-98.4	[25]
$a_{Al}^{liquid}$ , $a_{Cu}^{liquid}$	EMF	700, 800, 900	15.1-48.8	[26]
	EMF	1100	10-94.97	[27]
	Bubbling method	1100–1200	8.8-57.4	[28]
	EMF	1100	70-90	[29]
	Mass spectrometry	1100–1200	0–100	[30]
	Vapor pressure	1100	12-95	[31]
	Mass spectrometry	1000-1400	4.5-90.9	[32]
$\hat{H}_{Al}^{Liquid}$ ,	Isothermal calorimeter	1200	0-60	[33]
$ar{H}^{Liquid}_{Cu}$	Solution calorimeter	918	00.14-1.96	[34]
	Isoperibolic calorimetry	1317	0–70	[35]
$\mu_{Al}^{Solid}$	EMF	660-760	64.5-97.5	[36]
rAI	EMF	400-550	40-62	[37]
	EMF	500	0.4-98.4	[25]
	EMF	777	65–96	[38]
Phase equilibria <sup>a</sup>	DTA, EDS, XRD	400-1050	10-83	[9]

<sup>&</sup>lt;sup>a</sup> Only one detailed phase diagram study was found to be published after the basic assessment of the experimental phase diagram by MSIT [8].

temperature calorimeter were in agreement with [15,39] except the point at 60 at% Cu; they show no obvious temperature dependence from 700 °C to 1317 °C. Some partial enthalpy of mixing data of the liquid were measured by isothermal calorimetry [33,34]. All available activity data of Al in the liquid phase measured by electromotive force (EMF) and vapor pressure [26–32] are consistent with each other, except a few data points in [29,31] slightly deviate from the others. The enthalpy of formation of solid alloys was determined calorimetrically at room temperature [22–24] and at 500 °C [25] by using different calorimetric techniques. The chemical potential of Al in the solid phases was measured by EMF [25,36–38].

# 2.3. Phase equilibria

Murray [2] reviewed all experimental results prior to 1985 and plotted an evaluated Al–Cu phase diagram. Liu et al. [40] studied the Cu–rich Al–Cu phase equilibria and determined the  $\gamma_0/\gamma_1$  transition as a second order transition, and discard the existence of a high temperature  $\beta_0$  phase evaluated by [2]. Gröbner [8] updated the assessed experimental Al–Cu phase diagram which is mainly a combination of the review of [2] and the work of [40]. Ponweiser et al. [9] re–investigated the whole Al–Cu system, and mainly focused on the  $\eta_1$  phase, and they also provided their evaluated experimental phase diagram. The solubility range of the some

intermediate phases,  $\eta_1$ ,  $\eta_2$ ,  $\zeta_1$  and  $\zeta_2$ , in [9] are contradicting [2,8], while the temperatures of the invariant reactions show better agreement [8,9]. The  $\delta$  phase, also named as  $\gamma_2$  phase in some publications, with orthorhombic R3m structure and often reported as R-type brass [41,42,43], was assessed with a solubility range of 59.3-61.9 at% Cu in [2,8]. However, the phase boundaries are not well determined in the original sources [13,42-44], and a single phase region covering both the  $\delta$  and  $\gamma_1$  phases is marked as questionable area in [9]. The coexistence of  $\zeta_1$  and  $\gamma_1$  phase was also reported [45]. The two-phase region between  $\gamma_0$ -AlCu and  $\gamma_1$ -Al<sub>4</sub>Cu<sub>9</sub> was assessed by dashed lines in [2]. Liu et al. [40] clarified the disorder-order transition between  $\gamma_0$ -AlCu and  $\gamma_1$ -Al $\Delta$ Cu phase. The thermal analysis results of [9] further confirmed that as disorder-order transition. In the present work, we essentially accept the assessed phase diagram data in [8] amended with the experimental data in [9].

## 3. Thermodynamic modeling

All equations for the Gibbs energy functions and the notation are the same as given in Part I of this work [46]. The Gibbs energies of the pure components are accepted from the SGTE compilation by Dinsdale in the most recent electronic version (SGTE-Unary-v5.1b, 24 November 2010, www.sgte.org), which is identical with the print version [47], except that more figures for the parameters of Al are given. The Gibbs energies for liquid, the terminal Fcc phase, denoted as (Al) and (Cu), and the intermediate phase Bcc, denoted as  $\beta$  and  $\epsilon_1$ , are described by the substitutional solution model. The Gibbs energies of the two intermediate phases  $\eta_2$  (Al<sub>48</sub>Cu<sub>52</sub>),  $\delta_1$  (Al<sub>5</sub>Cu<sub>8</sub>) and  $\alpha_2$  (Al<sub>23</sub>Cu<sub>77</sub>), assumed to be stoichiometric, are described by a Gibbs energy of formation term relative to pure Al and Cu (Fcc). The Gibbs energies of all other intermediate phases are described by the compound energy formalism (CEF) [48,49], allowing for a certain solution range.

Proper sublattices have been selected for the phases according to their crystal structure and composition range shown in Table 1. The  $\theta$  phase is modeled as (Al)2(Al,Cu)1 and the  $\eta_1$  phase as (Al)1 (Al,Cu)1, which are the same as in [4]. Only one  $\eta$  phase was considered in [4], and the calculated  $\eta$  phase region extends to low Cu side at low temperature. However, the experimental works [8,9] show that the low temperature  $\eta_2$  phase contains higher Cu content than the high temperature  $\eta_1$  phase. In this work, a stoichiometric phase  $\eta_2$  was introduced with a model of (Al)48(Cu)52. As discussed in Section 2.3, the crystal structure and composition range of the  $\delta_1$  phase is still under dispute. In [2], the  $\delta_1$  phase was assessed as  $\gamma$ -brass structure together with the  $\gamma_0$ -AlCu,  $\gamma_1$ -Al\_4Cu\_9 phases, and with a composition range of 59.3–61.9% Cu. We modeled the  $\delta_1$  phase as stoichiometric (Al)5(Cu)8 phase for simplicity.

In the present work, the  $\gamma_0$ -AlCu phase is renamed as  $\gamma_2$  phase in order to obtain a consistent notation in combination with the Cu-Zn binary system. To reflect the crystal structure of the disordered  $\gamma_2$  phase in detail the four-sublattice model (Al,Cu)4(Cu)4 (Cu)6(Al,Cu)12 should be used. As the same simplification applied in the Cu-Zn binary system [46], we joined the first and forth sublattices to get the model (Cu)4(Cu)6(Al,Cu)16. For the ordered  $\gamma_1$  phase, a six-sublattice model (Al,Cu)2(Al,Cu)2(Cu)4(Cu)6(Al,Cu) 6(Al,Cu)6 should be used to model the order/disorder transition; it is obtained by splitting the first and forth sublattices of the detailed model of  $\gamma_2$ , thus allowing the ordering. In this work, the  $\gamma_1/\gamma_2$  transition is not modeled as order–disorder transition, but as first order transition;  $\gamma_1$  and  $\gamma_2$  are treated as two separate phases, and the model of the  $\gamma_1$  phase is also simplified by combining all the sublattices containing both Al and Cu as (Cu)4(Cu)6(Al,Cu)16. This model partially reflects the crystal structure of the  $\gamma$  phase and is just simplified sufficiently for a realistic joining with the  $\gamma$ -CuZn phase in the modeling of the Al-Cu-Zn ternary system, as discussed in [50].

#### Table 3

Phase names; name in tdb-parameters; models (sublattice formula), and parameters of the Gibbs energy equations developed in present work. Gibbs energy is given in J mol-formula<sup>-1</sup> and temperature (T) is given in Kelvin. The Gibbs energies for the pure elements are found in the SGTE compilation [47].

```
Liquid; Liq; (Al,Cu)<sub>1</sub>
L_{AlCu}^{0,Liq} = -67094 + 8.555 \times T
L_{Al,Cu}^{1,Liq} = 32148 - 7.118 \times T
L_{AlCu}^{2,Liq} = 5915 - 5.889 \times T
L_{AlCu}^{3,Liq} = -7290 + 5.5 \times T
(Al), (Cu); Fcc; (Al,Cu)<sub>1</sub>
L_{Al,Cu}^{0,Fcc} = -53520 + 2 \times T
L_{AlCu}^{1,Fcc} = -39005-2.368 \times T
L_{Al,Cu}^{2,Fcc} = 1210
\alpha_2; Alpha2; (Al)<sub>0.23</sub>(Cu)<sub>0.77</sub>
G_{Al+Cl}^{0,Alpha2} = -13904.9 - 3 \times T + 0.23 \times G_{Al}^{0,Fcc} + 0.77 \times G_{Cl}^{0,Fcc}
β, ε<sub>1</sub>; Bcc; (Al,Cu)<sub>1</sub>
L_{Al,Cu}^{0,Bcc} = -73537 + 3.98 \times T
L_{AICU}^{1,Bcc} = 51696-12 \times T
\gamma_1; Gam1; (Cu)<sub>4</sub>(Cu)<sub>6</sub>(Al,Cu)<sub>16</sub>
G_{Cu:Cu:Al}^{0,Gam1} = -94322.4 + 389.508 \times T + 16 \times G_{Al}^{0,Fcc} + 10 \times G_{Cu}^{0,Fcc}
G_{Cu:Cu:Cu}^{0,Gam1} = -2600 + 26 \times G_{Cu}^{0,Fcc}
L_{Cu:Cu:Al,Cu}^{0,Gam1} = -1613680 - 708 \times T
L_{Cu:Cu:Al,Cu}^{1,Gam1} = -1585780-100 \times T
\gamma_2; Gam2; (Cu)_4(Cu)_6(Al,Cu)_{16}
G_{Cu:Cu:Al}^{0,Gam2} = -94322.4 + 389.508 \times T + 16 \times G_{Al}^{0,Fcc} + 10 \times G_{Cu}^{0,Fcc}
G_{Cu:Cu:Cu}^{0,Gam2} = 2600 + 26 \times G_{Cu}^{0,Fcc}
L_{Cu:Cu:Al.Cu}^{0,Gam2} = -1451526.6-859.628 \times T
L_{Cu:Cu:Al.Cu}^{1,Gam2} = -1516652.4 - 92 \times T
\delta_1; Al5Cu8; (Al)<sub>5</sub>(Cu)<sub>8</sub>
G_{Al;Cu}^{0,Al5Cu8} = -278450 + 5.2 \times T + 5 \times G_{Al}^{0,Fcc} + 8 \times G_{Cu}^{0,Fcc}
ζ; Al9Cu11; (Al,Cu)<sub>9</sub>(Cu)<sub>11</sub>
G_{Al:Al}^{0,Al9Cu11}=~20\times G_{Al}^{0,Bcc}
G_{Al-Cij}^{0,Al9Cu11} = -424720 + 20 \times T + 9 \times G_{Al}^{0,Fcc} + 11 \times G_{Cij}^{0,Fcc}
L_{Al,Cu;Cu}^{0,Al9Cu11} = 0
L_{Al,Cu;Cu}^{1,Al9Cu11} = -413300 + 100 \times T
\varepsilon_2; Eps2; (Al,Cu)<sub>1</sub>(Cu)<sub>1</sub>
G_{Al-Cl}^{0,Eps2} = -40326 + 4.1 \times T + G_{Al}^{0,Fcc} + G_{Cl}^{0,Fcc}
G_{Cu:Cu}^{0,Eps2} = 14850 + 30.1 \times T + 2 \times G_{Cu}^{0,Fcc}
L_{Al,Cu:Cu}^{0,Eps2} = -56870 - 55 \times T
\eta_1; Eta1; (Al,Cu)<sub>1</sub>(Cu)<sub>1</sub>
G_{Al;Cu}^{0,Eta1} = -39431 + 1.9 \times T + G_{Al}^{0,Fcc} + G_{Cu}^{0,Fcc}
G_{Cu:Cu}^{0,Eta1} = 4770 + 9 \times T + 2 \times G_{Cu}^{0,Fcc}
L_{Al,Cu:Cu}^{0,Eta1} = 0
L_{Al,Cu:Cu}^{1,Eta1} = -43168 - 18.5 \times T
\eta_2; Eta2; (Al)<sub>0.48</sub>(Cu)<sub>0.52</sub>
G_{Al;CIJ}^{0,Eta2} = -20987 + 1.62 \times T + 0.48 \times G_{Al}^{0,Fcc} + 0.52 \times G_{CIJ}^{0,Fcc}
θ; Al2Cu; (Al)<sub>2</sub>(Al,Cu)<sub>1</sub>
G_{Al:Al}^{0,Al2Cu} = 3 \times G_{Al}^{0,Bcc}
G_{Al:Cu}^{0,Al2Cu} = -48361 + 8 \times T + 2 \times G_{Al}^{0,Fcc} + G_{Cu}^{0,Fcc}
L_{Al;Al,Cu}^{0,Al2Cu} = 2211
```

## 3.1. Parameter optimization

The thermodynamic parameters from [7] were used as a start point. Then the new models for the  $\gamma_1$ ,  $\gamma_2$ ,  $\delta_1$  and  $\eta_2$  phases were introduced, finally  $\alpha_2$ . At last, all the parameters of the phases were optimized based on all assessed experimental data. Many tries and iteration have been carried out during the optimization. Thermodynamic calculations and the parameter optimization are performed using the Pandat software package (www.computherm.com) with PanOptimizer [51]. All thermodynamic Gibbs energy parameters determined in the present work are given in Table 3.

## 4. Results and discussion

## 4.1. Phase diagram

Fig. 1 represents the calculated phase diagram of Al–Cu system in this work. This diagram is essential to clearly assign the current phase notation to all single phase regions of this complex system. Fig. 2 compares the calculated phase diagram with experimental data [8–11]. The plotted "experimental" phase diagram is taken from the most recent and comprehensive critical review performed by Gröbner [8], and the more recent experimental data points are shown in original from [9–11]. For better readability, "original" data points essentially found in [2] are not shown here, the reader is referred to the compilation in [8].

The calculated result shows good agreement with experimental phase diagram [8,9,11], especially for the important invariant reaction temperatures. In case of disagreement between the experimental assessment [8] and the more recent and meticulous experimental data of [9] the latter have been accepted for the present modeling, e.g. in the region of the  $\gamma_1/\gamma_2$  transition, and the eutectoid  $\beta \leftrightarrow \gamma_1 + (Cu)$ , etc. Table 4 shows the calculated invariant reactions compared with experimental data [2,8,9] and previous Calphad calculations [4,5,7]. Murray [2] provided an assessed phase diagram enlarged at Al-rich corner compared with all original experimental data. The calculated results, especially the solvus line of (Al), agree well with the experimental phase diagram [2] as shown in Fig. 3.

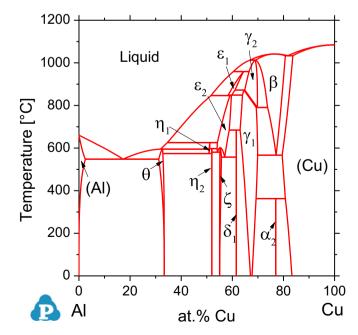
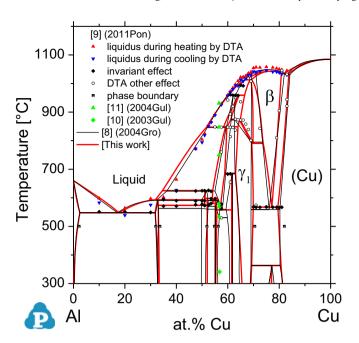


Fig. 1. Calculated phase diagram of the Al-Cu system from the present thermodynamic description.



**Fig. 2.** Comparison of the Al–Cu phase diagram calculated in this work with the assessed experimental phase diagram by MSIT [8] and more recent experimental data [9–11]. Additional short-hand notation indicating year and first author is given.

The DTA results in [9] and the assessed experimental phase diagrams [2,8] show the existence of the  $\eta_1/\eta_2$  phase transition within the  $\eta$  phase. It is most clearly seen in the almost degenerate reaction  $\eta_1 \leftrightarrow \eta_2 + \theta$  measured at 574 °C [9], 11 K above the 563 °C in the previous assessments [2, 8]. Due to the solution range in the  $\eta_1/\eta_2$  phases the second invariant temperature related to the  $\eta_1/\eta_2$ transition,  $\eta_1 + \zeta_{(2)} \leftrightarrow \eta_2$ , is measured at higher temperature, 580 °C, by a single sample [9]. We acknowledge the transition and treat  $\eta_1$  and  $\eta_2$  as separate phases, as opposed to all previous Calphad assessments treating it as one  $\eta$  phase [4,7]. However, due to the scarcity of data, the  $\eta_2$  phase is modeled as stoichiometric phase of Al<sub>48</sub>Cu<sub>52</sub>. The temperature of the reaction  $\eta_1 \leftrightarrow \eta_2 + \theta$  is calculated at 573.7 °C, close to the experimental result of 574 °C [9]. In addition, two other invariant reactions associate with  $\eta_2$ phase are calculated,  $\varepsilon_2 + \eta_1 \leftrightarrow \eta_2$  (597.9 °C) and  $\varepsilon_2 \leftrightarrow \eta_2 + \zeta$ (580.4 °C). These temperatures are very close to detected thermal events at 597 °C and 580 °C corresponding to invariant reactions [9], but were attributed to different reactions associated with the  $\zeta_2$  phase [9]. The  $\zeta_2(\zeta_1)$  phase is not well studied and assumed (or speculated) to possess a large solubility range as discussed below.

Gulay and Harbrecht [10, 11] clearly determined the crystal structures of  $\zeta_1$  (oF88, Fmm2) and  $\zeta_2$  (ol24, Imm2) phases and also carried out DSC measurements. The calculated results in this work agree well with five of the detected thermal signals of alloy Al43.4Cu56.6 ( $\zeta_1$ -Al<sub>3</sub>Cu<sub>4</sub>) at heating in the range 568–931 °C [11], see Fig. 2. The origin of the signal at 533 °C in that sample remains unclear because the previously assumed  $\zeta_1/\zeta_2$  transition near that temperature [2,8,9] is highly questionable based on the results of [10,11] as discussed below.

Gulay et al. [11] also found that the  $\zeta_1$ -phase resists a heat treatment at 400 °C for 10 days, which questions the currently assumed decomposition of this high-temperature phase  $\zeta_1$  at 530 °C based on thermal signals [8,2]. In their second work, Gulay et al. [10] showed that the  $\zeta_2$ -Al<sub>3</sub>Cu<sub>4- $\delta$ </sub> phase, currently perceived as stable low-temperature phase [2,8], decomposes during heat treatment at 400 °C into  $\zeta_1$  (plus some  $\eta_2$ ). Thus,  $\zeta_1$  even forms at 400 °C, which puts a further question mark on the previously assumed thermal stability ranges [2,8] of the " $\zeta$ " phases. They also

**Table 4**Comparison of calculated and experimental invariant reactions in the Al–Cu system.

ystem.						
Reaction	T/°C	Cu content in phases at%		ıt%	Reference/	
	, -	<b>r</b>			comments	
Congruent		L	β			
$L \leftrightarrow \beta$	1049	75	75		[2,8]	
	$1052 \pm 5$				[9] Exp.	
	1045.6	76.6	76.6		[4] Calphad	
	1045.6	76.6	76.6		[5] Calphad	
	1043.9	76.8	76.8		[7] Calphad	
	1043.3	76.7	76.7		[This work]	
Eutectic		L	(Cu)	β		
$L \leftrightarrow (Cu) + \beta$	1032	83.0	84.2	82.0	[2,8]	
	$1035 \pm 5$	$83.0 \pm 0.5$	$84.5 \pm 0.5$	$82.0 \pm 0.5$	[9] Exp.	
	1036.1	81.8	83.7	80.7	[4] Calphad	
	1036.1	81.7	83.7	80.7	[5] Calphad	
	1036.1	81.7	83.7	80.7	[7] Calphad	
Devisership	1035.1	81.8	83.7	80.8	[This work]	
Peritectic	1027	L	β	γ <sub>2</sub>	[0]	
$L+\beta\leftrightarrow\gamma_2$	1037	69.2	70.9	~ 69	[8]	
	$993 \pm 2$	$63.0 \pm 0.5$	$69.0 \pm 0.5$	$65.0 \pm 1.0$	[9] Exp.	
	1020.3	67.9	69.9	68.4	[4] Calphad	
	1019.0	67.9 67.9	69.9 69.9	68.4	[5] Calphad	
	1019.0 <b>1013.2</b>	66.9	<b>69.2</b>	68.4 <b>68.8</b>	[7] Calphad [This work]	
Peritectic	1015.2		L		[IIIIS WOIK]	
$\gamma_2 + L \leftrightarrow \varepsilon_1$	958	$\frac{\gamma_2}{62.9}$	59.8	ε <sub>1</sub> 62.1	[2,8]	
12+2 ↔ €1	$960 \pm 2$	$65.5 \pm 0.5$	$60.0 \pm 0.5$	$64.5 \pm 0.5$	[9] Exp.	
	959.2	66.7	60.4	64.4	[4] Calphad	
	959.1	66.7	60.4	64.4	[5] Calphad	
	959.1	66.7	60.4	64.4	[7] Calphad	
	959.3	66.5	60.4	64.4	[This work]	
Peritectoid	000.5	$\epsilon_1$	γ <sub>2</sub>	γ1	[1110 110111]	
$\varepsilon_1 + \gamma_2 \leftrightarrow \gamma_1$	873	61.4	66.0	63.9	[2]	
$(\gamma_2 \leftrightarrow \gamma_1,  \varepsilon_1)$	$874 \pm 2$	0	65.0	03.0	[9] Exp.	
$\varepsilon_1 + \gamma_2 \leftrightarrow \gamma_1$	877.4	62.5	65.6	64.7	[4] Calphad	
$\varepsilon_1 + \gamma_2 \leftrightarrow \gamma_1$	880.7	62.6	65.6	64.8	[5] Calphad	
$\varepsilon_1 + \gamma_2 \leftrightarrow \gamma_1$	876.7	62.5	65.6	64.7	[7] Calphad	
$\varepsilon_1 + \gamma_2 \leftrightarrow \gamma_1$	872.5	61.1	64.9	64.5	[This work]	
Deg., Ptd.,Cat. <sup>a</sup>		$\epsilon_1$	$\epsilon_2$	L	•	
$\varepsilon_1 \leftrightarrow \varepsilon_2$ , L	848	$\sim$ 59.4	$\sim$ 59.4	52.2	[2,8]	
$\varepsilon_1 \leftrightarrow \varepsilon_2$ , L	$847\pm1$	$59.5 \pm 0.5$	$59.5 \pm 0.5$	$52.5 \pm 0.5$	[9] Exp.	
$\epsilon_1\!+\!L\!\leftrightarrow\epsilon_2$	850.1	58.5	57.7	52.0	[4] Calphad	
$\epsilon_1 \! + \! L \! \leftrightarrow  \epsilon_2$	851.1	58.5	57.7	52.0	[5] Calphad	
$\epsilon_1\!+\!L \!\leftrightarrow \epsilon_2$	851.1	58.5	57.7	52.0	[7] Calphad	
$\varepsilon_1 \leftrightarrow \varepsilon_2 + L$	846.6	58.3	59.1	51.7	[This work]	
Ptd., Deg., Etd. <sup>a</sup>		$\gamma_1$	$\epsilon_1$	$\epsilon_2$		
$\gamma_1 + \epsilon_1 \leftrightarrow \epsilon_2$	850	62.5	~61.1	~61.1	[2,8]	
$\varepsilon_1 \leftrightarrow \varepsilon_2$ , $\gamma_1$	$847 \pm 1$	$64.0 \pm 0.5$	$62.5 \pm 0.5$	$62.5 \pm 0.5$	[9] Exp.	
$\varepsilon_1 \leftrightarrow \varepsilon_2 + \gamma_1$	836.2	63.6	60.2	58.1	[4] Calphad	
$\varepsilon_1 \leftrightarrow \varepsilon_2 + \gamma_1$	837.0	63.9	60.2	58.0	[5] Calphad	
$\varepsilon_1 \leftrightarrow \varepsilon_2 + \gamma_1$	836.4	63.5	60.2	58.0	[7] Calphad	
$\varepsilon_1 \leftrightarrow \varepsilon_2 + \gamma_1$	848.7	64.0	<b>59.8</b>	59.7	[This work]	
Deg., Etd. <sup>a</sup>	- 000	$\gamma_2$ 69.0	β	γ1	[0] Evp	
$(\gamma_2 \leftrightarrow \gamma_1, \beta)$	~ 800 781.1	68.9	73.2	68.7	[9] Exp.	
$\gamma_2 \leftrightarrow \beta + \gamma_1$	781.1 781.3	69.0	73.2 73.2	68.7 68.9	[4] Calphad [5] Calphad	
$\gamma_2 \leftrightarrow \beta + \gamma_1$ $\gamma_2 \leftrightarrow \beta + \gamma_1$	781.1	68.9	73.2	68.7	[7] Calphad	
$\gamma_2 \leftrightarrow \beta + \gamma_1$ $\gamma_2 \leftrightarrow \beta + \gamma_1$	<b>790.7</b>	69.8	73.8	<b>69.3</b>	[This work]	
Peritectoid	750.7	γ1	$\varepsilon_2$	$\delta_1$	[TIIIS WOTK]	
$\gamma_1 + \varepsilon_2 \leftrightarrow \delta_1$	686	62.8	59.2	61.9	[2,8]	
2 · · •1	$684 \pm 1$	$63.0 \pm 0.5$	$58.5 \pm 0.5$	$61.5 \pm 0.5$	[9] Exp.	
	686.6	62.8	57.2	60	[4] Calphad	
	690.2	63.7	57.2	60	[5] Calphad	
	686.6	62.8	57.2	60	[7] Calphad	
	684.6	63.0	58.8	61.5	[This work]	
Peritectic		$\epsilon_2$	L	$\eta_1$	-	
$\varepsilon_2 + L \leftrightarrow \eta_1$	624	55.0	36.3	51.8	[2,8]	
••	$625\pm2$	$54.5 \pm 0.5$	$\textbf{38.5} \pm \textbf{0.5}$	$52.0 \pm 0.5$	[9] Exp.	
	624.9	54.1	34.6	51.0	[4] Calphad	
	625.0	54.1	34.9	51.0	[5] Calphad	
	625.0	54.1	34.9	51.0	[7] Calphad	
	625.0	54.1	34.6	51.2	[This work]	
Peritectic		$\eta_1$	L	θ		
$\eta_1\!+\!L\!\leftrightarrow\!\theta$	591	49.8	32.2	32.8	[2,8]	
	$591 \pm 2$	$51.5 \pm 0.5$	$32.5 \pm 0.5$	$33.5 \pm 0.5$	[9] Exp.	
	595.9	50.6	31.9	32.9	[4] Calphad	

Table 4 (continued)

Reaction	T/°C	C Cu content in phases at%			Reference/ comments	
	595.7 595.7 <b>593.3</b>	50.6 50.6 <b>50.9</b>	32.2 32.2 <b>31.9</b>	32.9 32.9 <b>32.9</b>	[5] Calphad [7] Calphad [This work]	
Etd., Ptd. a	500	ε2	η <sub>1</sub>	ζ	[2,0]	
$\varepsilon_2 + \eta_1 \leftrightarrow \zeta_1$	590	56.5 56.5 + 0.5	52.4	56.2 55.5 + 0.5	[2,8]	
$\varepsilon_2 + \eta_1 \leftrightarrow \zeta_2$ $\varepsilon_2 \leftrightarrow \eta + \zeta$	$597 \pm 1$ 579.3	54.3	53.0 ± 0.5 51.1	55.5 ± 0.5	[9] Exp. [4] Calphad	
$\varepsilon_2 \leftrightarrow \eta + \zeta$	579.3	54.3	51.1	55	[5] Calphad	
$\varepsilon_2 \leftrightarrow \eta + \zeta$	579.3	54.3	51.1	55	[7] Calphad	
-2 11.5		ε2	η1	$\eta_2$	[.]F	
$\epsilon_2 + \eta_1 \leftrightarrow \eta_2$	597.9	54.0	51.2	52	[This work]	
Ptd., Etd. <sup>a</sup>		$\eta_1$	ζ	$\eta_2$		
$\eta_1 \leftrightarrow \eta_2 + \zeta_2$	560	~52.3	55.3	$\sim$ 52.3	[2,8]	
$\zeta_1 + \eta_1 \leftrightarrow \zeta_2$	570	52.3	55.2	50.5 . 0.5	[2,8]	
$\eta_1 + \zeta_2 \leftrightarrow \eta_2$ Eutectoid	$580 \pm 1$	52.5 ± 0.5	54.5 ± 0.5	$52.5 \pm 0.5$	[9] Exp.	
$\varepsilon_2 \leftrightarrow \eta_2 + \zeta$	580.4	ε <sub>2</sub> <b>54.1</b>	ζ <b>55.1</b>	η <sub>2</sub> <b>52</b>	[This work]	
Eutectoid	J00. <del>4</del>	<b>34.1</b> η <sub>1</sub>	θ	$\eta_2$	[IIIIS WOIK]	
$\eta_1 \leftrightarrow \theta + \eta_2$	563	49.8	33.0	49.8	[2,8]	
11 . 12	$574 \pm 3$	$52.0 \pm 0.5$	$33.5 \pm 0.5$	$52.5 \pm 0.5$	[9] Exp.	
	573.7	50.9	32.9	52	[This work]	
Eutectoid		β	(Cu)	$\gamma_1$		
$\beta \leftrightarrow (Cu) + \gamma_1$	567	76.1	80.3	69.0	[2]	
	559	76.1	80.3	69.0	[8]	
	$567 \pm 2$	$76.0 \pm 0.5$	$81.5 \pm 0.5$	$70.0 \pm 0.5$	[9] Exp.	
	564.7	77.1 77.1	79.5	69.1	[4] Calphad	
	566.9 564.7	77.1	79.5 79.5	70.5 69.1	[5] Calphad [7] Calphad	
	<b>567.5</b>	77.1	<b>79.5</b>	69.9	[This work]	
Eutectoid	007.0	ε <sub>2</sub>	ζ	$\delta_1$	[1110 110111]	
$\varepsilon_2 \leftrightarrow \zeta + \delta_1$	560	57.9	56.9	59.3	[2,8]	
	582.7	55.6	55	60	[4] Calphad	
	582.7	55.6	55	60	[5] Calphad	
	582.7	55.6	55	60	[7] Calphad	
Post of the	557.7	57.0	55.7	61.5	[This work]	
Eutectic $L \leftrightarrow \theta + (Al)$	548.2	L 17.1	θ 31.9	(Al) 2.48	[2.0]	
$L \leftrightarrow \theta + (AI)$	548.2 550 + 2	17.1 17.0 + 1.0	31.9 $32.0 + 0.5$	2.48 $2.5 + 0.5$	[2,8] [9] Exp.	
	547.6	17.5 _ 1.0	31.8	2.54	[4] Calphad	
	547.5	17.5	31.8	2.54	[5] Calphad	
	547.5	17.5	31.8	2.54	[7] Calphad	
	547.7	17.7	31.8	2.50	[This work]	
Peritectiod		γ1	(Cu)	$\alpha_2$		
$\gamma_1 + (Cu) \leftrightarrow \alpha_2$	363	69	80.3	77.25	[2,8]	
	363.0	69.3	80.7	77	[This work]	

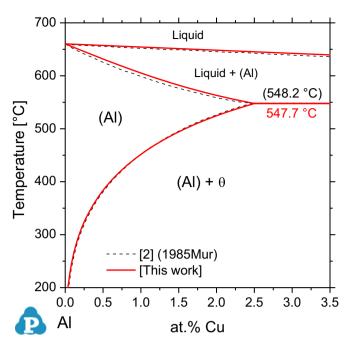
<sup>&</sup>lt;sup>a</sup> The type of these reactions differ in different references. Abbreviations used here are Ptd.: Peritectoid; Deg.: Degenerate; Cat.: Catatectic; Etd.: Eutectoid.

detected a small but reproducible exothermic effect on cooling single-phase  $\zeta_2$  in DSC at 341 °C [10], (see Fig. 2). That might be a kinetic effect of  $\zeta_2$ -sample transition/decomposition considering the metastability of  $\zeta_2$  observed in the same work at 400 °C.

The  $\zeta_1$  and  $\zeta_2$  phases had previously been simplified to a single stoichiometric Al<sub>9</sub>Cu<sub>11</sub> phase [4,7]. However, solution ranges exist even though not very well established. The  $\zeta_1/\zeta_2$  transition and thermal stability ranges are highly questionable [10,11]. Therefore, it is modeled in the present work as a single phase,  $\zeta$ , with a narrow solution range close to the Al<sub>9</sub>Cu<sub>11</sub> composition.

The phase labeled  $\delta_1$  is modeled in the present work as stoichiometric  $Al_5Cu_8$ , which is closer to the experimentally determined composition, see Fig. 2, as compared to the previously assumed model,  $Al_2Cu_3$  [4,7], see Fig. 4. The simplification as stoichiometric phase is considered reasonable because the solution range of that phase is not well established [9].

The  $\gamma_1/\gamma_2$  transition modeled as very narrow two-phase region, is very close to the experimentally determined order–disorder transition line [9]. Also the eutectoid  $\beta = \gamma_1 + (Cu)$  invariant



**Fig. 3.** Comparison of the Al–Cu phase diagram at Al-rich corner calculated in this work with the assessed experimental phase diagram [2].

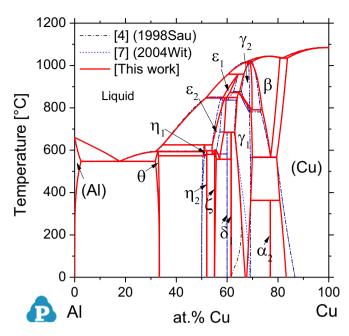
temperature is in excellent agreement with the experimental data [9], see Fig. 2. Moreover, the single-phase  $\varepsilon_2$  phase region in this work agrees better with the experimental phase diagram [2,8] than previous calculations [4,7].

The calculated liquidus line of  $\beta$  is slightly below the experimental data, whereas the liquidus of  $\epsilon_2$  is slightly above, see Fig. 2. It is hard to reduce this deviation while maintaining the very good agreement with the invariant reaction temperatures, see Table 4, and all other thermodynamic and phase equilibrium data. The same deviation is observed in previous Calphad calculations [4,7], as shown in Fig. 4. However, a slight improvement is obtained in the present work, especially for the  $\epsilon_2$  liquidus.

It is noted that the calculated phase diagrams of [4] and [7] virtually overlap, the only visible distinction in Fig. 4 is the  $(\delta_1 + \gamma_1)/\gamma_1$  phase boundary, which veers off to much lower Cu content with decreasing temperature in the diagram of Saunders [4]. That produces a very unrealistic shape of the  $(\delta_1 + \gamma_1)/\gamma_1$  phase boundary at low temperature and an artificial eutectoid  $(\delta_1 = \zeta + \gamma_1)$  at -60.4 °C. In addition, one of the liquid phase parameters was selected in a way making the enthalpy of mixing temperature dependent [4], a dependency that is not supported by any experimental data. Witusiewicz et al. [7] remediated both artifacts by revising the parameters of the liquid and the  $\gamma_1$  phase, keeping all other parameters from [4]. However, all invariant reaction temperatures of these two calculations [4,7] generally agree within 1 K, the largest difference is 1.3 K at the peritectic  $L+\beta=\gamma_2$ in Table 4, calculated at 1013.2 °C in the present work. One should also notice that the temperature of the invariant reaction of  $L+\beta \leftrightarrow \gamma_2$  is determined to be 1037 °C in [8], but 933  $\pm$  2 °C in [9], the present work and previous calculations cannot come close to either of them, thus an intermediate temperature between the reported values is accepted.

## 4.2. Thermodynamic properties of the liquid phase

Fig. 5 shows the reported experimental data [15–21,35] for the enthalpy of mixing of the liquid phase. Disregarding the data from [16,17], all other experimental data are agree well with each other,



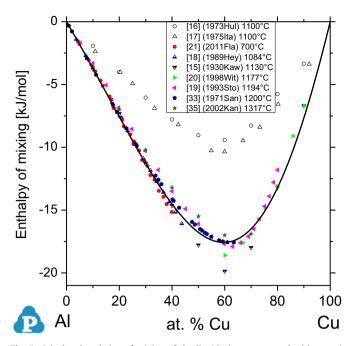
**Fig. 4.** Comparison of the calculated phase diagrams of the Al–Cu system in this work with previously published Calphad assessments [4,7].

and do not show any temperature dependence from 700 °C to 1317 °C. Thus the implicit temperature dependence in the modeling of [4], who did not show any experimental data for comparison at all, is not justified. Actually, the data in [16] were provided as private communication by Yazawa and Itagaki, who are the authors of [17]. Those data are much more positive than the later determined enthalpy of mixing data [18–21,35], and thus not accepted in this work. The calculated partial enthalpies of the components in the liquid phase are also in good agreement with the experimental data as shown in Fig. 6.

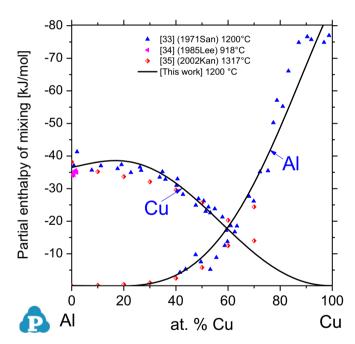
Fig. 7 represents the calculated activity of the components in liquid at 1100 °C in agreement with experimental data from [26–32]. For better resolution of the aluminum activity,  $a_{\rm Al}$ , in the Curich region Fig. 8 shows these data as relative chemical potential,  $\mu_{\rm Al}$ – $G_{Al}^{0,Liquid}$ =RT ln  $a_{\rm Al}$ . It reveals that the relative chemical potential data from [29,31], thus also  $a_{\rm Al}$ , are slightly above the other experimental results [27,30,32], which are in perfect agreement with the calculated curve.

## 4.3. Thermodynamic properties of solid phases

Fig. 9 represents the calculated enthalpy of formation of the solid phases at 500 °C compared with the experimental data and a calculation using the parameters of [7]. Fig. 10 shows the relative chemical potential of Al in the solid phases, the calculated results from this work are compared with experimental data and a calculation using the parameters of [7]. Some deviations are noted and we found that it is difficult to obtain a closer agreement with the experimental data in Fig. 9 and Fig. 10 while maintaining the good agreement with the phase diagram and liquid phase thermodynamic data. The largest deviation is between 63–69 at% Cu, corresponding to the  $\delta_1+\gamma_1$  and the  $\gamma_1$  phase regions, in which the phase boundary is still questionable as mentioned in [9]. Although different models have been applied for  $\delta_1$  and  $\gamma_1$  phases in this work compared to previous Calphad modeling [7], similar deviations between calculated and experimental data are observed.



**Fig. 5.** Calculated enthalpy of mixing of the liquid phase compared with experimental data [16–21,33,35]. Reference states: pure liquid elements.



**Fig. 6.** Calculated partial enthalpy of the elements in the liquid phase (at 1200 °C) compared with experimental data [33–35]. Reference states: pure liquid elements.

## 5. Conclusion

The key improvements of this thermodynamic re-assessment compared to previous Calphad descriptions are:

- For the first time all available experimental data of the Al-Cu system have been extensively compared with calculated results.
- The model selected for the  $\gamma_1$  and  $\gamma_2$  phases partially reflects the crystal structure of the  $\gamma$  phase while being sufficiently simple for a realistic joining with the  $\gamma$ -CuZn phase in the Al-Cu-Zn ternary system.
- The  $\gamma_1/\gamma_2$  transition, modeled as very narrow two-phase region, is very close to the experimentally determined order–disorder

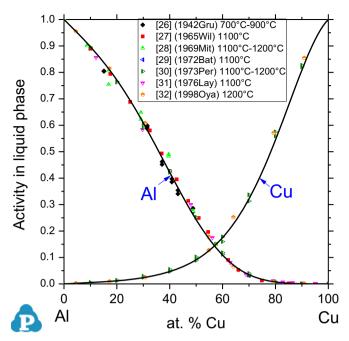
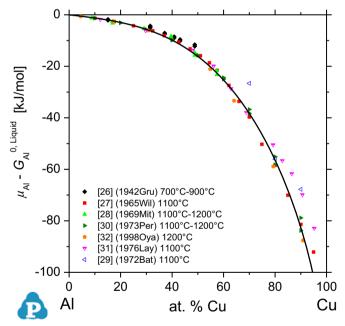


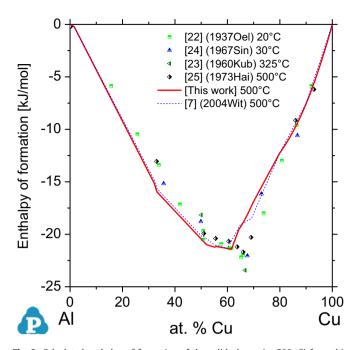
Fig. 7. Calculated activity of the elements in the liquid phase (at  $1100 \,^{\circ}$ C) agree well with the experimental data [26–32]. Reference states: pure liquid elements.



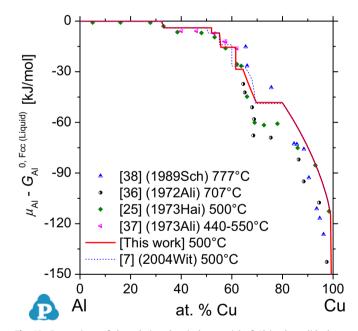
**Fig. 8.** Calculated relative chemical potential of Al (at  $1100\,^{\circ}$ C) in the liquid phase compared with experimental data [26–32]. Reference state: liquid Al.

transition line [9].

- The  $\eta_1/\eta_2$  phase transition has been introduced, and the temperature of the invariant reactions associated with those two phases agree well with experimental thermal signals.
- The extent of single-phase ε<sub>2</sub> phase in this work agrees better with experimental data.
- The stoichiometry Al<sub>5</sub>Cu<sub>8</sub> for the δ<sub>1</sub> phase is closer to the experimentally determined composition.
- The  $\alpha_2$  phase forming below 363 °C has been introduced.
- Very good agreement with the extensive thermodynamic properties of the liquid Al–Cu phase and reasonable agreement for solid phase thermodynamic properties is demonstrated in detail.



**Fig. 9.** Calculated enthalpy of formation of the solid phases (at  $500\,^{\circ}$ C) from this work and [7] compared with experimental data [22–25] Reference states: solid Al (Fcc) and Cu (Fcc).



**Fig. 10.** Comparison of the relative chemical potential of Al in the solid phases obtained by experiments [25,36-38] and calculated using the present thermodynamic description. Reference state changes from solid Al (Fcc) at 500 °C to liquid Al for the experimental data at 707 °C and 777 °C.

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# Appendix A. Supplementary material

Supplementary data associated with this article can be found in the online version at http://dx.doi.org/10.1016/j.calphad.2015.10.004.

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