

《量子信息基础》2022.3.1 随堂作业:

1. <即教材*问题 1.5>

Consider the wave function

$$\Psi(x, t) = Ae^{-\lambda|x|}e^{-i\omega t}$$

where A , λ , and ω are positive real constants.

(a) Normalize Ψ .

$$\begin{aligned}\int |\Psi(x, t)|^2 dx &= \int |Ae^{-\lambda|x|}e^{-i\omega t}|^2 dx = |A|^2 \int_{-\infty}^{\infty} e^{-2\lambda|x|} dx \\ &= 2|A|^2 \int_0^{\infty} e^{-2\lambda x} dx = \frac{|A|^2}{\lambda} = 1 \\ \therefore A &= \sqrt{\lambda}\end{aligned}$$

(b) Determine the expectation values of x and x^2 .

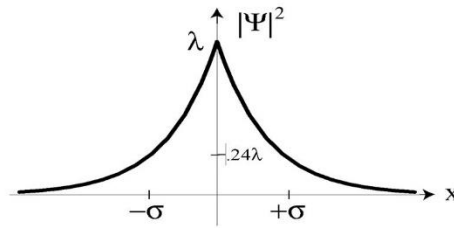
$$\begin{aligned}\langle x \rangle &= \int_{-\infty}^{+\infty} x |\Psi(x, t)|^2 dx = \int_{-\infty}^{\infty} x |Ae^{-\lambda|x|}|^2 dx = |A|^2 \int_{-\infty}^{\infty} x e^{-2\lambda|x|} dx \\ &= \lambda \int_{-\infty}^{\infty} x e^{-2\lambda|x|} dx \\ \because x e^{-2\lambda|x|} &\text{是奇函数} \quad \therefore \langle x \rangle = 0\end{aligned}$$

$$\begin{aligned}\langle x^2 \rangle &= \int_{-\infty}^{+\infty} x^2 |\Psi(x, t)|^2 dx = |A|^2 \int_{-\infty}^{\infty} x^2 e^{-2\lambda|x|} dx = 2\lambda \int_0^{\infty} x^2 e^{-2\lambda x} dx \\ &= - \int_0^{\infty} x^2 d e^{-2\lambda x} = -x^2 e^{-2\lambda x} \Big|_0^{\infty} + \int_0^{\infty} e^{-2\lambda x} d x^2 \\ &= \int_0^{\infty} 2x e^{-2\lambda x} dx = -\frac{1}{\lambda} \int_0^{\infty} x d e^{-2\lambda x} \\ &= -\frac{1}{\lambda} x e^{-2\lambda x} \Big|_0^{\infty} + \frac{1}{\lambda} \int_0^{\infty} e^{-2\lambda x} dx = -\frac{1}{2\lambda^2} \int_0^{\infty} d e^{-2\lambda x} = \frac{1}{2\lambda^2}\end{aligned}$$

(c) Find the standard deviation of x . Sketch the graph of $|\Psi|^2$, as a function of x , and mark the points $(\langle x \rangle - \sigma)$ and $(\langle x \rangle + \sigma)$, to illustrate the sense in which σ represents the “spread” in x .

$$\sigma = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = \frac{1}{\sqrt{2}\lambda}$$

$|\Psi|^2$ 的图像如下图所示:



- (d) What is the probability that the particle would be found outside this range (which means outside of $(\langle x \rangle - \sigma)$ and $(\langle x \rangle + \sigma)$)?

$$P = 2 \int_{\sigma}^{\infty} |\Psi(x, t)|^2 dx = 2\lambda \int_{\sigma}^{\infty} e^{-2\lambda x} dx = - \int_{\sigma}^{\infty} d e^{-2\lambda x} = e^{-2\lambda \sigma} = e^{-\sqrt{2}} = 0.2431$$

2. <PPT 最后一页>

A photon propagates in the z direction and passes a linear optical polarizer which is oriented in the x direction (see figure below). The state in Figure (a) is ψ_a while the state in Figure (b) is ψ_b .

- (a) Write down the formula of ψ_c , assuming the light beam is polarized with an angle of α to the x axis.

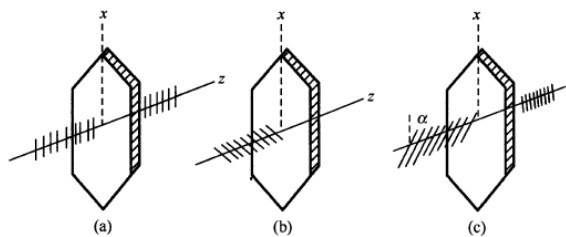
$$\psi_{\alpha} = \cos \alpha \cdot \psi_x + \sin \alpha \cdot \psi_y$$

- (b) How much is the probability that a single photon could pass the polarizer in (c)?

$$P = |\cos \alpha|^2$$

- (c) How does the system maintain the normalization condition?

$$|\cos \alpha|^2 + |\sin \alpha|^2 = 1$$



3. <即教材*问题 1.15>

Show that

$$\frac{d}{dt} \int_{-\infty}^{\infty} \psi_1^* \psi_2 dx = 0$$

for any two (normalizable) solutions to the Schrödinger equation (with the same $V(x)$), ψ_1 and ψ_2 .

$$\frac{d}{dt} \int_{-\infty}^{\infty} \psi_1^* \psi_2 dx = \int_{-\infty}^{\infty} \frac{\partial \psi_1^*}{\partial t} \psi_2 + \psi_1^* \frac{\partial \psi_2}{\partial t} dx$$

复共轭函数的薛定谔方程

$$-i\hbar \frac{\partial \Psi_1^*}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \Psi_1^*}{\partial x^2} + V\Psi_1^*$$

$$i\hbar \frac{\partial \Psi_2}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \Psi_2}{\partial x^2} + V\Psi_2$$

$$\frac{d}{dt} \int_{-\infty}^{\infty} \Psi_1^* \Psi_2 dx = \int_{-\infty}^{\infty} \left[\frac{\hbar}{2im} \frac{\partial^2 \Psi_1^*}{\partial x^2} - \frac{V}{i\hbar} \Psi_1^* \right] \Psi_2 + \Psi_1^* \left[-\frac{\hbar}{2im} \frac{\partial^2 \Psi_2}{\partial x^2} + \frac{V}{i\hbar} \Psi_2 \right] dx$$

$$\frac{d}{dt} \int_{-\infty}^{\infty} \Psi_1^* \Psi_2 dx = \int_{-\infty}^{\infty} \frac{\hbar}{2im} \frac{\partial^2 \Psi_1^*}{\partial x^2} \Psi_2 - \Psi_1^* \frac{\hbar}{2im} \frac{\partial^2 \Psi_2}{\partial x^2} dx$$

$$\frac{d}{dt} \int_{-\infty}^{\infty} \Psi_1^* \Psi_2 dx = \frac{\hbar}{2im} \int_{-\infty}^{\infty} \frac{\partial}{\partial x} \left(\frac{\partial \Psi_1^*}{\partial x} \Psi_2 - \Psi_1^* \frac{\partial \Psi_2}{\partial x} \right) dx$$

$$\frac{d}{dt} \int_{-\infty}^{\infty} \Psi_1^* \Psi_2 dx = \frac{\hbar}{2im} \left(\frac{\partial \Psi_1^*}{\partial x} \Psi_2 - \Psi_1^* \frac{\partial \Psi_2}{\partial x} \right) \Big|_{-\infty}^{\infty}$$

$$\frac{d}{dt} \int_{-\infty}^{\infty} \Psi_1^* \Psi_2 dx = 0$$

* David J. Griffiths, and Darrell F. Schroeter, Introduction to Quantum Mechanics (3rd Edition), Cambridge University Press (2018).