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高エネルギー加速器研究機構



原始ブラックホールと宇宙論的降着

Cosmological accretions on Primordial Black  
Holes (PBHs) Kaz Kohri

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# Why PBHs?

- We need  $30 M_{\odot}$  BHs to explain **LIGO/VIRGO** GW events
- We do not know the **origin of  $30 M_{\odot}$  BHs** to be astrophysical or cosmological.
- PBHs should be a good candidate for Cold Dark Matter (**CDM**), but we need to know the full cosmic history.
- Some **inflation** modes predict PBHs formed at small scales in the early Universe (before 1 sec)
- Scenarios have been constrained by **BBN, CMB anisotropies (with accretion), lensing, gamma-ray**, and so on.  
*Carr, Kohri, Sendouda, J.Yokoyama (2010)*
- In future, scenarios can be investigated further by **PIXIE** (CMB  $\mu$ -distortion), **SKA/Ominiscope** (21cm), **CTA** (gamma-ray), **DECIGO** (Gravitational Wave), ...

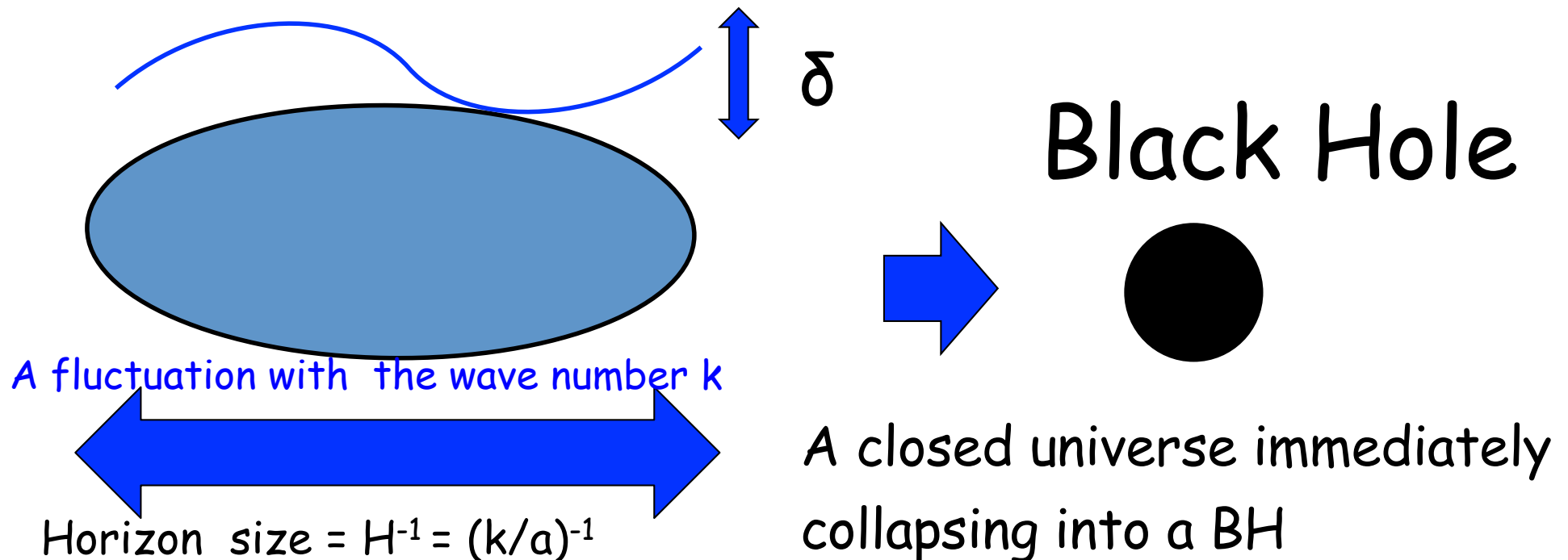
# Conditions for a PBH formation in Radiation dominated (RD) Universe

Zel'dovich and Novikov (1967), Hawking (1971), Carr (1975)

Harada, Yoo and KK (2013)

- Gravity could be stronger than pressure

$$\delta > \delta_c \sim p / \rho \sim c_s^2 = w = 1/3$$



# $P_\zeta(k)$ and PBH abundance $\beta(M)$

- Fraction of PBH to the total at its formation epoch with Gaussian fluctuation.

$$\beta(M) \equiv \frac{\rho_{\text{PBH}}(M)}{\rho_{\text{tot}}} = 2 \int_{\delta_{\text{th}}}^{\infty} d\delta \frac{1}{\sqrt{2\pi\sigma^2}} \exp\left(-\frac{\delta^2}{2\sigma^2}\right) = \text{erfc}\left(\frac{\delta_{\text{th}}}{\sqrt{2}\sigma}\right)$$

$\sim 1/3 - 0.5$

- Finally we have a relation between  $\beta$  and fluctuation  $\sigma$  (or  $\beta$  and  $\Omega$ )

$$\beta(M) \sim \text{erfc}\left(\frac{\delta_{\text{th}}}{\sqrt{2}\sigma}\right) \simeq \sqrt{\frac{2}{\pi}} \frac{\sigma}{\delta_{\text{th}}} \exp\left(-\frac{\delta_{\text{th}}^2}{2\sigma^2}\right)$$

$$= 1.5 \times 10^{-18} \left(\frac{m_{\text{PBH}}}{10^{15} \text{ g}}\right)^{1/2} \left(\frac{\Omega_{\text{PBH}} h^2}{0.1}\right)$$

$\sim P_\zeta$

# Typical quantities of PBHs in RD

- Mass (horizon mass =  $\rho(t_{\text{form}}) H(t_{\text{form}})^{-3}$ )

$$M_{\text{PBH}} \sim M_{\text{pl}}^2 t_{\text{form}} \sim \frac{M_{\text{pl}}^3}{T_{\text{form}}^2} \sim 10^{15} g \left( \frac{T_{\text{form}}}{3 \times 10^8 \text{ GeV}} \right)^{-2} \sim 30 M_{\odot} \left( \frac{T_{\text{form}}}{40 \text{ MeV}} \right)^{-2}$$

- Lifetime

$$\tau_{\text{PBH}} \sim \frac{M_{\text{PBH}}^3}{M_{\text{pl}}^4} \sim 4 \times 10^{17} \text{ sec} \left( \frac{M_{\text{PBH}}}{10^{15} g} \right)^3 \sim 3 \times 10^{68} \text{ yrs} \left( \frac{M_{\text{PBH}}}{30 M_{\odot}} \right)^3$$

- Hawking Temperature

$$T_{\text{PBH}} \sim \frac{M_{\text{pl}}^2}{M_{\text{PBH}}} \sim 0.1 \text{ MeV} \left( \frac{M_{\text{PBH}}}{10^{15} g} \right)^{-1} \sim 3 \times 10^{-11} \text{ K} \left( \frac{M_{\text{PBH}}}{30 M_{\odot}} \right)^{-1}$$

- Wave number of horizon length

$$k = aH \sim 10^5 \text{ Mpc}^{-1} \left( \frac{M_{\text{PBH}}}{10^{15} M_{\odot}} \right)^{-1/2} \sim 10^5 \text{ Mpc}^{-1} \left( \frac{T_{\text{form}}}{\text{MeV}} \right)^{+1}$$

- Fraction to CDM

$$f_{\text{fraction}} \equiv \frac{\Omega_{\text{PBH}}}{\Omega_{\text{CDM}}} \sim \left( \frac{\beta}{10^{-18}} \right) \left( \frac{M_{\text{PBH}}}{10^{15} g} \right)^{-1/2} \sim \left( \frac{\beta}{10^{-8}} \right) \left( \frac{M_{\text{PBH}}}{30 M_{\odot}} \right)^{-1/2} \sim 10^8 \left( \frac{M_{\text{PBH}}}{30 M_{\odot}} \right)^{-1/2} \sqrt{P_{\delta}} \exp \left[ -\frac{1}{18 P_{\delta}} \right]$$



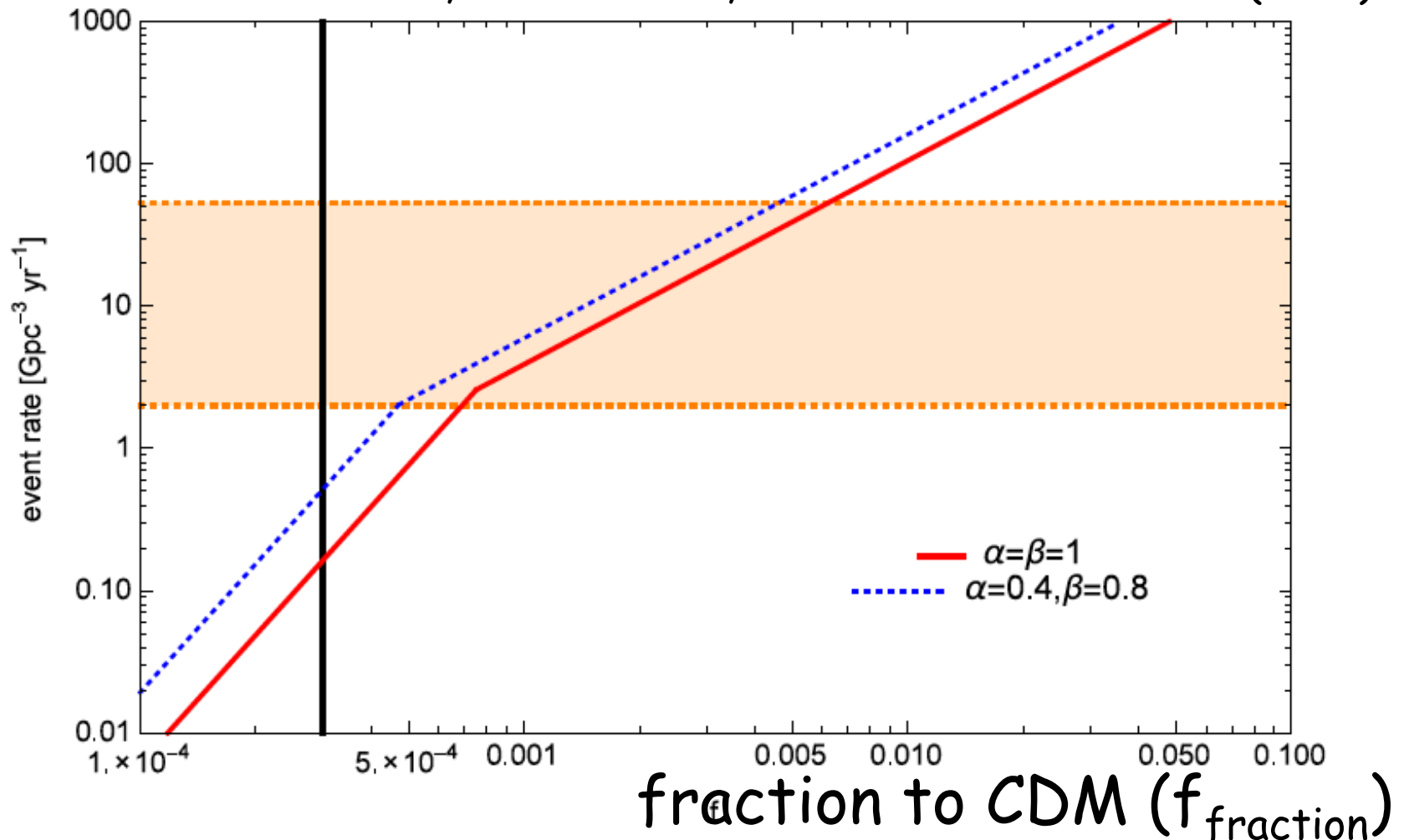
# GW150914 and its merger rate

M. Sasaki, T. Suyama, T. Tanaka and S. Yokoyama (2016).

3-body is important for the  
BBH formations

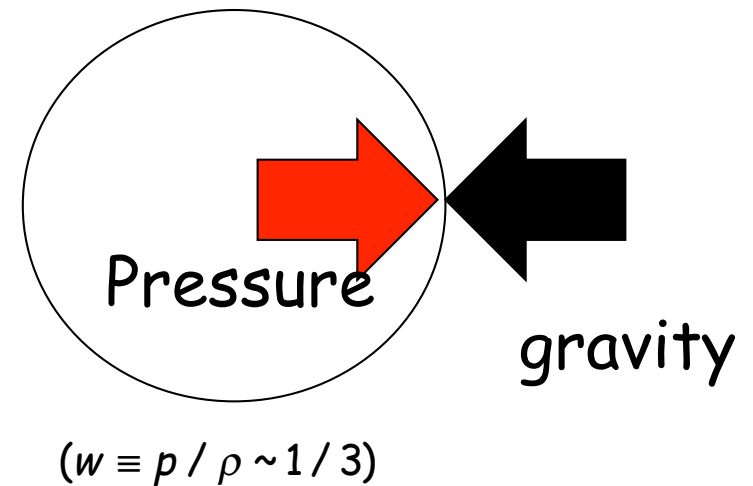
See also, Ali-Haïmoud, Kovetz and Kamionkowski (2017)

Rate of GW140914



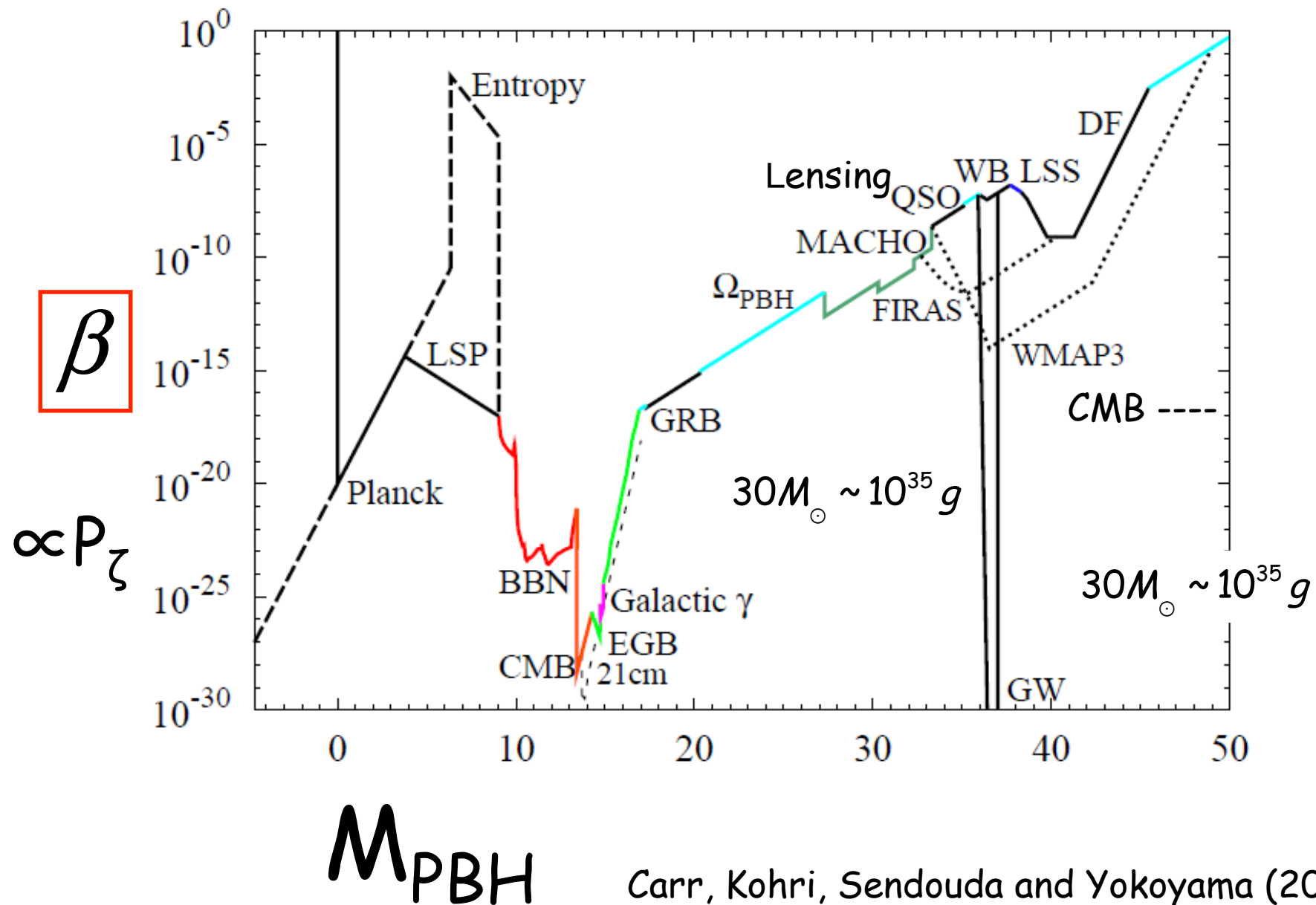
# Features of PBH formations in RD

- Perfectly spherical due to radiation pressure



- No evolutions of density perturbations
- Small angular momentum

# Upper bounds on $\beta$ for PBH



Carr, Kohri, Sendouda and Yokoyama (2010)  
Josan, Green and Malik (2009)

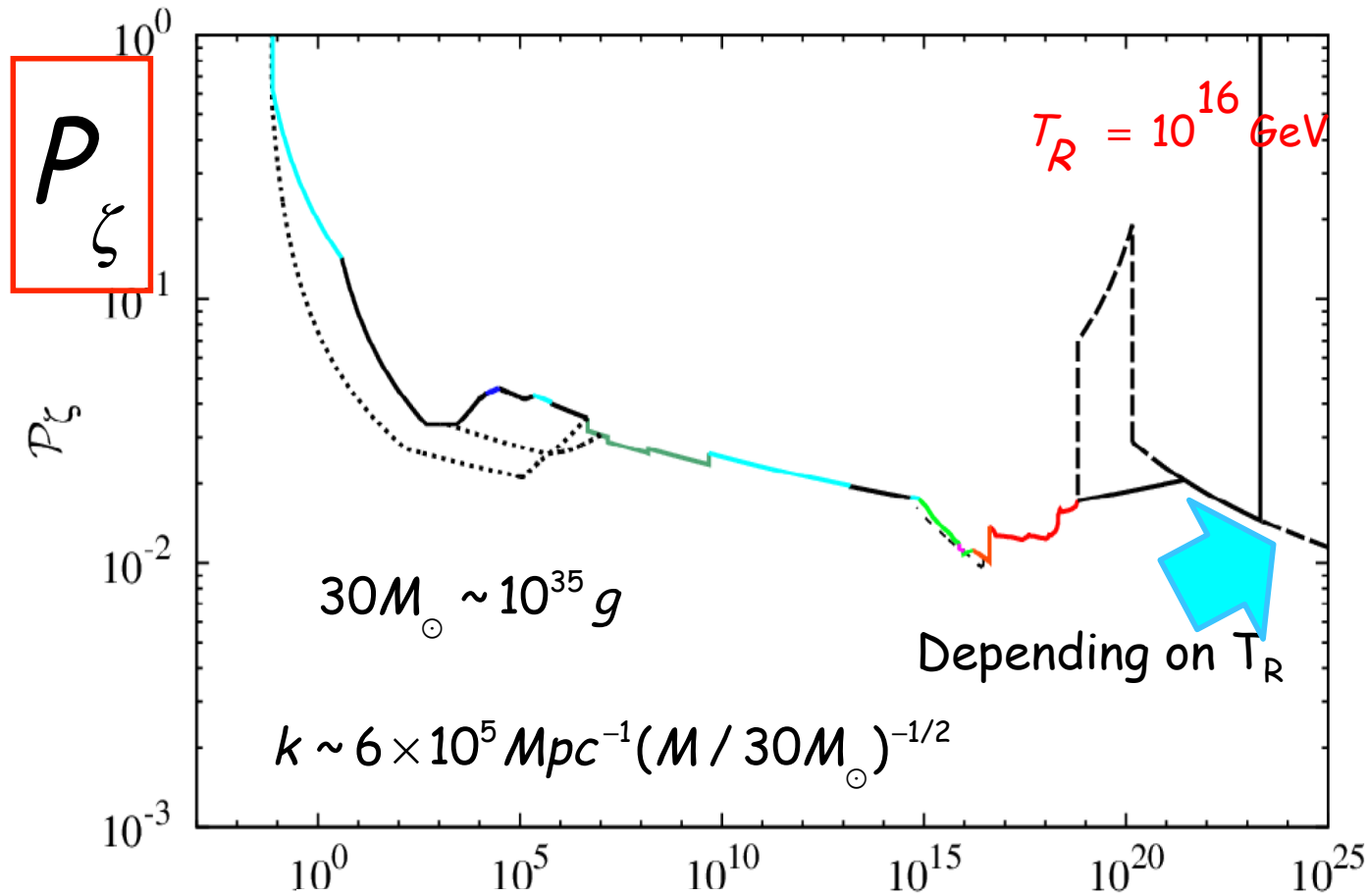


# Upper bounds on $P_\zeta$ for PBH

Carr, KK, Sendouda, J.Yokoyama (2010)

Alabidi, KK, Sasaki, Sendouda (2012)

Amplitude of curvature perturbation



$k \text{ (Mpc}^{-1}\text{)}$  wave number  $k = p \times a$

# CMB bound on PBHs by disk-accretion in the late MD epoch

Poulin, Serpico, Calore, Clesse, KK (2017)

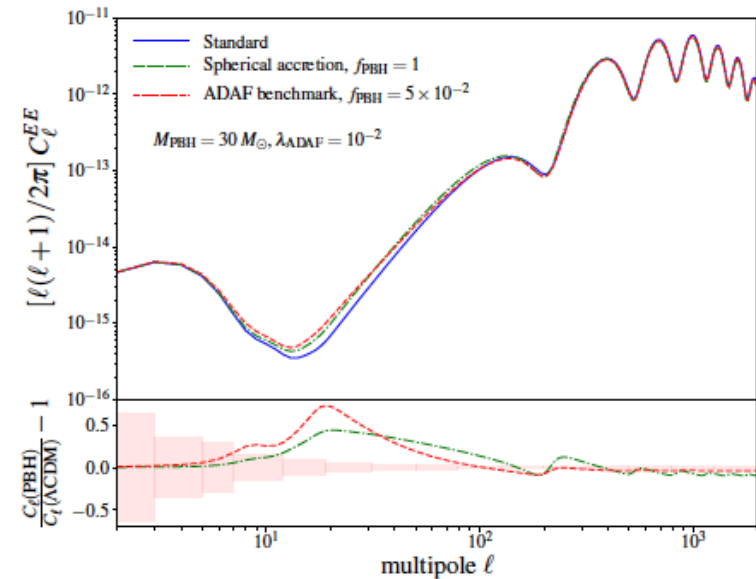
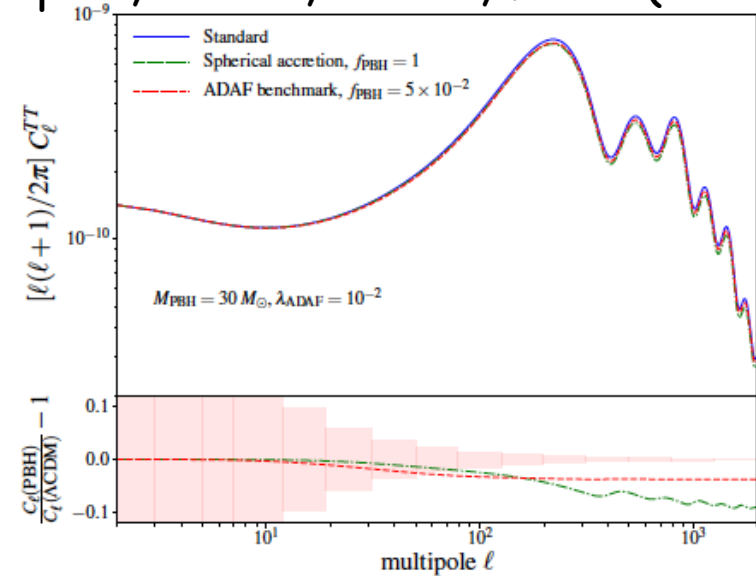
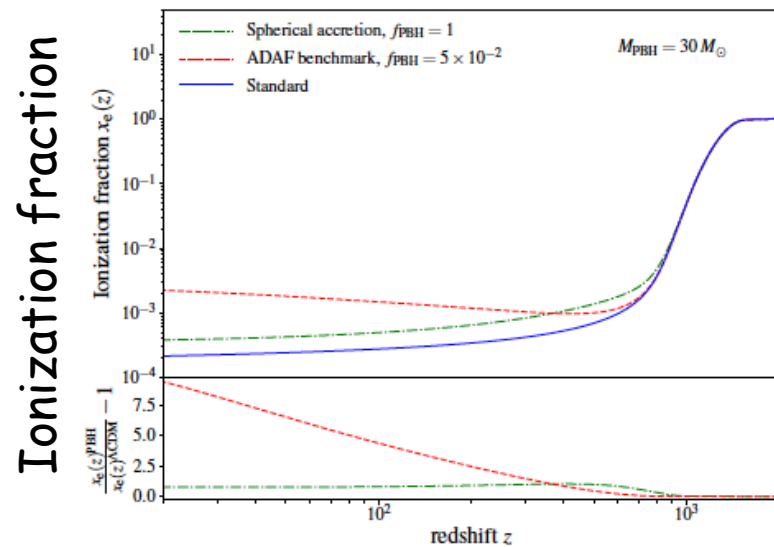
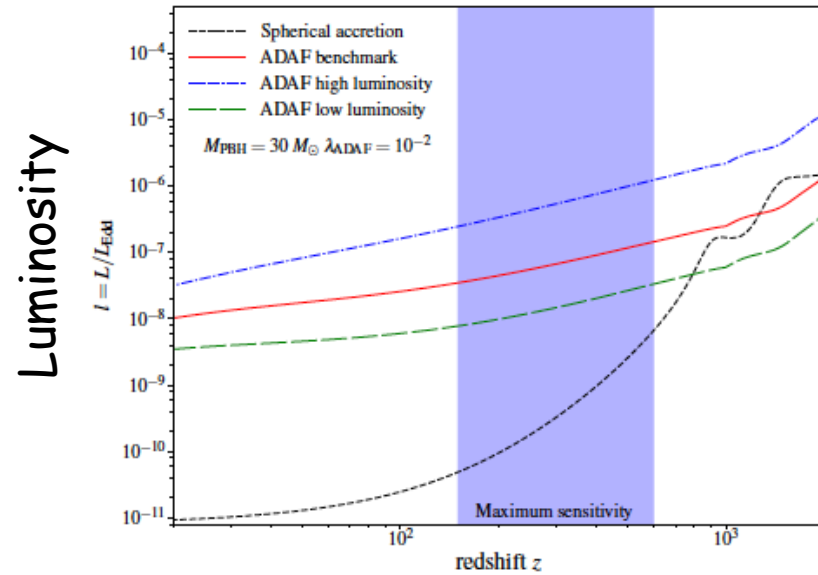
- A non-spherical accretion disk around a PBH caused by an angular momentum emits radiation

$$\dot{M}_{\text{HB}} \equiv 4\pi\lambda\rho_{\infty}v_{\text{eff}}r_{\text{HB}}^2 \equiv 4\pi\lambda\rho_{\infty}\frac{(GM)^2}{v_{\text{eff}}^3}$$
$$l \simeq \omega r_{\text{HB}}^2 \simeq \left(\frac{\delta\rho}{\rho} + \frac{\delta v}{v_{\text{eff}}}\right)v_{\text{eff}}r_{\text{HB}}$$

- CMB anisotropies are affected
- From observations, we can constrain the number density of PBHs.

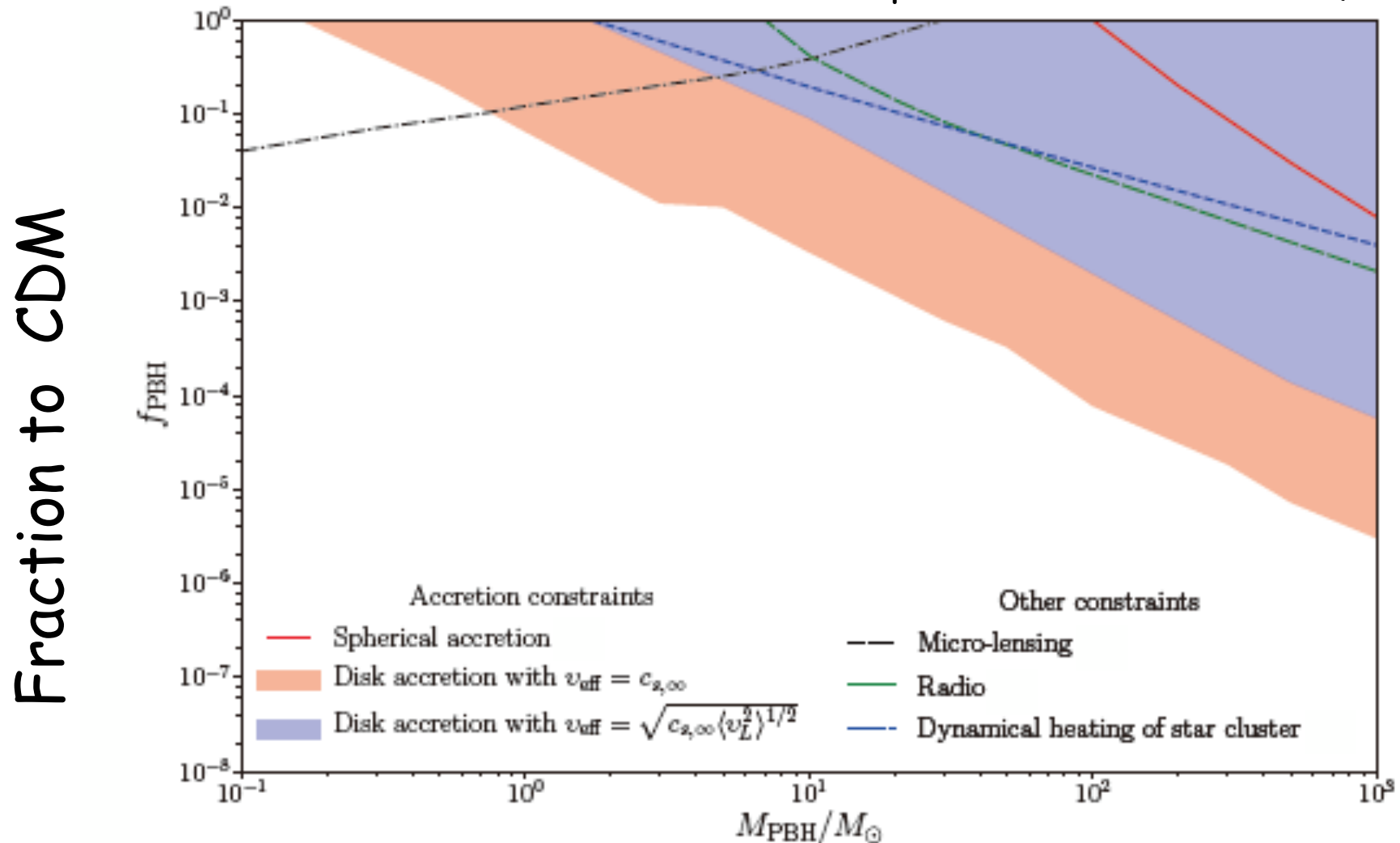
# Modified CMB anisotropy

Poulin, Serpico, Calore, Clesse, Kohri (2017)



# CMB bound by disk-accretion in the latest MD epoch

Poulin, Serpico, Calore, Clesse, KK (2017)



# PBH formation at the (early) matter dominated (MD) Universe

Polnarev and Khlopov (1982)

Harada, Yoo, KK, Nakao, Jhingan (2016)

- **Pressure is zero**, which could induce an immediate collapse and producing more PBHs?
- **Density perturbations can evolve**, which produces non-spherical objects and cannot be included inside the Horizon radius. That means less PBHs can be produced?

# Matter domination

- Three radius in Lagrangian coordinate  $q_i$

$$r_1 = (a - \alpha b)q_1 \quad \text{Zel'dovich Approximation}$$

$$r_2 = (a - \beta b)q_2$$

$$r_3 = (a - \gamma b)q_3$$

- Eccentricity
- Hoop

$$e^2 = 1 - \left( \frac{r_2(t_c)}{r_3(t_c)} \right)^2 = 1 - \left( \frac{\alpha - \beta}{\alpha - \gamma} \right)^2$$

$$\mathcal{C} = 16 \left( 1 - \frac{\gamma}{\alpha} \right) E \left( \sqrt{1 - \left( \frac{\alpha - \beta}{\alpha - \gamma} \right)^2} \right) r_f,$$

- Hoop conjecture for PBH production

$$\mathcal{C} \lesssim 2\pi r_g.$$



# Abundance of PBHs formed in MD

- Probability distribution by peak statistics (BBKS)

Doroshkevich (1970)

$$\begin{aligned}
 & w(\alpha, \beta, \gamma) d\alpha d\beta d\gamma \\
 &= -\frac{27}{8\sqrt{5}\pi\sigma_3^6} \exp \left[ -\frac{1}{10\sigma_3^2}(\alpha + \beta + \gamma)^2 - \frac{1}{4\sigma_3^2} \{(\alpha - \beta)^2 + (\beta - \gamma)^2 + (\gamma - \alpha)^2\} \right] \\
 & \cdot (\alpha - \beta)(\beta - \gamma)(\gamma - \alpha) d\alpha d\beta d\gamma.
 \end{aligned}$$

$$\sigma_H = \sqrt{5}\sigma_3$$

- Probability

$$\beta_0 = \int_0^\infty d\alpha \int_{-\infty}^\alpha d\beta \int_{-\infty}^\beta d\gamma \theta(1 - h(\alpha, \beta, \gamma)) w(\alpha, \beta, \gamma)$$

$$\begin{aligned}
 h(\alpha, \beta, \gamma) &= \frac{2}{\pi} \frac{\alpha - \gamma}{\alpha^2} E \left( \sqrt{1 - \left( \frac{\alpha - \beta}{\alpha - \gamma} \right)^2} \right) \\
 h(\alpha, \beta, \gamma) &:= \mathcal{C} / (2\pi r_g)
 \end{aligned}$$

# Effects by finite angular momentum

Harada, Yoo, KK, Nakao (2017)

- Probability distribution

$$a_* := L/(GM^2/c)$$

$$f_{\text{BH}(2)}(a_*) da_* \propto \frac{1}{a_*^{5/3}} \exp \left( -\frac{1}{2\sigma_H^{2/3}} \left( \frac{2}{5} \mathcal{I} \right)^{4/3} \frac{1}{a_*^{4/3}} \right) da_*$$

- Probability

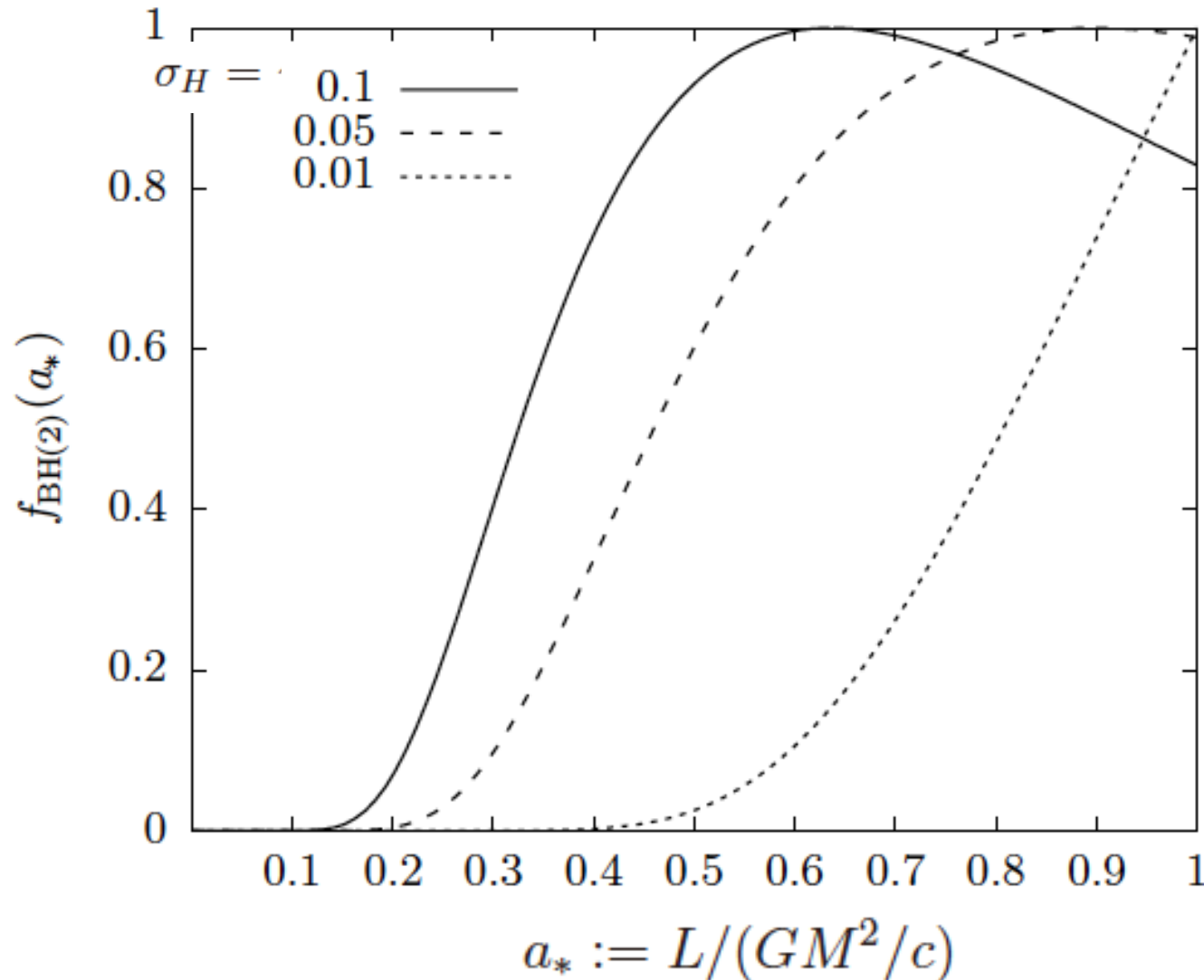
$$\beta_0 \simeq \int_0^\infty d\alpha \int_{-\infty}^\alpha d\beta \int_{-\infty}^\beta d\gamma \theta[\delta_H(\alpha, \beta, \gamma) - \delta_{\text{th}}] \theta[1 - h(\alpha, \beta, \gamma)] w(\alpha, \beta, \gamma)$$

$$\delta_H(\alpha, \beta, \gamma) = \alpha + \beta + \gamma \quad \delta_{\text{th}} := \left( \frac{2}{5} \mathcal{I} \sigma_H \right)^{2/3}$$

# Spin distribution

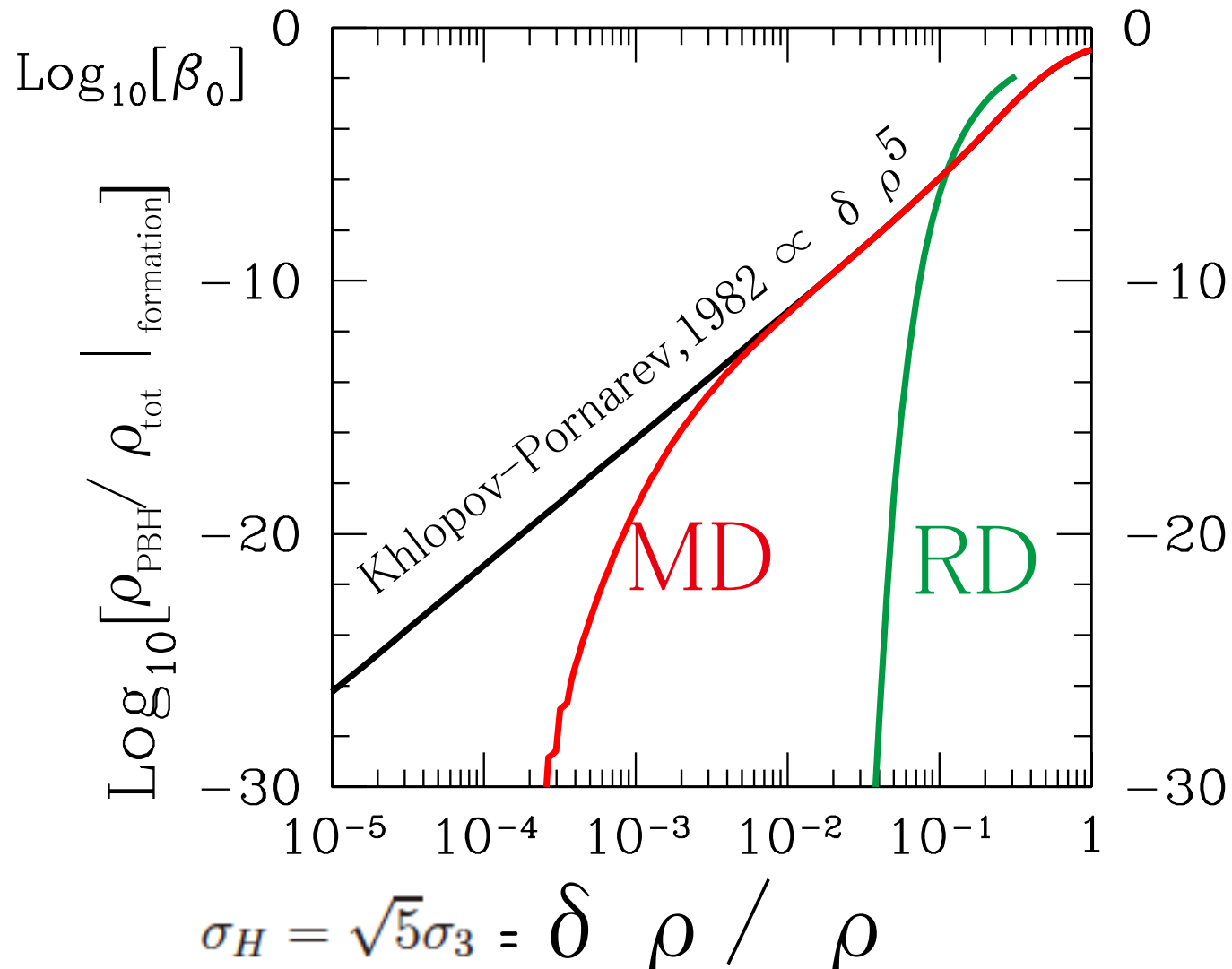
More highly-spinning halos cannot collapse into PBHs, which means  
that the PBHs produced tend to have high spins in MD

Harada, Yoo, KK, Nakao (2017)



# Beta in matter-domination

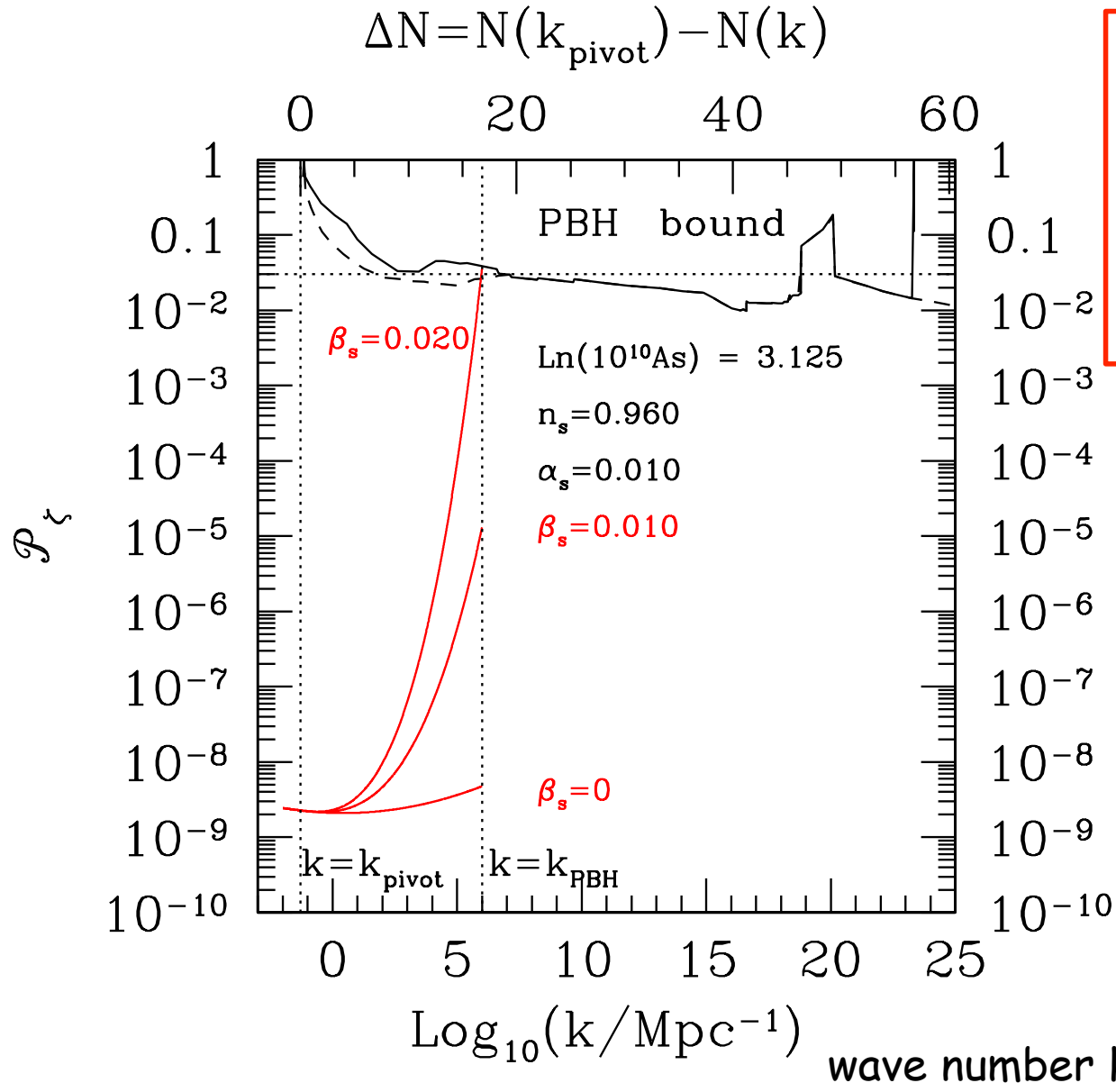
Harada, Yoo, KK, Nakao (2017)



# $P_\zeta$ vs $k$

KK and T.Terada, 2018

Amplitude of curvature perturbation



Planck (2015)

$$n_s = 0.9586 \pm 0.0056,$$

$$\alpha_s = 0.009 \pm 0.010,$$

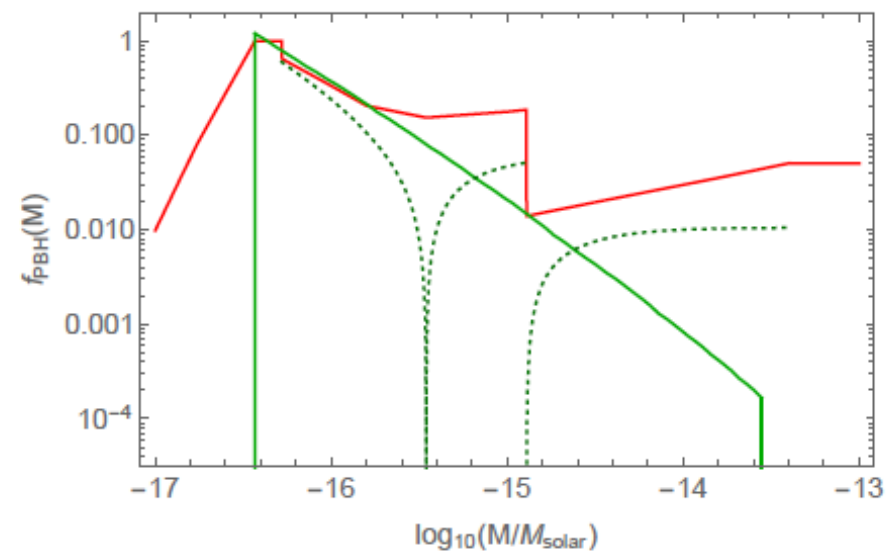
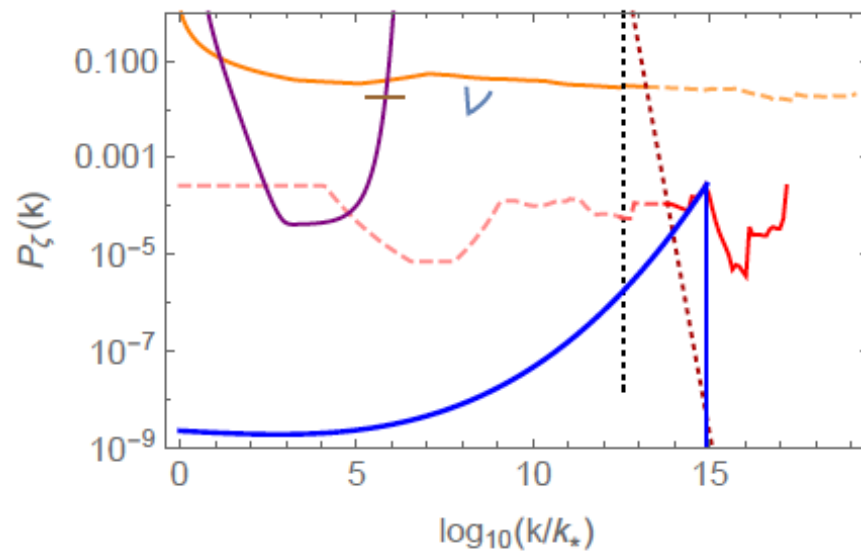
$$\beta_s = 0.025 \pm 0.013.$$

at 68% C.L.

# 100 % Dark Matter by PBHs

KK and T.Terada, 2018

$$n_s = 0.96, \alpha_s = 0, \beta_s = 0.0019485.$$



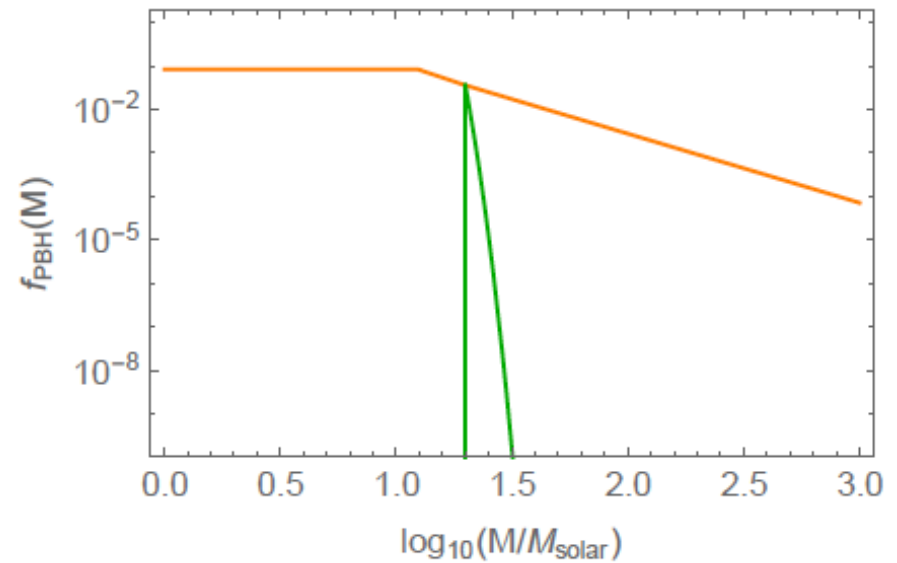
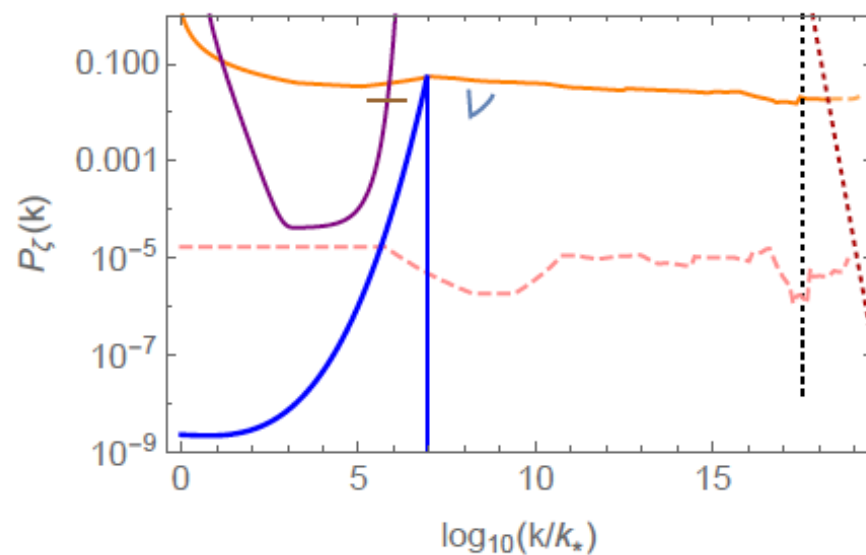
black dotted line shows  
 $T_R = 10^4 \text{ GeV}$ .



# LIGO/VIRGO event

KK and T.Terada, 2018

$$n_s = 0.96, \alpha_s = 0, \beta_s = 0.026.$$



black dotted line shows the reheating at  $T_R = 10^9 \text{ GeV}$

# Summary

- PBH can be formed at small scales even in both radiation and matter dominated epochs
- More PBHs are produced in MD
- We may detect gravitational wave signals secondarily-induced by large SCALAR fluctuations at small scales by e.g., [aLIGO](#), [KAGRA](#), [DECIGO](#) ...
- We will be able to distinguish a model from others by using future small-scale probes such as [PIXIE-like satellite](#) (CMB  $\mu$ -distortion), [SKA/Ominiscope](#) (21cm, Pulsar timing), [CTA](#) (gamma-ray), [DECIGO](#) (GW)...