# Formulas Sheet

## Physical Constants:

| Quantity | Symbol, equation | Value |
|---|---|---|
| Speed of light | $c$ | $2.9979 \times 10^{8}$ m s$^{-1}$ |
| Electron charge | $e$ | $1.602 \times 10^{-19}$C |
| Planck constant | $h$ | $6.626 \times 10^{-34}$Js |
| Planck constant, reduced | $\hbar = h/2\pi$ | $1.055 \times 10^{-34}$ J s |
| Conversion constant | $\hbar c$ | 197.327MeVfm = 197.327eVnm |
| Electron mass | $m_e$ | $9.109 \times 10^{-31}$ kg = 0.511MeV/c$^2$ |
| Proton mass | $m_p$ | $1.673 \times 10^{-27}$ kg$^2$ = 938.272MeV/c$^2$ |
| Neutron mass | $m_n$ | $1.675 \times 10^{-27}$ kg = 939.566MeV/c$^2$ |
| Fine structure constant | $\alpha = e^2/\hbar c$ | 1/137.036 |
| Classical electron radius | $r_e = e^2/m_e c^2$ | $2.818 \times 10^{-15}$ m |
| Electron Compton wavelength | $\lambda = h/m_e c = r_e/\alpha$ | $2.426 \times 10^{-12}$ m |
| Proton Compton wavelength | $\lambda = h/m_p c$ | $1.321 \times 10^{-15}$ m |
| Bohr radius | $a_0 = r_e/\alpha^2$ | $0.529 \times 10^{-10}$ m |
| Rydberg energy | $\mathcal{R} = m_e c^2 \alpha^2/2$ | 13.606eV$^{-11}$MeVT$^{-1}$ |
| Bohr magneton | $\mu_B = e\hbar/2m_e$ | $5.788 \times 10^{-11}$ |
| Nuclear magneton | $\mu_N = e\hbar/2m_p$ | $3.152 \times 10^{-14}$MeVT$^{-1}$ |
| Avogadro number | $N_A$ | $6.022 \times 10^{23}$ mol$^{-1}$ |
| Boltzmann constant | $k$ | $1.381 \times 10^{-23}$ J K$^{-1}$ |
| | | $= 8.617 \times 10^{-5}$eVK$^{-1}$ |
| Gas constant | $R = N_A k$ | 8.31 J mol$^{-1}$ K$^{-1}$ |
| Gravitational constant | $G$ | $6.673 \times 10^{-11}$ m$^3$ kg$^{-1}$ s$^{-2}$ |
| Permittivity of free space | $\epsilon_0 = 1/\mu_0 c^2$ | $8.854 \times 10^{-12}$Fm$^{-1}$ |
| Permeability of free space | $\mu_0$ | $4\pi \times 10^{-7}$ N A$^{-2}$ |

## Conversion of units:

1fm = $10^{-15}$ m, 1fm = $10^{-15}$ m, 1 barn = $10^{-28}$ m$^2$ = 100fm$^2$, 1 atmosphere = 101325 Pa, Thermal energy at $T = 300$ K : $kT = [38.682]^{-1}$eV 0°C = 273.15 K, 1eV = $1.602 \times 10^{-19}$ J, 1eV/c$^2$ = $1.783 \times 10^{-36}$ kg

## Properties of the Solution of 1D Stationary Schrodinger Equation:

1. For 1D potential, all stationary solutions are non-degenerate.
2. Stationary square integrable solution exist only for $E > min(V(x))$
3. If V(x) is real, then $\Psi$(x) can be taken to be real.
4. Eigenvalues of a Hermitian Hamiltonian are all real.
5. The eigenfunctions of a Hermitian operator form a complete orthogonal basis set, for smooth potentials.
6. 1D Schrodinger equation Solution is real up to an over all phase.
7. For a given 1D even potential the stationary states are either even or odd.
8. The wave function and its first order space derivative are continuous all over space and in particular at the boundaries of a finite potential.
9. At boundaries with Dirac delta function potential, the wave function is still continuos but the first space derivative of the wavefunction is discontinuous.
10. Due to square integrability, physical solution should be finite all over space, no blow ups, in particular at infinity.
11. The number of nodes (zeros) of the eigenfunction increases by one unit between adjacent states as we move from the ground state (zero nodes) to higher excited states.
12. Bound states exist only for confining potential (classically between turning points of the potential).

## The Origin of Quantum Physics:

$$E = h\omega \qquad p = \frac{h}{\lambda} = \hbar k \tag{1}$$

$$\lambda = \frac{h}{p} = \frac{2\pi}{k} \qquad \lambda_C = \frac{h}{mc} \tag{2}$$

## Blackbody Radiation:

Plank energy spectral density: $\rho(\nu) = \frac{8\pi h\nu^3}{c^3}\left[\frac{1}{\exp\left(\frac{h\nu}{k_B T}\right)-1}\right]$

## Photoelectric Effect:

$W$ is the work function of irradiated metal. $V_s$ is the stopping potential.

$$K = \frac{1}{2}\, mv^2 = h\nu - W = \frac{hc}{\lambda} - W \ ; \quad K = \frac{1}{2}\, mv^2 = |e|V_s \leftrightarrow V_s = \frac{K}{|e|}$$

$$\nu \geq \frac{W}{h} = \nu_{\min}; \quad \nu = \frac{1}{T} = \frac{c}{\lambda}; \quad \nu = \frac{\omega}{2\pi}$$

## de Broglie Formula:

$$\lambda = \frac{h}{p} = \frac{h}{mv}$$

## Bohr Hydrogen-like Atom:

$$E_n = -\frac{Z^2}{n^2}R \qquad R = 13.6\,\mathrm{eV}$$

$$r_n = \frac{a_0}{Z}n^2 \qquad a_0 = 0.53 \times 10^{-10}\,\mathrm{m}$$

$$h\nu = E_n - E_m = Z^2 R\left(\frac{1}{m^2} - \frac{1}{n^2}\right) \qquad n > m$$

## The Wave Function:

$\Psi(x,t)$ obeys Schrodinger's equation, and the normalization condition $\int_{-\infty}^{\infty} |\Psi(x,t)|^2 dx = 1$:

$$\langle x|x\rangle = \int_{-\infty}^{\infty} x|\Psi(x,t)|^2 dx; \langle p|p\rangle = \int_{-\infty}^{\infty} \Psi^*(x,t)\frac{\hbar}{i}\frac{\partial}{\partial x}\Psi(x,t)dx; \langle Q(\hat{x},\hat{p})|Q(\hat{x},\hat{p})\rangle = \int_{-\infty}^{\infty} \Psi^*(x,t)Q\left(x, \frac{h}{i}\frac{\partial}{\partial x}\right)\Psi \tag{3}$$

$$i\hbar\frac{\partial}{\partial t}\Psi(x,t) = H\Psi(x,t); \qquad \Psi(x,t) = \psi(x)e^{-iEt/\hbar}; \qquad H\psi(x) = E\psi(x) \tag{4}$$

$$\rho(x,t) = |\Psi(x,t)|^2; \qquad \frac{\partial}{\partial t}\rho(\mathbf{x},t) + \nabla\cdot\mathbf{J}(\mathbf{x},t) = 0; \qquad J(x,t) = \frac{i\hbar}{2m}\left(\Psi\frac{\partial\Psi^*}{\partial x} - \Psi^*\frac{\partial\Psi}{\partial x}\right) \tag{5}$$

$$\mathbf{J}(\mathbf{x},t) = \frac{h}{2im}\left(\psi^*\nabla\psi - \psi\nabla\psi^*\right) \quad p = \frac{\hbar}{i}\nabla; \qquad [x_i, p_j] = i\hbar\ \delta_{i,j} \tag{6}$$

Hermitian conjugate $A^\dagger$ is defined by: $\int (A\psi(x))^*\psi(x), dx = \int \psi(x)A^\dagger\psi(x), dx$.

$$H = -\frac{\hbar^2}{2m}\frac{d}{dx^2} + V(x)$$

## Fourier Transform & Wavepackets

$$\Psi(x) = \frac{1}{\sqrt{2\pi}} \int dk \Phi(k) e^{ikx}, \qquad \Phi(k) = \frac{1}{\sqrt{2\pi}} \int dx \Psi(x) e^{-ikx} \qquad \int dx |\Psi(x)|^2 = \int dk |\Phi(k)|^2 = 1 \tag{7}$$

$$\Psi(x) = \frac{1}{(2\pi)^{\frac{3}{2}}} \int d^3k \Phi(\mathbf{k}) e^{i\mathbf{k}\cdot\mathbf{x}}, \quad \Phi(k) = \frac{1}{(2\pi)^{\frac{3}{2}}} \int d^3x \Psi(\mathbf{x}) e^{-i\mathbf{k}\cdot\mathbf{x}} \quad \int d^x |\Psi(\mathbf{x})|^2 = \int d^k |\Phi(\mathbf{k})|^2 = 1 \tag{8}$$

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} e^{ikx} dx = \delta(k) \qquad \frac{1}{(2\pi)^3} \int_{-\infty}^{\infty} e^{i\mathbf{k}\cdot\mathbf{x}} d^3x = \delta^{(3)}(k) \qquad \Psi(x,t) = \frac{1}{(2\pi)^3} \int \phi(k) e^{i(\mathbf{k}\cdot\mathbf{x} - \omega(\mathbf{k})t)}\, dk^3 \tag{9}$$

$$v_{group} = \frac{d\omega}{dk}; \;\; \Delta k \Delta x \simeq 1 \tag{10}$$

## Complete Basis Set:

Given that $H\psi_n(x) = E_n\psi_n(x)$; $\int \phi_n^*(x)\phi_m(x)dx = \delta_{nm}$, where $\{\phi_n\}$ is a complete set, then:

$$\psi(x) = \sum_n c_n \phi_n(x); \quad c_n = \int \psi_n^*(x)\psi(x)\, dx \tag{11}$$

$$\int \psi^*(x)\psi(x)dx = \sum_n |c_n|^2 = 1 \tag{12}$$

$$E = \int \psi_n^*(x) H \psi_m(x) dx = \sum_n |c_n|^2 E_n \tag{13}$$

$$\Psi(x,0) = \psi(x) = \sum_n c_n \phi_n(x) \implies \Psi(x,t) \begin{cases} = \sum_n c_n e^{-iE_nt/\hbar} \phi_n(x) \\ = e^{i\hat{H}t/\hbar} \Psi(x,0) \end{cases} \tag{14}$$

$$c_n = \int \phi_n^* \Psi(x,0) dx \tag{15}$$

## Commutator Properties:

$$[A,A] = 0; \quad [A,B] = -[B,A]; \quad [A+B,C] = [A,C] + [B,C] \tag{16}$$

$$[AB,C] = [A,C]B + A[B,C]; \quad [A,BC] = [A,B]C + B[A,C]; \tag{17}$$

$$[A,[B,C]] + [B,[C,A]] + [C,[A,B]] = 0 \tag{18}$$

## Uncertainty Principle:

$$(\Delta Q)^2 = \langle Q^2 | Q^2 \rangle - \langle Q|Q \rangle^2 = \langle (Q - \langle Q|Q \rangle)^2 | (Q - \langle Q|Q \rangle)^2 \rangle \qquad \Delta x \Delta p \geq \frac{\hbar}{2} \tag{19}$$

Where $\Delta Q$ is the uncertainty for the Hermitian Operator Q.

$$(\Delta Q)^2 (\Delta R)^2 \geq \frac{1}{2} \left| \langle [Q, R] \rangle \right|^2 \tag{20}$$

## Operators:

For the operator $\hat{A}$, $\hat{A}\psi = a\psi$. $a$ in an eigenvalue and $\psi$ is an eigenfunction of $A$. Then, the following properties hold:

- $\hat{A}^n \psi = a^n \psi, \quad \hat{A}^{-1}\psi = a^{-1}\psi, \quad e^{i\hat{A}}\psi = e^{ia}\psi, \quad F(\hat{A})\psi = F(a)\psi$
- If $\hat{A}^\dagger = A,\ \hat{A}\,|\phi_n\rangle = a_n\,|\phi_n\rangle \quad \Longrightarrow \ a_n \in \mathbb{R}, \langle \phi_m | \phi_n | \phi_m | \phi_n \rangle = \delta_{mn}$
- If $\{\phi_n\}$ is a complete and orthonormal for a Hermitian operator, then the operator is diagonal in the eigenbasis, $\{\phi_n\}$, with eigenvalues,, $\{a_n\}$, as the diagonal elements. The basis set is unique iff there are no degenerate eigenvalues.
- If two Hermitian operators, $\hat{A}$ and $\hat{B}$, commute and have no degenerate eigenvalues. Then each eigenvector of $\hat{A}$ is also an eigenvector of $\hat{B}$. A common orthonormal basis can be made of the joint eigenvectors of $\hat{A}$ and $\hat{B}$.

## 1D Infinite Square Well:

$$H\psi_n(x) = E_n \psi_n(x) \qquad \int \psi_n^*(x)\psi_n(x)dx = \delta_{n,m} \tag{21}$$

$$\phi_n(x) = \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi}{a}x\right);\ n = 1, 2, \ldots \qquad \phi_n(x,t) = \phi_n(x)e^{-iE_n t/\hbar} \tag{22}$$

$$V(x) = \begin{cases} 0, 0 \leq x \leq a \\ \infty, \text{otherwise} \end{cases} \qquad E_n = \frac{\hbar^2 k_n^2}{2m} = \frac{n^2\pi^2\hbar^2}{2ma^2} \tag{23}$$

$$\Psi(x,t) = \sum_{n=1}^{\infty} c_n \phi_n(x) e^{-iE_n t/\hbar} \qquad c_n = \int_0^a \phi_n(x)\Psi(x,0)\,dx \tag{24}$$

## Particle on a Ring:

$$\psi_\pm(\theta) = \frac{1}{\sqrt{2\pi}} \exp\left\{ \left( \pm i \frac{R\theta}{\hbar} \sqrt{2mE} \right) \right\} = \frac{1}{\sqrt{2\pi}} e^{\pm ikx} \qquad x = R\theta; L = 2\pi R; k = \frac{2\pi n}{L} = \frac{n}{R} \tag{25}$$

$$\psi(\theta) = \frac{1}{\sqrt{2\pi}} \exp\{\pm in\theta\} \qquad E_n = \frac{n^2\hbar^2}{2mR^2}, \quad n = 0 \pm 1, \pm 2, \pm 3, ... \tag{26}$$

## Harmonic Oscillator:

$$V(x) = \frac{1}{2}kx^2 = \frac{1}{2}m(\omega x)^2 \quad \left(\omega \equiv \sqrt{k/m}\right) \qquad E_n = \hbar\omega\left(n + \frac{1}{2}\right); n = 0, 1, 2, \ldots \tag{27}$$

$$H = \frac{1}{2m}[p^2 + (m\omega x)^2] = \hbar\omega\left(N + \frac{1}{2}\right) \qquad N = a_+ a_- \quad (= a^\dagger a) \tag{28}$$

$$N\psi_n = n\psi_n;\ a_\pm = \sqrt{\frac{m\omega}{2\hbar}}(\hat{x} \mp i\frac{\hat{p}}{m\omega}) \qquad N(a_+\psi_n) = [N, a_+]\psi_n \tag{29}$$

$$[N, a_\pm] = \pm a_\pm;\ [a_-, a_+] = 1 \qquad a_\pm \equiv \frac{1}{\sqrt{2\hbar m\omega}}(\mp ip + m\omega x) \tag{30}$$

$$x = \sqrt{\frac{\hbar}{2m\omega}}(a_+ + a_-) \qquad p = i\sqrt{\frac{m\omega\hbar}{2}}(a_+ - a_-) \tag{31}$$

$$a_+\psi_n = \sqrt{n+1}\psi_{n+1} \qquad a_-\psi_n = \sqrt{n}\psi_{n-1} \tag{32}$$

$$\psi_0(x) = \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \exp\left(-\frac{m\omega}{2\hbar}x^2\right) \qquad \psi_n = \frac{1}{\sqrt{n!}}(a_+)^n\psi_0 \tag{33}$$

$$\xi \equiv \sqrt{\frac{m\omega}{\hbar}}x \qquad \mathcal{H}_n(\xi) = (-1)^n e^{\xi^2}\left(\frac{d}{d\xi}\right)^n e^{-\xi^2} \tag{34}$$

$$\psi_n(x) = \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \frac{1}{\sqrt{2^n n!}}\mathcal{H}_n(\xi)e^{-\xi^2/2} \qquad \int_{-\infty}^{+\infty} \mathcal{H}_n\mathcal{H}_m e^{-x^2}dx = 2^n n!\sqrt{\pi}\delta_{n,m} \tag{35}$$

## Models of Dirac Delta Distribution $\delta(x)$:

$$\delta(x) = \lim_{\alpha\to\infty}\frac{\sin(\alpha x)}{\pi x};\ \delta(x) = \lim_{\epsilon\to 0^+}\frac{1}{2\pi}\int_{-\infty}^{\infty} e^{-ikx}e^{-\epsilon|k|}\,dk = \lim_{\epsilon\to 0^+}\frac{\epsilon}{\pi(x^2+\epsilon^2)};\ \delta(x) = \lim_{\epsilon\to 0}\frac{\Theta(x+\epsilon)-\Theta(\epsilon)}{\epsilon} \tag{36}$$

where $\Theta(x)$ is Heaviside or step function,

$$\Theta(x) = \begin{cases} 1, & x > 0 \\ 0, & x < 0 \end{cases} \tag{37}$$

## Bound State of Single $\delta$-Potential: $\boxed{E < 0}$

$$V = -\alpha\delta(x),\ \alpha > 0 \qquad \psi(x) = \sqrt{\frac{m\alpha}{\hbar^2}}\, e^{-\frac{m\alpha}{\hbar^2}|x|} \qquad E = -\frac{m\alpha^2}{2\hbar^2} \tag{38}$$

## Scattering State: $\boxed{E > 0}$

$$V(x) = -\alpha\delta(x);\ k = \frac{\sqrt{2mE}}{\hbar} \qquad \psi(x) = \begin{cases} Ae^{ikx} + Be^{-ikx}, & x < 0 \\ Fe^{ikx}, & x > 0 \end{cases} \tag{39}$$

$$T = \frac{|F|^2}{|A|^2} = \frac{1}{1+\beta^2} \qquad R = \frac{|B|^2}{|A|^2} = \frac{\beta}{1+\beta} \qquad (\beta = m\alpha/\hbar^2 k) \tag{40}$$

## Miscellaneous:

$$\int_{-\infty}^{\infty} e^{-ax^2}\, dx = \sqrt{\frac{\pi}{a}}; \qquad \int_{-\infty}^{\infty} e^{-(ax^2+bx)}\, dx = e^{b^2/4a}\sqrt{\frac{\pi}{a}}; \qquad \delta_{ij} = \begin{cases} 1, & i = j \\ 0, & i \neq j \end{cases} \tag{41}$$

## Matrix Algebra:

Let A be a 2X2 matrix defined as: $A = \begin{bmatrix} a & b \\ c & d \end{bmatrix}$ then:

$$A^{-1} = \frac{1}{|A|}\begin{bmatrix} d & -b \\ -c & a \end{bmatrix};\ |A| = ad - bc \tag{42}$$

$$\sigma_x = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \qquad \sigma_y = \begin{bmatrix} 0 & -1 \\ i & 0 \end{bmatrix} \qquad \sigma_z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \tag{43}$$

$$\{\sigma_i, \sigma_j\} = 2\delta_i j \qquad [\sigma_i, \sigma_j] = 2i\epsilon_{ijk} \qquad \sigma_i\sigma_j = \begin{cases} 1, & i = j \\ -\sigma_j\sigma_i, & i \neq j \end{cases} \tag{44}$$

## Orbital Angular Momentum:

$$\hat{L}_x = \hat{y}\hat{p}_z - \hat{z}\hat{p}_y, \quad \hat{L}_y = \hat{z}\hat{p}_x - \hat{x}\hat{p}_z, \quad \hat{L}_z = \hat{x}\hat{p}_y - \hat{y}\hat{p}_x \tag{45}$$

$$\left[\hat{L}_x, \hat{L}_y\right] = i\hbar\hat{L}_z, \quad \left[\hat{L}_y, \hat{L}_z\right] = i\hbar\hat{L}_x, \quad \left[\hat{L}_z, \hat{L}_x\right] = i\hbar\hat{L}_y. \tag{46}$$

$$\hat{L}^2 \equiv \hat{L}_x\hat{L}_x + \hat{L}_y\hat{L}_y + \hat{L}_z\hat{L}_z, \quad \left[\hat{L}^2, \hat{L}_i\right] = 0 \tag{47}$$

$$\nabla = \hat{r}\frac{\partial}{\partial r} + \hat{\theta}\frac{1}{r}\frac{\partial}{\partial\theta} + \hat{\phi}\frac{1}{r\sin\theta}\frac{\partial}{\partial\phi} \tag{48}$$

$$\nabla^2 = \frac{1}{r^2}\frac{\partial}{\partial r}\left(r^2\frac{\partial}{\partial r}\right) + \frac{1}{r^2\sin\theta}\frac{\partial}{\partial\theta}\left(\sin\theta\frac{\partial}{\partial\theta}\right) + \frac{1}{r^2\sin^2\theta}\left(\frac{\partial^2}{\partial\phi^2}\right) \tag{49}$$

$$\nabla^2 = \frac{1}{r}\frac{\partial^2}{\partial r^2}r + \frac{1}{r^2}\left(\frac{\partial^2}{\partial\theta^2} + \cot\theta\frac{\partial}{\partial\theta} + \frac{1}{\sin^2\theta}\frac{\partial^2}{\partial\phi^2}\right) \tag{50}$$

$$\hat{L}^2 = -\hbar^2\left(\frac{\partial^2}{\partial\theta^2} + \cot\theta\frac{\partial}{\partial\theta} + \frac{1}{\sin^2\theta}\frac{\partial^2}{\partial\phi^2}\right) \tag{51}$$

$$\hat{L}_z = \frac{\hbar}{i}\frac{\partial}{\partial\phi}; \quad \hat{L}_\pm = \hbar e^{\pm i\phi}\left(\pm\frac{\partial}{\partial\theta} + i\cot\theta\frac{\partial}{\partial\phi}\right) = \hat{L}_x \pm i\hat{L}_y \tag{52}$$

## General Momentum Operator $\boxed{\vec{L},\ \vec{S},\ \vec{J}}$

$$[J_x, J_y] = i\hbar J_z \qquad [J_z, J_x] = i\hbar J_y \qquad [J_y, J_z] = i\hbar J_x \tag{53}$$

$$J^2\left|Jm\right\rangle = \hbar^2 j(j+1)\left|Jm\right\rangle \qquad J_z\left|Jm\right\rangle = \hbar m\left|Jm\right\rangle \qquad J^2 = J_x^2 + J_y^2 + J_z^2 \tag{54}$$

$$J^2 = J_\pm J_\mp + {J_z}^2 \mp \hbar J_z \tag{55}$$

$\left|Jm\right\rangle$ are common eigenstates of the $J^2, J_z$;

Thus both $J^2$ and $J_z$ are diagonal matrices for any fixed value of J, the dimension of these matrices is $(2J+1)\times(2J+1)$ since $m = -J, -J+1, \ldots, J-1, J$, (2J+1 Values).

$J = \frac{1}{2} \implies 2J+1 = 2 \implies 2\times 2$ matrices.

$J = 1 \implies 2J+1 = 3 \implies 3\times 3$ matrices.

$$J_\pm = J_x \pm iJ_y \qquad J_\pm\left|Jm\right\rangle = \sqrt{J(J+1) - m(m\pm 1)}\left|Jm\pm 1\right\rangle \tag{56}$$

so that $J_\pm$ are off diagonal matrices.

We order the eigenvectors for a fixed J-value as the following:

$\left|J, J\right\rangle \quad \left|J, J-1\right\rangle \quad \left|J, J-2\right\rangle \quad \ldots \quad \left|J, -J\right\rangle$

For a fixed J here is the expectation of operator A:

$$A = \begin{pmatrix} \langle J|\,A\,|J\rangle & \langle J|\,A\,|J-1\rangle & \ldots & \langle J|\,A\,|-J\rangle \\ \langle J-1|\,A\,|J\rangle & \langle J-1|\,A\,|J-1\rangle & \ldots & \langle J-1|\,A\,|-J\rangle \\ \vdots & \vdots & \vdots & \vdots \\ \langle -J|\,A\,|J\rangle & \langle -J|\,A\,|J-1\rangle & \ldots & \langle -J|\,A\,|-J\rangle \end{pmatrix} \tag{57}$$

in the above matrix, for a fixed J we denote $|Jm\rangle = |m\rangle$.

## Spin Angular Momentum Operator

For $j = \frac{1}{2} \implies |\uparrow\rangle = \left|\frac{1}{2}, \frac{1}{2}\right\rangle ; |\downarrow\rangle = \left|\frac{1}{2}, -\frac{1}{2}\right\rangle$ *then* :

$$S_x = \frac{\hbar}{2}\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \frac{\hbar}{2}\sigma_1 \qquad S_y = \frac{\hbar}{2}\begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = \frac{\hbar}{2}\sigma_2 \qquad S_z = \frac{\hbar}{2}\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \frac{\hbar}{2}\sigma_3 \tag{58}$$

$$[S_x, S_y] = i\hbar S_z \qquad [S_z, S_x] = i\hbar S_y \qquad [S_y, S_z] = i\hbar S_x \tag{59}$$

$$S_+ = S_x + iS_y \qquad S_- = S_x - iS_y \tag{60}$$

$$S_+ = \frac{\hbar}{2}\begin{pmatrix} 0 & 2 \\ 0 & 0 \end{pmatrix} \qquad S_- = \frac{\hbar}{2}\begin{pmatrix} 0 & 0 \\ 2 & 0 \end{pmatrix} \tag{61}$$

## Spherical Harmonics:

$$\langle\theta\,\phi|l\,m\rangle = Y_{\ell,m}(\theta,\phi) \equiv \mathcal{N}_{\ell,m} P_\ell^m(\cos\theta)e^{im\phi} \tag{62}$$

$$\hat{L}_z Y_{\ell m} = \hbar m Y_{\ell m} \tag{63}$$

$$\hat{L}^2 Y_{\ell m} = \hbar^2 \ell(\ell+1) Y_{\ell m} \tag{64}$$

$$\langle l'\,m'|l\,m\rangle = \int d\Omega Y^*_{\ell' m'}(\theta,\phi) Y_{\ell m}(\theta,\phi) = \delta_{\ell',\ell}\delta_{m',m}, \qquad \int d\Omega = \int_0^{2\pi} d\phi \int_{-1}^{1} d(\cos\theta) \tag{65}$$

$$Y_{0,0}(\theta,\phi) = \frac{1}{\sqrt{4\pi}}; \quad Y_{1,\pm1}(\theta,\phi) = \mp\sqrt{\frac{3}{8\pi}}\sin\theta\exp(\pm i\phi); \quad Y_{1,0}(\theta,\phi) = \sqrt{\frac{3}{4\pi}}\cos\theta \tag{66}$$

$$\int |\theta\,\phi\rangle\langle\theta\,\phi|\,d\Omega = 1 \qquad \sum_{lm} |l\,m\rangle\langle l\,m| = 1 \tag{67}$$

## Central Potentials $V(\mathbf{r}) = V(r)$:

$$\left[\frac{-\hbar^2}{2m}\nabla^2 + V(r)\right]\psi(r,\theta,\phi) = E\psi(r,\theta,\phi) \tag{68}$$

$$-\frac{\hbar^2}{2m}\left[\frac{1}{r}\frac{\partial^2}{\partial r^2}r - \frac{L^2(\theta,\phi)}{\hbar^2 r^2} + V(r)\right]\psi(r,\theta,\phi) = E\psi(r,\theta,\phi) \tag{69}$$

$$\psi(r,\theta,\phi) = R(r)Y_{l,m} * (\theta\phi) \tag{70}$$

$$\psi(r,\theta,\phi) = \frac{u(r)}{r}Y_{\ell m}(\theta,\phi) \tag{71}$$

$$\left(-\frac{\hbar^2}{2m}\frac{d^2}{dr^2} + V(r) + \frac{\hbar^2\ell(\ell+1)}{2mr^2}\right)u(r) = Eu(r) \tag{72}$$

$$u(r) \sim r^{\ell+1},\ \text{as}\, r \to 0 \tag{73}$$

$$\left[-\frac{\hbar^2}{2m}\frac{d^2}{dr^2} + V_{\text{eff}}(r)\right]u(r) = E\,u(r); \qquad V_{\text{eff}}(r) = V(r) + \frac{\hbar^2\ell(\ell+1)}{2mr^2} \tag{74}$$

$$\nabla^2 = \frac{1}{r}\frac{\partial^2}{\partial r^2}r + \frac{1}{r^2}\left(\frac{\partial^2}{\partial\theta^2} + \cot\theta\frac{\partial}{\partial\theta} + \frac{1}{\sin^2\theta}\frac{\partial^2}{\partial\phi^2}\right) \tag{75}$$

$$\hat{L}^2 = -\hbar^2\left(\frac{\partial^2}{\partial\theta^2} + \cot\theta\frac{\partial}{\partial\theta} + \frac{1}{\sin^2\theta}\frac{\partial^2}{\partial\phi^2}\right) \tag{76}$$

$$\hat{L}_z = \frac{\hbar}{i}\frac{\partial}{\partial\phi}; \quad \hat{L}_\pm = \hbar e^{\pm i\phi}\left(\pm\frac{\partial}{\partial\theta} + i\cot\theta\frac{\partial}{\partial\phi}\right) = \hat{L}_x \pm i\hat{L}_y \tag{77}$$

## Hydrogen Atom (Z=1):

$$V(r) = -\frac{Ze^2}{r} \qquad V_{\text{eff}}(r) = -\frac{Ze^2}{r} + \frac{\hbar^2\ell(\ell+1)}{2mr^2} \tag{78}$$

$$H = \frac{\mathbf{p}^2}{2m} - \frac{Ze^2}{r} \tag{79}$$

$$E_n = -\frac{Z^2e^2}{2a_0}\frac{1}{n^2}, \quad a_0 = \frac{\hbar^2}{me^2} \simeq 0.529\times 10^{-10}\text{ m}, \quad \frac{e^2}{2a_0} \simeq 13.6\text{eV} \tag{80}$$

$$\psi_{n\ell m}(r,\theta,\phi) = A\left(\frac{r}{a_0}\right)^{\ell}\left(\text{Polynomial in } \frac{r}{a_0} \text{ of degree } n-(\ell+1)\right)e^{-\frac{Zr}{na_0}}Y_\ell^m(\theta,\phi) \tag{81}$$

$$\psi_{n\ell m}(r,\theta,\phi) = R_{n\ell}(r)Y_\ell^m(\theta,\phi) = \sqrt{\left(\frac{2}{na}\right)^3\frac{(n-\ell-1)!}{2n(n+\ell)!}}e^{\frac{-r}{na}}\left(\frac{2r}{na}\right)^{\ell}[L_{n-\ell-1}^{2\ell+1}\left(\frac{2r}{na}\right)]Y_\ell^m(\theta,\phi) \tag{82}$$

$$\int \psi^*_{n\ell m}(\vec{r})\psi_{n'\ell'm'}(\vec{r})d^3\vec{r} = \delta_{nn'}\delta_{mm'}\delta_{\ell\ell'} \tag{83}$$

$$\hat{H}\psi_{n\ell m}(\vec{r}) = E_n\psi_{n\ell m}(\vec{r}) \tag{84}$$

$$n = 1,2,\ldots, \quad \ell = 0,1,\ldots,n-1, \quad m = -\ell,\ldots,\ell \tag{85}$$

$$\psi_{n\ell m}(r,\theta,\phi) = \frac{u_{n\ell}(r)}{r}Y_\ell^m(\theta,\phi) \tag{86}$$

$$u_{1,0}(r) = \frac{2r}{a_0^{3/2}}\exp\left(-r/a_0\right) \tag{87}$$

$$u_{2,0}(r) = \frac{2r}{(2a_0)^{3/2}}\left(1-\frac{r}{2a_0}\right)\exp\left(-r/2a_0\right) \tag{88}$$

$$u_{2,1}(r) = \frac{1}{\sqrt{3}}\frac{1}{(2a_0)^{3/2}}\frac{r^2}{a_0}\exp\left(-r/2a_0\right) \tag{89}$$

$$\Psi_{100}(r,\theta,\phi) = \frac{2r}{\sqrt{4\pi a_0^3}}e^{-r/a_0} \tag{90}$$

## Transfer Matrix

$$\begin{bmatrix} F \\ G \end{bmatrix} = M\begin{bmatrix} A \\ B \end{bmatrix} \qquad M = \begin{bmatrix} M_{11} & M_{12} \\ M_{12}^* & M_{11}^* \end{bmatrix} \qquad \det\{M\} = 1 \tag{91}$$

$$T = \frac{1}{|M_{22}|^2}; \qquad R + T = 1 \tag{92}$$

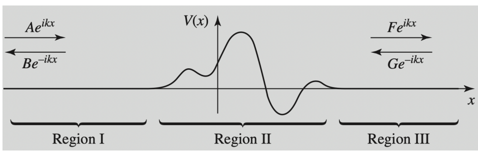


Figure 1: Scattering from an arbitrary localized potential ($V(x) = 0$ except for Region II)

## Spherical Coordinates:

$$\hat{\mathbf{r}} = \sin\theta\cos\phi\hat{\mathbf{x}} + \sin\theta\sin\phi\hat{\mathbf{y}} + \cos\theta\hat{\mathbf{z}}$$

$$\hat{\boldsymbol{\theta}} = \cos\theta\cos\phi\hat{\mathbf{x}} + \cos\theta\sin\phi\hat{\mathbf{y}} - \sin\theta\hat{\mathbf{z}}$$

$$\hat{\boldsymbol{\phi}} = -\sin\phi\hat{\mathbf{x}} + \cos\phi\hat{\mathbf{y}}$$

$$\hat{\mathbf{x}} = \sin\theta\cos\phi\hat{\mathbf{r}} + \cos\theta\cos\phi\hat{\boldsymbol{\theta}} - \sin\phi\hat{\boldsymbol{\phi}}$$

$$\hat{\mathbf{y}} = \sin\theta\sin\phi\hat{\mathbf{r}} + \cos\theta\sin\phi\hat{\boldsymbol{\theta}} + \cos\phi\hat{\boldsymbol{\phi}}$$

$$\hat{\mathbf{z}} = \cos\theta\hat{\mathbf{r}} - \sin\theta\hat{\boldsymbol{\theta}}$$

## Identical Particles:

Particles that share the same intrinsic properties: mass, charge, spin, magnetic moment, ...etc. Identical particles are indistinguishable

$$\Psi(1,2,\ldots,i,\ldots,j,\ldots,N) = \pm\Psi(1,2,\ldots,j,\ldots,i,\ldots,N) \tag{93}$$

are either symmetric for Bosons (integer spin particles) or antisymmetric for Fermions (half integer spin particles) under exchange of any two particles.

$\Longrightarrow$ Pauli Exclusion Principle: Two identical Fermions cannot occupy the same quantum state.

$$\Psi_{Fermions}(R_1,R_2,\ldots,R_N) = \frac{1}{\sqrt{N!}}\begin{pmatrix} \Psi_1(R_1) & \Psi_2(R_1) & \ldots & \Psi_N(R_1) \\ \Psi_1(R_2) & \Psi_2(R_2) & \ldots & \Psi_N(R_2) \\ \vdots & \vdots & \vdots & \vdots \\ \Psi_1(R_N) & \Psi_2(R_N) & \ldots & \Psi_N(R_N) \end{pmatrix} \tag{94}$$

For Bosons use the same Slater-determinant by replacing all signs be $+$ . Spin $\frac{1}{2}$ particles can be in any of the following states:

Singlet State (Antisymmetric state):

$$\chi_s(s1,s2) = \frac{1}{\sqrt{2}}(\chi_\uparrow(s1)\chi_\downarrow(s2) - \chi_\downarrow(s1)\chi_\uparrow(s2)) \tag{95}$$

with total spin s=0.

Triplet States (Symmetric state): with total spin $S_z$ =1, 0, -1:

$$\chi_T(s_1,s_2)\begin{cases} \chi_{1,1}(s_1,s_2) = \chi_\uparrow(s_1)\chi_\uparrow(s_2) \\ \chi_{-1,-1}(s_1,s_2) = \chi_\downarrow(s_1)\chi_\downarrow(s_2) \\ \chi_{1,0}(s_1,s_2) = \frac{1}{\sqrt{2}}(\chi_\uparrow(s_1)\chi_\downarrow(s_2) + \chi_\downarrow(s_1)\chi_\uparrow(s_2)) \end{cases} \tag{96}$$

## Electron Gas in Volume V:

Each eigenvalue in k-space occupy a volume $\frac{(2\pi)^3}{V}$ so that 2 for spin states:

$$\sum_k f(k) = 2\frac{V}{(2\pi)^3}\int d^3k f(k) \tag{97}$$

Highest occupied wave vector is called Fermi wave vector $k_F$

$$k_F = (3\pi^2 n)^{1/3};\ n = \frac{N}{V}\ \text{Electronic Density} \tag{98}$$

## Ground State of N Electrons in Volume V:

$$E_F = \frac{\hbar^2 k^2}{2m} \qquad E = 2\sum_k \frac{\hbar^2 k^2}{2m} = \frac{3}{5}NE_F \tag{99}$$

$$2\frac{\Omega_k}{\frac{(2\pi)^3}{V}} = N = 2\frac{\frac{4}{3}\pi k_F^3}{\frac{(2\pi)^3}{V}} \qquad k_F = (3\pi^2 n)^{1/3} \tag{100}$$

$$n = \frac{N}{V} = \frac{1}{v} = \frac{1}{\frac{4}{3}\pi r_s^3} \qquad r_s = \left(\frac{3}{4\pi n}\right)^{1/3} \tag{101}$$

$$k_F = (3\pi^2 n)^{1/3} = \frac{(9\pi/4)^{1/3}}{r_s} = \frac{1.92}{r_s} \tag{102}$$

## Bloch Theorem in Solids:

$$\psi(R+x) = e^{ik\cdot R}\psi(x) \tag{103}$$

$$\underline{\text{In } 3D\text{:}}\quad k_i = \frac{2\pi n_i}{L};\quad i = x,y,z;\quad n = 0,\pm1,\pm2,... \Rightarrow dn = \frac{V}{(2\pi)^3}dk^3 \tag{104}$$

$$\sum_{k\sigma} F(k,\sigma) = \frac{V}{(2\pi)^3}\sum_\sigma \int dk^3 F(k,\sigma) \tag{105}$$

## Probability Distributions:

We have three distributions, one for classical particles (Maxwell-Boltzmann) and two for undistinguishable particles (Fermi-Dirac for fermions and Bose-Einstein for bosons). Their equations are the following:

$$P(E) = \frac{1}{e^{\beta(E-E_f)}+1} \qquad \text{Fermi–Dirac Dist.} \tag{106}$$

$$P(E) = \frac{1}{e^{\beta(E-\mu)}-1} \qquad \text{Bose–Einstein Dist.} \tag{107}$$

$$P(E) = \frac{1}{e^{\beta(E-\mu)}} \qquad \text{Maxwell–Boltzmann Dist.} \tag{108}$$

$$\text{Where } \beta = \frac{1}{k_B T} \tag{109}$$

The plots of these distributions for different temperatures:

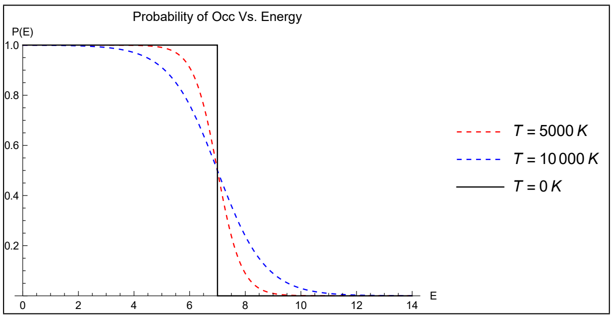


**Figure 1:** Fermi–Dirac Dist. when $E_f = 7eV$. For fermions

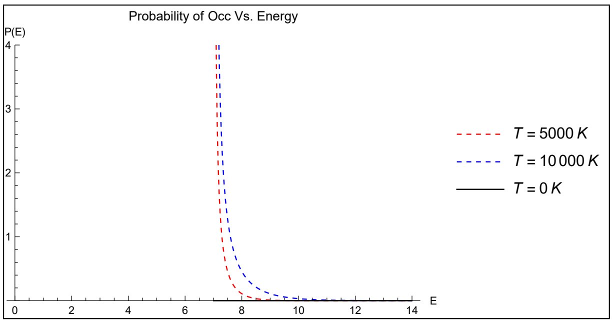


**Figure 2:** Bose–Einstein Dist. when $\mu = 7eV$. For bosons

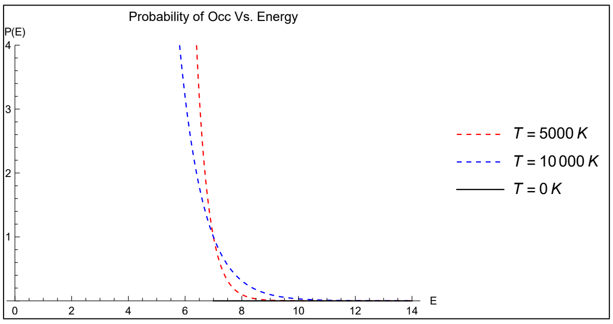


**Figure 3:** Maxwell–Boltzmann Dist. when $\mu = 7eV$. For bosons

## Non-Degenerate Perturbation Theory

$$H = H^0 + \lambda H' = H(\lambda) \tag{110}$$

$$H(\lambda)\psi_n(\lambda) = E_n(\lambda)\psi_n(\lambda) \tag{111}$$

$$\psi_n(\lambda) = \psi_n^0 + \lambda\psi_n^1 + \lambda^2\psi_n^2 + \cdots \tag{112}$$

$$E_n(\lambda) = E_n^0 + \lambda E_n^1 + \lambda^2 E_n^2 + \cdots \tag{113}$$

$$\text{Assume } \psi_n(\lambda),\ E_n(\lambda) \text{ converges} \tag{114}$$

$$\tag{115}$$

$$\lambda^0 : \quad H^0|\psi_n^0\rangle = E_n^0|\psi_n^0\rangle \tag{116}$$

$$\lambda^1 : \quad H^0|\psi_n^1\rangle + H'\psi_n^0 = E_n^0|\psi_n^1\rangle + E_n^1|\psi_n^0\rangle \tag{117}$$

$$\lambda^2 : \quad H^0|\psi_n^2\rangle + H'\psi_n^1 = E_n^0|\psi_n^2\rangle + E_n^1|\psi_n^1\rangle + E_n^2|\psi_n^0\rangle \tag{118}$$

$$E_n^1 = \langle\psi_n^0|H'|\psi_n^0\rangle \tag{119}$$

$$|\psi_n^1\rangle = \sum_{m\neq n} \frac{\langle\psi_m^0|H'|\psi_n^0\rangle}{E_n^0 - E_m^0}|\psi_m^0\rangle \tag{120}$$

$$E_n^2 = \langle\psi_n^0|H'|\psi_n^1\rangle \tag{121}$$

$$E_n^2 = \sum_{m\neq n} \frac{|\langle\psi_m^0|H'|\psi_n^0\rangle|^2}{E_n^0 - E_m^0} \tag{122}$$

Let $|\psi_1^0\rangle$ and $|\psi_2^0\rangle$ be degenerate states of $H^0$, then

$$H' = \begin{bmatrix} H'_{11} & H'_{12} \\ H'_{21} & H'_{22} \end{bmatrix}; \quad \text{where} \quad H'_{ij} = \langle\psi_i^0|H'|\psi_j^0\rangle \tag{123}$$

$$|H' - E^1 I| = \begin{bmatrix} H'_{11} - E^1 & H'_{12} \\ H'_{21} & H'_{22} - E^1 \end{bmatrix} \begin{bmatrix} \alpha_1 \\ \alpha_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \tag{124}$$

$$E_\pm^1 = \frac{1}{2}\left[H'_{11} + H'_{22} \pm \sqrt{(H'_{11} - H'_{11})^2 + 4|H'_{12}|^2}\right] \quad \text{and} \quad |\psi_\pm\rangle = \alpha_1|\psi_1^0\rangle + \alpha_2|\psi_2^0\rangle \tag{125}$$

## WKB Approximation

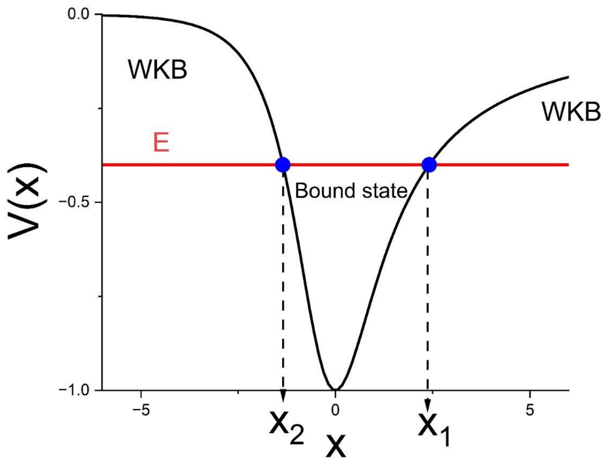


Figure 2: WKB Bound States

$$p(x) \equiv \sqrt{2m[E - V(x)]} \tag{126}$$

$$\psi(x) \approx \frac{C}{\sqrt{p(x)}}\, e^{\pm \frac{i}{\hbar}\int p(x)\,dx} \qquad ; \quad E > V \tag{127}$$

$$\psi(x) \approx \begin{cases} \frac{C_1}{\sqrt{p(x)}} \sin\left[\frac{1}{\hbar}\int_{x_1}^{x} p(x')\,dx' + \frac{\pi}{4}\right], & x_1 < x; \\ \frac{C_2}{\sqrt{p(x)}} \sin\left[\frac{1}{\hbar}\int_{x}^{x_2} p(x')\,dx' + \frac{\pi}{4}\right], & x < x_2 \end{cases} \tag{128}$$

Potential well with no vertical walls

$$\int_{x_1}^{x_2} p(x)\,dx = \left(n - \frac{1}{2}\right)\pi\hbar \tag{129}$$

Potential well with one vertical wall

$$\int_{0}^{x_2} p(x)\,dx = \left(n - \frac{1}{4}\right)\pi\hbar \tag{130}$$

Potential well with two vertical walls

$$\int_{0}^{a} p(x)\,dx = n\pi\hbar \tag{131}$$

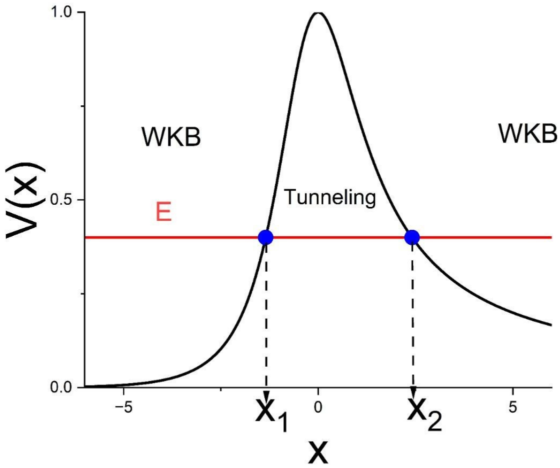


Figure 3: WKB Tunneling

$$K(x) \equiv \sqrt{2m[V(x) - E]} \tag{132}$$

$$\psi(x) \approx \frac{C}{\sqrt{K(x)}}\, e^{\pm\frac{1}{\hbar}\int K(x)\,dx} \qquad ; \quad E < V \tag{133}$$

$$\psi(x) \approx \begin{cases} \frac{C_1}{\sqrt{K(x)}} \exp\Big\{\Big[-\frac{1}{\hbar}\int_{x_1}^{x} K(x')\,dx'\Big]\Big\}, & x > x_1; \\ \frac{C_2}{\sqrt{K(x)}} \exp\big\{\big[-\frac{1}{\hbar}\int_{x}^{x_2} K(x')\,dx'\big]\big\}, & x < x_2. \end{cases} \tag{134}$$

$$T \sim e^{-2\gamma}, \qquad \gamma \equiv \frac{1}{\hbar}\int_0^a K(x)\,dx \tag{135}$$

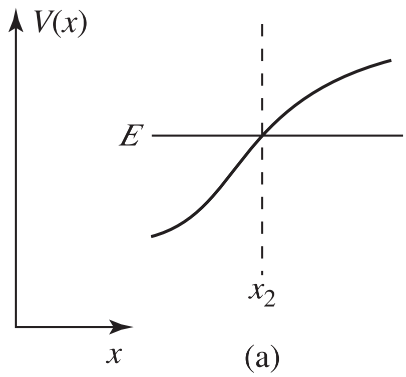


Figure 4: Right Turning Point

$$\psi(x) \approx \begin{cases} \frac{2D}{\sqrt{p(x)}} \sin\left[\frac{1}{\hbar}\int_x^{x_2} p(x')\,dx' + \frac{\pi}{4}\right], & x < x_2; \\ \frac{D}{\sqrt{K(x)}} \exp\left\{\left[-\frac{1}{\hbar}\int_{x_2}^{x} K(x')\,dx'\right]\right\}, & x > x_2. \end{cases} \tag{136}$$

Table 9.1: *Some properties of the Airy functions.*

*Differential Equation*: $\dfrac{d^2y}{dz^2} = zy.$

*Solutions*: Linear combinations of Airy functions, Ai($z$) and Bi(z).

*Integral Representation*: $\mathrm{Ai}(z) = \dfrac{1}{\pi}\displaystyle\int_0^\infty \cos\left(\frac{s^3}{3} + sz\right) ds,$

$$\mathrm{Bi}(z) = \frac{1}{\pi}\int_0^\infty \left[e^{-\frac{s^3}{3} + sz} + \sin\left(\frac{s^3}{3} + sz\right)\right] ds.$$

*Asymptotic Forms*:

$$\left.\begin{aligned} \mathrm{Ai}(z) &\sim \frac{1}{2\sqrt{\pi}z^{1/4}} e^{-\frac{2}{3}z^{3/2}} \\ \mathrm{Bi}(z) &\sim \frac{1}{\sqrt{\pi}z^{1/4}} e^{\frac{2}{3}z^{3/2}} \end{aligned}\right\} z \gg 0; \qquad \left.\begin{aligned} \mathrm{Ai}(z) &\sim \frac{1}{\sqrt{\pi}(-z)^{1/4}} \sin\left[\frac{2}{3}(-z)^{3/2} + \frac{\pi}{4}\right] \\ \mathrm{Bi}(z) &\sim \frac{1}{\sqrt{\pi}(-z)^{1/4}} \cos\left[\frac{2}{3}(-z)^{3/2} + \frac{\pi}{4}\right] \end{aligned}\right\} z \ll 0.$$

## Time Dependant Perturbation:

$$H(t) = H_o + H'(t); \quad H_0 |\psi_n\rangle = E_n |\psi_n\rangle$$
$$|\Psi(t)\rangle = \sum_n C_n(t) e^{-iE_n t/\hbar} |\psi_n\rangle$$

Schrodinger equation

$$i\hbar \frac{\partial}{\partial t} |\Psi(t)\rangle = H(t) |\Psi(t)\rangle$$

gives:

$$\begin{cases} i\hbar \dot{C}_k(t) = \sum_n H'_{kn}(t) e^{i\omega_{kn} t} C_n(t) \\ H'_{kn}(t) = \langle\psi_k| \langle H'(t)\rangle |\psi_n\rangle \\ \omega_{kn} = (E_k - E_n)/\hbar \end{cases}$$

with initial condition $|\psi(t_o)\rangle = |\Phi_o\rangle$ so that:

$$C_n(t_o) = \langle\psi_n|\Phi_o\rangle e^{-iE_n t_o/\hbar}$$

as initial condition for the above coupled equation in $C_n(t)$

## Two Level System:

$$|\Psi(t)\rangle = C_a(t) e^{-iE_a t/\hbar} |\psi_a\rangle + C_b(t) e^{-iE_b t/\hbar} |\psi_b\rangle$$
$$H_o |\psi_a\rangle = E_a |\psi_a\rangle ; \qquad H_o |\psi_b\rangle = E_b |\psi_b\rangle$$

Gives, assuming $H_{aa} = H_{bb} = 0$,

$$\begin{cases} \dot{C}_a(t) = -\frac{i}{\hbar} H'_{ab}(t) e^{-i\omega_o t} C_b(t) \\ \dot{C}_b(t) = -\frac{i}{\hbar} H'_{ba}(t) e^{i\omega_o t} C_a(t) \\ \omega_o = \frac{E_b - E_a}{\hbar} \end{cases}$$

Initial Condition $C_a(0) = 1; \quad C_b(0) = 0$

## First Order Perturbation:

$$C_b^{(1)} = -\frac{i}{\hbar}\int_{t_o}^{t} H'_{ba}(t')e^{i\omega_o t'}dt'$$

Transmission probability:

$$P_{ab} = \left|C_b^{(1)}\right|^2 = \frac{1}{\hbar^2}\left|\int_{t_o}^{t} H'_{ba}(t')e^{i\omega_o t'}dt'\right|^2$$

# Sinusoidal Perturbation:

$$\hat{H}'(\mathbf{r},t) = V(\mathbf{r})\cos(\omega t)$$

$$P_{ab}(t) = |C_b(t)|^2 \approx \frac{|V_{ab}|^2}{\hbar^2}\frac{\sin^2[(\omega_o-\omega)t/2]}{(\omega_o-\omega)^2}$$

$$\frac{dN_b}{dt} = -N_b A - N_b B_{ba}\rho(\omega_o) + N_a B_{ab}\rho(\omega_o) = 0$$