

# **Part 8. Partial Differential Equations**

## **Chapter 29. Finite Difference: Parabolic Equations**

### **Lecture 31 & 32**

**The Heat Conduction Equations (30.1)**

**Explicit & Implicit Methods (30.2, 30.3)**

**The Crank-Nicolson Method (30.4)**

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# Finite Difference: Parabolic Equations

- Parabolic equations are employed to characterize time-variable (*unsteady-state*) problems.
- Conservation of energy can be used to develop an *unsteady-state* energy balance for the differential element in a long, thin insulated rod.

# Introduction: Parabolic PDEs

- General form for 2<sup>nd</sup> order linear PDE with 2 independent and 1 dependent variables:

$$A \frac{\partial^2 u}{\partial x^2} + B \frac{\partial^2 u}{\partial x \partial y} + C \frac{\partial^2 u}{\partial y^2} + D = 0$$

- Criteria for parabolic Eq:  $B^2 - 4AC = 0$

- e.g. in heat-conduction equation:  $\alpha \frac{\partial^2 T}{\partial x^2} = \frac{\partial T}{\partial t}$

$$A = \alpha, \quad B = 0$$

$$C = 0, \quad D = -1$$

$$B^2 - 4AC = 0 - 4(\alpha)(0) = 0$$

**Parabolic**

## Example. Heat Conduction Equation for Metal Rod

- Energy balance together with Fourier's law of heat conduction yields **heat-conduction equation**:

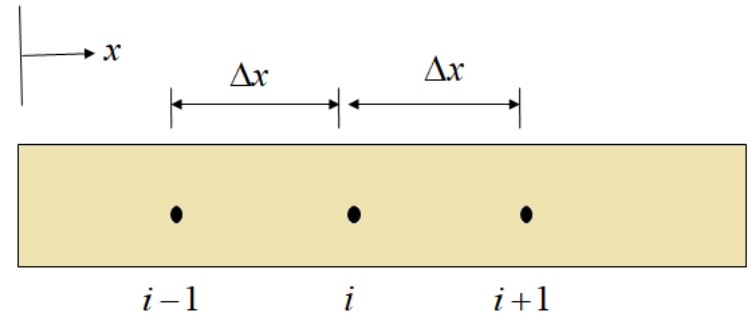
$$k \frac{\partial^2 T}{\partial x^2} = \frac{\partial T}{\partial t}$$

- Just as elliptic PDEs, parabolic equations can be solved by substituting finite divided differences for the partial derivatives.
- In contrast to elliptic PDEs, we must now consider changes in time as well as in space.

# Example of Parabolic PDE: Heat Conduction Equation for Metal Rod

The internal temperature of a metal rod exposed to two different temperatures at each end can be found using the heat conduction equation.

$$\alpha \frac{\partial^2 T}{\partial x^2} = \frac{\partial T}{\partial t}$$



**Discretizing Parabolic PDE**

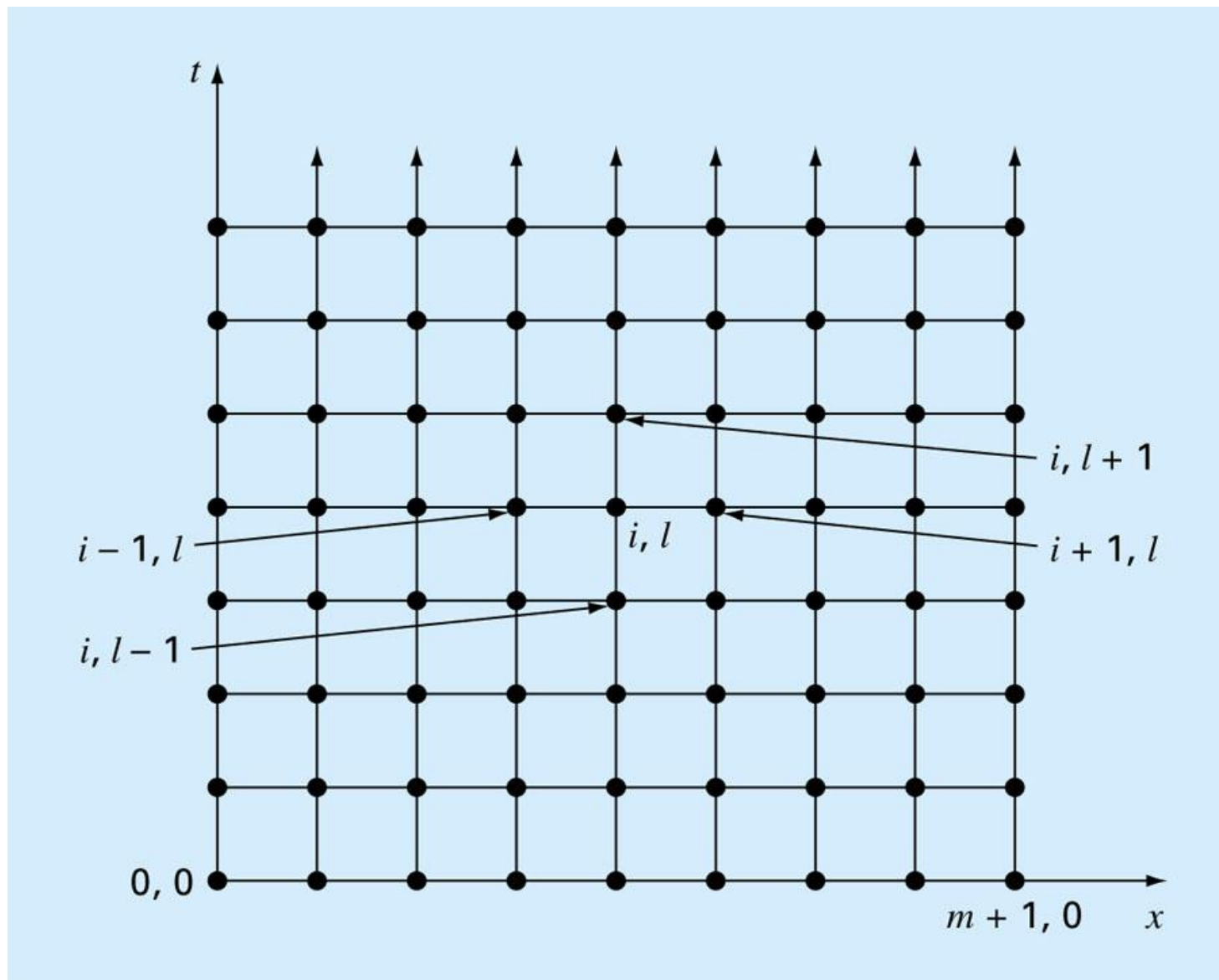
For a rod of length  $L$  divided into  $n+1$  nodes:  $\Delta x = L / n$

The time is similarly broken into time steps of  $\Delta t$

$$x = (i)(\Delta x)$$

Then:

$$t = (j)(\Delta t)$$



# Methods for Solution of Parabolic PDE

Explicit Method

Implicit Method

Crank-Nicolson Method (implicit)

# Explicit Method



# Explicit Method

$$\left. \frac{\partial^2 T}{\partial x^2} \right|_{i,j} \cong \frac{T_{i+1}^j - 2T_i^j + T_{i-1}^j}{(\Delta x)^2}$$

$$\left. \frac{\partial T}{\partial t} \right|_{i,j} \cong \frac{T_i^{j+1} - T_i^j}{\Delta t}$$

$$\alpha \frac{\partial^2 T}{\partial x^2} = \frac{\partial T}{\partial t} \quad \longrightarrow \quad \alpha \frac{T_{i+1}^j - 2T_i^j + T_{i-1}^j}{(\Delta x)^2} = \frac{T_i^{j+1} - T_i^j}{\Delta t}$$

$$T_i^{j+1} = T_i^j + \alpha \frac{\Delta t}{(\Delta x)^2} (T_{i+1}^j - 2T_i^j + T_{i-1}^j) \quad \lambda = \alpha \frac{\Delta t}{(\Delta x)^2}$$

$$T_i^{j+1} = T_i^j + \lambda (T_{i+1}^j - 2T_i^j + T_{i-1}^j)$$

# Explicit Method

$$T_i^{j+1} = T_i^j + \lambda(T_{i+1}^j - 2T_i^j + T_{i-1}^j)$$

- Eq. can be solved explicitly  $\rightarrow$  can be written for each internal location node of the rod for time node  $j$  in terms of the temperature at time node  $j+1$  .
- If the temperature at node  $j=0$  , and the boundary temperatures are known, the temperature at the next time step can be found.
- The process is continued by first finding the temperature at all nodes of  $j=1$ , and using these to find the temperature at the next time node,  $j=2$ . This process continues until the time at which we are interested in finding the temperature is reached.

# Implicit Method

# Implicit Method

Why ?

- Using the explicit method, we were able to find the temperature at each node, one equation at a time.
- However, the temperature at a specific node was only dependent on the temperature of the neighboring nodes from the previous time step. This is contrary to what we expect from the physical problem.
- The implicit method allows us to solve this and other problems by developing a system of simultaneous linear equations for the temperature at all interior nodes at a particular time.

# Implicit Method

$$\left. \frac{\partial T}{\partial t} \right|_{i,j+1} \approx \frac{T_i^{j+1} - T_i^j}{\Delta t} \quad \left. \frac{\partial^2 T}{\partial x^2} \right|_{i,j+1} \approx \frac{T_{i+1}^{j+1} - 2T_i^{j+1} + T_{i-1}^{j+1}}{(\Delta x)^2}$$

The second derivative on the left hand side of the equation is approximated by the Central divided difference scheme at time level  $j+1$  at node  $(i)$  as

$$\alpha \frac{\partial^2 T}{\partial x^2} = \frac{\partial T}{\partial t} \quad \alpha \frac{T_{i+1}^{j+1} - 2T_i^{j+1} + T_{i-1}^{j+1}}{(\Delta x)^2} = \frac{T_i^{j+1} - T_i^j}{\Delta t}$$

$$\lambda = \alpha \frac{\Delta t}{(\Delta x)^2}$$

$$-\lambda T_{i-1}^{j+1} + (1 + 2\lambda) T_i^{j+1} - \lambda T_{i+1}^{j+1} = T_i^j$$

**Example.** Solve the 1D heat conduction in rod equation using **explicit** and **implicit** method, if  $L=10$  cm and one side is  $100$  C and one side of the rod is  $0$  C .  $T(x, t=0)=20$  C.

$$\Delta x = 2cm \quad \Delta t = 1s \quad \lambda = 0.243$$

Solve using new time step,  $j=1$

1.486	-0.243	0	0
-0.243	1.486	-0.243	0
0	-0.243	1.486	-0.243
0	0	-0.243	1.486



Solve using new time step,  $j=2$

1.486	-0.243	0	0
-0.243	1.486	-0.243	0
0	-0.243	1.486	-0.243
0	0	-0.243	1.486





# Convergence and Stability

- Convergence means that as  $\Delta x$  and  $\Delta t$  approach zero, the results of the finite difference method approach the true solution.
- Stability means that errors at any stage of the computation are not amplified but are attenuated as the computation progresses.
- The explicit method is both convergent and stable if

$$\lambda \leq 1/2$$

*or*

$$\Delta t \leq \frac{1}{2} \frac{\Delta x^2}{k}$$

# Pros and Cons of Implicit Method

## Advantages

- Unconditionally stable
- Can use larger time step values
- More accurate than explicit method

## Challenges

- Computationally intense
- First order accurate in time

# **Crank-Nicolson Method**

## **(implicit)**

# Crank-Nicolson Method (implicit)

WHY:

Using the implicit method:

approximation of  $\frac{\partial^2 T}{\partial x^2}$  is of  $O(\Delta x)^2$  accuracy,

approximation of  $\frac{\partial T}{\partial t}$  is of  $O(\Delta t)$  accuracy.

# Crank-Nicolson Method (implicit)

approximating the second derivative at the midpoint of the time step.

$$\left. \frac{\partial^2 T}{\partial x^2} \right|_{i,j} \approx \frac{\alpha}{2} \left[ \frac{T_{i+1}^j - 2T_i^j + T_{i-1}^j}{(\Delta x)^2} + \frac{T_{i+1}^{j+1} - 2T_i^{j+1} + T_{i-1}^{j+1}}{(\Delta x)^2} \right] \quad \left. \frac{\partial T}{\partial t} \right|_{i,j} \approx \frac{T_i^{j+1} - T_i^j}{\Delta t}$$

$$\frac{\alpha}{2} \left[ \frac{T_{i+1}^j - 2T_i^j + T_{i-1}^j}{(\Delta x)^2} + \frac{T_{i+1}^{j+1} - 2T_i^{j+1} + T_{i-1}^{j+1}}{(\Delta x)^2} \right] = \frac{T_i^{j+1} - T_i^j}{\Delta t} \quad \lambda = \alpha \frac{\Delta t}{(\Delta x)^2}$$

$$-\lambda T_{i-1}^{j+1} + 2(1 + \lambda)T_i^{j+1} - \lambda T_{i+1}^{j+1} = \lambda T_{i-1}^j + 2(1 - \lambda)T_i^j + \lambda T_{i+1}^j$$

**Example.** Solve the 1D heat conduction in rod equation using **Crank-Nicolson** method, if  $L=10$  cm and one side is 100 C and one side of the rod is 0 C .  $T(x, t=0)=20$  C.

$$\Delta x = 2cm \quad \Delta t = 1s \quad \lambda = 0.243$$

## Comparison of 3 methods

	<b>Explicit</b>	<b>Implicit</b>	<b>Crank Nicolson</b>
$T_1^1$	39.4	33.4	27.89
$T_2^1$	20	22.2	20.8
$T_3^1$	20	19.8	19.9
$T_4^1$	15.1	16.7	16.1

# Group Problem Solving

Consider a steel rod that is subjected to a temperature of 100 C on the left end and 25 C on the right end. If the rod is of length 0.05 m, find the temperature distribution in the rod from  $t = 0$  and  $t = 9$  seconds using three methods blow. Use  $\Delta x = 0.01$  m and  $\Delta t = 3$  s. The initial temperature of the rod is 20C. By knowing the following constants:

Using:

- Explicit method (Group 1)
- Implicit Method (Group 2)
- Crank-Nicolson Method (Group 3)

$$\left\{ \begin{array}{l} k = 54 \frac{W}{m - K} \\ \rho = 7800 \frac{kg}{m^3} \\ C = 490 \frac{J}{kg - K} \end{array} \right.$$

