

Contents

1	Methods	1
1.1	Abstract	1
1.2	Technical details	1
1.3	Convergence criteria	2
1.3.1	Simulations with finite bond dimension	2
1.3.2	Simulations with finite system size	2

1

Methods

1.1 Abstract

We describe the technical details of the algorithms used to compute quantities of interest. We report the convergence behaviour of the algorithms and discuss validity and sources of error.

1.2 Technical details

For the models treated in this thesis, the corner transfer matrix \mathcal{A} and the row-to-row transfer matrix T are symmetric. But due to the accumulation of machine-precision sized errors in the matrix multiplication and singular value decomposition, this will, after many algorithm steps, no longer be the case. In order for results to remain valid, we manually enforce symmetricity after each step.

The tensor network contractions at each algorithm step will cause the elements of \mathcal{A} and T to tend to infinity, which means that they will at some point exceed the maximum value of a floating point number as it can be stored in memory. But because the elements of \mathcal{A} and T represent Boltzmann weights, they can be scaled by a constant factor, which allows us to prevent this overflow if we use a suitable scaling. For example by requiring that

$$\mathrm{Tr} \mathcal{A}^4 = 1, \tag{1.1}$$

so that the interpretation of \mathcal{A}^4 as a reduced density matrix of an effective one-dimensional quantum is valid.

1.3 Convergence criteria

1.3.1 Simulations with finite bond dimension

The convergence of the CTMRG algorithm with fixed bond dimension m (the infinite system algorithm) can be defined in multiple ways (*cite*). In this thesis, the convergence after step i of the algorithm is defined as

$$c_i = \sum_{\alpha=1}^m |s_{\alpha}^{(i)} - s_{\alpha}^{(i-1)}|, \quad (1.2)$$

where s_{α} are the singular values of the corner transfer matrix \mathcal{A} . If the convergence falls below some threshold ϵ , the algorithm terminates.

The assumption is that once the singular values stop changing to some precision, the optimal projection is sufficiently close to its fixed point and the transfer matrices \mathcal{A} and T represent an environment only limited by the length scale given by m , i.e.

$$\xi(m) \ll N \quad (1.3)$$

is satisfied.

Convergence of quantities at the critical point

The convergence of the various quantities at the critical point of the Ising model is shown in ??.

Cross check with correlation length, report on boundary conditions

1.3.2 Simulations with finite system size

In the finite-system algorithm, we require

$$N \ll \xi(m), \quad (1.4)$$

so the question becomes at which m the results are sufficiently converged in m . Throughout this work, we have used the residual probability (also called truncation error)

$$P(m)^{(i)} = \frac{\sum_{\alpha=m+1}^{dm} (s_{\alpha}^{(i)})^2}{\sum_{\alpha=1}^{dm} (s_{\alpha}^{(i)})^2}, \quad (1.5)$$

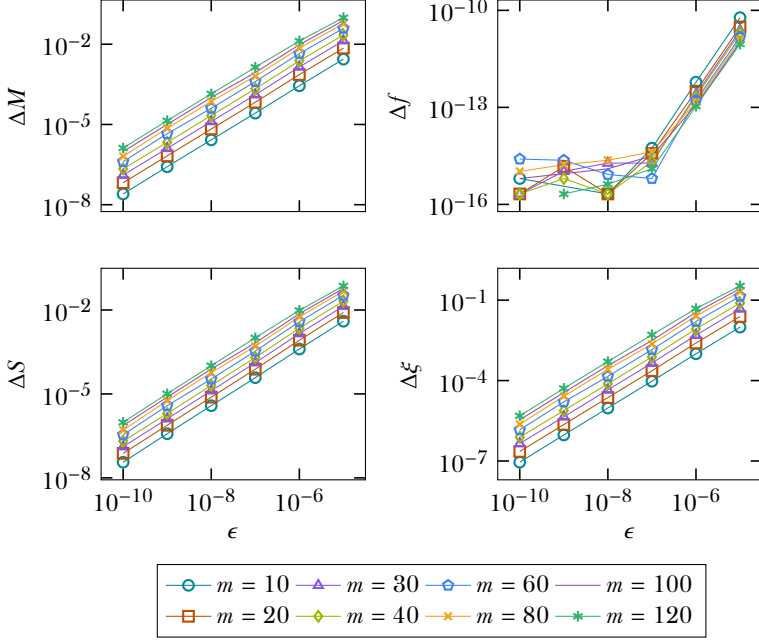


Figure 1.1: hallootjes

which quantifies the fraction of the spectrum of the corner transfer matrix that is thrown away, as a measure of how accurate the transfer matrices represent the finite system of size N . Here, d is the dimension of the local tensors ($d = 2$ for the Ising model).

If, for a given m , we have

$$P(m) < P_{\max} \quad (1.6)$$

we deem the result accurate enough.

In the limit $m \rightarrow d^n$, with

$$n = \frac{N-1}{2} \quad (1.7)$$

the number of algorithm steps, we obtain the exact result for the transfer matrices and hence the partition function, i.e. $P(m) \rightarrow 0$.

To justify that for small enough P_{\max} , we obtain good results for a wide range of N , figure bla bla.