

Lecture 09 – chapter 32

Inductance

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Before Lecture

Week 10

The General Form of Faraday's Law

Why no negative sign?

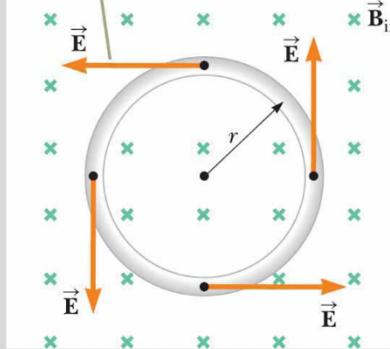
$$\Delta V = \oint \vec{E} \cdot d\vec{s}$$

$$E \cdot \oint d\vec{s} = -\frac{d}{dt}(BA)$$

$$\Rightarrow E(2\pi r) = -\frac{dB}{dt}(\pi r^2)$$

$$\Rightarrow E = -\frac{r}{2} \frac{dB}{dt}$$

If \vec{B} changes in time, an electric field is induced in a direction tangent to the circumference of the loop.



CJLI

Week 3

Electric Potential (電位)

- A charged particle moving in an electric field will experience a change in potential

$$\Delta V \equiv \frac{\Delta U}{q} = - \int_A^B \vec{E} \cdot d\vec{s}$$

A negative sign

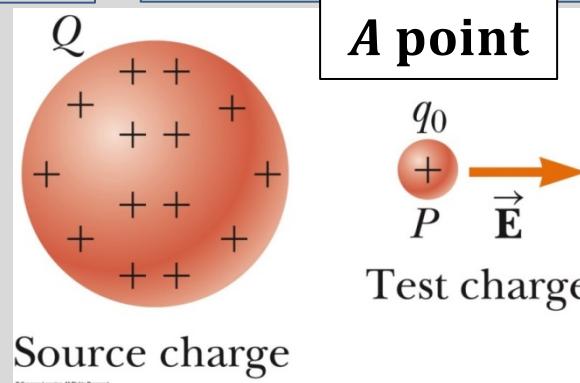
(25.2)
(25.3)

It indicates that the $V_B < V_A$.

- The potential energy per unit charge, U/q_0
- Has a value at every point in an electric field
- Scalar quantity, Units: **V** (voltage, volt, 伏特)
- Electron-Volts
- Unit of energy commonly used in atomic and nuclear physics
- $1 \text{ eV} = 1.60 \times 10^{-19} \text{ C} \cdot \text{V} = 1.60 \times 10^{-19} \text{ J}$

Section 25.1

20



Electric field lines always point in the direction of decreasing electric potential

Before Lecture

Week 10

The General Form of Faraday's Law

$$\Delta V = \oint \vec{E} \cdot d\vec{s}$$

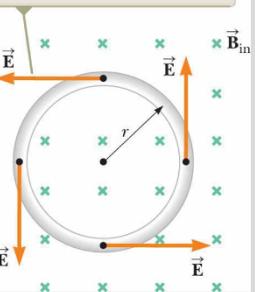
If $A=B$, a closed loop, $\Delta V \neq 0$

$$E \cdot \oint d\vec{s} = -\frac{d}{dt}(BA)$$

$$\Rightarrow E(2\pi r) = -\frac{dB}{dt}(\pi r^2)$$

$$\Rightarrow E = -\frac{r}{2} \frac{dB}{dt}$$

If \vec{B} changes in time, an electric field is induced in a direction tangent to the circumference of the loop.



CJLI

Induced emf and Electric Fields

- Changing magnetic flux \rightarrow An electric field in the conductor
 - Physical loop built up by conductor

$$\varepsilon = -\frac{d\Phi_B}{dt}$$

- Changing magnetic field \rightarrow An electric field in empty space
 - For any closed path, the emf can be expressed as the line integral of electrical field over the path

$$\varepsilon = \oint \vec{E} \cdot d\vec{s}$$

- Faraday's law written in a general form:

$$\oint \vec{E} \cdot d\vec{s} = -\frac{d\Phi_B}{dt}$$

- The induced electric field by magnetic field, not by the charges

- The induced field cannot be an electrostatic field and is non-conservative

Section 31.4

53

Week 3

Electric Potential (電位)

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$$\Delta V \equiv \frac{\Delta U}{q} = -\oint_A^B \vec{E} \cdot d\vec{s} \quad (25.2)$$

(25.3)

- The potential energy per unit charge, U/q_0
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If $A=B$, $\Delta V = 0$

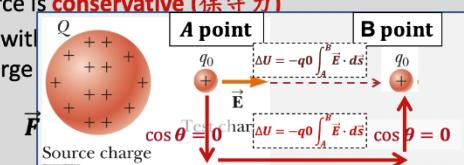
Section 25.1

20

Electric Potential Energy

- Electrostatic force is **conservative** (保守力)

- The work done with field on the charge



- The potential energy (電位能) of the charge-field system

$$\Delta U = -q \int_A^B \vec{E} \cdot d\vec{s}$$

Because the force is conservative, the line integral does not depend on the path taken by the charge

Introduction

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The induced electric field is different in nature from the electric field from the stationary charge

Previous Lecture

- Some History

$$\varepsilon = - \frac{d\Phi_B}{dt}$$

- Faraday's Law of Induction

- Motional emf

- A moving conductor in the \vec{B} field
 - If the above is part of a closed circuit,

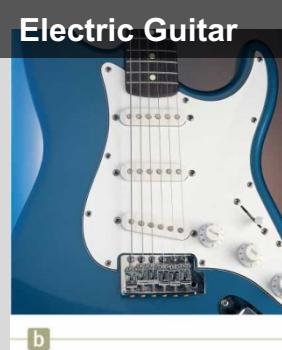
- Lenz's law

$$\oint \vec{E} \cdot d\vec{s} = - \frac{d\Phi_B}{dt}$$

- The General Form of Faraday's Law

- Generators & Motors

- Eddy Currents



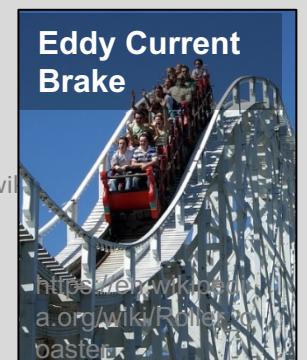
Electric guitars



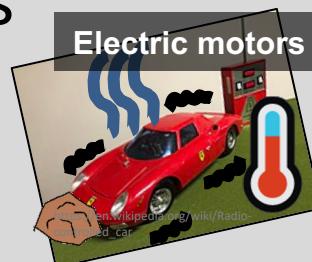
Electric Generator



Electric motors



Eddy Current Brake



Electric motors



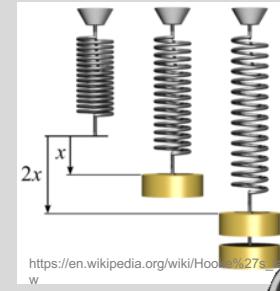
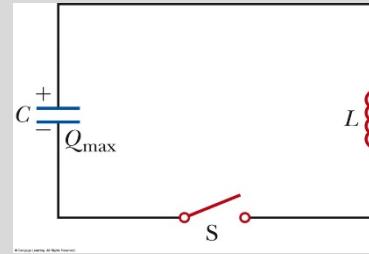
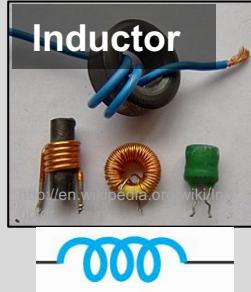
Induction cooking



台灣高鐵

This Lecture

- Inductance
 - Self-Inductance
- *RL* Circuit
 - Energy in a Magnetic Field
- Energy Storage Summary
 - Charged capacitor, Inductor, Resistor
- Mutual Inductance
- *LC* Circuits
- *RLC* Circuit



Inductance (電感)

- A time-varying current in a circuit or a coil produces a changing magnetic flux
 - Self-inductance : an induced emf on self **opposing** initial emf **自感**
 - Mutual induction : an induced emf on another coil **互感**
- The electrical circuit element called an **inductor**
 - Energy is stored in the magnetic field
 - Electrical field in capacitor

Self-Inductance

自感

- When the switch is closed
 - Before reaching maximum, the increasing flux creates an induced emf **opposite** the direction of the emf of the battery by **Lenz's law**

- L is called the **inductance** (電感) of the coil

$$L = \frac{\Phi_B}{i} \Rightarrow \varepsilon_L = -\frac{d\Phi_B}{dt} = -L \frac{di}{dt} \quad (32.1)$$

- SI unit of inductance is the **henry (H)**
- Example of solenoid (電磁閥, 螺管線圈)
 - N turns and length ℓ
 - A common circuit element act like capacitor

$$\Phi_B = NBA = N\mu_0 niA = N\mu_0 \frac{N}{l} iA = \frac{\mu_0 N^2 i A}{l}$$

$$L = \frac{\Phi_B}{i} = \mu_0 \frac{N^2}{l} A = \mu_0 n^2 V$$

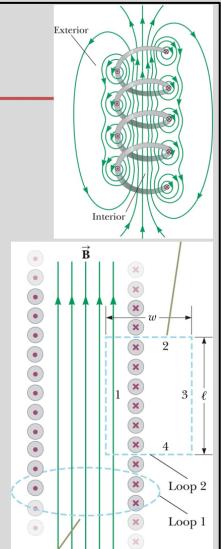
Solenoid (電磁閥) Week 8

- As the length increases
 - The interior field becomes more uniform
 - The exterior field becomes weaker
- Ideal Solenoid
 - The turns are closely spaced
 - The length \gg radius of the turns

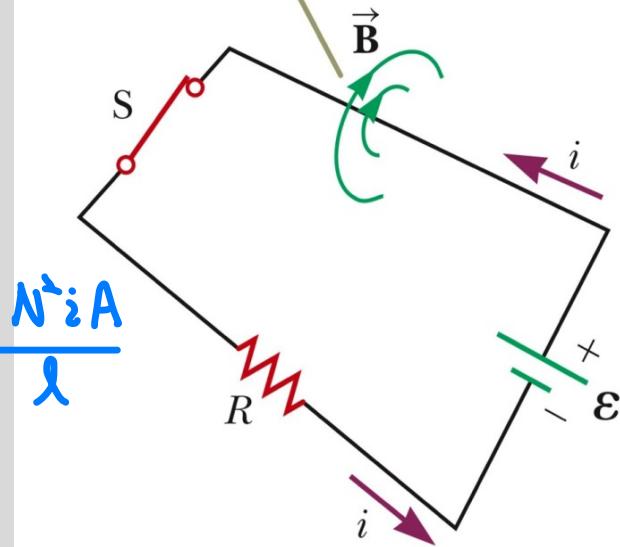
✓ Ampere's Law for loop 2

$$\oint \vec{B} \cdot d\vec{s} = B\ell = \mu_0 NI \Rightarrow B = \mu_0 \frac{N}{\ell} I = \mu_0 nI$$

$n = N/\ell$ is the number of turns per unit length



Section 30.4



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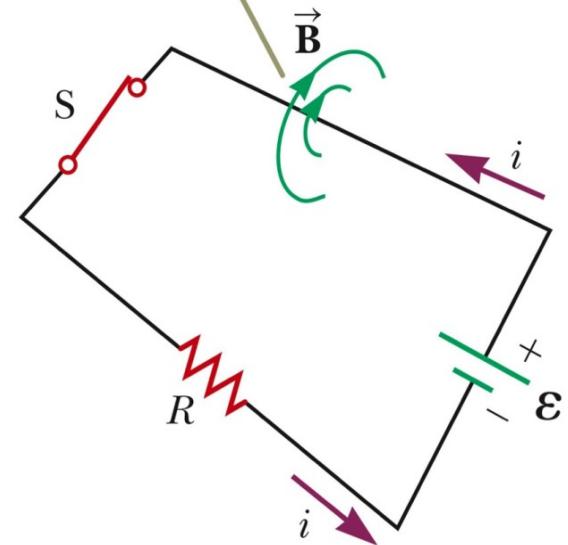
Self-Inductance

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 - N turns and length ℓ
 - A common circuit element act like capacitor

$$\Phi_B = NBA = N\mu_0 n i A = N\mu_0 \frac{N}{l} i A$$

$$L = \frac{\Phi_B}{i} = \mu_0 \frac{N^2}{l} A = \mu_0 n^2 V$$

After the switch is closed, the current produces a magnetic flux through the area enclosed by the loop. As the current increases toward its equilibrium value, this magnetic flux changes in time and induces an emf in the loop.



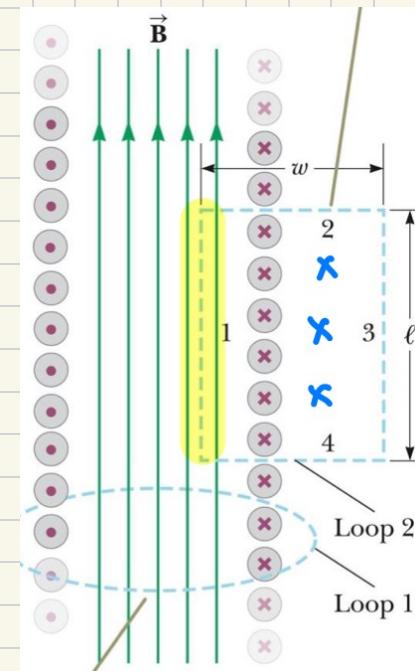
Self Inductance of a Solenoid

1° B in a solenoid (Ampere's Law)

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 I$$

$$\Rightarrow B \cdot l = \mu_0 (NI)$$

$$\Rightarrow B = \mu_0 \frac{N}{l} I = \mu_0 n I$$



2° $\phi_B = N(BA) = N(\mu_0 n i)A$

3° $L = \frac{\phi_B}{i} = N \mu_0 n A = \mu_0 \frac{N^2}{l^2} (Al) = \mu_0 n^2 V_{\#}$

RL Circuit

- Inductor (電感) 
- Usually self-inductance of the rest of the circuit is negligible compared to the inductor

$v_L(t) = L \frac{di(t)}{dt}$ Because the inductance of a coil produces a back emf, an inductor in a circuit opposes any change in current and prevents the current from changing instantaneously.

- When switch S_1 close and S_2 set to a (at $t = 0$)

$$\mathcal{E} - i(t)R - L \frac{di(t)}{dt} = 0$$

$$\frac{\mathcal{E}}{R} - i - \frac{L}{R} \frac{di}{dt} = 0 \quad \text{Let } x = (\mathcal{E}/R) - i \quad \text{and} \quad dx = -di$$

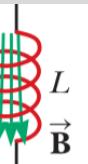
Rearrange this $x = -\frac{L}{R} \frac{dx}{dt}$ $-\frac{R}{L} dt = \frac{dx}{x}$

$$\Rightarrow x + \frac{L}{R} \frac{dx}{dt} = 0 \quad \int_{x_0}^x \frac{dx}{x} = -\frac{R}{L} \int_0^t dt \rightarrow \ln \frac{x}{x_0} = -\frac{R}{L} t$$

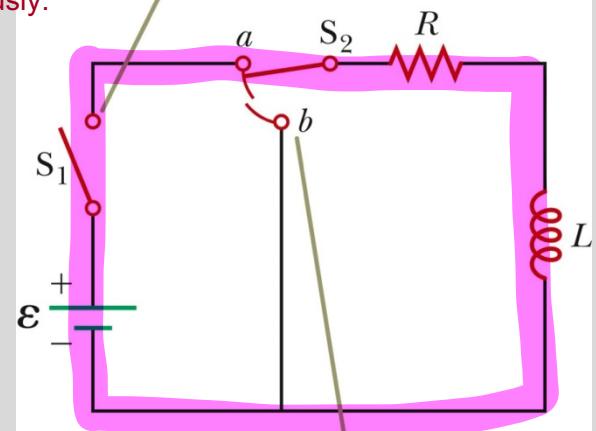
$$\Rightarrow \frac{\mathcal{E}}{R} - i = \frac{\mathcal{E}}{R} e^{-Rt/L} \quad x = x_0 e^{-Rt/L}$$

$x = (\mathcal{E}/R) - i$
 $i = 0 \text{ at } t = 0$

$$\rightarrow i = \frac{\mathcal{E}}{R} (1 - e^{-Rt/L}) \quad x_0 = \mathcal{E}/R$$



When switch S_1 is thrown closed, the current increases and an emf that opposes the increasing current is induced in the inductor.



When the switch S_2 is thrown to position b, the battery is no longer part of the circuit and the current decreases.

* We cannot use KVL here! As there is an inductor in the loop.

$$\oint \vec{E} \cdot d\vec{l} = - \frac{d\phi_B}{dt} \quad (\text{Faraday's Law})$$

$$\Rightarrow - \oint \vec{E} \cdot d\vec{l} = \frac{d\phi_B}{dt}$$

1. $- \oint \vec{E} \cdot d\vec{l} = \Delta V$

$$\Rightarrow \epsilon - iR = L \frac{di}{dt}$$

2. $L = \frac{\phi_B}{i} \Rightarrow \phi_B = iL$

$$\Rightarrow \frac{\epsilon}{R} - i = \frac{L}{R} \frac{di}{dt}$$

Let $\frac{\epsilon}{R} - i = x$

$$\Rightarrow x = - \frac{L}{R} \frac{dx}{dt}$$

$\Rightarrow dx = - di$

$$\Rightarrow \frac{dx}{dt} = -\frac{R}{L}x \quad (\text{exponential decay equation})$$

$$\Rightarrow x = x_0 e^{-Rt/L}$$

1. $x_0 = \frac{\epsilon}{R} - i(0) = \frac{\epsilon}{R}$

$$\Rightarrow \frac{\epsilon}{R} - i = \frac{\epsilon}{R} e^{-Rt/L}$$

2. $x = \frac{\epsilon}{R} - i$

$$\Rightarrow i = \frac{\epsilon}{R} (1 - e^{-Rt/L})$$

τ (time constant) = $\frac{L}{R}$

$$\Rightarrow i = \frac{\epsilon}{R} (1 - e^{-\frac{t}{\tau}})$$

RL Circuit

- Inductor (電感) 

– Usually self-inductance of the rest of the circuit is negligible compared to the inductor

$$v_L(t) = L \frac{di(t)}{dt}$$

- When switch S_1 close and S_2 set to a (at $t = 0$)

$$\epsilon - i(t)R - L \frac{di(t)}{dt} = 0 \quad = \frac{\epsilon}{R} (1 - e^{-\frac{t}{\tau}})$$

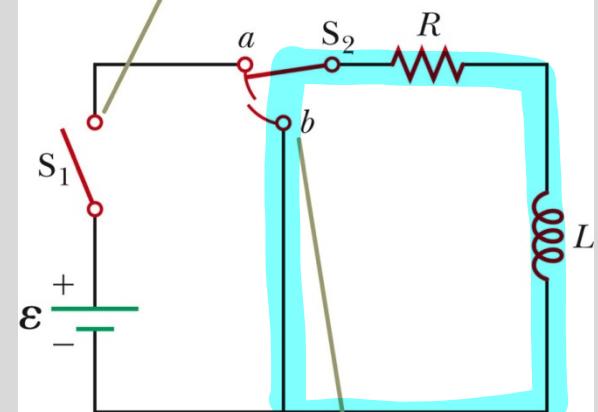
$$i(t) = \frac{\epsilon}{R} - \left(\frac{\epsilon}{R} - i(0) \right) e^{-Rt/L} = \frac{\epsilon}{R} (1 - e^{-Rt/L})$$

- Now set S_2 to position b

$$i(t)R + L \frac{di(t)}{dt} = 0 \Rightarrow i(t) = i(0)e^{-Rt/L} = \frac{\epsilon}{R} e^{-Rt/L}$$

$$\tau = L / R$$

When switch S_1 is thrown closed, the current increases and an emf that opposes the increasing current is induced in the inductor.



When the switch S_2 is thrown to position b, the battery is no longer part of the circuit and the current decreases.

$$\oint \vec{E} \cdot d\vec{l} = - \frac{d\phi_B}{dt}$$

$$\Rightarrow - \oint \vec{E} \cdot d\vec{l} = \frac{d\phi_B}{dt} \quad \boxed{L = \frac{\phi_B}{i} \Rightarrow \phi_B = Li}$$

$$\Rightarrow -iR = L \frac{di}{dt}$$

$$\Rightarrow \frac{di}{dt} = -\frac{R}{L} i \quad (\text{Exponential decay equation})$$

$t=0$, steady state (DC)

For an inductor L ,

$$V = L \frac{di}{dt} \Rightarrow V = \frac{di}{dt} = 0$$

$\Rightarrow L$ behaves like a short circuit

$$\Rightarrow i_0 = \frac{\epsilon}{R}$$

$$\Rightarrow i = \frac{\epsilon}{R} e^{-Rt/L}$$

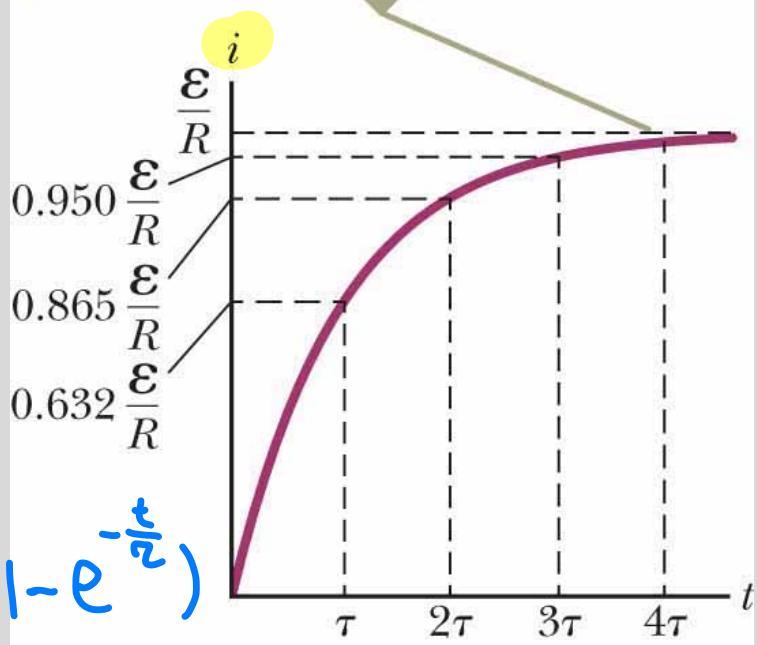
τ (time constant) = $\frac{L}{R}$

$$\Rightarrow i = \frac{\epsilon}{R} e^{-t/\tau}$$

RL Circuit: i-t plot

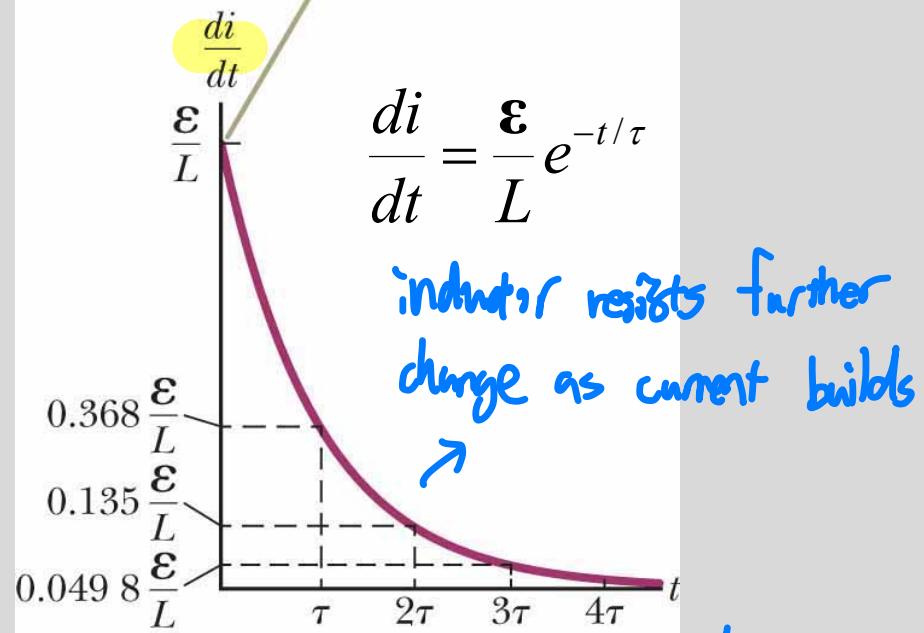
$$\frac{d(i(t))}{dt} = \frac{\epsilon}{R} (-e^{-\frac{t}{\tau}}) (-\frac{1}{L}) \\ = \frac{\epsilon}{L} e^{-\frac{t}{\tau}}$$

After switch S_1 is thrown closed at $t = 0$, the current increases toward its maximum value ϵ/R .



$$(1 - e^{-1}) = 0.632 = 63.2\%$$

The time rate of change of current is a maximum at $t = 0$, which is the instant at which switch S_1 is thrown closed.



$$i(t) = \frac{\epsilon}{R} e^{-\frac{t}{\tau}}$$

Inductor

網路資訊，同學要記得獨立思考，自主查證



已由「The science works」設為置頂

@Thescienceworks 1年前

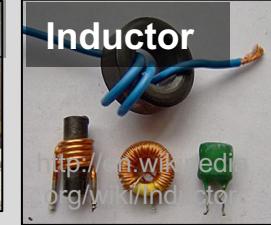
The direction of magnetic field is incorrect in the wire field demonstration, sorry for this silly mistake.

322 回覆

CJLI



https://en.wikipedia.org/wiki/Induction_cooking



<http://en.wikipedia.org/wiki/Inductor>

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Energy in a Magnetic Field

power P

- For the energy rate in watt (J/s)

$$i\epsilon = i^2 R + L i \frac{di}{dt}$$

- $i\epsilon$: energy supplied by the battery
- $i^2 R$: energy delivered to the resistor
- $L i (di/dt)$: energy is being stored in the magnetic field

- For the energy (J)

$$\frac{dU_B}{dt} = L i \frac{di}{dt} \quad U = L \int_0^i i \, di = \frac{1}{2} L i^2$$

② Energy Stored in an Inductor's Magnetic Field U_B

$$1^{\circ} \quad \mathcal{E} = -\frac{d\phi_B}{dt} \text{ (Faraday's Law)} \quad \boxed{L = \frac{\phi_B}{i}}$$

$$\Rightarrow \mathcal{E} = -L \frac{di}{dt} \quad \Rightarrow \phi_B = Li$$

$$\Rightarrow V_L = L \frac{di}{dt} \quad \begin{matrix} \text{(we often drop the minus sign)} \\ \text{by the passive sign convention} \end{matrix}$$

$$2^{\circ} \quad P = IV = i \left(L \frac{di}{dt} \right) = Li \frac{di}{dt}$$

3° Power is the rate of energy

$$\frac{dU_B}{dt} = P = L \cdot i \cdot \frac{di}{dt}$$

multiply both
sides by dt

$$\Rightarrow dU_B = L \cdot i \cdot di$$

$$\Rightarrow \int dU_B = \int_0^i L \cdot i \cdot di$$

$$\Rightarrow U_B = L \int_0^i i \cdot di$$

$$\Rightarrow U_B = \frac{1}{2} L i^2 *$$

For solenoid

- Energy (J)

$$L = \mu_0 n^2 V \quad B = \mu_0 n i$$

$$U_B = \frac{1}{2} L i^2$$

Previous slide

Self-Inductance

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- L is called the **inductance (電感)** of the coil

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$$L = \frac{\Phi_B}{i} = \mu_0 \frac{N^2}{l} A = \mu_0 n^2 V$$

Solenoid (電磁閥)
Week 8

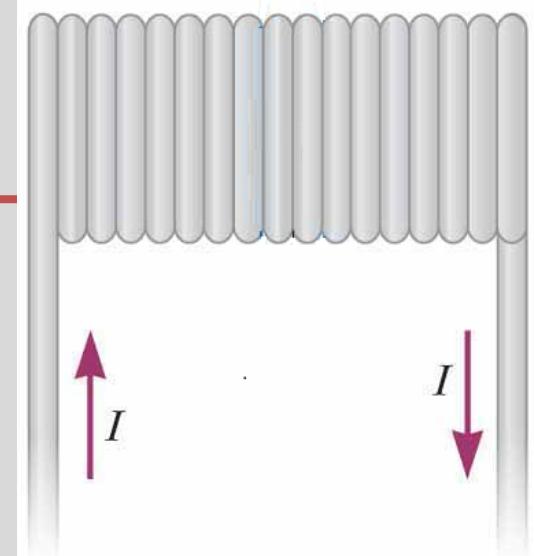
- As the length increases
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 - The exterior field becomes weaker
- Ideal Solenoid
 - The turns are closely spaced
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Ampere's Law for loop 2

$$\oint \vec{B} \cdot d\vec{s} = B\ell = \mu_0 NI \Rightarrow B = \mu_0 \frac{N}{\ell} = \mu_0 nI$$

$n = N/\ell$ is the number of turns per unit length

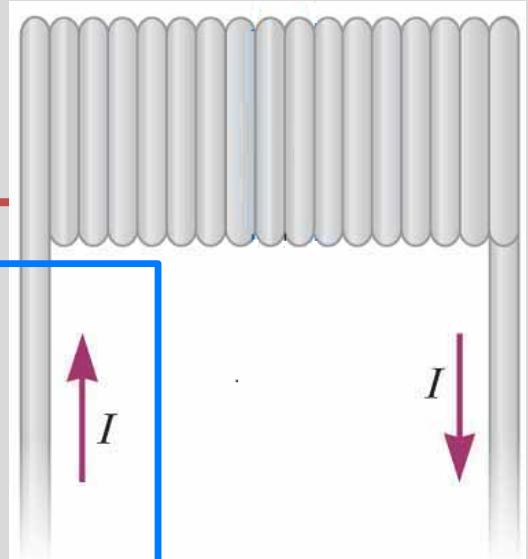
Section 32.4



For solenoid

- Energy (J)

$$L = \mu_0 n^2 V \quad B = \mu_0 n i$$



$$U_B = \frac{1}{2} L i^2 = \frac{1}{2} \mu_0 n^2 V \left(\frac{B}{\mu_0 n} \right)^2 = \frac{B^2}{2 \mu_0} V$$

- Energy density (J/m³)

$$u_B = \frac{U_B}{V} \Rightarrow u_B = \frac{B^2}{2 \mu_0}$$

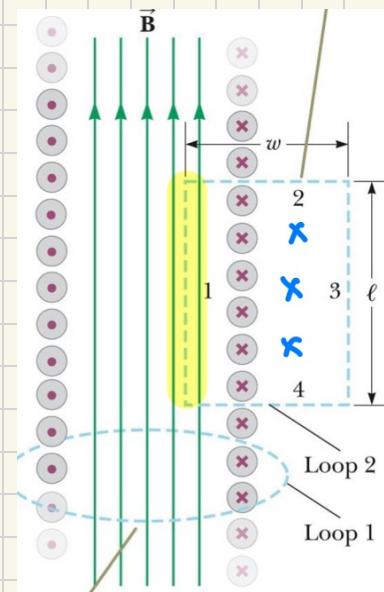
② Energy and Energy Density

1' Ampere's Law around Loop 2

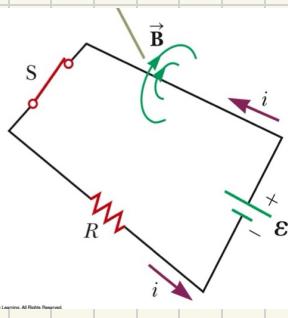
$$\oint \vec{B} \cdot d\vec{s} = \mu_0 I$$

$$\Rightarrow Bl = \mu_0 Ni$$

$$\Rightarrow B = \mu_0 \frac{N}{l} i = \mu_0 n i$$



$$\begin{aligned}
 2' \quad L &= \frac{\phi_B}{i} = \frac{NBA}{i} = \frac{N(\mu_0 n i) A}{i} \\
 &= N \mu_0 n A = \mu_0 \frac{N^2}{l^2} (Al) = \mu_0 n^2 V
 \end{aligned}$$



$$3^{\circ} \text{ Energy } U_B = \frac{1}{2} L i^2$$

$$= \frac{1}{2} (\mu_0 n^2 V) \left(\frac{B}{\mu_0 n} \right)^2$$

$$= \frac{B^2}{2 \mu_0} V \quad \#$$

$$4^{\circ} \text{ Energy Density } U_B = \frac{U_B}{V} = \frac{B^2}{2 \mu_0} \#$$

$$1. \text{ From } 1^{\circ}, \quad B = \mu_0 n i$$

$$\Rightarrow i = \frac{B}{\mu_0 n}$$

$$2. \text{ From } 2^{\circ}, \quad L = \mu_0 n^2 V$$

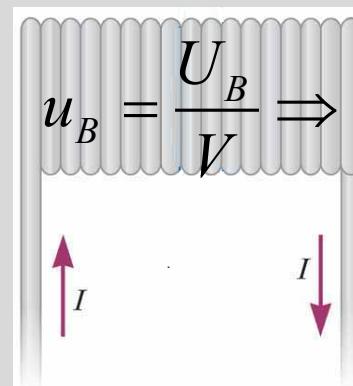


Quick Quiz 31.3

- You are performing an experiment that requires the highest-possible **magnetic energy density** in the interior of a very long current-carrying solenoid. Which of the following adjustments increases the energy density? (More than one choice may be correct.)
 - (a) increasing the number of turns per unit length on the solenoid $n \uparrow$
 - (b) increasing the cross-sectional area of the solenoid
 - (c) increasing only the length of the solenoid while keeping the number of turns per unit length fixed
 - (d) increasing the current in the solenoid $i \uparrow$

highest-possible
magnetic energy density

Answers: (a), (d)



$$i \uparrow$$

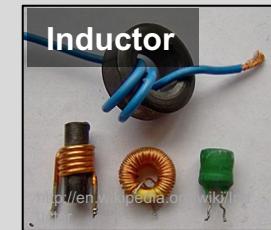
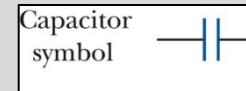
$$B = \mu_0 n i$$

$$u_B = \frac{B^2}{2\mu_0}$$

$$= \frac{\mu_0 n^2 i^2}{2}$$

Energy Storage Summary

- A resistor, inductor and capacitor all store energy through different mechanisms
- Charged capacitor
 - Stores energy as **electric potential energy** (Electrical field)
- Inductor
 - When it carries a current, stores energy as **magnetic potential energy** (magnetic field)
- Resistor
 - Energy delivered is transformed into **internal energy**



Resistor



Section 32.3

the heat increases the random motion of atoms in a resistor

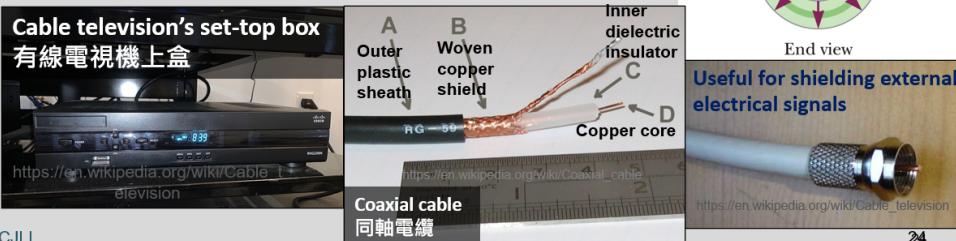
Example: The Coaxial Cable

- Coaxial cables are often used to connect electrical devices, such as your video system, and in receiving signals in television cable systems

Week5 Example 27.3

- Coaxial cables are used extensively for cable television and other electronic applications
 - Two concentric cylindrical conductors
 - The region between the conductors is completely filled with polyethylene plastic
 - Unwanted current leakage through the plastic, in the *radial* direction
- $a = 0.500 \text{ cm}$, $b = 1.75 \text{ cm}$, $L = 15.0 \text{ cm}$. The resistivity of the plastic is $1.0 \times 10^{13} \Omega \cdot \text{m}$.

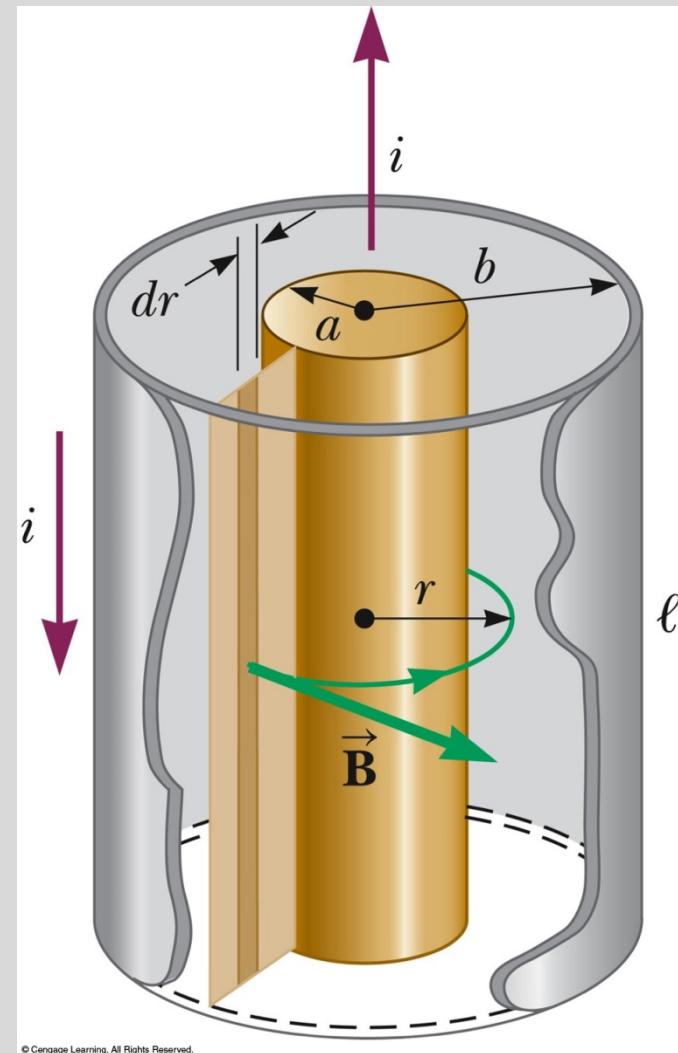
Week4



CJLI

Section 32.3

24



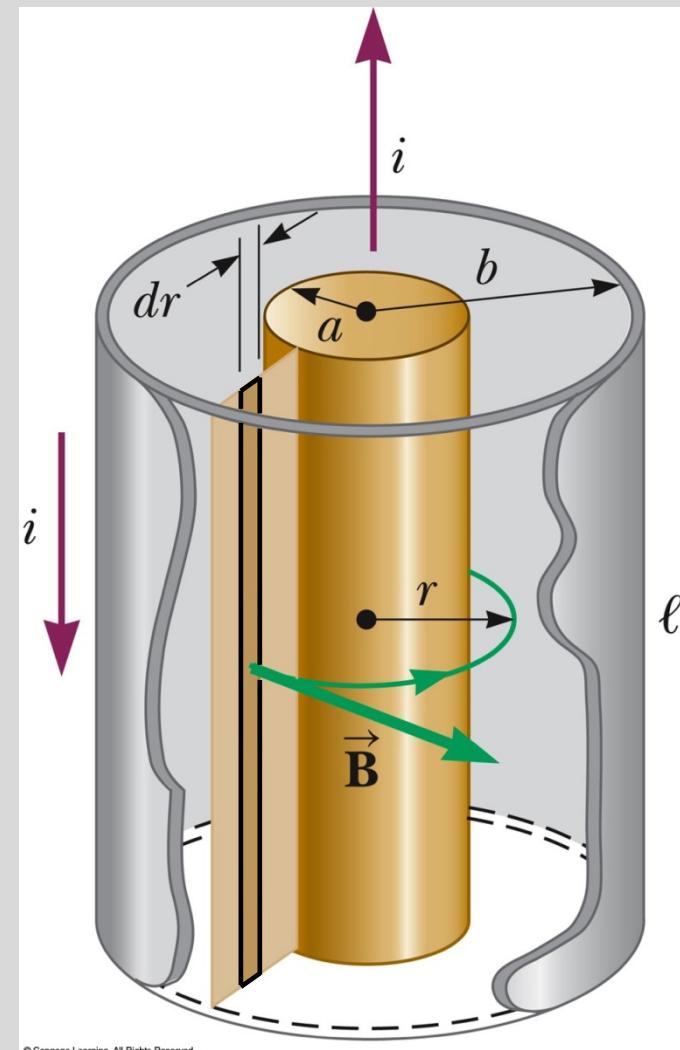
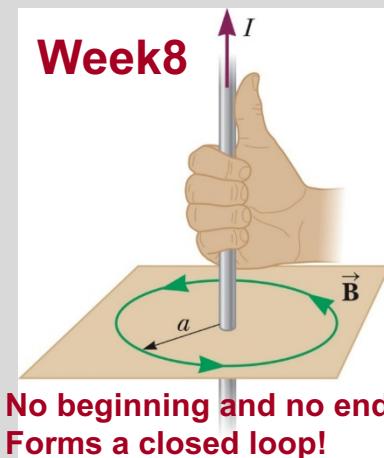
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29

Example: The Coaxial Cable

- Coaxial cables are often used to connect electrical devices, such as your video system, and in receiving signals in television cable systems
- Model a long coaxial cable as a thin, cylindrical conducting shell of radius b concentric with a solid cylinder of radius a

$$d\Phi_B = BdA = B\ell dr$$



Example: The Coaxial Cable

A Long, Straight Week8 Conductor

$$d\vec{s} \times \hat{r} = |d\vec{s} \times \hat{r}| \hat{k} = \left[dx \sin\left(\frac{\pi}{2} - \theta\right) \right] \hat{k} = (dx \cos \theta) \hat{k}$$

$$(1) d\vec{B} = (dB) \hat{k} = \frac{\mu_0 I}{4\pi} \frac{dx \cos \theta}{r^2} \hat{k}$$

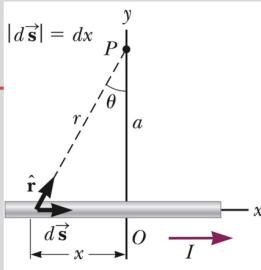
$$(2) r = \frac{a}{\cos \theta}$$

$$x = -a \tan \theta$$

$$(3) dx = -a \sec^2 \theta d\theta = -\frac{a d\theta}{\cos^2 \theta}$$

$$(4) dB = -\frac{\mu_0 I}{4\pi} \left(\frac{a d\theta}{\cos^2 \theta} \right) \left(\frac{\cos^2 \theta}{a^2} \right) \cos \theta = -\frac{\mu_0 I}{4\pi a} \cos \theta d\theta$$

$$B = -\frac{\mu_0 I}{4\pi a} \int_{\theta_1}^{\theta_2} \cos \theta d\theta = \frac{\mu_0 I}{4\pi a} (\sin \theta_1 - \sin \theta_2) \quad (30.4)$$

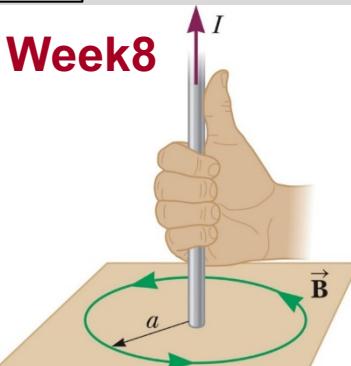


connect
video
n television
hin.

An infinitely long,
straight wire,
 $\theta_1 = \pi/2$ and $\theta_2 = -\pi/2$

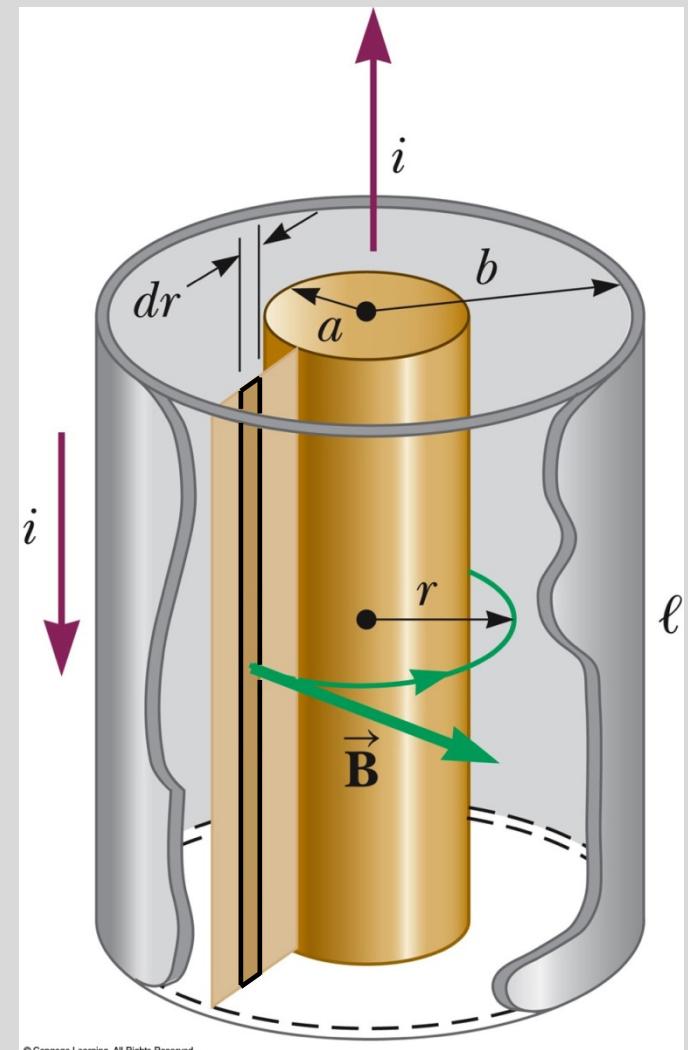
**Same magnitude
at a distance of a
radius a**

$$d\Phi_B = BdA = B\ell dr$$



**No beginning and no end
Forms a closed loop!**

Section 32.3



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Example: The Coaxial Cable

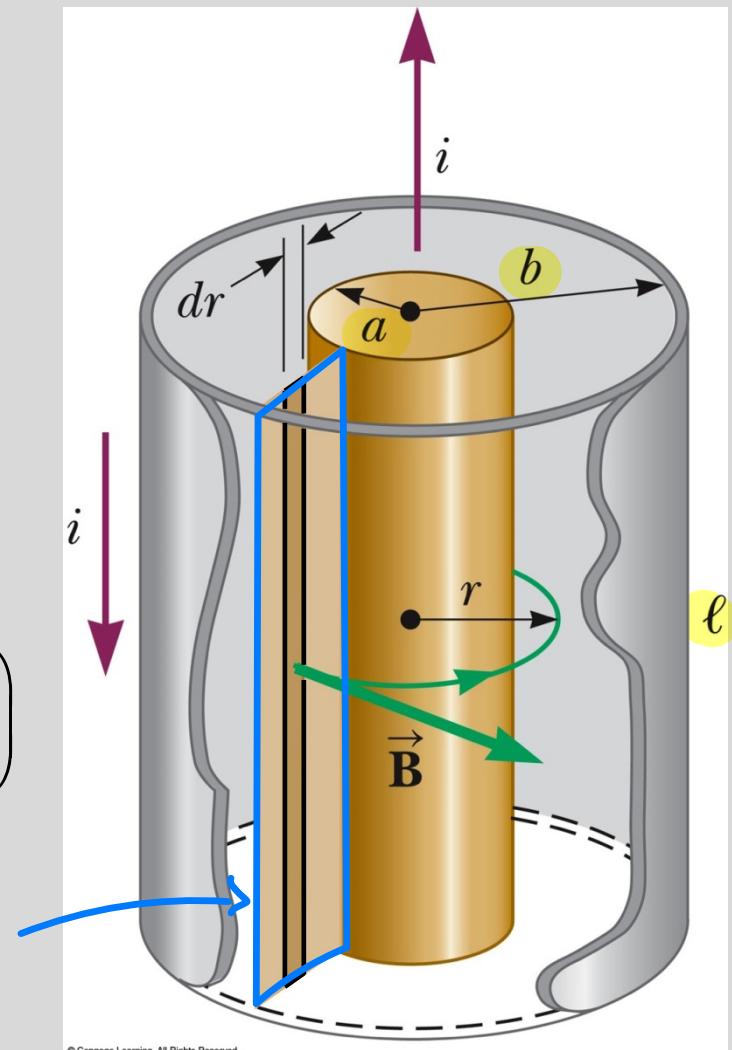
- Coaxial cables are often used to connect electrical devices, such as your video system, and in receiving signals in television cable systems
- Model a long coaxial cable as a thin, cylindrical conducting shell of radius b concentric with a solid cylinder of radius a

$$d\Phi_B = BdA = B\ell dr$$

$$\Phi_B = \int_a^b \frac{\mu_0 i}{2\pi r} \ell dr = \frac{\mu_0 i \ell}{2\pi} \int_a^b \frac{dr}{r} = \frac{\mu_0 i \ell}{2\pi} \ln\left(\frac{b}{a}\right)$$

$$L = \frac{\Phi_B}{i} = \frac{\mu_0 \ell}{2\pi} \ln\left(\frac{b}{a}\right)$$

chosen surface



② Inductance in a Coaxial Cable

$$L = \frac{\Phi_B}{i}$$

$$= \frac{\int B \cdot dA}{i}$$

1. $dA = l dr$

2. Obtain B using Ampere's Law

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 I \Rightarrow B = \frac{\mu_0 i}{2\pi r}$$

$$= \frac{\int_a^b \frac{M_0 i}{2\pi r} l dr}{i}$$

$$= \int_a^b \frac{M_0 l}{2\pi r} dr$$

$$= \frac{M_0 l}{2\pi} \int_a^b \frac{dr}{r}$$

$$= \frac{M_0 l}{2\pi} \ln \frac{b}{a} *$$

1. Why do the currents flow in opposite directions in the inner and outer conductors?

This is **intentional and necessary** for coaxial cables.

- The **inner wire** carries current $+i$ (say, upward).
- The **outer cylindrical shell** carries **return current** $-i$ (downward).

Why this matters:

- **Net external magnetic field = 0** outside the coaxial cable.
- The opposing currents **cancel each other's magnetic field** outside the cable.
- This makes coaxial cables highly resistant to **EM interference, signal radiation, and cross-talk.**

| ✓ This current configuration confines the magnetic field **between** the conductors — where inductance is calculated.

Mutual Inductance

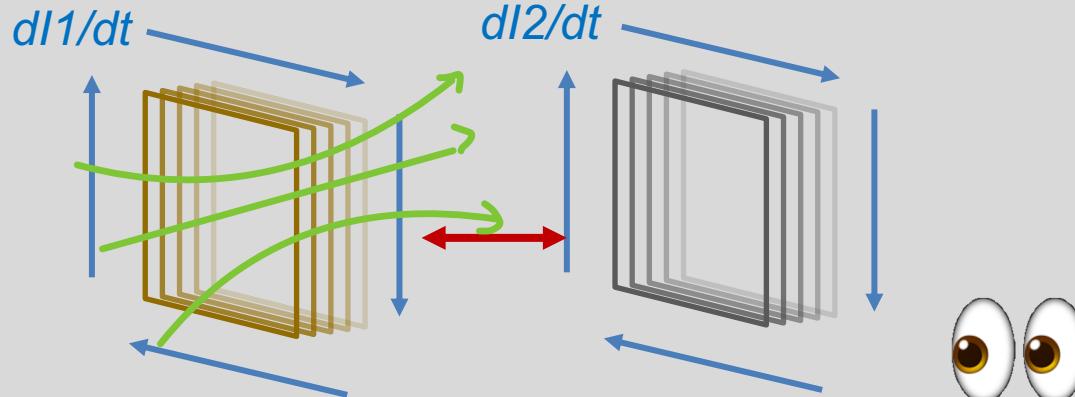
互感

- If current i_1 varies with time, the emf induced

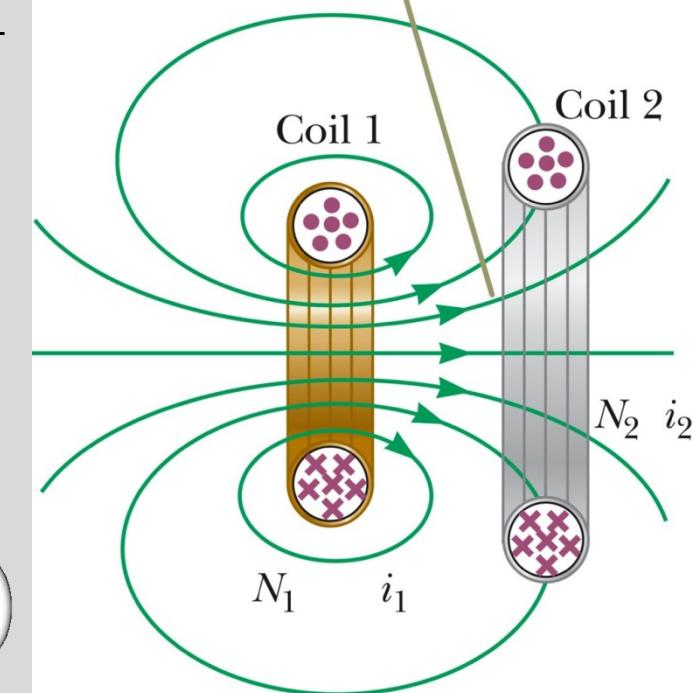
$$\mathbf{\epsilon}_2 = -N_2 \frac{d\Phi_{12}}{dt} = -M_{12} \frac{di_1}{dt} \quad (32.15)$$

- Mutual inductance M_{12} of coil 2 with respect to coil 1

$$M_{12} = \frac{N_2 \Phi_{12}}{i_1} \quad (32.16)$$



A current in coil 1 sets up a magnetic field, and some of the magnetic field lines pass through coil 2.



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Mutual Inductance

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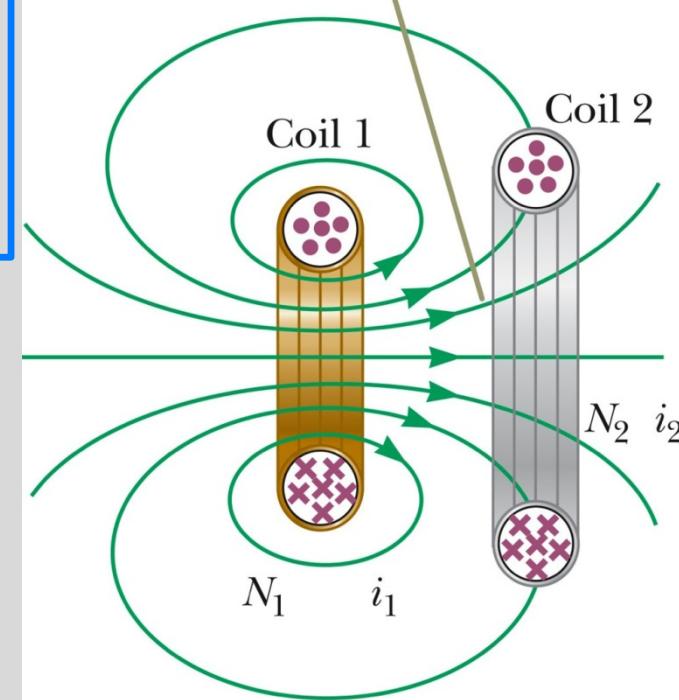
- If current 2 varies with time, the emf induced


$$\mathbf{\epsilon}_1 = -M_{21} \frac{di_2}{dt}$$

- $M_{12} = M_{21} = M$ (mutual inductance)

$$\mathbf{\epsilon}_1 = -M \frac{di_2}{dt} \quad \text{and} \quad \mathbf{\epsilon}_2 = -M \frac{di_1}{dt}$$

A current in coil 1 sets up a magnetic field, and some of the magnetic field lines pass through coil 2.



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① $i = \text{constant}$

$$N_2 \phi_{B_2} = M i_1 \Rightarrow M = \frac{N_2 \phi_{B_2}}{i_1}$$

② $i \neq \text{constant}$

$$\left\{ \begin{array}{l} N_2 \frac{d\phi_{B_2}}{dt} = M \frac{di_1}{dt} \\ \hline \end{array} \right.$$

$$\left\{ \begin{array}{l} \mathcal{E}_2 = - N_2 \frac{d\phi_{B_2}}{dt} \\ \hline \end{array} \right.$$

differentiate both sides

$$\Rightarrow \mathcal{E}_2 = - M \frac{di_1}{dt}$$

(Faraday's Law of Induction)

What is this formula actually doing?

This equation **does not imply** that mutual inductance depends on i_1 in real time.

Instead, it's a **definition** that calculates M_{12} based on a **test value** of current i_1 and the resulting flux Φ_{12} through coil 2:

M_{12} = Flux through coil 2 caused by unit current in coil 1

Important clarification:

- Mutual inductance M is a **property of the geometry and material** (e.g., number of turns, coil spacing, area, core permeability).
 - It is **constant** for a fixed setup.
-

So why is $M \propto \frac{1}{i_1}$ in this equation?

Because you're using this as a **one-time measurement**:

- You apply current i_1 ,
- You **measure** the resulting flux linkage $N_2\Phi_{12}$,
- You calculate:

$$M_{12} = \frac{N_2\Phi_{12}}{i_1}$$

If you **double** i_1 , the flux Φ_{12} will also double → and the ratio stays **constant**.

In short:

The formula looks like $M \propto \frac{1}{i_1}$, but it's not actually a variable dependence — it's just a way to **define** M from one instance of i_1 and the resulting flux.

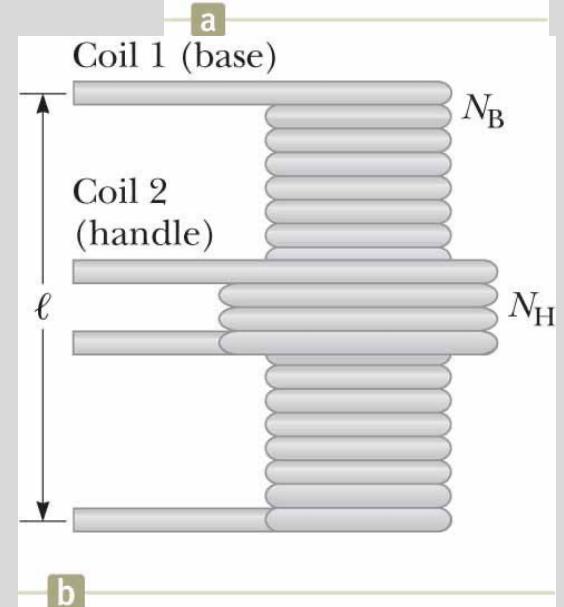
Example: “Wireless” Battery Charger

- An electric toothbrush has a base designed to hold the toothbrush handle when not in use. As shown in the photo, the handle has a cylindrical hole that fits loosely over a matching cylinder on the base. When the handle is placed on the base, a changing current in a solenoid inside the base cylinder induces a current in a coil inside the handle. This induced current charges the battery in the handle.

$$B = \mu_0 \frac{N_B}{\ell} i$$

$$M = \frac{N_H \Phi_{BH}}{i} = \frac{N_H B A}{i} = \boxed{\mu_0 \frac{N_B N_H}{\ell} A}$$

(
i : current in the base
A : cross-sectional area of the base coil)



@ Mutual Inductance of a "Wireless" Battery Charger

$$M = \frac{N_H \Phi_{BH}}{i}$$

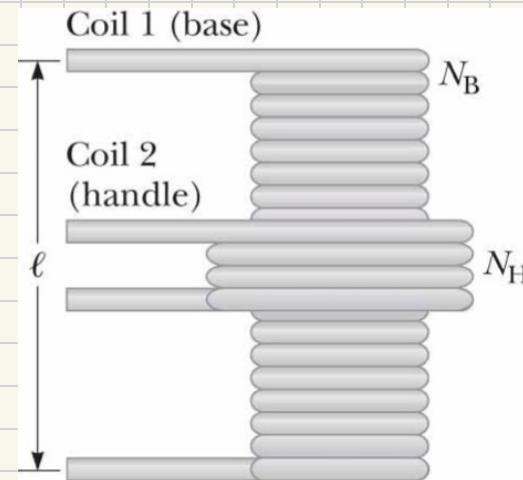
↓

$$= \frac{N_H B A}{i}$$

←

$$\Phi_{BH} = BA \cos 0$$

$$= BA$$

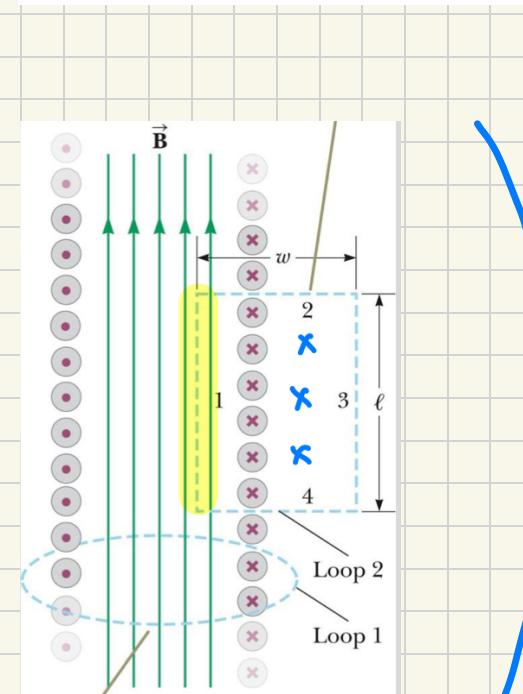


Obtain B using Ampere's Law

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 I$$

$$\Rightarrow B l = \mu_0 N i$$

$$\Rightarrow B = \mu_0 \frac{N}{l} i = \mu_0 n i$$



$$= \frac{N_H (\mu_0 n_i) A}{i}$$

$$= \mu_0 N_H n A$$

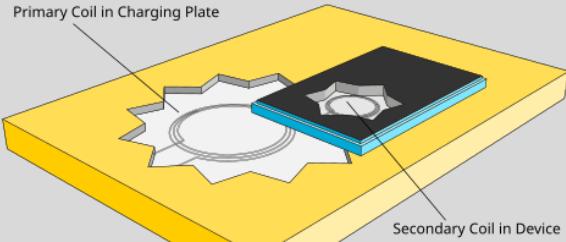
[

$$= \mu_0 \frac{N_B N_H}{l} A$$

← #

$$n = \frac{N_B}{l}$$

Inductive charging



https://en.wikipedia.org/wiki/Inductive_charging



https://en.wikipedia.org/wiki/Inductive_charging



The primary coil in the charger induces a current in the secondary coil of the device being charged.



https://en.wikipedia.org/wiki/Inductive_charging

LC Circuits

- Assume the capacitor is initially charged and then the switch is closed

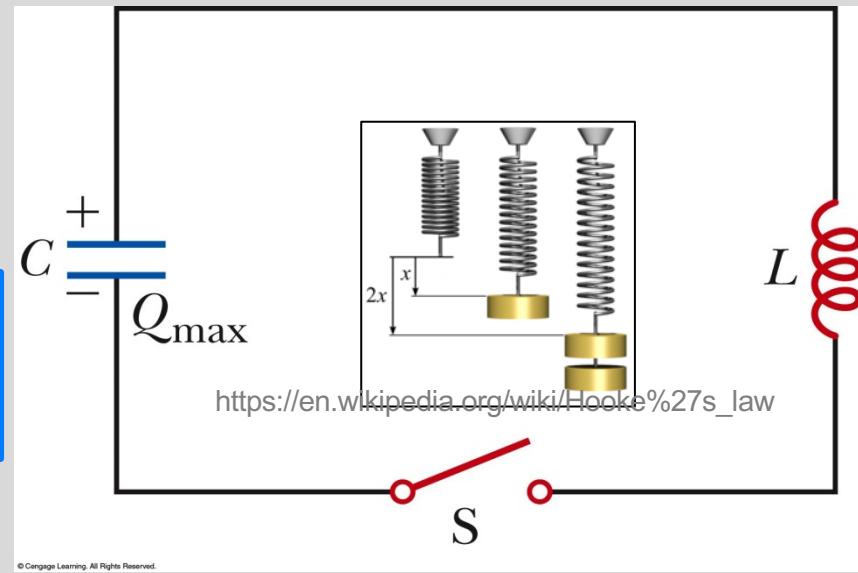
$$U = U_E + U_B = \frac{q^2}{2C} + \frac{1}{2} Li^2 = \frac{Q_{\max}^2}{2C}$$

Energy conservation

KVL: $v_C + v_L = 0 \rightarrow$

$$\frac{q}{C} = -L \frac{di}{dt} = -L \frac{d^2q}{dt^2}$$

$$\frac{d^2q}{dt^2} = -\frac{1}{LC} q = -w^2 q$$



Hooke's Law , k: spring constant

A spring system move as SHM

$$F = -kx = ma$$

$$\frac{d^2x}{dt^2} = -\frac{k}{m} x = -w^2 x$$

Time Functions of an *LC* Circuit

- The function of charge
 $q = Q_{\max} \cos(\omega t + \phi)$ $\frac{d^2 q}{dt^2} = -\frac{1}{LC} q = -\omega^2 q$
(When = 0, q=Q_{max}) (32.21)

- The function of current

$$i = \frac{dq}{dt} = -\omega Q_{\max} \sin(\omega t + \phi) = -I_{\max} \sin(\omega t + \phi) \quad (32.23)$$

- The function of total energy

$$U = U_E + U_B = \frac{Q_{\max}^2}{2C} \cos^2 \omega t + \frac{1}{2} L I_{\max}^2 \sin^2 \omega t \rightarrow$$

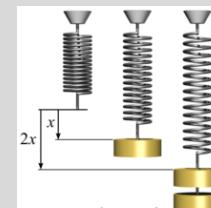
$$U = U_E + U_B = \frac{q^2}{2C} + \frac{1}{2} L i^2 = \frac{Q_{\max}^2}{2C}$$

Omitting ϕ assumes the system starts at maximum charge, $\phi = 0$. (32.26)

- The angular frequency

$$\omega = \frac{1}{\sqrt{LC}} \quad \begin{aligned} &\text{A spring system move as SHM} \\ &F = -kx = ma \\ &\frac{d^2 x}{dt^2} = -\frac{k}{m} x = -\omega^2 x \end{aligned} \quad (32.22)$$

- It is the natural frequency of oscillation of the circuit.



The function of charge on the capacitor C

$$q = Q_{\max} \cos(\omega t + \phi)$$

$$\Rightarrow \frac{dq}{dt} = i = -\omega Q_{\max} \sin(\omega t + \phi)$$

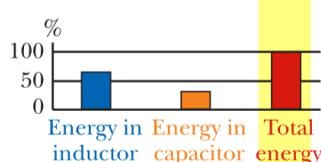
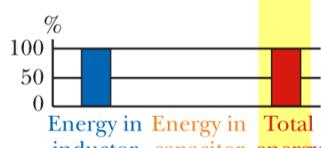
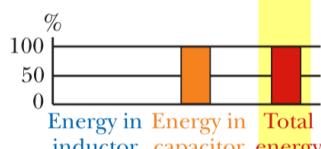
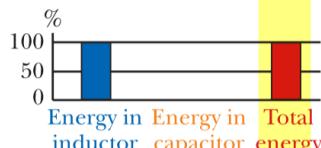
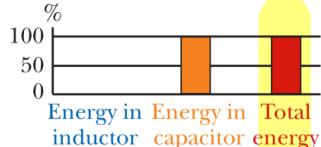
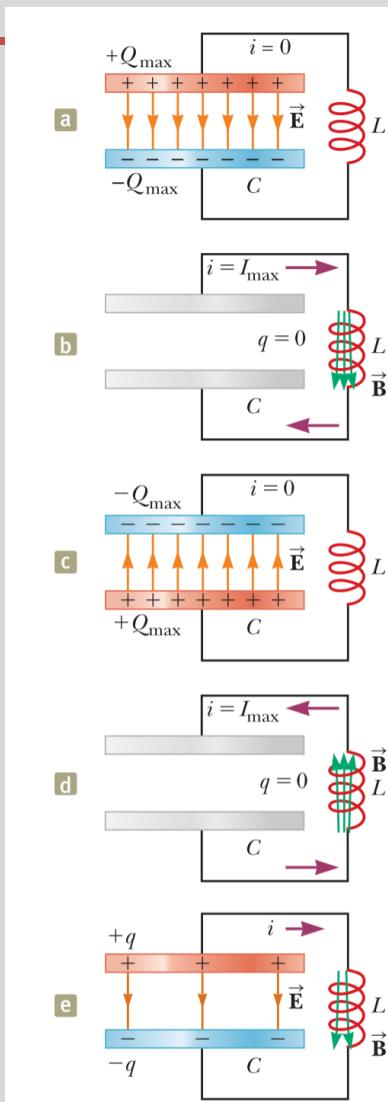
$$\begin{aligned}\Rightarrow \frac{d^2q}{dt^2} &= -\omega^2 Q_{\max} \cos(\omega t + \phi) \\ &= -\omega^2 q \quad \boxed{\quad} \\ &= -\frac{q}{LC} \quad \boxed{\quad}\end{aligned}$$

$$\omega = \frac{1}{\sqrt{LC}}$$

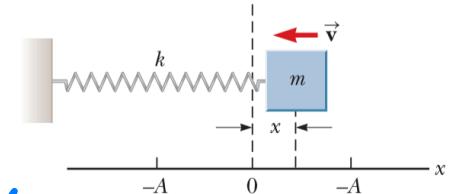
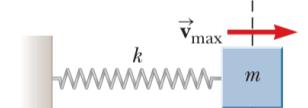
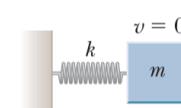
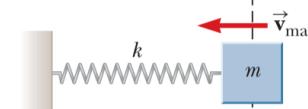
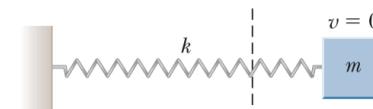
LC Circuits

spring

LC-Grauit



energy
conservation

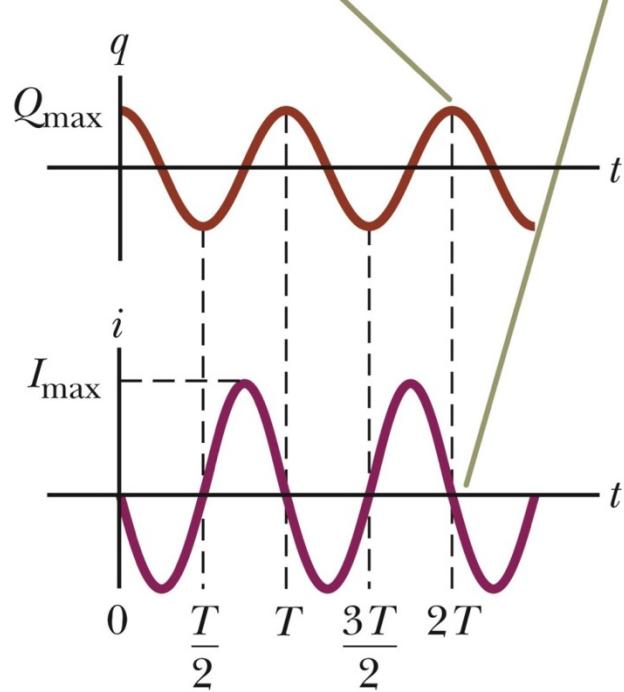


$$\frac{1}{2} kx^2 \Leftrightarrow \frac{Q^2}{2C}$$

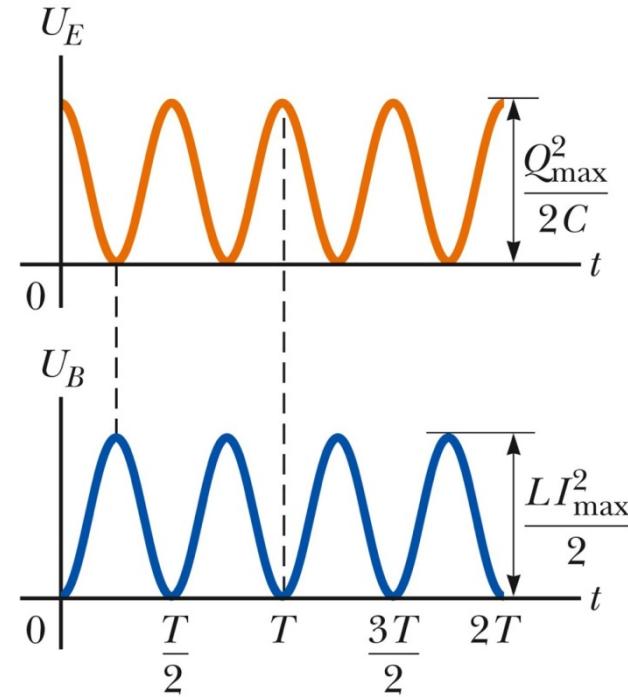
$$\frac{1}{2} mv^2 \Leftrightarrow \frac{1}{2} Li^2$$

Charge, Current, Energy

The charge q and the current i are 90° out of phase with each other.

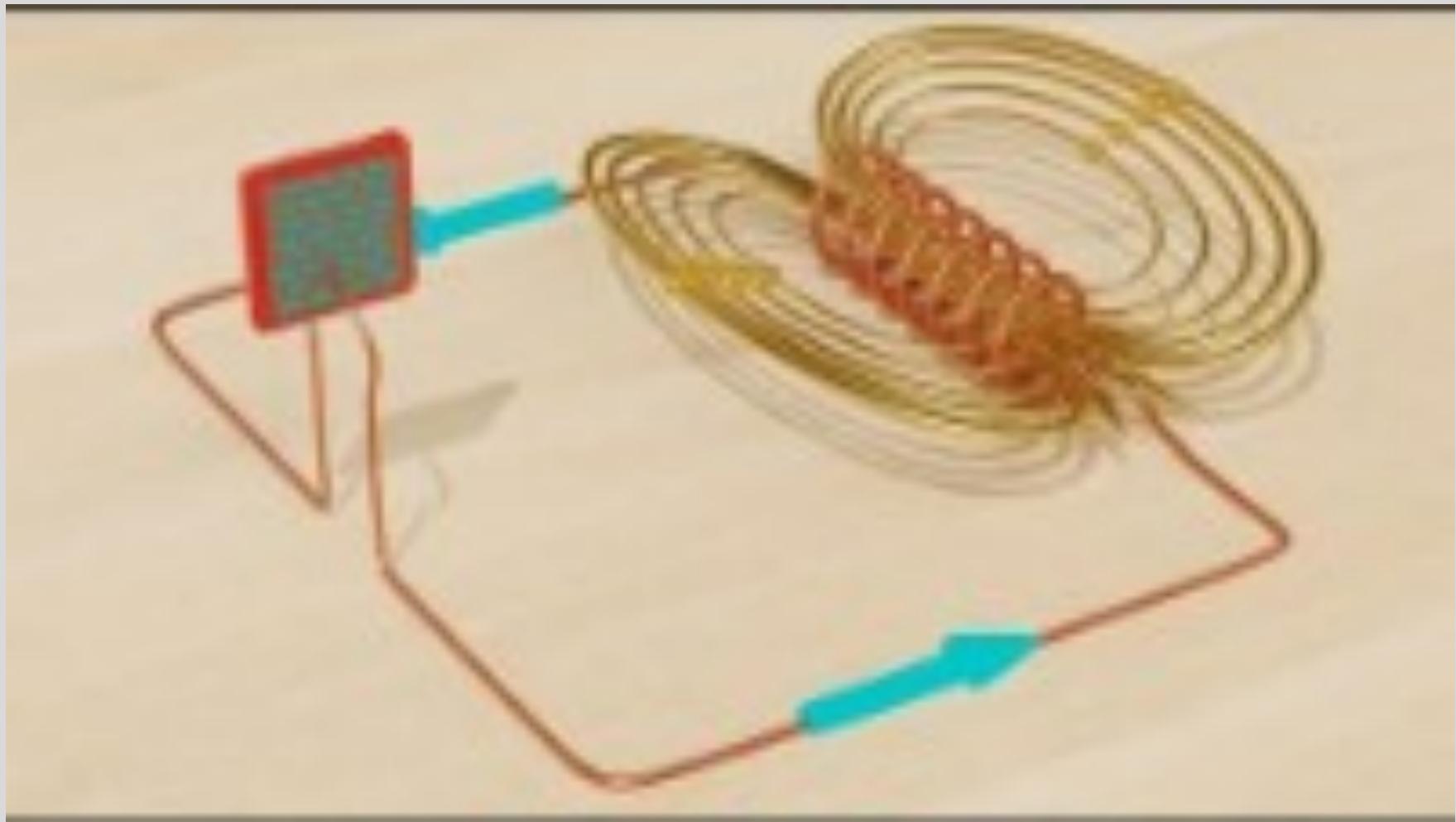


The sum of the two curves is a constant and is equal to the total energy stored in the circuit.



LC Circuits

網路資訊，同學要記得獨立思考，自主查證



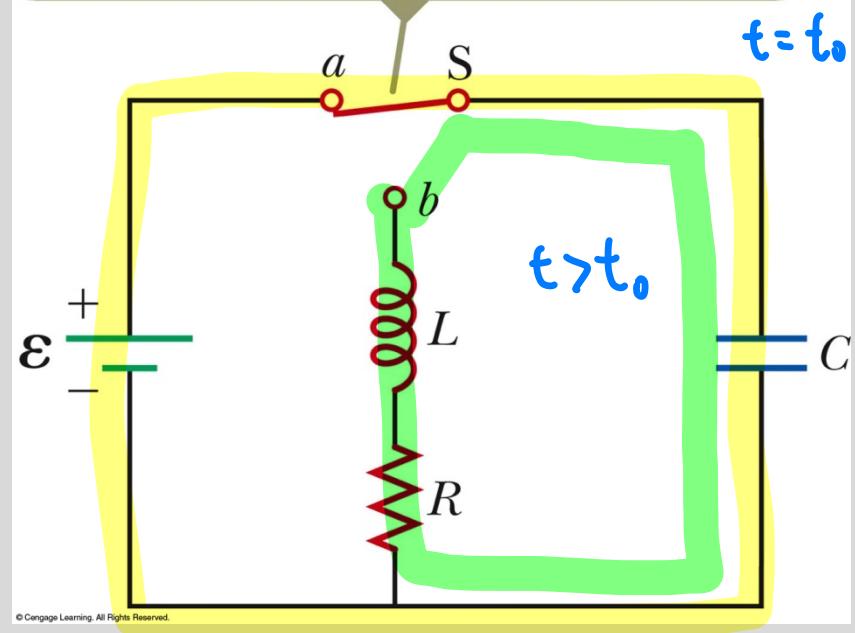
https://www.youtube.com/watch?v=2_y_3_3V-so

2nd - order linear
differential equation

The RLC Circuit

- The total energy is not constant
 - Transformation to internal energy in the resistor at the rate of $dU/dt = -I^2R$
- Kirchhoff's loop laws
$$L \frac{d^2q}{dt^2} + R \frac{dq}{dt} + \frac{q}{C} = 0$$
- The RLC circuit is analogous to a damped harmonic oscillator
 - Small R
 - Large R
 - $R=R_c$

The switch is set first to position *a*, and the capacitor is charged. The switch is then thrown to position *b*.



Damped RLC Circuit

causes the amplitude to

shrink over time

- $q = Q_{\max} e^{-Rt/2L} \cos \omega_d t$

$$\omega_d = \left[\frac{1}{LC} - \left(\frac{R}{2L} \right)^2 \right]^{1/2}$$

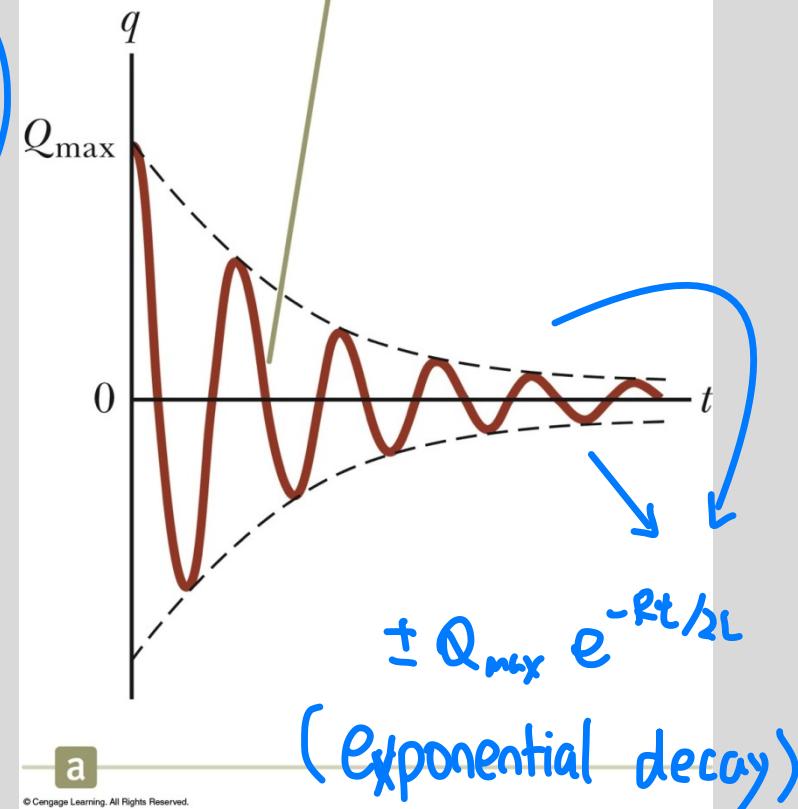
(damped angular velocity)

$$R_C = \sqrt{4L/C}$$

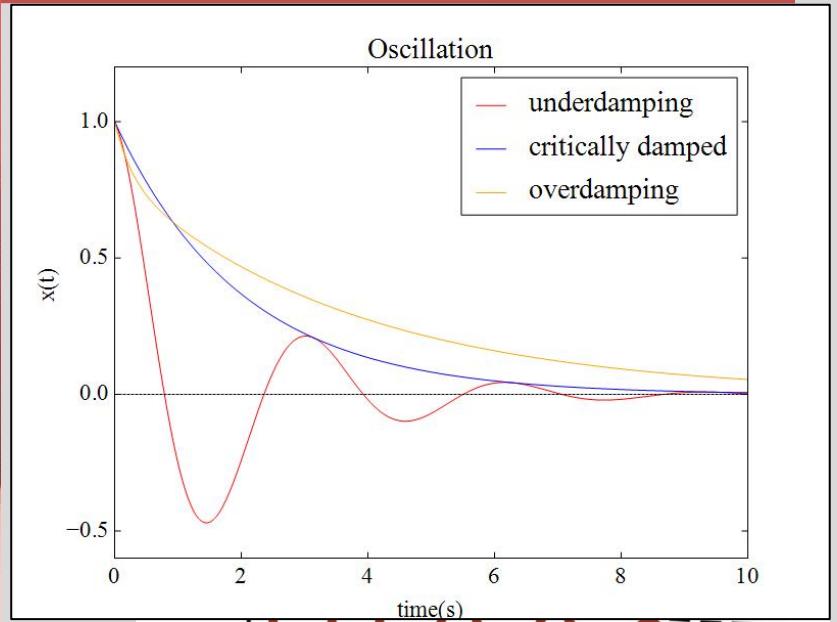
(critical resistance)

- $R < R_C$: damped oscillation
- $R = R_C$: critically damped
- $R > R_C$: overdamped

The q -versus- t curve represents a plot of Equation 32.31.

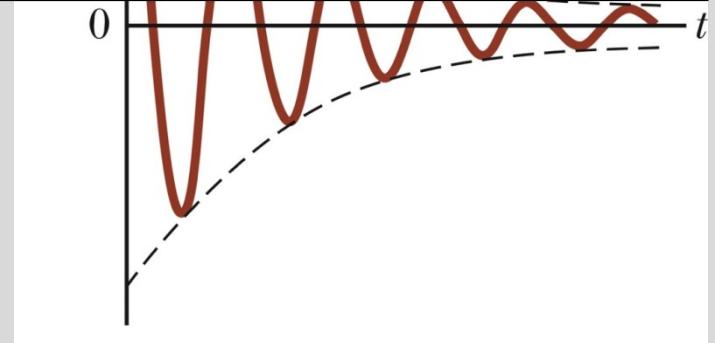


Damped *RLC* Circuit



- $R < R_C$: damped oscillation
- $R = R_C$: critically damped
- $R > R_C$: overdamped

油壓關門裝置，避免撞擊噪音，
但也希望它能儘快回到平衡位置



Damped RLC Circuit

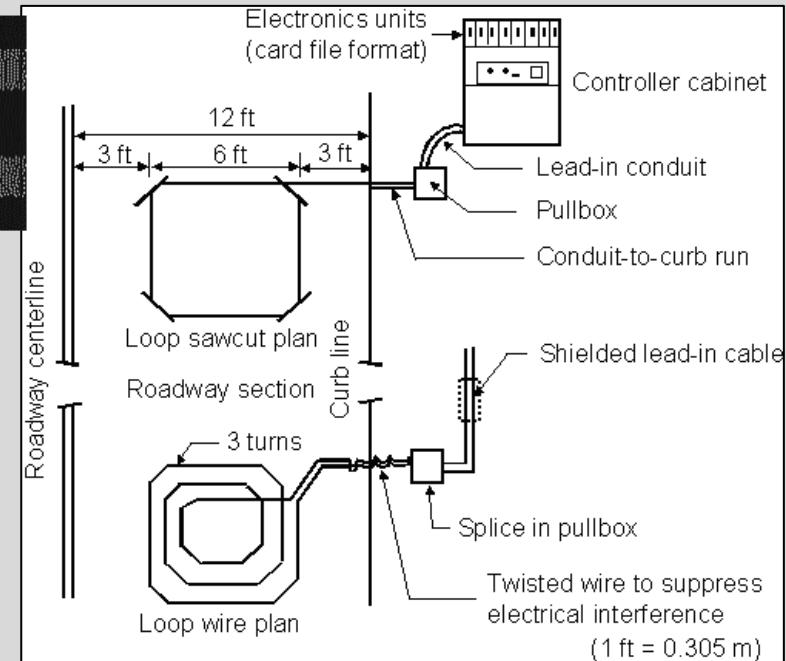
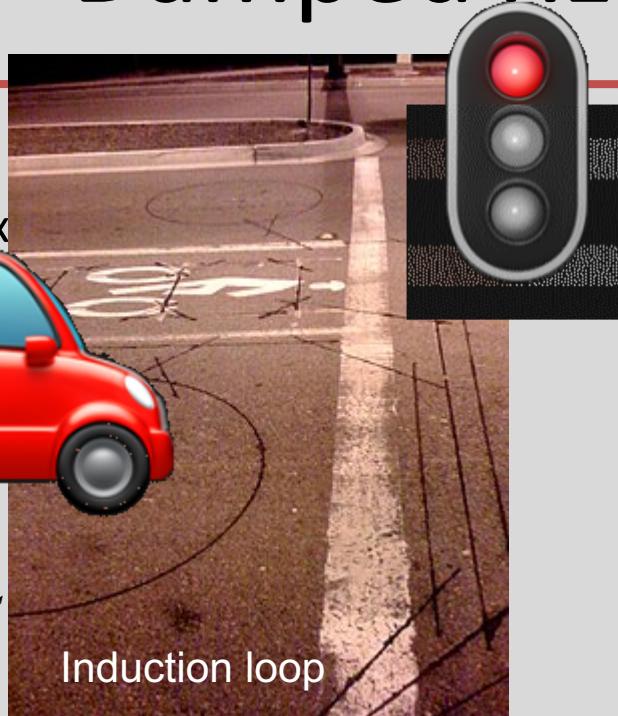
- $q = Q_{\max}$



R_C

Induction loop

- $R < R_C$: damped oscillation
- $R = R_C$: critically damped
- $R > R_C$: overdamped



Damped RLC Circuit

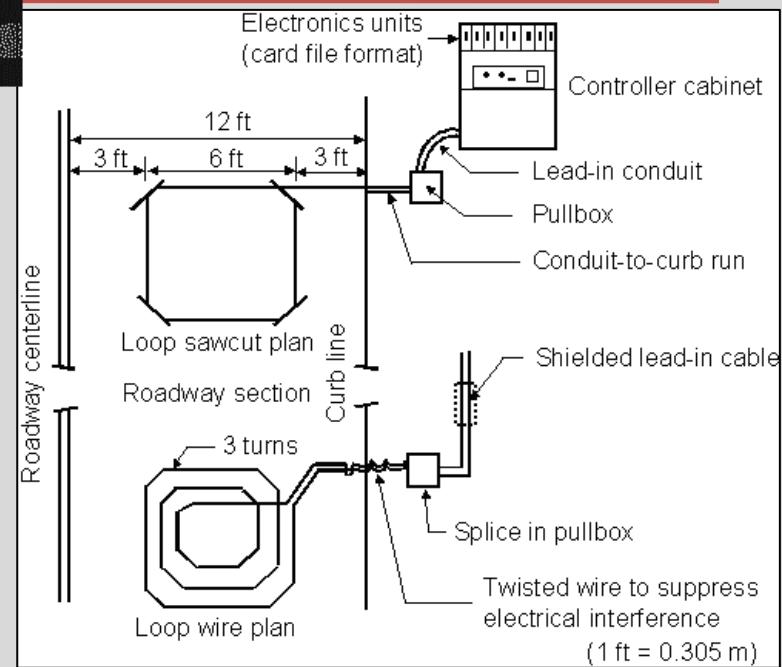
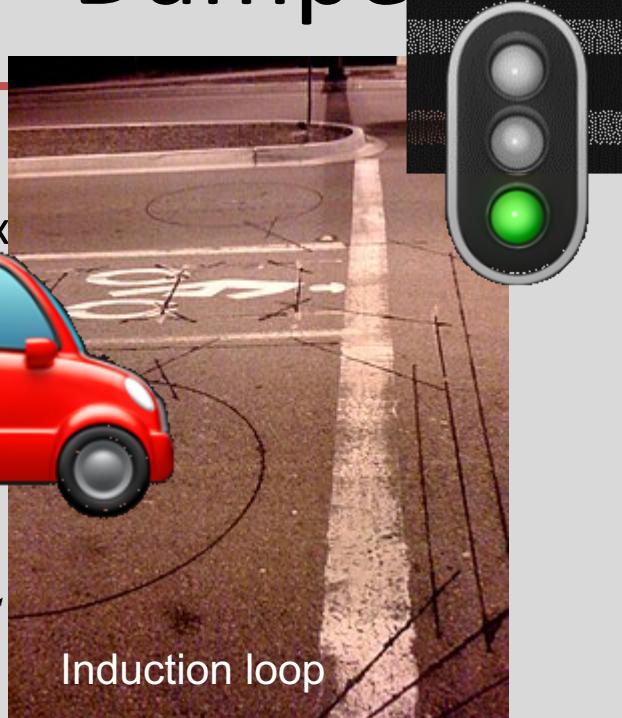
- $q = Q_{\max}$



R_C

Induction loop

- $R < R_C$: damped oscillation
- $R = R_C$: critically damped
- $R > R_C$: overdamped



Summary: Analogies Between Electrical and Mechanic Systems

Table 32.1 Analogies Between the *RLC* Circuit and the Particle in Simple Harmonic Motion

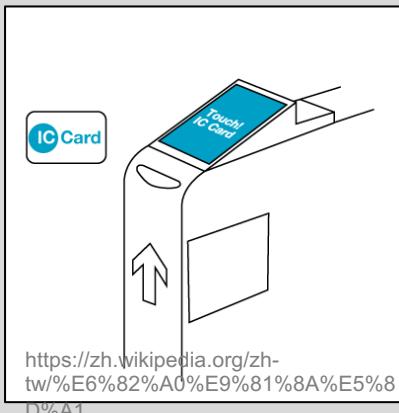
<i>RLC</i> Circuit		One-Dimensional Particle in Simple Harmonic Motion
Charge	$q \leftrightarrow x$	Position
Current	$i \leftrightarrow v_x$	Velocity
Potential difference	$\Delta V \leftrightarrow F_x$	Force
Resistance	$R \leftrightarrow b$	Viscous damping coefficient
Capacitance	$C \leftrightarrow 1/k$	(k = spring constant)
Inductance	$L \leftrightarrow m$	Mass
Current = time derivative of charge	$i = \frac{dq}{dt} \leftrightarrow v_x = \frac{dx}{dt}$	Velocity = time derivative of position
Rate of change of current = second time derivative of charge	$\frac{di}{dt} = \frac{d^2q}{dt^2} \leftrightarrow a_x = \frac{dv_x}{dt} = \frac{d^2x}{dt^2}$	Acceleration = second time derivative of position
Energy in inductor	$U_B = \frac{1}{2}Li^2 \leftrightarrow K = \frac{1}{2}mv^2$	Kinetic energy of moving object
Energy in capacitor	$U_E = \frac{1}{2} \frac{q^2}{C} \leftrightarrow U = \frac{1}{2}kx^2$	Potential energy stored in a spring
Rate of energy loss due to resistance	$i^2R \leftrightarrow bv^2$	Rate of energy loss due to friction
<i>RLC</i> circuit	$L \frac{d^2q}{dt^2} + R \frac{dq}{dt} + \frac{q}{C} = 0 \leftrightarrow m \frac{d^2x}{dt^2} + b \frac{dx}{dt} + kx = 0$	Damped object on a spring

法拉第電磁感應定律



無線射頻辨識系統 RFID:
悠遊卡、信用卡感應支付、停車票卡、國道ETC

每張悠遊卡皆內建獨立的電子標籤，當卡片進入磁場感應範圍內，透過電磁感應，電子標籤內的線圈產生感應電流，以供應電子標籤運作，並將資訊傳送至讀卡機，以解讀晶片資料。



讀卡機 reader、電子標籤 tag

近場通訊 NFC: 手機支付



<https://pansci.asia/archives/359263>

Summary

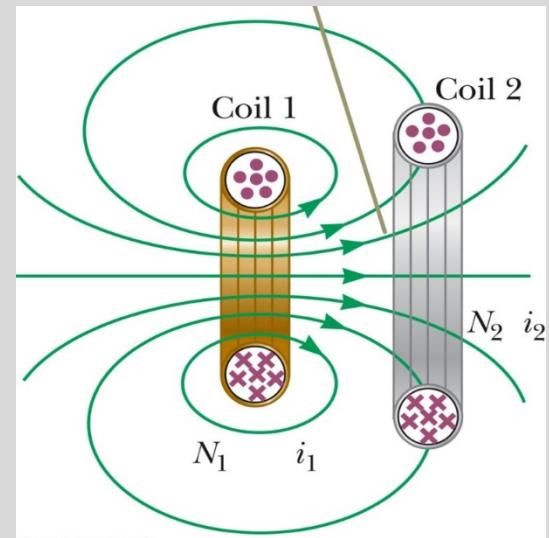
- Inductance (L)
$$\boldsymbol{\epsilon}_L = -L \frac{di}{dt}$$
- Mutual inductance

$$M = \frac{N_2 \Phi_{12}}{i_1} = \frac{N_1 \Phi_{21}}{i_2}$$

$$\boldsymbol{\epsilon}_1 = -M \frac{di_2}{dt} \quad \text{and} \quad \boldsymbol{\epsilon}_2 = -M \frac{di_1}{dt}$$

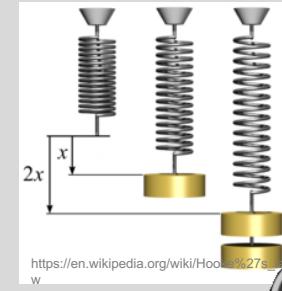
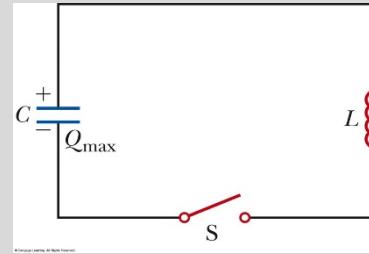
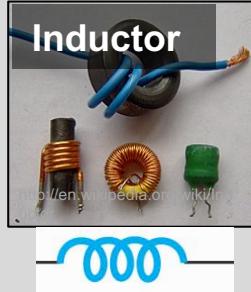
- Inductor 

$$v_L(t) = L \frac{di(t)}{dt}$$



This Lecture

- Inductance
 - Self-Inductance
- *RL* Circuit
 - Energy in a Magnetic Field
- Energy Storage Summary
 - Charged capacitor, Inductor, Resistor
- Mutual Inductance
- *LC* Circuits
- *RLC* Circuit



Announcement

- Homework [W11](#) has been assigned on Moodle.
 - One-week submission: 100% credit.
 - Afterwards: 70% credit.