

1. ρ is density, and v_i is flow velocity. Prove the following two equations (hint: the continuity equation is used):

$$(a) \quad \frac{\partial(\rho Q_{lm..})}{\partial t} + \frac{\partial(\rho v_j Q_{lm..})}{\partial x_j} = \rho \frac{DQ_{lm..}}{Dt}$$

$$(b) \quad v_i \frac{\partial(\rho v_i)}{\partial t} + v_i \frac{\partial(\rho v_j v_i)}{\partial x_j} = \frac{\partial(\frac{1}{2} \rho v^2)}{\partial t} + \frac{\partial(\frac{1}{2} \rho v_j v^2)}{\partial x_j}$$

The proof of equation (a) begins by the definition of the material derivative, as

$$\frac{D(\cdot)}{Dt} = \frac{\partial(\cdot)}{\partial t} + \mathbf{v} \cdot \nabla(\cdot).$$

We shall take into the above equation, as

$$\rho \frac{DQ}{Dt} = \rho \frac{\partial Q}{\partial t} + \rho \mathbf{v} \cdot \nabla Q,$$

and

$$\left[\frac{\partial \rho Q}{\partial t} \right]_{lm} = \frac{\partial(\rho Q_{lm})}{\partial t} \quad \text{and} \quad [\nabla \cdot (\rho Q \otimes \mathbf{v})]_{lm} = \frac{\partial(\rho Q_{lm} v_j)}{\partial x_j}.$$

$$\frac{\partial(\rho Q)}{\partial t} = \frac{\partial \rho}{\partial t} Q + \rho \frac{\partial Q}{\partial t} \quad \text{and} \quad \nabla \cdot (Q \otimes \rho \mathbf{v}) = (\nabla \cdot (\rho \mathbf{v})) Q + \nabla Q \cdot \rho \mathbf{v}.$$

$$\frac{\partial(\rho Q)}{\partial t} + \nabla \cdot (Q \otimes \rho \mathbf{v}) = \rho \frac{\partial Q}{\partial t} + \left(\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) \right) Q + \nabla Q \cdot \rho \mathbf{v} = \rho \frac{DQ}{Dt},$$

with the continuity equation used. So that

$$\rho \frac{DQ_{lm}}{Dt} = \frac{\partial(\rho Q_{lm})}{\partial t} + \frac{\partial \rho v_j Q_{lm}}{\partial x_j}.$$

The proof of equation (b) begins by

$$\frac{\partial(\frac{1}{2} \rho v^2)}{\partial t} = \frac{1}{2} \rho \frac{\partial v^2}{\partial t} + |\mathbf{v}|^2 \frac{\partial(\frac{1}{2} \rho)}{\partial t} = v_i \rho \frac{\partial v_i}{\partial t} + |\mathbf{v}|^2 \frac{\partial(\frac{1}{2} \rho)}{\partial t}$$

$$v_i \frac{\partial(\rho v_i)}{\partial t} = |\mathbf{v}|^2 \frac{\partial \rho}{\partial t} + \rho v_i \frac{\partial v_i}{\partial t}$$

$$\frac{\partial(\frac{1}{2} \rho v_j v^2)}{\partial x_j} = \nabla \cdot \left(\frac{1}{2} \rho \mathbf{v} |\mathbf{v}|^2 \right) = |\mathbf{v}|^2 \nabla \cdot \left(\frac{1}{2} \rho \mathbf{v} \right) + \mathbf{v} \cdot \nabla \left(\frac{1}{2} \rho |\mathbf{v}|^2 \right) = |\mathbf{v}|^2 \nabla \cdot \left(\frac{1}{2} \rho \mathbf{v} \right) + \mathbf{v} \cdot \nabla(\rho \mathbf{v}) \cdot \mathbf{v}.$$

$$v_i \frac{\partial(\rho v_j v_i)}{\partial x_j} = \mathbf{v} \cdot \nabla \cdot (\rho \mathbf{v} \otimes \mathbf{v}) = \mathbf{v} \cdot (\nabla(\rho \mathbf{v}) \cdot \mathbf{v} + \mathbf{v} \nabla \cdot (\rho \mathbf{v})) = \mathbf{v} \cdot \nabla(\rho \mathbf{v}) \cdot \mathbf{v} + |\mathbf{v}|^2 \nabla \cdot (\rho \mathbf{v}).$$

Combine them, we get

$$\frac{\partial(\frac{1}{2} \rho v^2)}{\partial t} + \frac{\partial(\frac{1}{2} \rho v_j v^2)}{\partial x_j} = v_i \rho \frac{\partial v_i}{\partial t} + |\mathbf{v}|^2 \frac{\partial(\frac{1}{2} \rho)}{\partial t} + |\mathbf{v}|^2 \nabla \cdot \left(\frac{1}{2} \rho \mathbf{v} \right) + \mathbf{v} \cdot \nabla(\rho \mathbf{v}) \cdot \mathbf{v}.$$

$$v_i \frac{\partial(\rho v_i)}{\partial t} + v_i \frac{\partial(\rho v_j v_i)}{\partial x_j} = \rho v_i \frac{\partial v_i}{\partial t} + |\mathbf{v}|^2 \frac{\partial \rho}{\partial t} + |\mathbf{v}|^2 \nabla \cdot (\rho \mathbf{v}) + \mathbf{v} \cdot \nabla(\rho \mathbf{v}) \cdot \mathbf{v}.$$

The purple parts are the continuity equation, which are vanished. So that

$$\frac{\partial(\frac{1}{2} \rho v^2)}{\partial t} + \frac{\partial(\frac{1}{2} \rho v_j v^2)}{\partial x_j} = v_i \frac{\partial(\rho v_i)}{\partial t} + v_i \frac{\partial(\rho v_j v_i)}{\partial x_j}.$$

2. The volume flow Q of a centrifugal pump is a function of the input power P , impeller diameter D , rotational rate Ω , and the density ρ and viscosity μ of the fluid:

$$Q = f(P, D, \Omega, \rho, \mu)$$

Rewrite this as a dimensionless relationship as below with Ω , ρ , and D as repeating variables.

$$\frac{Q}{\Omega D^3} = f\left(\frac{P}{\rho \Omega^3 D^5}, \frac{\mu}{\rho \Omega D^2}\right)$$

Use the dimensional analysis, as

$$[Q] = \text{m}^3/\text{s}, \quad [P] = \text{W} = \text{J}/\text{s} = \text{kg} \cdot \text{m}^2/\text{s}^3, \quad [D] = \text{m}, \quad [\Omega] = \text{s}^{-1}, \quad [\rho] = \text{kg}/\text{m}^3, \quad [\mu] = \text{kg}/(\text{m} \cdot \text{s})$$

Choose three unrelated quantities (the dimensions are unrelated) as basis, as

$$[D] = \text{m}, \quad [\Omega] = \text{s}^{-1}, \quad [\rho] = \text{kg}/\text{m}^3.$$

There exist three other quantities, we can construct the dimension equation respectively, as

$$[Q] = [D]^\alpha [\Omega]^\beta [\rho]^\gamma \Rightarrow \text{m}^3/\text{s} = \text{m}^\alpha \text{s}^{-\beta} (\text{kg}/\text{m}^3)^\gamma \Rightarrow \alpha = 3, \beta = 1, \gamma = 0.$$

So that $Q/(D^3\Omega)$ is one dimensionless quantity.

$$[P] = [D]^\alpha [\Omega]^\beta [\rho]^\gamma \Rightarrow \text{kg} \cdot \text{m}^2/\text{s}^3 = \text{m}^\alpha \text{s}^{-\beta} (\text{kg}/\text{m}^3)^\gamma \Rightarrow \alpha = 5, \beta = 3, \gamma = 1.$$

So that $P/(\rho D^5 \Omega^3)$ is one dimensionless quantity.

$$[\mu] = [D]^\alpha [\Omega]^\beta [\rho]^\gamma \Rightarrow \text{kg}/(\text{m} \cdot \text{s}) = \text{m}^\alpha \text{s}^{-\beta} (\text{kg}/\text{m}^3)^\gamma \Rightarrow \alpha = 2, \beta = 1, \gamma = 1.$$

So that $\mu/(\rho D^2 \Omega)$ is one dimensionless quantity. We can construct the dimensionless relation, as

$$\frac{Q}{\Omega D^3} = f\left(\frac{P}{\rho \Omega^3 D^5}, \frac{\mu}{\rho \Omega D^2}\right).$$

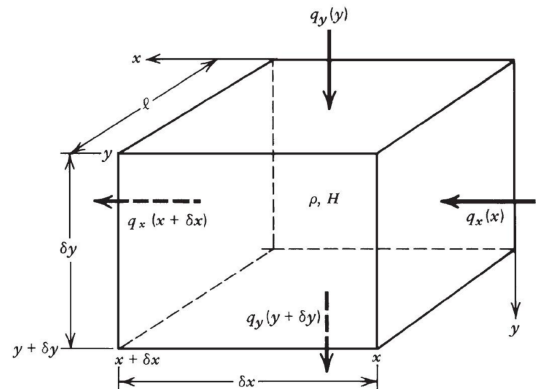
3. One important mathematical physical equations is the thermal conduction equation (a type of parabolic equation) that describes temperature distribution due to thermal conduction inside a solid body. The thermal conduction equation can be derived by neglecting the advection from the energy conservation equation:

$$\rho c_p \frac{DT}{Dt} = -\frac{\partial q_i}{\partial x_i} + \tau_{ij} \dot{\epsilon}_{ij} + \alpha T \frac{Dp}{Dt} + \rho H \quad \rightarrow \quad \rho c_p \frac{\partial T}{\partial t} = -\frac{\partial q_i}{\partial x_i} + \rho H$$

Please derive the thermal conduction equation

$$\rho c_p \frac{\partial T}{\partial t} = k \nabla^2 T + \rho H$$

by conducting thermal balance analysis for a element on the right:



The heat change in this element is

$$\rho c_p \frac{\partial T}{\partial t} \delta x \delta y \delta z = \rho c_p \frac{\partial T}{\partial t} \delta V.$$

The energy change of the element due to the outflow is

$$\left(\frac{q_x(x) - q_x(x + \delta x)}{\delta x} + \frac{q_y(y) - q_y(y + \delta y)}{\delta y} + \frac{q_z(z) - q_z(z + \delta z)}{\delta z} \right) \delta x \delta y \delta z + \rho H \delta x \delta y \delta z = -(q_{x,x} + q_{y,y} + q_{z,z} + \rho H)$$

So that, the energy conservation reads

$$\rho c_p \frac{\partial T}{\partial t} \delta V = (-\nabla \cdot \mathbf{q} + \rho H) \delta V.$$

The Fourier's law reads

$$\mathbf{q} = -k \nabla T.$$

The final conservation becomes

$$\rho c_p \frac{\partial T}{\partial t} = k \nabla^2 T + \rho H.$$