Fourier Transforms!
PHYS 250 (Autumn 2019) – Lecture 11

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Outline

Moving towards physics data analysis

As we discussed last time, I would like to take the direction of this quarter more towards practical physics data analysis and algorithms. We will start with **Fourier Transforms** and **Neural Networks** and then analyze data from the Large Hadron Collider.

Fourier Analysis and Neural Networks

- **Fourier Series and Analysis:**
  - Discussing the basics of Fourier Series
  - Evaluate, computationally, the coefficients of a simple series for both a square and sawtooth wave
  - Extend discussion to the **Fast Fourier Transform**
- **Neural Networks:**
  - Training computers to discover, identify, and analyze patterns of interest in datasets
  - Modeling perspective on what a neural network achieves
  - Structure and function of a neuron
  - Mathematical properties of a neural network

This will be a mixture of Python Notebooks and Lecture Slides

Outline

Square Wave

Let's start with the Fourier Series for the **square wave** function.

As you may have heard in a variety of contexts (e.g. PHYS 133!) we can decompose any periodic function or periodic signal into the weighted sum of a (possibly infinite) set of simple oscillating functions, namely sines and cosines (or, equivalently, complex exponentials).

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos(nx) + b_n \sin(nx) \tag{1}$$

This is possible because the trigonometric functions for a **set of complete, orthogonal basis vectors** that span the space.

Now open up the `Fourier-Series.ipynb` jupyter notebook and we will discuss this more deeply!

Fourier series for a square wave

We may determine the coefficients of a sine and cosine expansion to be:

$$a_n = \frac{2}{n\pi} \sin\left(n\omega_0 \frac{\pi}{2}\right) \tag{2}$$

which yields a discrete function:

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos(nx) \tag{3}$$

Now open up the `Fourier-Transforms-Analysis.ipynb` jupyter notebook so that we can look at this in more detail!

Outline

Euler's Formula

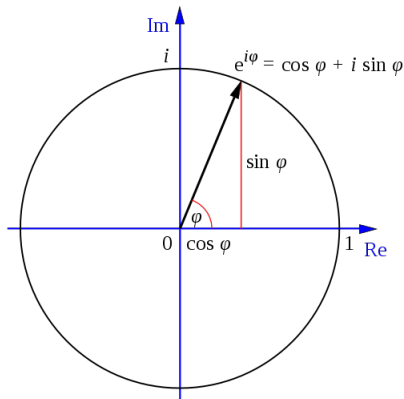We can generalize this by making use of Euler's Formula.

Euler's formula states that for any real number ϕ:

$$e^{i\phi} = \cos(\phi) + i\sin(\phi) \qquad (4)$$

When $\phi = \pi$, Euler's formula evaluates to

$$e^{i\pi} + 1 = 0, \qquad (5)$$

which is known as Euler's identity.

The implication is that it is possible to recover the amplitude of each wave in a Fourier series using an integral, which has many useful properties (in particular, that it's then continuous).

Fourier Transforms (I)

I will use the following definitions for the Fourier transform $\hat{f}(\xi)$ of a function $f(x)$, where x typically represents either a **spatial or time domain**, and ξ typically represents a corresponding inverse notion of **spatial or time frequency**.

$$\hat{f}(\xi) = \int_{-\infty}^{\infty} f(x)e^{-2\pi i \xi x} dx \tag{6}$$

The **inverse transform** is then obtained via

$$f(x) = \int_{-\infty}^{\infty} \hat{f}(\xi)e^{2\pi i \xi x} d\xi \tag{7}$$

In the case of spatial coordinates, x denotes length and ξ denotes inverse wavelength: $\xi = \frac{1}{\lambda}$. In the time domain, x denotes time and ξ denotes frequency. In the case that $x = t$ is in seconds, but ξ is **angular** frequency ω then a factor of 2π appears to get the normalization correct.

Fourier Transforms (II)

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$$
\begin{align}
\hat{f}(\omega) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(t) e^{-i\omega t} dt \tag{8} \\
f(t) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(\omega) e^{i\omega t} d\omega \tag{9}
\end{align}
$$

Since $\omega = 2\pi\xi = \frac{2\pi}{\lambda}$.

The $\frac{1}{\sqrt{2\pi}}$ factor in both these integrals is a common normalization in quantum mechanics but maybe not in engineering where only a single $\frac{1}{2\pi}$ factor is often used.

Discrete Fourier Transforms (I)

If $\hat{f}(\omega)$ or $f(t)$ are known analytically or numerically, the Fourier transform integrals can be evaluated using the integration techniques studied earlier.

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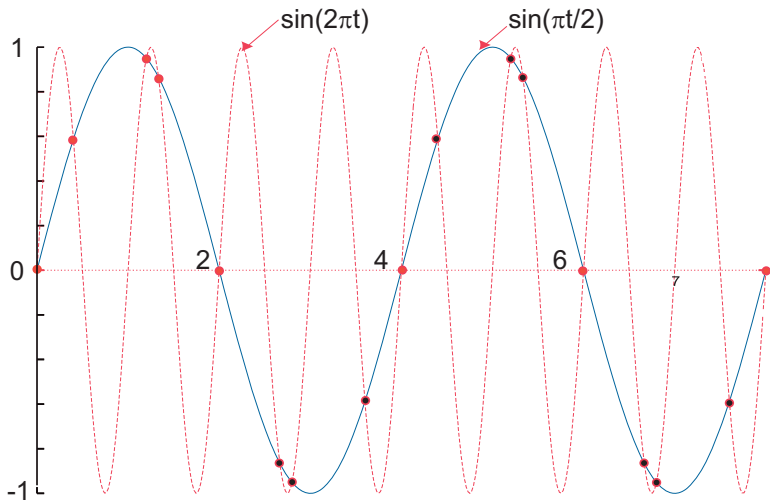
In this way the **DFT can be thought of as a technique for interpolating, compressing, and extrapolating data**.

Do you see any issues with this "sampling"?

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sin(2πt) sin(πt/2)

Discrete Fourier Transforms (II)

The DFT algorithm results from evaluating the integral not from 1 to $+1$ but rather from time 0 to time T over which the signal is measured, and from approximating the integration of the integral by computing a discrete sum:

$$
\begin{align}
\hat{f}(\omega_n) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(t) e^{-i\omega_n t} dt \tag{10}\\
&\simeq \frac{1}{\sqrt{2\pi}} \int_{0}^{T} f(t) e^{-i\omega_n t} dt \tag{11}\\
&\simeq \frac{1}{\sqrt{2\pi}} \sum_{k=1}^{N} h f(t_k) e^{-i\omega_n t_k} \qquad (h \equiv \text{stepsize}) \tag{12}\\
&\simeq \frac{h}{\sqrt{2\pi}} \sum_{k=1}^{N} f_k e^{-2\pi i k n / N} \tag{13}\\
\hat{f}_n \equiv \frac{\hat{f}(\omega_n)}{h} &= \frac{1}{\sqrt{2\pi}} \sum_{k=1}^{N} f_k e^{-2\pi i k n / N} \tag{14}
\end{align}
$$

Discrete Fourier Transforms (III)

We then need the inverse as well, which we can obtain with $d\omega \to 2\pi/Nh$ we invert the $\hat{f}_n$

$$f_k = \frac{1}{\sqrt{2\pi}} \sum_{n=1}^{N} \frac{2\pi}{Nh} \hat{f}_n e^{i\omega_n t} \tag{15}$$

Once we know the N values of the transform $\hat{f}_n$, we can use this expression to evaluate $f(t)$ for any time t. The frequencies ωn are determined by the number of samples taken and by the total sampling time $T = Nh$ as

$$\omega_n = n\frac{2\pi}{Nh} \tag{16}$$

Clearly, the larger we make the time $T = Nh$ over which we sample the function, the smaller will be the frequency steps or resolution. Accordingly,if you want a smooth frequency spectrum, you need to have a small frequency step $2\pi/T$.

Lastly, we can simplify this expression to yield a clear computational approach:

$$f_k = \frac{\sqrt{2\pi}}{N} \sum_{n=1}^{N} Z^{-nk} \hat{f}_n \qquad (Z = e^{-2\pi i/N}) \tag{17}$$

$$\hat{f}_n = \frac{1}{\sqrt{2\pi}} \sum_{k=1}^{N} Z^{nk} f_k \qquad (n = 0, 1, \cdots, N) \tag{18}$$

With this formulation, the computer needs to compute only powers of Z.