

Fourier Transforms!
PHYS 250 (Autumn 2018) – Lecture 13

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Outline

- 1 *Reminders*
 - Reminders from Lecture 12
- 2 *Follow-up with the Square Wave*
 - Square Wave
- 3 *Extension to the Fourier Transform*
 - Euler's Formula
 - Fourier Transforms

Reminders from last time

We discussed both Poisson's equation and Fourier Series in the last lecture.

PDEs and Fourier Series

- **PDEs: Poisson's Equation**

- We discussed and then wrote a function for computing the solution to Poisson's equation
- Specifically, we elaborated on how we could structure the (relatively simple) function in order to take various constraints (i.e. sources) for the solution into account more easily

- **Fourier Series:**

- Started discussing the basics of Fourier Series
- Evaluated, computationally, the coefficients of a simple series for both a square and sawtooth wave

Today we will go much more in depth with Fourier Transforms and Analysis!
This will be a mixture of Python Notebooks and Lecture Slides

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Square Wave

First, I want to do a little follow-up with the Fourier Series for the **square wave** function that we discussed last time.

Recall that we determined the coefficients of a sine and cosine expansion to be:

$$a_n = \frac{2}{n\pi} \sin\left(n\omega_0 \frac{\pi}{2}\right) \quad (1)$$

which yields a discrete function:

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos(nx) \quad (2)$$

Now open up the `Fourier-Transforms-Analysis.ipynb` jupyter notebook so that we can look at this in more detail!

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Euler's Formula

We can generalize this by making use of Euler's Formula.

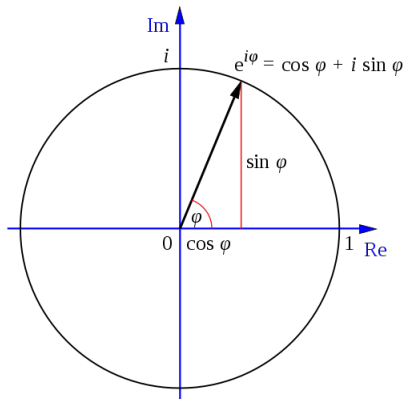
Euler's formula states that for any real number ϕ :

$$e^{i\phi} = \cos(\phi) + i \sin(\phi) \quad (3)$$

When $\phi = \pi$, Euler's formula evaluates to

$$e^{i\pi} + 1 = 0, \quad (4)$$

which is known as Euler's identity.



The implication is that it is possible to recover the amplitude of each wave in a Fourier series using an integral, which has many useful properties (in particular, that it's then continuous).

Fourier Transforms (I)

I will use the following definitions for the Fourier transform $\hat{f}(\xi)$ of a function $f(x)$, where x typically represents either a **spatial or time domain**, and ξ typically represents a corresponding inverse notion of **spatial or time frequency**.

$$\hat{f}(\xi) = \int_{-\infty}^{\infty} f(x) e^{-2\pi i \xi x} dx \quad (5)$$

The **inverse transform** is then obtained via

$$f(x) = \int_{-\infty}^{\infty} \hat{f}(\xi) e^{2\pi i \xi x} d\xi \quad (6)$$

In the case of spatial coordinates, x denotes length and ξ denotes inverse wavelength: $\xi = \frac{1}{\lambda}$. In the time domain, x denotes time and ξ denotes frequency. In the case that $x = t$ is in seconds, but ξ is **angular** frequency ω then a factor of 2π appears to get the normalization correct.

Fourier Transforms (II)

In the case of spatial coordinates, x denotes length and ξ denotes inverse wavelength: $\xi = \frac{1}{\lambda}$. In the time domain, x denotes time and ξ denotes frequency. In the case that $x = t$ is in seconds, but ξ is **angular** frequency ω then a factor of 2π appears to get the normalization correct.

$$\hat{f}(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(t) e^{-i\omega t} dt \quad (7)$$

$$f(t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(\omega) e^{i\omega t} d\omega \quad (8)$$

Since $\omega = 2\pi\xi = \frac{2\pi}{\lambda}$.

The $\frac{1}{\sqrt{2\pi}}$ factor in both these integrals is a common normalization in quantum mechanics but maybe not in engineering where only a single $\frac{1}{2\pi}$ factor is often used.

Discrete Fourier Transforms (I)

If $\hat{f}(\omega)$ or $f(t)$ are known analytically or numerically, the Fourier transform integrals can be evaluated using the integration techniques studied earlier.

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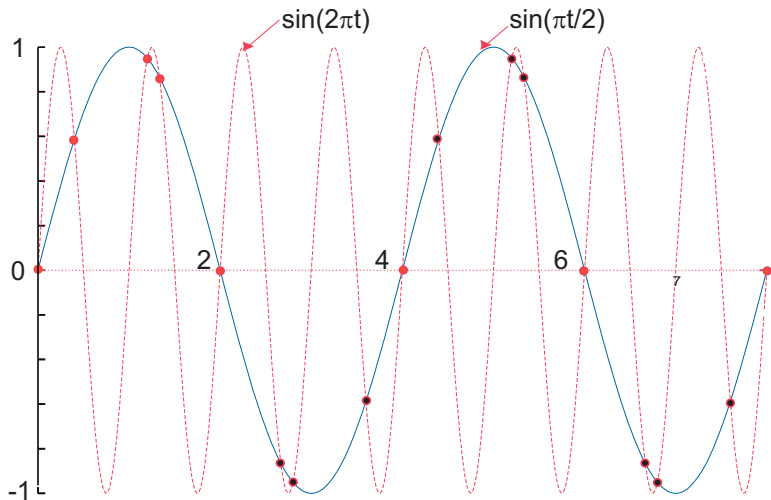
In this way the **DFT can be thought of as a technique for interpolating, compressing, and extrapolating data.**

Discussion

Do you see any issues with this “sampling”?

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Discrete Fourier Transforms (II)

The DFT algorithm results from evaluating the integral not from -1 to $+1$ but rather from time 0 to time T over which the signal is measured, and from approximating the integration of the integral by computing a discrete sum:

$$\hat{f}(\omega_n) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(t) e^{-i\omega_n t} dt \quad (9)$$

$$\simeq \frac{1}{\sqrt{2\pi}} \int_0^T f(t) e^{-i\omega_n t} dt \quad (10)$$

$$\simeq \frac{1}{\sqrt{2\pi}} \sum_{k=1}^N h f(t_k) e^{-i\omega_n t_k} \quad (h \equiv \text{stepsize}) \quad (11)$$

$$\simeq \frac{h}{\sqrt{2\pi}} \sum_{k=1}^N f_k e^{-2\pi i k n / N} \quad (12)$$

$$\hat{f}_n \equiv \frac{\hat{f}(\omega_n)}{h} = \frac{1}{\sqrt{2\pi}} \sum_{k=1}^N f_k e^{-2\pi i k n / N} \quad (13)$$

Discrete Fourier Transforms (III)

We then need the inverse as well, which we can obtain with $d\omega \rightarrow 2\pi/Nh$ we invert the \hat{f}_n

$$f_k = \frac{1}{\sqrt{2\pi}} \sum_{n=1}^N \frac{2\pi}{Nh} \hat{f}_n e^{i\omega_n t} \quad (14)$$

Once we know the N values of the transform \hat{f}_n , we can use this expression to evaluate $f(t)$ for any time t . The frequencies ω_n are determined by the number of samples taken and by the total sampling time $T = Nh$ as

$$\omega_n = n \frac{2\pi}{Nh} \quad (15)$$

Clearly, the larger we make the time $T = Nh$ over which we sample the function, the smaller will be the frequency steps or resolution. Accordingly, if you want a smooth frequency spectrum, you need to have a small frequency step $2\pi/T$.

Discrete Fourier Transforms (IV)

Lastly, we can simplify this expression to yield a clear computational approach:

$$f_k = \frac{\sqrt{2\pi}}{N} \sum_{n=1}^N Z^{-nk} \hat{f}_n \quad (Z = e^{-2\pi i/N}) \quad (16)$$

$$\hat{f}_n = \frac{1}{\sqrt{2\pi}} \sum_{k=1}^N Z^{nk} f_k \quad (n = 0, 1, \dots, N) \quad (17)$$

With this formulation, the computer needs to compute only powers of Z .