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SIMULATED QUANTUM ANISOTROPIC ANNEALING APPLIED TO ARTIFICIAL
NEURAL NETWORK WEIGHT SELECTION

THESIS
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THESIS

Presented to the Faculty of the Electrical and Computer Engineering
of the Air Force Institute of Technology

Air University

In Partial Fulfillment of the
Requirements for the Degree of
Master of Science in Computer Science

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June, 2016

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Preface

Justin Fletcher

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ANN	

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Abstract

SIMULATED QUANTUM ANISOTROPIC ANNEALING APPLIED TO ARTIFICIAL NEURAL NETWORK WEIGHT SELECTION

I. Introduction

In chapter three the traversal a error manifold in the problem configuration space is discussed at length. Several traversal methodologies are proposed and evaluated.

II. Background

This chapter serves as a comprehensive review of the physical and computational concepts material to the topic of this thesis. A broad overview of artificial neural networks and the application and history thereof is presented. Next, simulated annealing is described, along with a summary of some related works and a description of the physical inspiration for the algorithm. The chapter concludes with a very brief overview of the quantum mechanics, with emphasis placed on those concepts which will be employed throughout this thesis document. Finally, the notation and terminology conventions adopted in this document are established.

2.1 Artificial Neural Networks

It has long been recognized (?) that the capacity of biological information processing systems to flexibly and quickly process large quantities of data greatly exceeds that of sequential computing machinery. This information processing capability arises from the complex, nonlinear, parallel nature of biological information processors (Source? Hayken?). The family of models designed to replicate this powerful information processing architecture are collectively called artificial neural networks (ANN). In the most general sense, artificial neural networks are parallel distributed information processors (Hayken, or 1) comprising many simple processing elements. Networks store information about experienced stimuli and can make that information available. In such a network, interneuron connection strengths are used to encode information, and are modified via a learning strategy. Artificial neural networks are characterized by three features: a network topology or architecture, an activation function, and a learning strategy. Each will be discussed in the following sections. I will also review the history of artificial neural networks, and the biological inspiration for the computational model.

2.1.1 Biological Inspiration. Fig[an image of a neuron, mapped to a schematic of a neuron, mapped to a processing element]

Integration of magnitude-encoded, rather than frequency-encoded, signals. Threshold functions relationship to the biological shape, size... Papers needed.

Biological neural networks are many orders of magnitude slower than those based in ... It is not the size of the network or the number of interconnections alone which confer upon the human brain its remarkable efficiency (Faggin, 1991). The

2.1.2 Historical Overview. (McCulloch and Pitts, 1943)

Hebbian stuff with (Rochester, Holland, Haibt, and Duda, 1956)

2.1.3 Network Topology.

2.1.4 Activation Functions.

2.1.5 Learning Strategies. General discussion.

Reference (Mendel and McClaren, 1970) and (Haykin, pg 50)

2.1.5.1 Back Propagation Training. Discuss (Rumelhart, Hinton, Williams, 1986)

Derive back prop. (Haykin, pg 161)

Discuss the implication of local minima.

Fig[Error surface with backprop]

2.1.5.2 Simulated Annealing. (Ackley, Hinton, and Sejnowski 1985)

...This Boltzmann machine is also of historical importance, as it was the first successfully-implemented multilayered neural network. (Ref: Haykin Book)

2.2 Related Works in Simulated Annealing

Read (Haykin pg 556)

Read (Haykin pg 560)

Simulated annealing (SA) is a stochastic optimization algorithm which can be used to find the global minimum of a cost function mapped from the configurations of a combinatorial optimization problem. The concept of simulated annealing was introduced in by

Kirkpatrick et al. in (Kirkpatrick) as an application of the methods of statistical mechanics to the problem of discrete combinatorial optimization. Specifically, simulated annealing is an extension of the Metropolis-Hastings (Metropolis) algorithm which can be used to estimate the ground energy state of a many-body systems at thermal equilibrium. Kirkpatrick et al. applied the Metropolis-Hastings algorithm sequentially, with decreasing temperature values in order to approximate a solid slowly cooling to low temperatures. Later work by Geoff (Geoff) and Cortana et al. (Cortana) extended simulated annealing to the continuous domain. The basic simulated annealing algorithm, as it will be used in this paper, is presented in [algo1].

[algo1]

The physical inspiration for simulated annealing. See (Haykin pg 546)

In the parlance of simulated annealing (Kirkpatrick) a system at its maximum temperature is said to be *melted*. In the melted state, most perturbations of the system configuration are accepted by the algorithm. Analogously, a system that has a temperature of zero, which indicates that the algorithm cannot move to any higher-error state, is said to be *frozen*. Note that a frozen system may still be perturbed into a lower-energy state.

When considering only the influence of classical thermal fluctuations in particle energy levels, the probability of a particle traversing a barrier of height ΔV at a temperature T is on the order of:

$$\mathcal{P}_t = e^{-\frac{\Delta V}{T}} \quad (2.1)$$

2.2.1 Reheating.

2.2.2 Application of Simulated Annealing to ANN Synaptic Weight Selection.

2.3 Related Works in Quantum Mechanics

Quantum mechanics is the branch of physics concerned with the physical laws of nature at very small scales. Many aspects of physical reality are observable only at these scales. Several techniques described in this document are either inspired by, or are simple

models of quantum mechanical processes. These concepts are very briefly reviewed in this section.

2.3.1 Quantum Tunneling. One of the quantum phenomena for which there is no classical analog is potential barrier penetration, also known as quantum tunneling. This phenomenon arises from the probabilistic and wavelike behavior of particles in quantum physics. Tunneling plays a significant role in the behavior of bound and scattering quantum mechanical systems.

A particle with energy E incident upon a potential energy barrier of height $\Delta V > E$ has a non-zero probability of being found in, or past, the barrier. Classically, this behavior is forbidden. The probability of tunneling, \mathcal{P}_t , through a step barrier of height ΔV is described by:

$$\mathcal{P}_t = e^{-\frac{w\sqrt{\Delta V}}{\Gamma}} \quad (2.2)$$

where Γ is the tunneling field strength [Ref: Multivariable Opt: QAC - Mukherjee]. Figure [1] depicts a one-dimensional example of quantum tunneling.

[Figure 1]

2.3.2 Quantum Annealing. Quantum annealing is the use of quantum, rather than thermal fluctuations to traverse the free energy landscape of a system. This is accomplished by introducing an additional Hamiltonian term that does not commute with the classical Hamiltonian. This non-commutation implies that [What does it mean?... The non-commutative causes the quantum effects...but how?]. The term is introduced to account for the presence of a tunneling field which controls the frequency with which quantum fluctuations occur in the system. In effect, this term controls the relative importance of quantum effects on the behavior of the modeled system. This term, much like thermal energy in simulated annealing, is gradually reduced over the course of the simulation. [Par Source: Quantum annealing in a kinetically constrained system] The time dependent

Schrödinger equation ¹ for such a system has the form:

$$[\lambda(t)H' + H_0]\psi = i\hbar \frac{\partial \psi}{\partial t} \quad (2.3)$$

[Eq from Mult Opt... Murherjee] where $\lambda(t)$ is the time-variance function of the tunneling field, H' is the Hamiltonian term describing the tunneling field, and H_0 is the classical Hamiltonian.

The fluctuations induced by the tunneling field are tunneling events, which transition the system from one configuration to a different, lower-energy configuration directly, without assuming any of the higher energy configurations between the two. Said differently, the quantum tunneling field enables the penetration of energy barriers. The addition of these quantum fluctuations also ensures that each possible state of the system can be reached (KCS Paper).

2.4 Notation and Terminology Conventions

There is a great deal of academic writing describing quantum annealing in the language of physics, but very little writing describing the concept from an algorithmic perspective. For this reason a new, more specific term is introduced in this document. Simulated quantum annealing (SQA) is the quantum mechanical counterpart of simulated thermal annealing.

The term neuron will be used in this document to describe the information processing elements of a neural network. This convention is selected both for conciseness and for the useful adjectival form, neural, which will be of great explanatory utility in the coming chapters.

Read (Haykin pg 561) Table 11.1

¹Note that the presence of the Schrödinger equation in section does not imply that quantum annealing requires the annealed system must be an approximation to a wavefunction. It merely serves as an exposition of the properties of physical system which is modeled.

III. Methodology

3.1 Simulated Quantum Annealing

It is instructive to contrast equations 2.1 and 2.2. Both describe the same value, but the importance of the width and height of the traversed barrier in the two equations is considerably different. For systems in which quantum tunneling is possible, the probability of penetrating a barrier of height ΔV is increased by a factor of approximately $e^{\Delta V}$, for large values of ΔV . This relationship is depicted graphically in Fig. 2 which shows the probability of barrier traversal for a system which allows quantum fluctuations, divided by the same probability for a system which only considers thermal fluctuations. Therefore, physical models which considers quantum effects are much more likely predict penetration of tall, thin energy barriers than those which only include classical thermal effects.

[Figure 2]

3.2 Traversing the Error Manifold

3.2.1 Unidimensional Weight Perturbation.

3.2.2 Omni-dimensional Weight Perturbation.

3.2.3 Constrained Step-Size Omni-dimensional Weight Perturbation.

3.2.4 Quantum Weight Perturbation. It is shown in Proof [n] that this algorithm is certain to eventually find the minimum possible

3.2.5 Stochastic-Anisotropic Quantum Weight Perturbation.

3.2.6 Quantum Annealing. [After discussing the way in which the algorithm is implemented] ...The net effect of this design is to allow the algorithm to move from a local minima configuration, to a different, lower-error configuration, without requiring the evaluation of intervening, higher-error configurations. This means that the probability of "tunneling" to a state

IV. Results

Is my algorithm computationally efficient as in Haykin, pg 229?

V. Conclusion

Appendix A. First appendix title

A.1 In an appendix

This is appendix section A.1.

Note: I highly recommend you create each chapter in a separate file including the `\chapter` command and `\include` the file. Then you can use `\includeonly` to process selected chapters and you avoid having to latex/preview/print your entire document every time.

Bibliography

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