

《量子信息基础》第三章第二部分：

1. Derive the following equation of Bloch function in periodic potentials from the time-independent Schrodinger equation:

$$\frac{d^2u}{dx^2} + 2ik \frac{du}{dx} + \left(\frac{2m}{\hbar^2} (E - V) - k^2 \right) u = 0$$

将布洛赫定理

$$\psi(x) = e^{ikx}u(x)$$

代入周期性方势阱的薛定谔方程

$$\frac{d^2\psi}{dx^2} + \frac{2m}{\hbar^2} (E - V(x))\psi = 0$$

的左面，得到

$$\begin{aligned} \frac{d^2\psi}{dx^2} + \frac{2m}{\hbar^2} (E - V)\psi &= \frac{d^2e^{ikx}u}{dx^2} + \frac{2m}{\hbar^2} (E - V)e^{ikx}u \\ &= \frac{d}{dx} \left(ike^{ikx}u + e^{ikx} \frac{du}{dx} \right) + \frac{2m}{\hbar^2} (E - V)e^{ikx}u \\ &= -k^2e^{ikx}u + 2ike^{ikx} \frac{du}{dx} + e^{ikx} \frac{d^2u}{dx^2} + \frac{2m}{\hbar^2} (E - V)e^{ikx}u \\ &= e^{ikx} \left[\frac{d^2u}{dx^2} + 2ik \frac{du}{dx} + \left(\frac{2m}{\hbar^2} (E - V) - k^2 \right) u \right] = 0 \end{aligned}$$

所以

$$\frac{d^2u}{dx^2} + 2ik \frac{du}{dx} + \left(\frac{2m}{\hbar^2} (E - V) - k^2 \right) u = 0$$

2. (Text book* Problem 7.8)

Let the two “good” unperturbed states be

$$\psi_{\pm}^0 = \alpha_{\pm}\psi_a^0 + \beta_{\pm}\psi_b^0$$

where α_{\pm} and β_{\pm} are determined (up to normalization) by Equation 7.27 (or Equation 7.29). Show explicitly that

(a) ψ_{\pm}^0 are orthogonal ($\langle \psi_+^0 | \psi_-^0 \rangle = 0$);

$$\begin{aligned} \langle \psi_+^0 | \psi_-^0 \rangle &= \langle \alpha_+\psi_a^0 + \beta_+\psi_b^0 | \alpha_-\psi_a^0 + \beta_-\psi_b^0 \rangle \\ &= \alpha_+^*\alpha_- \langle \psi_a^0 | \psi_a^0 \rangle + \alpha_+^*\beta_- \langle \psi_a^0 | \psi_b^0 \rangle + \beta_+^*\alpha_- \langle \psi_b^0 | \psi_a^0 \rangle + \beta_+^*\beta_- \langle \psi_b^0 | \psi_b^0 \rangle \\ &= \alpha_+^*\alpha_- + \beta_+^*\beta_- \end{aligned}$$

由 7.27 可得

$$\beta_{\pm} = \frac{\alpha_{\pm}(E_{\pm}^1 - W_{aa})}{W_{ab}}$$

$$\langle \psi_+^0 | \psi_-^0 \rangle = \alpha_+^*\alpha_- + \frac{\alpha_+^*\alpha_-}{|W_{ab}|^2} [E_+^1 E_-^1 - W_{ab}(E_+^1 + E_-^1) + |W_{ab}|^2]$$

由于

$$E_{\pm}^1 = \frac{1}{2} \left[W_{aa} + W_{bb} \pm \sqrt{(W_{aa} - W_{bb})^2 + 4|W_{ab}|^2} \right]$$

代入可得

$$\langle \psi_+^0 | \psi_-^0 \rangle = 0$$

$$(b) \langle \psi_+^0 | H' | \psi_-^0 \rangle = 0;$$

$$\begin{aligned} \langle \psi_+^0 | H' | \psi_-^0 \rangle &= [\alpha_+^* \quad \beta_+^*] \begin{bmatrix} W_{aa} & W_{ab} \\ W_{ba} & W_{bb} \end{bmatrix} \begin{bmatrix} \alpha_- \\ \beta_- \end{bmatrix} \\ &= \alpha_+^* \alpha_- W_{aa} + \alpha_+^* \beta_- W_{ab} + \beta_+^* \alpha_- W_{ba} + \beta_+^* \beta_- W_{bb} \\ &= \alpha_+^* \alpha_- W_{aa} + \alpha_+^* \alpha_- (E_-^1 - W_{aa}) + \alpha_+^* \alpha_- (E_+^1 - W_{aa}) \\ &\quad - \alpha_+^* \alpha_- W_{bb} = \alpha_+^* \alpha_- (E_+^1 + E_-^1 - W_{aa} - W_{bb}) = 0 \end{aligned}$$

$$(c) \langle \psi_{\pm}^0 | H' | \psi_{\pm}^0 \rangle = E_{\pm}^1, \text{ where } E_{\pm}^1 \text{ given by Equation 7.33.}$$

$$\begin{aligned} \langle \psi_{\pm}^0 | H' | \psi_{\pm}^0 \rangle &= [\alpha_{\pm}^* \quad \beta_{\pm}^*] \begin{bmatrix} W_{aa} & W_{ab} \\ W_{ba} & W_{bb} \end{bmatrix} \begin{bmatrix} \alpha_{\pm} \\ \beta_{\pm} \end{bmatrix} \\ &= |\alpha_{\pm}|^2 W_{aa} + \alpha_{\pm}^* \beta_{\pm} W_{ab} + \beta_{\pm}^* \alpha_{\pm} W_{ba} + |\beta_{\pm}|^2 W_{bb} \end{aligned}$$

$$\beta_{\pm} W_{ab} = \alpha_{\pm} (E_{\pm}^1 - W_{aa})$$

$$\alpha_{\pm} W_{ba} = \beta_{\pm} (E_{\pm}^1 - W_{bb})$$

$$\langle \psi_{\pm}^0 | H' | \psi_{\pm}^0 \rangle = (|\alpha_{\pm}|^2 + |\beta_{\pm}|^2) E_{\pm}^1 = E_{\pm}^1$$

3. (Text book* Problem 7.1)

Suppose we put a delta-function bump in the center of the infinite square well:

$$H' = \alpha \delta(x - a/2)$$

where α is a constant.

(a) Find the first-order correction to the allowed energies. Explain why the energies are not perturbed for even n .

$$\psi_n^0 = \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi}{a}x\right)$$

$$E_n^1 = \langle \psi_n^0 | H' | \psi_n^0 \rangle = \frac{2\alpha}{a} \int_0^a \sin^2\left(\frac{n\pi}{a}x\right) \delta\left(x - \frac{a}{2}\right) dx = \frac{2\alpha}{a} \sin^2\left(\frac{n\pi}{2}\right)$$

For even n

$$E_n^1 = 0$$

- (b) Find the first three nonzero terms in the expansion (Equation 7.13) of the correction to the ground state, ψ_1^1 .

$$\psi_1^1 = \sum_{m=2}^{\infty} \frac{\langle \psi_m^0 | H' | \psi_1^0 \rangle}{E_1^0 - E_m^0} \psi_m^0$$

$$\langle \psi_m^0 | H' | \psi_1^0 \rangle = \frac{2\alpha}{a} \int_0^a \sin\left(\frac{m\pi}{a}x\right) \sin\left(\frac{\pi}{a}x\right) \delta\left(x - \frac{a}{2}\right) dx = \frac{2\alpha}{a} \sin\left(\frac{m\pi}{2}\right)$$

$$E_1^0 - E_m^0 = \frac{\hbar^2 \pi^2}{2Ma^2} (1 - m^2)$$

The first three non-zero terms are $m = 3, 5, 7$.

$$\begin{aligned} \psi_1^1 &= \frac{2\alpha}{a} \frac{2Ma^2}{\hbar^2 \pi^2} \left[\frac{1}{8} \sqrt{\frac{2}{a}} \sin\left(\frac{3\pi}{a}x\right) - \frac{1}{24} \sqrt{\frac{2}{a}} \sin\left(\frac{5\pi}{a}x\right) + \frac{1}{48} \sqrt{\frac{2}{a}} \sin\left(\frac{7\pi}{a}x\right) \cdots \right] \\ &= \frac{M\alpha}{\hbar^2 \pi^2} \sqrt{\frac{a}{2}} \left[\sin\left(\frac{3\pi}{a}x\right) - \frac{1}{3} \sin\left(\frac{5\pi}{a}x\right) + \frac{1}{6} \sin\left(\frac{7\pi}{a}x\right) \cdots \right] \end{aligned}$$

* David J. Griffiths, and Darrell F. Schroeter, Introduction to Quantum Mechanics (3rd Edition), Cambridge University Press (2018).