QComp Homework 1

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1 Problem 3

For an arbitrary 1-qbit gate W, 2-qbit matrix for cW_{01} is obtained using

$$cW_{01} = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \otimes I + \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \otimes W.$$

Also, 2-qbit matrix for W_0 (which only applies W to the first qbit) is obtained using

$$W_0 = W \otimes I$$
.

1.1 1

$$HXH = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 2 & 0 \\ 0 & -2 \end{pmatrix} = Z$$

1.2 2

$$cZ_{01} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} = cZ_{10}$$

1.3 3

1.4 4

$$c(e^{i\alpha})_{01} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & e^{i\alpha} & 0 \\ 0 & 0 & 0 & e^{i\alpha} \end{pmatrix} = U_1(\alpha)_{(0)}$$

2 Problem 4

Separable states should allow representation in the form

$$(a|0\rangle + b|1\rangle) \otimes (c|0\rangle + d|1\rangle) = ac|00\rangle + ad|01\rangle + bc|10\rangle + bd|11\rangle.$$

2.1 1

If the state is separable, then

$$ac = \frac{2}{3}$$
, $ad = \frac{1}{3}$, $bc = 0$, $bd = -\frac{2}{3}$.

It is clear that none of a, b, c, d are 0, so $\frac{c}{d} = \frac{ac}{ad} = 2$. Thus $bc = bd\frac{c}{d} = 2bd = -\frac{4}{3}$. Contradiction, so the state is entangled.

2.2 2

$$\frac{1}{2}(|00\rangle - i|01\rangle + i|10\rangle + |11\rangle) = \frac{1}{\sqrt{2}}(|0\rangle + i|1\rangle) \otimes \frac{1}{\sqrt{2}}(|0\rangle - i|i\rangle)$$

Separable state.

2.3 3

If the state is separable, then (omitting the normalization coefficient)

$$ac = 1$$
, $ad = -1$, $bc = 1$, $bd = 1$.

It is clear that none of a, b, c, d are 0, so $\frac{c}{d} = \frac{ac}{ad} = -1$. Thus $bc = bd\frac{c}{d} = -bd = -1$. Contradiction, so the state is entangled.

3 Problem 6

Let l = n - m - k. We can represent the states $|\Phi_x\rangle$ as

$$|\Phi_x\rangle_{n-m} = \sum_{y} b_{x,y} |y\rangle_k |\Omega_{x,y}\rangle_l.$$

3.1 a

After the measurement of m qbits, the state is $|x\rangle_m |\Phi_x\rangle_{n-m} = \sum_y b_{x,y} |x\rangle_m |y\rangle_k |\Omega_{x,y}\rangle_l$ with probability $|a_x|^2$. When another k qbits are measured, any vector y, for which $b_{x,y} \neq 0$, can be result of measurement, and the state collapses to $|x\rangle_m |y\rangle_k |\Omega_{x,y}\rangle$ with probability $|b_{x,y}|^2$. So the possible final states are $|x\rangle_m |y\rangle_k |\Omega_{x,y}\rangle_l$ with probability $|a_x|^2 |b_{x,y}|^2$.

3.2 b

In general, after the measurement of m+k qbits, the state could be one of $\sum_{z} c_{z} |x\rangle_{m} |y\rangle_{k} |z\rangle_{l}$ for some x and y. The probability amplitudes c_{z} are given by

$$\begin{split} c_z &= \left\langle x \right|_m \left\langle y \right|_k \left\langle z \right|_l \left| \Psi \right\rangle_n = \left\langle x \right|_m \left\langle y \right|_k \left\langle z \right|_l \sum_{x'} a_{x'} \left| x' \right\rangle_m \left| \Phi_{x'} \right\rangle_{n-m} = \\ &\left\langle x \right|_m \left\langle y \right|_k \left\langle z \right|_l \sum_{x'} a_{x'} \sum_{y'} b_{x',y'} \left| x' \right\rangle_m \left| y' \right\rangle_k \left| \Omega_{x',y'} \right\rangle_l = \\ &\sum_{x'} \sum_{y'} a_{x'} b_{x',y'} \left\langle x \right|_m \left\langle y \right|_k \left\langle z \right|_l \left| x' \right\rangle_m \left| y' \right\rangle_k \left| \Omega_{x',y'} \right\rangle_l. \end{split}$$

It is clear that the only possibly non-zero summand is when x'=x, and y=y', so $c_z=a_xb_{x,y}\langle z|_{l}|\Omega_{x,y}\rangle_{l}$. Now,

$$\sum_{z} c_{z} |x\rangle_{m} |y\rangle_{k} |z\rangle_{l} = a_{x}b_{x,y} \sum_{z} \langle z|_{l} |\Omega_{x,y}\rangle_{l} |z\rangle_{l} = a_{x}b_{x,y} |x\rangle_{m} |y\rangle_{k} |\Omega_{x,y}\rangle_{l}$$

Thus, the possible final states are $|x\rangle_m |y\rangle_k |\Omega_{x,y}\rangle_l$ with probability $|a_x|^2 |b_{x,y}|^2$, the same result as in (a).