QComp Homework 1

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1 Problem 3

For an arbitrary 1-qbit gate W, 2-qbit matrix for cW_{01} is obtained using

$$cW_{01} = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \otimes I + \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \otimes W.$$

Also, 2-qbit matrix for W_0 (which only applies W to the first qbit) is obtained using

$$W_0 = W \otimes I$$
.

1.1 1

$$HXH = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 2 & 0 \\ 0 & -2 \end{pmatrix} = Z$$

1.2 2

$$cZ_{01} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} = cZ_{10}$$

1.3 3

1.4 4

$$c(e^{i\alpha})_{01} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & e^{i\alpha} & 0 \\ 0 & 0 & 0 & e^{i\alpha} \end{pmatrix} = U_1(\alpha)_{(0)}$$

2 Problem 4

Separable states should allow representation in the form

$$(a |0\rangle + b |1\rangle) \otimes (c |0\rangle + d |1\rangle) = ac |00\rangle + ad |01\rangle + bc |10\rangle + bd |11\rangle.$$

$2.1 \quad 1$

If the state is separable, then

$$ac = \frac{2}{3}$$
, $ad = \frac{1}{3}$, $bc = 0$, $bd = -\frac{2}{3}$.

It is clear that none of a, b, c, d are 0, so $\frac{c}{d} = \frac{ac}{ad} = 2$. Thus $bc = bd\frac{c}{d} = 2bd = -\frac{4}{3}$. Contradiction, so the state is entangled.

$2.2 \quad 2$

$$\frac{1}{2}(|00\rangle - i|01\rangle + i|10\rangle + |11\rangle) = \frac{1}{\sqrt{2}}(|0\rangle + i|1\rangle) \otimes \frac{1}{\sqrt{2}}(|0\rangle - i|i\rangle)$$

Separable state.

2.3 3

If the state is separable, then (omitting the normalization coefficient)

$$ac = 1$$
, $ad = -1$, $bc = 1$, $bd = 1$.

It is clear that none of a, b, c, d are 0, so $\frac{c}{d} = \frac{ac}{ad} = -1$. Thus $bc = bd\frac{c}{d} = -bd = -1$. Contradiction, so the state is entangled.

3 Problem 5

We can represent $|\psi\rangle$ using Shmidt decomposition:

$$|\psi\rangle = \sum_{i=1}^{2} \lambda_i |\varphi_i^A\rangle |\varphi_i^B\rangle.$$

To get λ_i , φ_i^A , φ_i^B , we can apply SVD to matrix A:

$$A = \begin{pmatrix} a_{00} & a_{01} \\ a_{10} & a_{11} \end{pmatrix} = U\Sigma V^T,$$

where U and V are unitary matrices, and Σ is diagonal with non-negative values:

$$\Sigma = \begin{pmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{pmatrix}, \quad U = \begin{pmatrix} \varphi_{1,0}^A & \varphi_{2,0}^A \\ \varphi_{1,1}^A & \varphi_{2,1}^A \end{pmatrix}, \quad V = \begin{pmatrix} \varphi_{1,0}^B & \varphi_{2,0}^B \\ \varphi_{1,1}^B & \varphi_{2,1}^B \end{pmatrix}, \quad \text{where } |\varphi_i^{\alpha}\rangle = \begin{pmatrix} \varphi_{i,0}^{\alpha} \\ \varphi_{i,1}^{\alpha} \end{pmatrix}.$$

Since U and V are unitary, they can be used as 1-qbit gate; we can also construct a 1-qbit gate B as follows:

$$B = \begin{pmatrix} \lambda_1 & \lambda_2 \\ \lambda_2 & -\lambda_1 \end{pmatrix}$$

It is unitary because $\lambda_i \geq 0$ and $\lambda_1^2 + \lambda_2^2 = 1$.

Now,

$$U\left|0\right\rangle = \left|\varphi_{1}^{A}\right\rangle, \quad U\left|1\right\rangle = \left|\varphi_{2}^{A}\right\rangle, \quad V\left|0\right\rangle = \left|\varphi_{1}^{B}\right\rangle, \quad V\left|1\right\rangle = \left|\varphi_{2}^{B}\right\rangle,$$

thus

$$U_1V_0cX_{01}B_0\left|00\right\rangle = U_1V_0cX_{01}\left(\lambda_1\left|00\right\rangle + \lambda_2\left|01\right\rangle\right) = U_1V_0\left(\lambda_1\left|00\right\rangle + \lambda_2\left|11\right\rangle\right) = \left(\lambda_1\left|\varphi_1^A\right\rangle\left|\varphi_1^B\right\rangle + \lambda_2\left|\varphi_2^A\right\rangle\left|\varphi_2^B\right\rangle\right).$$

So to prepare an arbitrary state, one needs to obtain Shmidt decomposition and then compute $U_1V_0cX_{01}B_0|00\rangle$.

3.1 a

$$\begin{split} |\Phi^{\pm}\rangle &= \frac{1}{\sqrt{2}}(|00\rangle \pm |11\rangle) = \frac{1}{\sqrt{2}}|0\rangle |0\rangle + \frac{1}{\sqrt{2}}|1\rangle |1\rangle \,. \\ B &= \frac{1}{\sqrt{2}}\begin{pmatrix} 1 & 1\\ 1 & -1 \end{pmatrix} = H, \quad U = I, \quad B = I \end{split}$$

3.2 b

$$\begin{split} |\Psi^{\pm}\rangle &= \frac{1}{\sqrt{2}}(|01\rangle \pm |10\rangle) = \frac{1}{\sqrt{2}}|0\rangle |1\rangle + \frac{1}{\sqrt{2}}|1\rangle |0\rangle \,. \\ B &= \frac{1}{\sqrt{2}}\begin{pmatrix} 1 & 1\\ 1 & -1 \end{pmatrix} = H, \quad U = I, \quad B = \begin{pmatrix} 0 & 1\\ 1 & 0 \end{pmatrix} = X \end{split}$$

3.3 c

Let
$$A = \sqrt{\frac{1}{2} + \frac{1}{\sqrt{5}}}$$
, $B = \sqrt{\frac{1}{2} - \frac{1}{\sqrt{5}}}$.

$$|\Psi^{\pm}\rangle = \frac{1}{\sqrt{12}}(3|00\rangle + |01\rangle + |10\rangle - |11\rangle) = \frac{1}{\sqrt{6}}\sqrt{3 + \sqrt{5}}(A|0\rangle + B|1\rangle)(A|0\rangle + B|1\rangle) + \frac{1}{\sqrt{6}}\sqrt{3 - \sqrt{5}}(A|0\rangle - B|1\rangle)(-A|0\rangle + B|1\rangle).$$

$$B = \begin{pmatrix} \frac{1}{\sqrt{6}}\sqrt{3+\sqrt{5}} & \frac{1}{\sqrt{6}}\sqrt{3-\sqrt{5}} \\ \frac{1}{\sqrt{6}}\sqrt{3-\sqrt{5}} & -\frac{1}{\sqrt{6}}\sqrt{3+\sqrt{5}} \end{pmatrix}, \quad U = \begin{pmatrix} A & B \\ B & -A \end{pmatrix}, \quad B = \begin{pmatrix} A & -B \\ B & A \end{pmatrix}$$

The Qiskit code is in .ipynb part.

4 Problem 6

Let l = n - m - k. We can represent the states $|\Phi_x\rangle$ as

$$|\Phi_x\rangle_{n-m} = \sum_{y} b_{x,y} |y\rangle_k |\Omega_{x,y}\rangle_l.$$

4.1 a

After the measurement of m qbits, the state is $|x\rangle_m |\Phi_x\rangle_{n-m} = \sum_y b_{x,y} |x\rangle_m |y\rangle_k |\Omega_{x,y}\rangle_l$ with probability $|a_x|^2$. When another k qbits are measured, any vector y, for which $b_{x,y} \neq 0$, can be result of measurement, and the state collapses to $|x\rangle_m |y\rangle_k |\Omega_{x,y}\rangle$ with probability $|b_{x,y}|^2$. So the possible final states are $|x\rangle_m |y\rangle_k |\Omega_{x,y}\rangle_l$ with probability $|a_x|^2|b_{x,y}|^2$.

4.2 b

In general, after the measurement of m+k qbits, the state could be one of $\sum_{z} c_z |x\rangle_m |y\rangle_k |z\rangle_l$ for some x and y. The probability amplitudes c_z are given by

$$\begin{split} c_z &= \left\langle x \right|_m \left\langle y \right|_k \left\langle z \right|_l \left| \Psi \right\rangle_n = \left\langle x \right|_m \left\langle y \right|_k \left\langle z \right|_l \sum_{x'} a_{x'} \left| x' \right\rangle_m \left| \Phi_{x'} \right\rangle_{n-m} = \\ &\left\langle x \right|_m \left\langle y \right|_k \left\langle z \right|_l \sum_{x'} a_{x'} \sum_{y'} b_{x',y'} \left| x' \right\rangle_m \left| y' \right\rangle_k \left| \Omega_{x',y'} \right\rangle_l = \\ &\sum_{x'} \sum_{y'} a_{x'} b_{x',y'} \left\langle x \right|_m \left\langle y \right|_k \left\langle z \right|_l \left| x' \right\rangle_m \left| y' \right\rangle_k \left| \Omega_{x',y'} \right\rangle_l. \end{split}$$

It is clear that the only possibly non-zero summand is when x'=x, and y=y', so $c_z=a_xb_{x,y}$ $\langle z|_l |\Omega_{x,y}\rangle_l$. Now,

$$\sum_{z} c_{z} |x\rangle_{m} |y\rangle_{k} |z\rangle_{l} = a_{x} b_{x,y} \sum_{z} \langle z|_{l} |\Omega_{x,y}\rangle_{l} |z\rangle_{l} = a_{x} b_{x,y} |x\rangle_{m} |y\rangle_{k} |\Omega_{x,y}\rangle_{l}$$

Thus, the possible final states are $|x\rangle_m |y\rangle_k |\Omega_{x,y}\rangle_l$ with probability $|a_x|^2 |b_{x,y}|^2$, the same result as in (a).