

§3 电磁感应

§3.1 基本规律: Faraday 电磁感应定律

一、Faraday 电磁感应定律

(1) 电磁感应现象:

闭合回路
+
磁通量变化 \longrightarrow 感应电动势

(2) 定量描述 (Neumann, Ernest)

Thm - Faraday 电磁感应定律:

单匝闭合回路磁通量 Φ , 感应电动势 $\varepsilon \propto -\frac{d\Phi}{dt}$

在 SI 下, 比例系数取 1, 即 $\varepsilon = -k \frac{d\Phi}{dt}$

(3) 全回路

① 设 N 匝线圈磁通量 Φ_i , 全磁通 $\Psi = \sum_{i=1}^N \Phi_i$

② 全回路电磁感应定律 $\varepsilon = -\frac{d\Psi}{dt}$

二、Lenz 定律

(1) 感应电动势定向问题

① 选定回路绕行方向, 以右手定则确定其正法线方向 \hat{n} , $\Phi = \vec{B} \cdot S\hat{n}$

② 电动势正负 (产生感应电流与绕行方向相同为正) 与 $\frac{d\Phi}{dt}$ 正负总相反



(2) 感应电流定向

Thm - Lenz 定律:

闭合回路激发感应电流产生磁场, 总是阻碍引起电磁感应外磁场造成磁通量变化.

(3) 本质: 能量守恒

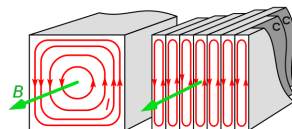
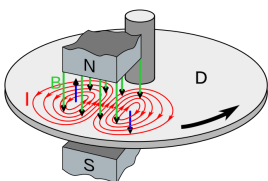
感应电流 $\left\{ \begin{array}{l} \text{产生焦耳热, } Q = \int i^2 R dt \\ \text{阻碍磁通量变化} \end{array} \right.$

\rightarrow 为使 Φ 变化必须克服阻力, 付出功 W

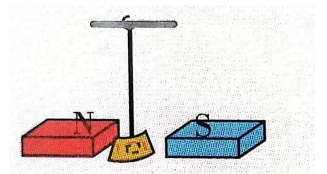
\Rightarrow 有阻碍才能解释 W , 从而 Q 不是无中生有

(4) 应用

① 涡流



② 电磁阻尼



§3.2 动生、感生电动势

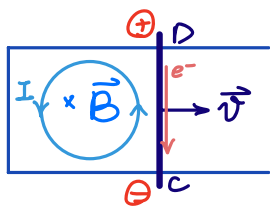
一、动生电动势

(1) 单棒切割模型.

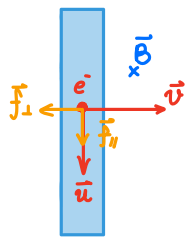
① 推导: e^- 受 Lorentz 力 $\vec{F} = -e\vec{v} \times \vec{B}$

非静电力 $\vec{K} = \frac{\vec{F}}{e} = \vec{v} \times \vec{B}$

动生电动势 $\mathcal{E} = \int_c^D \vec{K} \cdot d\vec{l} = \int_c^D (\vec{v} \times \vec{B}) \cdot d\vec{l}$

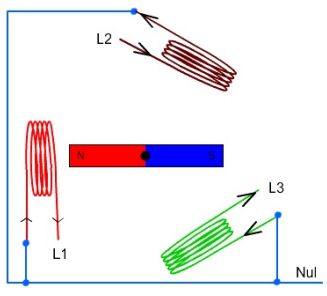
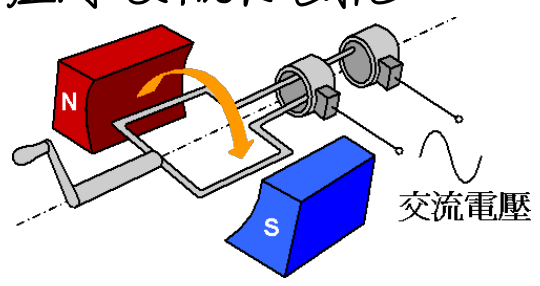


② 能量分析: $\begin{cases} dW_{\perp} = -e v B v dt \\ dW_{\parallel} = e v B u dt \end{cases} \Rightarrow dW_{\perp} + dW_{\parallel} = 0$



(2) 回路动生电动势 $\mathcal{E} = \oint_L (\vec{v} \times \vec{B}) \cdot d\vec{r}$

(3) 应用: 交流发电机



二、感生电动势

(1) 感生电动势

① 产生: 固定回路 $L = \partial \Sigma$, 磁场变化引起磁通量变化, 产生感生电动势.

② 推导: $\mathcal{E} = -\frac{d}{dt} \iint_{\Sigma} \vec{B} \cdot d^2\vec{r} = -\iint_{\Sigma} \frac{\partial \vec{B}}{\partial t} \cdot d^2\vec{r} = -\oint_L \frac{\partial \vec{A}}{\partial t} \cdot d\vec{r}$

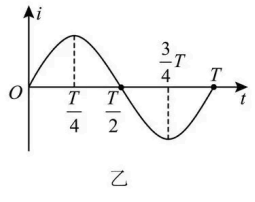
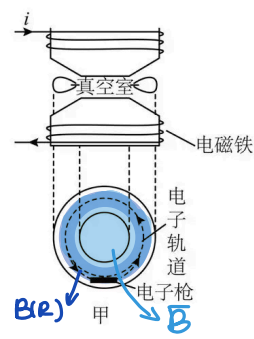
(2) 感生电场

① 非静电力 $\vec{K} = -\frac{\partial \vec{A}}{\partial t}$

② 涡旋电场力提供非静电力. $\vec{E}_{\text{涡}} = -\frac{\partial \vec{A}}{\partial t}$

③ 电场拓展: 静电场 + 涡旋电场 $\Rightarrow \vec{E} = -\nabla\phi - \frac{\partial \vec{A}}{\partial t}$

(3) 应用: 电子感应加速器



e^- 旋转周期 $\tau \ll T$

$m v = e R B(R) = \int e \frac{\pi R^2 \frac{dB}{dt}}{2\pi R} dt = \frac{1}{2} e \bar{B} R$

$\Rightarrow B(R) = \frac{1}{2} \bar{B}$

§3.3 自感与互感

一、互感

(1) 逻辑: 线圈1电流 $I_1 \Rightarrow$ 线圈2磁通 $\Phi_{12}(\psi_{12}) \Rightarrow$ 感应电动势 \mathcal{E}_2

(2) 定量关系: 1对2互感系数 M_{12} : $\psi_{12} = M_{12} I_1 \Rightarrow \psi_{12} \propto I_1$
感应电动势 $\mathcal{E}_{12} = -M_{12} \frac{dI_1}{dt}$
Biot-Savart

(3) 影响因素: 两线圈匝数、大小、形状、相对位置

(4) 两线圈间互感系数

$$\psi_{12} = M_{12} I_1 \quad \& \quad \psi_{21} = M_{21} I_2$$

$$M_{12} = \frac{\psi_{12}}{I_1} = \frac{N_2}{I_1} \oint_{L_2} \vec{A}_0 \cdot d\vec{L}_2 = \frac{N_2}{I_1} \oint_{L_2} \left(\frac{\mu_0 N_1 I_1}{4\pi} \oint_{L_1} \frac{d\vec{L}_1}{r_{12}} \right) \cdot d\vec{L}_2$$

$$= \frac{\mu_0 N_1 N_2}{4\pi} \oint_{L_1} \oint_{L_2} \frac{d\vec{L}_1 \cdot d\vec{L}_2}{r_{12}}$$

可交换 \Rightarrow 统一互感系数 $M_{12} = M_{21} = M = \frac{\mu_0 N_1 N_2}{4\pi} \oint_{L_1} \oint_{L_2} \frac{d\vec{L}_1 \cdot d\vec{L}_2}{r_{12}}$

(5) 单位: 亨利 (H), $1H = 1Wb/A$

二、自感

(1) 逻辑: 因线圈电流变化 $\frac{dI}{dt} \Rightarrow$ 线圈自身感应电动势 \mathcal{E}

(2) 定量关系: 自感系数 L : $\psi = LI$

自感电动势 $\mathcal{E} = -L \frac{dI}{dt}$

(3) 影响因素: 无磁介质, 与 I 无关, 仅取决于线圈形状、尺寸

三、多线圈耦合

(1) 两线圈互感

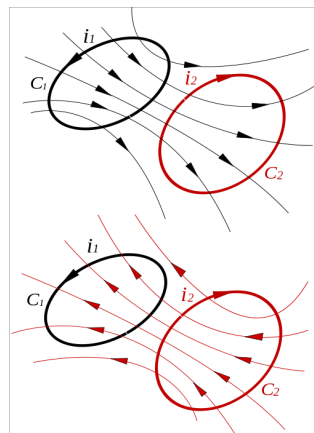
$$\begin{cases} L_1 = \frac{N_1 \Phi_{11}}{I_1} \\ L_2 = \frac{N_2 \Phi_{22}}{I_2} \end{cases} \quad \begin{matrix} \Phi_{21} \geq \Phi_{21} \\ \Phi_{12} \geq \Phi_{12} \end{matrix} \quad \begin{matrix} M = \frac{N_1 \Phi_{21}}{I_2} \Rightarrow \mathcal{E}_1 = \mathcal{E}_{11} + \mathcal{E}_{21} \\ M = \frac{N_2 \Phi_{12}}{I_1} \Rightarrow \mathcal{E}_2 = \mathcal{E}_{22} + \mathcal{E}_{12} \end{matrix}$$

$$\text{漏磁} \Rightarrow \begin{cases} \Phi_{12} = k_1 \Phi_{11}, k_{1,2} \in [0, 1] \\ \Phi_{21} = k_2 \Phi_{22} \end{cases}$$

$$\Rightarrow M^2 = \frac{N_1 N_2}{I_1 I_2} k_1 k_2 \Phi_{11} \Phi_{22} = k_1 k_2 L_1 L_2$$

$$\Rightarrow M = \sqrt{k_1 k_2} \sqrt{L_1 L_2} \triangleq K \sqrt{L_1 L_2}, K \in [0, 1] \text{ 称为耦合系数}$$

$$\text{完全无漏磁} \Rightarrow M = \sqrt{L_1 L_2}$$



(2) 两线圈串联

① 顺接: \mathcal{E}_{i1} 与 \mathcal{E}_{i2} 同向 ($i,j=1,2$)

$$\mathcal{E}_i = -L_i \frac{dI}{dt} - M \frac{dI}{dt} \quad (i=1,2)$$

$$\Rightarrow \mathcal{E} = -(L_1 + L_2 + 2M) \frac{dI}{dt}$$

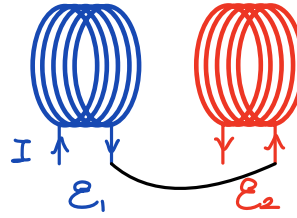
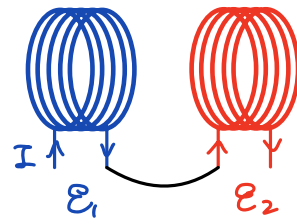
$$\Rightarrow L_{\text{顺}} = L_1 + L_2 + 2M$$

② 反接: \mathcal{E}_{i1} 与 \mathcal{E}_{i2} 反向 ($i,j=1,2$)

$$\mathcal{E}_i = -L_i \frac{dI}{dt} + M \frac{dI}{dt} \quad (i=1,2)$$

$$\Rightarrow \mathcal{E} = -(L_1 + L_2 - 2M) \frac{dI}{dt}$$

$$\Rightarrow L_{\text{反}} = L_1 + L_2 - 2M$$



(3) 两线圈并联

① 同名端并联: \mathcal{E}_{i1} 与 \mathcal{E}_{i2} 同向

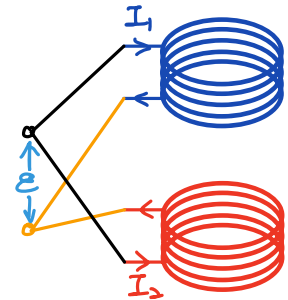
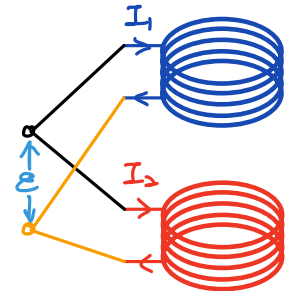
$$\begin{cases} \mathcal{E}_1 = -L_1 \frac{dI_1}{dt} - M \frac{dI_2}{dt} \\ \mathcal{E}_2 = -L_2 \frac{dI_2}{dt} - M \frac{dI_1}{dt} \end{cases}$$

$$\mathcal{E} = -L_{\text{同}} \left(\frac{dI_1}{dt} + \frac{dI_2}{dt} \right) \Rightarrow L_{\text{同}} = \frac{L_1 L_2 - M^2}{L_1 + L_2 - 2M}$$

② 异名端并联: \mathcal{E}_{i1} 与 \mathcal{E}_{i2} 反向

$$\begin{cases} \mathcal{E}_1 = -L_1 \frac{dI_1}{dt} + M \frac{dI_2}{dt} \\ \mathcal{E}_2 = -L_2 \frac{dI_2}{dt} + M \frac{dI_1}{dt} \end{cases}$$

$$\mathcal{E} = -L_{\text{异}} \left(\frac{dI_1}{dt} + \frac{dI_2}{dt} \right) \Rightarrow L_{\text{异}} = \frac{L_1 L_2 - M^2}{L_1 + L_2 + 2M}$$



(4) 零耦合

① 串联 $L = \sum L_i$ ② 并联 $\frac{1}{L} = \sum \frac{1}{L_i}$

四. 磁储能

(1) 自感磁能

$$dA = -\mathcal{E} i dt = L i di \Rightarrow A = L \int_0^I i di = \frac{1}{2} L I^2$$

$$\text{电感 } L \text{ 电流 } I, \text{ 自能 } W_L = \frac{1}{2} L I^2$$

(2) 互感磁能

$$dA = -\mathcal{E}_{12} i_2 dt - \mathcal{E}_{21} i_1 dt = M d(i_1 i_2) \Rightarrow A = M I_1 I_2$$

$$\text{互感 } M, \text{ 电流 } I_{1,2}, \text{ 互能 } W_m = M I_1 I_2$$

(3) N 线圈, $W = \frac{1}{2} \sum_{i=1}^N L_i I_i^2 + \sum_{1 \leq i < j \leq N} M_{ij} I_i I_j$

电磁感应 $\mathcal{E} = - \overset{\text{Lenz}}{\frac{d\psi}{dt}} \overset{\text{Faraday}}{\quad}$

$$= - \int_{\Sigma} \frac{\partial \vec{B}}{\partial t} \cdot d^2\vec{r} + \oint_{\partial\Sigma} (\vec{v} \times \vec{B}) \cdot d\vec{r}$$

$$= - \underbrace{\oint_{\partial\Sigma} \frac{\partial \vec{A}}{\partial t} \cdot d\vec{r}}_{\text{感生}} + \underbrace{\oint_{\partial\Sigma} (\vec{v} \times \vec{B}) \cdot d\vec{r}}_{\text{动生}}$$

电势

$$\left\{ \begin{array}{l} \text{自感: } \psi = LI \Rightarrow \mathcal{E} = -L \frac{dI}{dt} \\ \text{互感: } \psi_{ij} = M_{ij} I_j \Rightarrow \mathcal{E}_{ij} = -M_{ij} \frac{dI_j}{dt} \\ M_{ij} = M_{ji} = \frac{\mu_0 N_i N_j}{4\pi} \oint_{\gamma_i} \oint_{\gamma_j} \frac{d\vec{l}_i \cdot d\vec{l}_j}{r_{ij}} \end{array} \right.$$

耦合: $M_{ij} = K \sqrt{L_i L_j}, K \in [0, 1]$



$$\left\{ \begin{array}{l} \text{串} \left\{ \begin{array}{l} \text{顺 } L = L_1 + L_2 + 2M \\ \text{反 } L = L_1 + L_2 - 2M \end{array} \right. \Leftrightarrow \text{并} \left\{ \begin{array}{l} \text{同 } L = \frac{L_1 L_2 - M^2}{L_1 + L_2 - 2M} \\ \text{异 } L = \frac{L_1 L_2 - M^2}{L_1 + L_2 + 2M} \end{array} \right. \end{array} \right.$$

储能 $\left\{ \begin{array}{l} \text{自 } W_L = \frac{1}{2} L I^2 \\ \text{互 } W_M = M I_1 I_2 \end{array} \right. \Rightarrow W_0 = \sum_{i=1}^N \frac{1}{2} L_i I_i^2 + \sum_{1 \leq i < j \leq N} M_{ij} I_i I_j$