



CENTRO DE CIENCIAS
MATEMÁTICAS

FEA of Linear Vibrations in FEniCs

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FEniCs Installation



Strong Form to Weak Form

$$\rho \ddot{\mathbf{u}} = \operatorname{div} \boldsymbol{\sigma}, \quad \text{in } \Omega \quad (1)$$

Strong Form to Weak Form

$$\rho \ddot{\mathbf{u}} = \text{div } \boldsymbol{\sigma}, \quad \text{in } \Omega \quad (2)$$

To obtain the weak form we multiply the above equation with a test function \mathbf{v} and integrate over the domain Ω

$$\int_{\Omega} \boldsymbol{\sigma}(\mathbf{u}) : \boldsymbol{\epsilon}(\mathbf{v}) \, dV + \int_{\Omega} \rho \ddot{\mathbf{u}} \cdot \mathbf{v} \, dV = 0, \quad \forall \mathbf{v} \text{ in } \Omega, \quad (3)$$

Strong Form to Weak Form

$$\rho \ddot{\mathbf{u}} = \text{div } \boldsymbol{\sigma}, \quad \text{in } \Omega \quad (4)$$

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$$\int_{\Omega} \boldsymbol{\sigma}(\mathbf{u}) : \boldsymbol{\epsilon}(\mathbf{v}) \, dV + \int_{\Omega} \rho \ddot{\mathbf{u}} \cdot \mathbf{v} \, dV = 0, \quad \forall \mathbf{v} \text{ in } \Omega, \quad (5)$$

Assuming \mathbf{u} takes periodic form: $\mathbf{u}(\mathbf{x}, t) = \mathbf{u}_A(\mathbf{x}) e^{i\omega t}$;

Substituting this $\mathbf{u}(\mathbf{x}, t)$ in in the above eqn, we get:

$$\int_{\Omega} \boldsymbol{\sigma}(\mathbf{u}) : \boldsymbol{\epsilon}(\mathbf{v}) \, dV - \omega^2 \int_{\Omega} \rho \mathbf{u} \cdot \mathbf{v} \, dV = 0, \quad \forall \mathbf{v} \text{ in } \Omega, \quad (6)$$

Strong Form to Weak Form

$$\rho \ddot{\mathbf{u}} = \text{div } \boldsymbol{\sigma}, \quad \text{in } \Omega \quad (7)$$

To obtain the weak form we multiply the above equation with a test function \mathbf{v} and integrate over the domain Ω

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Eigen-value problem: $\mathbf{A}\mathbf{u} - \omega^2 \mathbf{M}\mathbf{u} = 0,$ (10)

A & **M** are matrices assembled from the governing equations

Strong Form to Weak Form

$$\rho \ddot{\mathbf{u}} = \text{div } \boldsymbol{\sigma}, \quad \text{in } \Omega \quad (11)$$

To obtain the weak form we multiply the above equation with a test function \mathbf{v} and integrate over the domain Ω

$$\int_{\Omega} \boldsymbol{\sigma}(\mathbf{u}) : \boldsymbol{\epsilon}(\mathbf{v}) \, dV + \int_{\Omega} \rho \ddot{\mathbf{u}} \cdot \mathbf{v} \, dV = 0, \quad \forall \mathbf{v} \text{ in } \Omega, \quad (12)$$

Assuming \mathbf{u} takes periodic form: $\mathbf{u}(\mathbf{x}, t) = \mathbf{u}_A(\mathbf{x})e^{i\omega t}$

Substituting this $\mathbf{u}(\mathbf{x}, t)$ in in the above eqn, we get:

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Eigen-value problem: $\mathbf{A}\mathbf{u} - \omega^2 \mathbf{M}\mathbf{u} = 0,$ (14)

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Strong Form to Weak Form

$$\rho \ddot{\mathbf{u}} = \text{div } \boldsymbol{\sigma}, \quad \text{in } \Omega \quad (15)$$

To obtain the weak form we multiply the above equation with a test function \mathbf{v} and integrate over the domain Ω

$$\int_{\Omega} \boldsymbol{\sigma}(\mathbf{u}) : \boldsymbol{\epsilon}(\mathbf{v}) \, dV + \int_{\Omega} \rho \ddot{\mathbf{u}} \cdot \mathbf{v} \, dV = 0, \quad \forall \mathbf{v} \text{ in } \Omega, \quad (16)$$

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Eigen-value problem: $\mathbf{A}\mathbf{u} - \omega^2 \mathbf{M}\mathbf{u} = 0,$ (18)

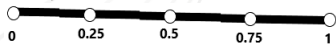
A & **M** are matrices assembled from the governing equations



FENICS
PROJECT

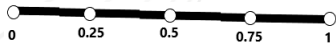
Matrix Formation Simplified

1D space Ω with a mesh of 5 equally spaced nodes:



Matrix Formation Simplified

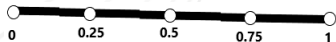
1D space Ω with a mesh of 5 equally spaced nodes:



- Linear Basis Functions: $V = \text{fe.FunctionSpace}(\text{mesh}, 1)$

Matrix Formation Simplified

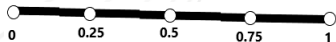
1D space Ω with a mesh of 5 equally spaced nodes:



- Linear Basis Functions: $V = fe.FunctionSpace(mesh, 1)$
- Basis functions: $\phi_0(x), \phi_1(x), \phi_2(x), \phi_3(x), \phi_4(x)$;
 $\phi_i(x) = 1$ at i^{th} node & zero at others and $\phi_i(x)$ is **linearly** interpolated between neighbouring nodes

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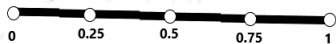
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- $A = \int_{\Omega} \sigma(u) : \epsilon(v) \, dx$

Matrix Formation Simplified

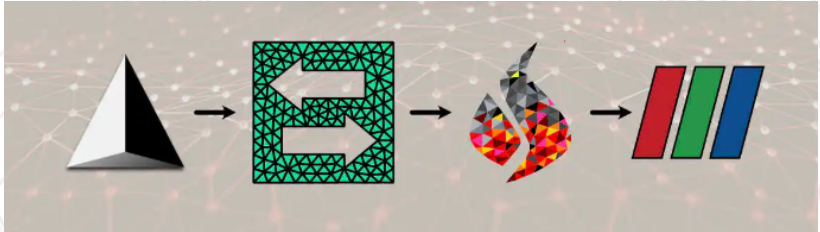
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- $A = \int_{\Omega} \sigma(u) : \epsilon(v) \, dx$

- Matrix $\mathbf{A} = \begin{bmatrix} 2 & -1 & 0 & 0 & 0 \\ -1 & 2 & -1 & 0 & 0 \\ 0 & -1 & 2 & -1 & 0 \\ 0 & 0 & -1 & 2 & -1 \\ 0 & 0 & 0 & -1 & 2 \end{bmatrix}$

Workflow

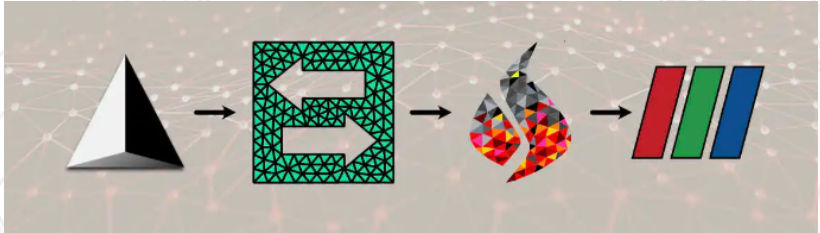


Workflow



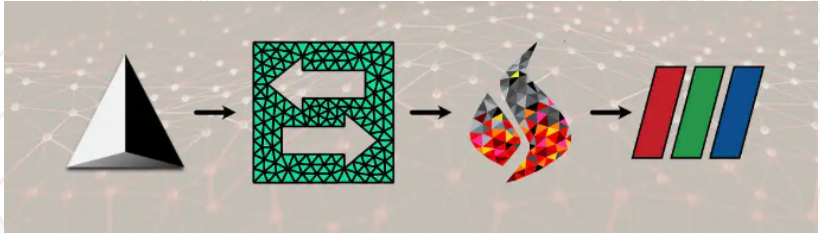
- GMSH: Geometry creation & Meshing

Workflow



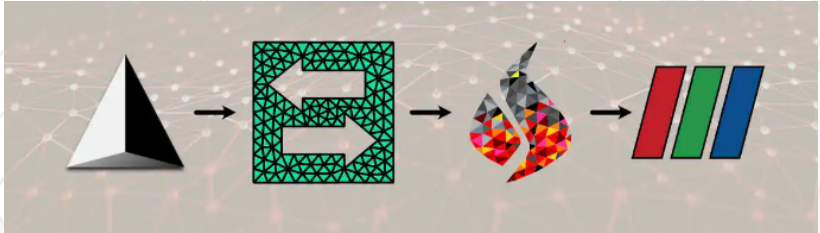
- GMSH: Geometry creation & Meshing
- Meshio: *file.msh* \rightleftharpoons *file.xdmf*

Workflow



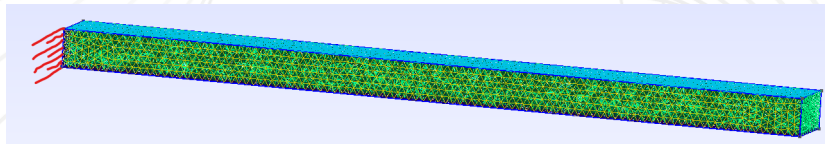
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- FEniCs: PDE Solver

Workflow



- GMSH: Geometry creation & Meshing
- Meshio: $file.msh \rightleftharpoons file.xdmf$
- FEniCs: PDE Solver
- Paraview: Visualisation & Post-processing

Uniform Cantilever Beam

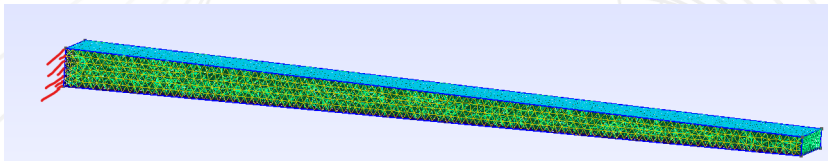


First 6 Modes:

Mode No.	ω - FEniCs [Hz]	ω - Beam theory [Hz]	Mode Participation Factor	Effective Mass	Relative Contribution
1	2.04991	2.01925	-7.83E-02	6.13E-03	61.30 %
2	4.04854	4.0385	7.88E-04	6.21E-07	0.01 %
3	12.81504	12.65443	4.34E-02	1.89E-03	18.85 %
4	25.12717	25.30886	-4.40E-04	1.94E-07	0.00 %
5	35.74168	35.43277	2.55E-02	6.49E-04	6.49 %
6	66.94816	70.86554	-3.71E-04	1.38E-07	0.00 %

Total relative mass of the first 6 modes: 86.65 %

Tapering Cantilever Beam

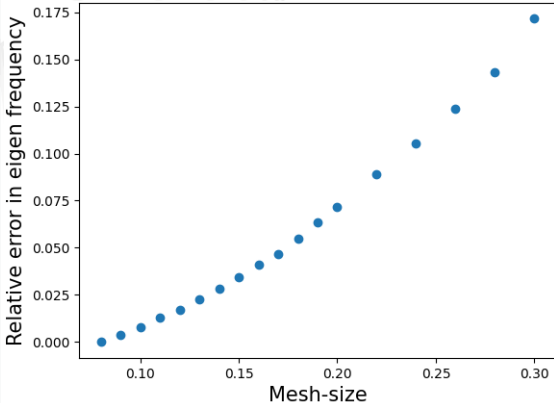


First 6 Modes:

Mode No.	ω - FEniCs [Hz]	Mode Participation Factor	Effective Mass	Relative Contribution
1	4.76505	-1.25E-01	1.57E-02	52.44 %
2	9.98477	-1.66E-05	2.76E-10	0.00 %
3	23.24731	7.78E-02	6.05E-03	20.17 %
4	52.46917	-4.62E-06	2.13E-11	0.00 %
5	59.93993	-5.02E-02	2.52E-03	8.39 %
6	76.63406	6.46E-05	4.18E-09	0.00 %

Total relative mass of the first 6 modes: 81.00 %

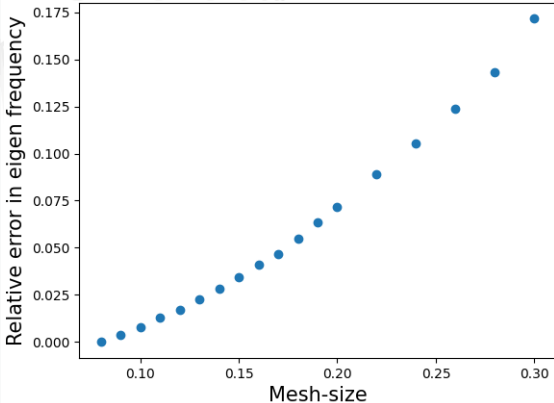
Tapering Beam Result Verification



- Mesh size limit (dl_0): 0.08m {Solver limit}

Response of Mode-1 eigen frequency upon refining the mesh

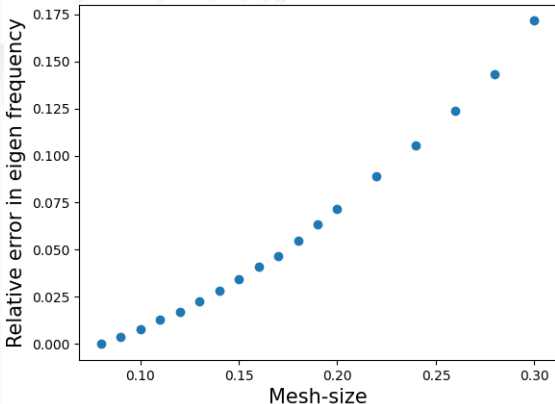
Tapering Beam Result Verification



- Mesh size limit (dl_0): 0.08m {Solver limit}
- Assuming the dl_0 mesh to give accurate output

Response of Mode-1 eigen frequency upon refining the mesh

Tapering Beam Result Verification

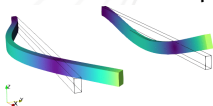


- Mesh size limit (dl_0): 0.08m {Solver limit}
- Assuming the dl_0 mesh to give accurate output
- Observing depletion of error upon mesh refinement

Response of Mode-1 eigen frequency upon refining the mesh

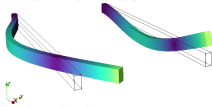
Prospects

- 1 Visualisation and post-processing in Paraview [Spivak, 2018]

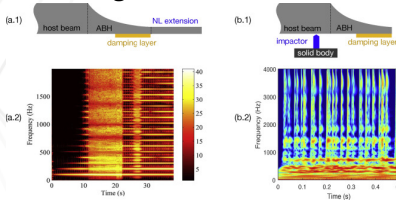


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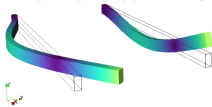
- 2 Simulating Acoustic Black Hole in FEniCs



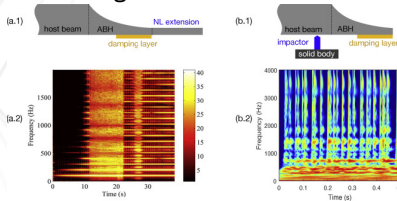
[Adrien et al., 2020]

Prospects

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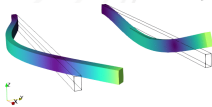
[Adrien et al., 2020]

- Power law: $h(x) = \epsilon x^\mu + h_{tip}$

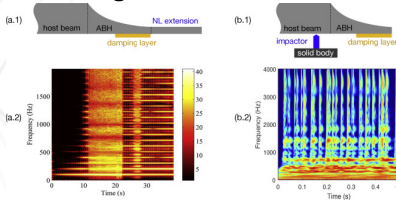
[Hook et al., 2019]

Prospects

1 Visualisation and post-processing in Paraview [Spivak, 2018]



2 Simulating Acoustic Black Hole in FEniCs



[Adrien et al., 2020]

- Power law: $h(x) = \epsilon x^\mu + h_{tip}$ [Hook et al., 2019]
- Geometry creation in GMSH using "BSpline" function

References



Adrien, P., François, G., Stephen, C., and Fabio, S. (2020).

The acoustic black hole: A review of theory and applications.

Journal of Sound and Vibration, 476.



Hook, K., Cheer, J., and Daley, S. (2019).

A parametric study of an acoustic black hole on a beam.

The Journal of the Acoustical Society of America, 145.



Spivak, M. (2018).

Calculus on manifolds: a modern approach to classical theorems of advanced calculus.

CRC press.



THANK YOU!