

Supporting Information: The Transport Mechanisms of Polar Solutes in a Cross-linked H_{II} Phase Lyotropic Liquid Crystal Membrane

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S1 Water content equilibration

We initially attempted to equilibrate our system with water by allowing water molecules to naturally penetrate the membrane from a water bath separating periodic images of the system in the z -direction (See Figure S1a). We allowed a dry, previously equilibrated system to further equilibrate in coexistence with a 3 nm-thick (in the z -direction) layer of water. Water readily enters the tail region where the density of monomers is low. About 3 times more water molecules occupy the tail region after 1000 ns of equilibration (See Figure S1b). Although the water level in the pore appears to plateau in this system, it is clear that equilibration of this system is kinetically limited since water does not fill the pores uniformly. The density of water along the pore axis, averaged over the last 50 nanoseconds of simulation, is close to zero at the membrane center. Therefore, we required a different equilibration technique in order to overcome the kinetic limitation.

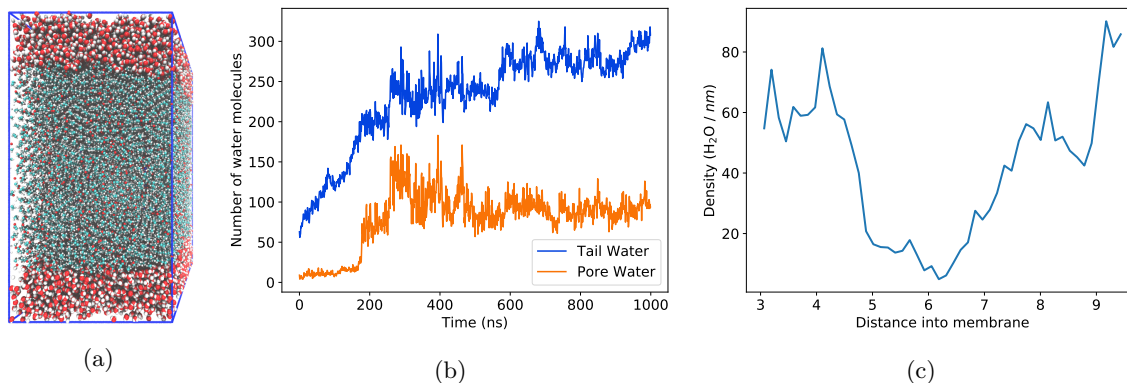


Figure S1: (a) Using an equilibrated dry configuration, we inserted a layer of water between periodic copies of the system in the z -direction. (b) Water slowly enters the membrane. Most water enters the tail region where the density of monomers is lowest. Water entering the pore plateaus after 500 ns. (c) Although the water content of the pore appears equilibrated in (b), the density of water throughout the pores is not uniform, with almost no water close to the pore center. Note that the density is not shown below 3 nm or above 9.5 nm, because the the system enters the water layer at those points.

We equilibrated 4 systems where we initially placed water in the pores and in the tails in addition to a water reservoir between periodic images. We filled the pores with water by running `gmx solvate` on our initial configuration and removing any water molecules placed outside of the pore region. The GROMACS command `gmx solvate` places water molecules in proximity to other atoms based on their van der waals radii and therefore does a decent job of preventing equilibration issues. This also means that the initial pore radius dictates the water content of the pore. We can however, put arbitrary amounts of water into the

tail region. We chose to test systems with initial pore radii of 5, 6, 7 and 8 Å with tail and pore water compositions given in Table S1.

Pore Radius (Å)	wt % water tails	wt % water pores
5	5.67	1.09
6	2.88	2.38
7	1.91	4.12
8	2.78	6.00

Table S1: We chose a diverse combination of initial pore and tail water contents in order to study its effect on equilibrium water content.

Systems appear to be most stable when there is more water in the pores than tails. In systems started with more water in the tails (Figures S2a and S2b), the pore water content tends to increase over time, while that of the tails decreases or stays stable. Filling the pores with water is likely a very long process since it requires monomers to make space for water molecules. Systems started with a higher pore water content (Figures S2c and S2d) tend to plateau relatively quickly, with about one third of the water staying in the tails. We used this ratio in order to construct the initial configurations used for the studies in the main text. Clearly, a more complicated methodology is needed in order to predict the equilibrium water content of a given LLC membrane. However, getting the value exactly right is not required for our study. Instead, we can observe mechanisms as a function of water content in each region.

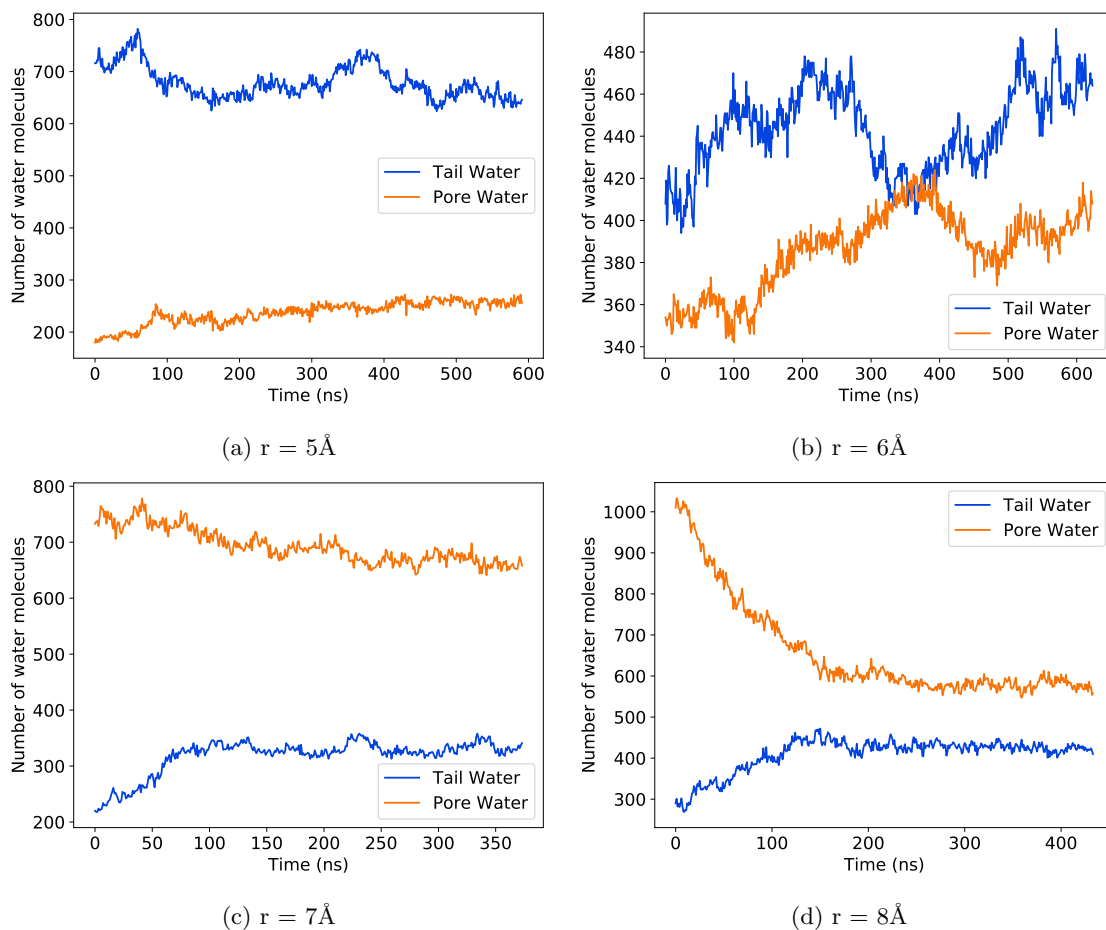


Figure S2

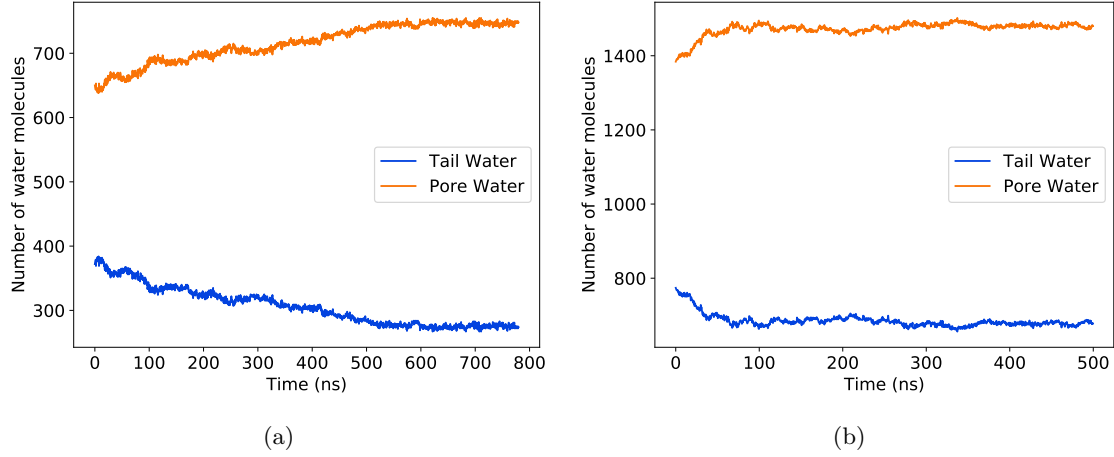


Figure S3: We created solvated systems with one third of the total water initially placed in the tail region. (a) With 5 wt % total water, the water content equilibrates after 600 ns, with $\sim 72\%$ of the total water in the pores. (b) With 10 wt % total water, the water content equilibrates after 100 ns, with $\sim 69\%$ the total water in the pores.

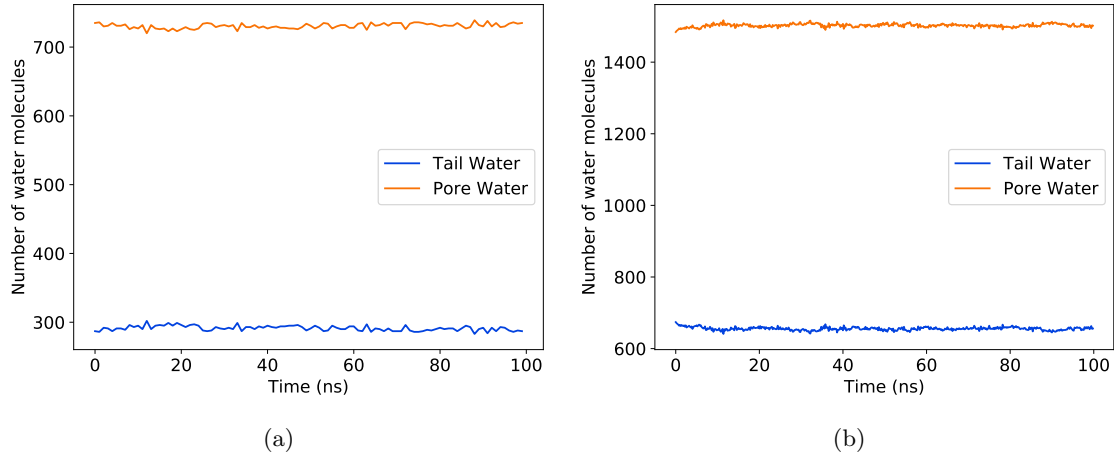


Figure S4

Since the equilibrium water content is unclear based on the previous simulations, we elected to choose and study systems with two different water contents. We removed the water reservoir and allowed the pore and tail water contents to equilibrate with 5 and 10 wt % total water. We placed one third of the total water needed in the tails, based on the intuition gained in Figure S2. We considered the water content equilibrated once the water contents plateaued. The 10 wt % water system equilibrated within the first 100 ns of simulation, while the 5 wt% system did not plateau until ~ 600 ns (See Figure S4). The pores contain 72 % and 69 % of the total water in the 5 and 10 wt % systems respectively.

We cross-linked the equilibrated solvated systems, then allowed them to equilibrate further for 100 ns.

S2 Choosing a transport model

We used the toolbox created by Meroz and Sokolov in order to justify our choice of transport model.[1] The solutes in our systems exhibit anomalous transport properties characteristic of a Continuous Time Random Walk (CTRW).

Mean Squared Displacement

The general form of a mean squared displacement (MSD) curve is:

$$\langle x^2(t) \rangle \sim t^\alpha \quad (1)$$

For brownian motion, $\alpha = 1$ and the MSD is linear. When $\alpha \neq 1$, the particle of interest exhibits anomalous diffusion. Values of α greater than 1 give rise to superdiffusion, while values of α less than 1 give rise to subdiffusion.

We can calculate the ensemble-averaged MSD curve by averaging the MSDs of each particle trajectory, where each MSD is calculated using:

$$\delta^2(t) = \|\mathbf{r}(t) - \mathbf{r}(0)\|^2 \quad (2)$$

where $\|\cdot\|$ represents the Euclidean norm.

The mean squared displacement of solutes in our model is a non-linear function of time, with $\alpha < 1$ which is indicative of anomalous subdiffusion. Figure S5a plots the ensemble-averaged MSD curve for 24 ethanol molecules diffusing in a 10 wt% water H_{II} LLC membrane system. We fit a power law of the form Ae^α to the MSD curve. We performed 2000 bootstrap trials by randomly sampling 24 MSD curves with replacement from the 24 total ethanol MSD curves. The bootstrapped average value of α is 0.75 for this system.

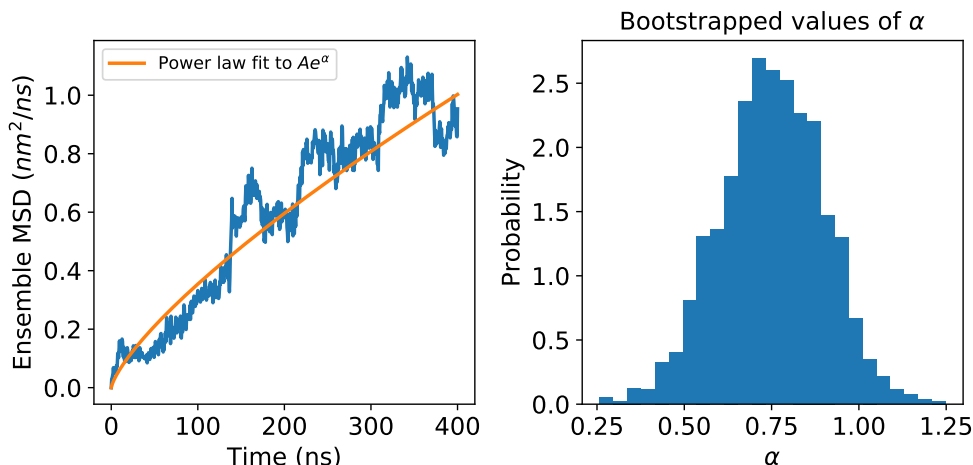


Figure S5: (a) We fit a curve with the form of Equation 1 to the ensemble-averaged MSD curve. (b) The average value of α , obtained using fits to MSDs calculated from bootstrapped ensembles, is less than 1 suggesting that ethanol molecules in our model exhibit subdiffusive behavior.

Ergodicity

The ergodicity of a system can help us narrow down the possible anomalous diffusion mechanisms. In an ergodic system, the time-averaged behavior of an observable should yield the same result as the ensemble average of the same observable. Examples of anomalous diffusion processes that are ergodic include random walks on fractals (RWF) and fractional brownian motion (FBM). Non-ergodic systems generally give rise to CTRWs with the possibility of combination with a RWF and/or FBM.[1]

We tested the ergodicity of our system by comparing the ensemble-averaged and time-averaged MSD curves. We calculated the MSD of each ethanol trajectory using Equation 2 and a time-averaged algorithm:

$$\delta^2(t) = \frac{1}{N-t} \sum_{i=0}^{N-t-1} \|\mathbf{r}(i+t) - \mathbf{r}(i)\|^2 \quad (3)$$

where N is the total number of simulation frames, and t represents the length of subinterval or number of frames per subinterval. We averaged the MSD curves from each trajectory in order to create final MSD plots.

The ethanol molecules exhibit non-ergodic behavior because their time-averaged and ensemble-averaged MSDs do not agree with each other (Figure S6a). We validated our analysis using a 1 ns simulation of a box of tip3p water molecules. As expected, since the particles exhibit Brownian motion, the time-averaged and ensemble-averaged MSDs agree with each within error (Figure S6b).

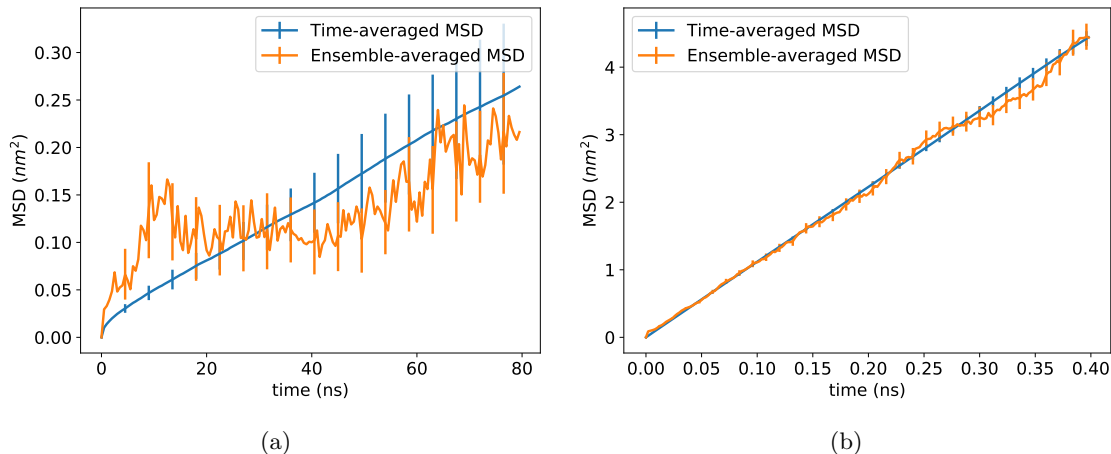


Figure S6: (a) The time-averaged and the ensemble-averaged MSDs for ethanol in an H_{II} nanopore are not in agreement, implying non-ergodicity. (b) A box of tip3p water molecules is expected to be ergodic and it is shown to be true here because both MSDs are in agreement.

Autocorrelation of steps

Based on the previous two sections, our model can likely be studied as a CTRW. However, it is still possible that our CTRW model might also be convoluted with an FBM or a RWF process. In a pure CTRW, the steps are uncorrelated. Both FBM and RWF exhibit anti-correlated steps.

The steps in our system are not correlated. We showed this by calculating the autocorrelation function (ACF) of the step lengths in the z -direction. The ACF of a representative trajectory is shown in Figure S7.

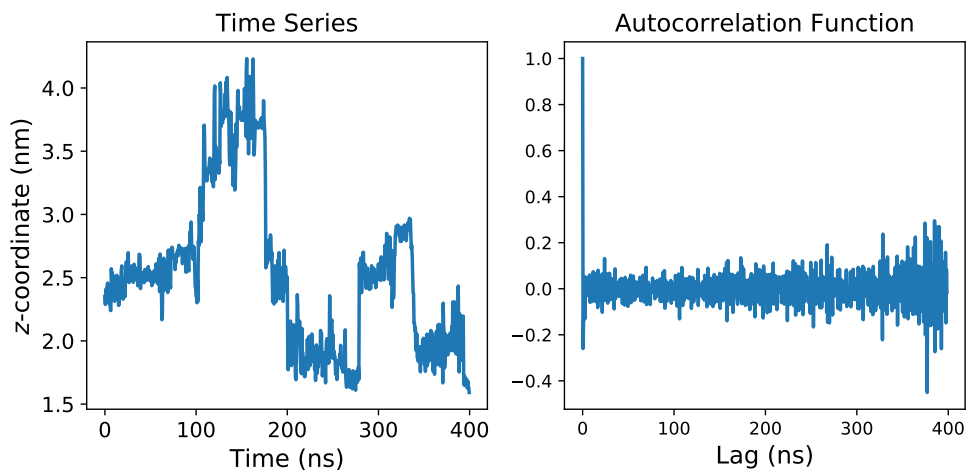


Figure S7: The autocorrelation function (right) of a representative ethanol center of mass z -coordinate trajectory (left) almost immediately decays to zero, indicating a complete loss of memory of its previous position. Noise increases at large time lags due to decreased sampling.

References

- [1] Y. Meroz and I. M. Sokolov, “A Toolbox for Determining Subdiffusive Mechanisms,” *Physics Reports*, vol. 573, pp. 1–29, Apr. 2015.