

Unmanned Autonomous Systems - 31390

Modeling of Unmanned Aerial Vehicles

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Two major areas of modeling

- **Kinematics**: studies of the motion of a mechanical system, taking into account the constraints of a kinematic structure, without taking into account for forces and torques that govern the motion of the system.
- **Dynamics**: studies the motion of a system as a function of the forces and torques that act on a kinematic structure

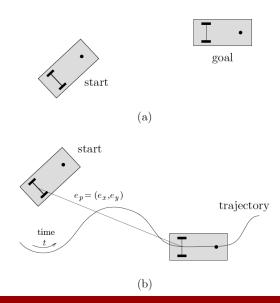


In order to study any physical system, we need to create a representation of the system

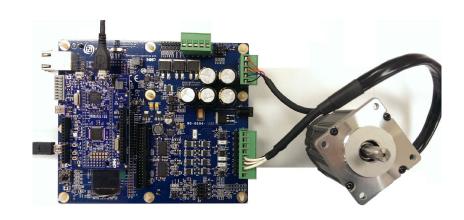
A **model** is a representation, through mathematical equations, of a system.

Example:

wheeled robot motion

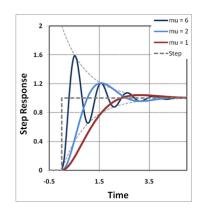


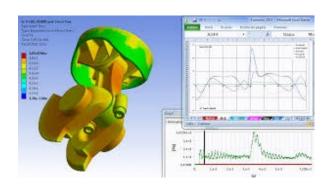
Motor Control

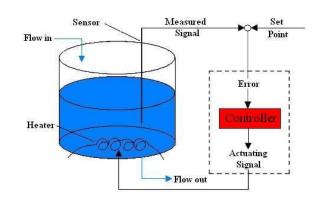




- What is, in practice, the purpose of a model?
 - understanding the behavior of a system (e.g. source of vibrations in a precision machine)
 - predicting the behavior of a system, given an external agent, without need to observe it in the real world
 - design of mechanical components based on the actions acting on elements (e.g. design a rod-crank mechanism for a combustion motor of a car
 - synthesis of the controller of a system and off-line regulation of control parameter









Kinematics



In order to study the dynamic behavior of mechanical systems, we need to first model their kinematics

Reminder: the kinematics is the study of the motion of a body, or a set of interconnected bodies without considering the actions acting on the system.

Studying the kinematics of a system means to define the relations among the position, velocity and acceleration of the bodies interconnected in the system

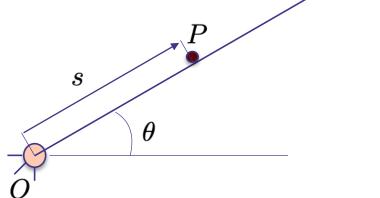


- interest is on defining quantities such as **position**, **velocity and acceleration**
- need to specify a reference frame (and a coordinate system in it to actually write the vector expressions)
- Velocity and acceleration depend of the choice of the reference frame
- Only when we go to laws of motion, the reference frame needs to be the inertial frame



Kinematics of a point in the plane:

a first exercise

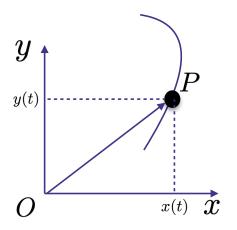


A bar BO can rotate around O The motion of the bar is described by the variable $\; \theta = \theta(t) \;$

A point P can slide on the bar BO, and its coordinate on the bar is described by the variable s=s(t)



Kinematics of a point in the plane



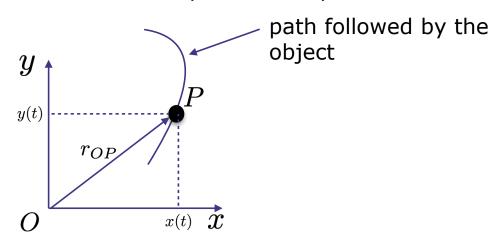
How can we describe the kinematics of a point moving on a trajectory?

Or more generally, how can we describe the motion of a system that consists of interconnected bodies? (e.g. a robot, a car, a production line)

We need a mathematical representation of the position of the bodies and their derivatives (velocity, acceleration, jerk, ...)



Kinematics of a point in the plane



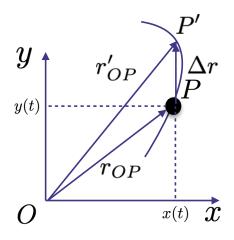
 r_{OP} position vector (specifies the position P of the object w.r.t. O)

Note
$$\overrightarrow{r_{OP}} = \overrightarrow{r_{OP}}(t)$$
 $x = x(t)$ $y = y(t)$

$$t = g(x)$$
$$y = y(g(x)) = f(x)$$



Kinematics of a point in the plane



Velocity vector:

$$\overrightarrow{v_P} = \frac{d}{dt} \overrightarrow{r_{OP}}(t) = \lim_{\Delta t \to 0} \frac{\Delta \overrightarrow{r_{OP}}}{\Delta t}$$

Acceleration vector:

$$\overrightarrow{a_P} = \frac{d}{dt} \overrightarrow{v_P}(t) = \frac{d^2}{dt^2} \overrightarrow{r_{OP}}(t)$$



Speed (scalar)
$$v_P = \|v_P\|$$

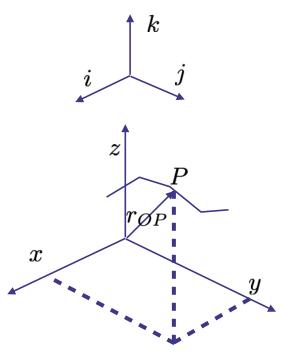
Magnitude of acceleration
$$a_P = \|\vec{a}_P\|$$

Note: the time derivatives have been considered w.r.t. a reference frame

We can describe them in various coordinate systems:

- Cartesian
- Cylindrical





$$\overrightarrow{i}$$
 \overrightarrow{j} \overrightarrow{k} unit vectors

NOTE: from now on, we call these unit vectors: $i \ j \ k$

Cartesian coordinate system

i j k are an orthogonal set

$$i \times j = k$$
 $j \times k = i$
 $k \times i = j$

$$\overrightarrow{r_{OP}} = x(t)i + y(t)j + z(t)k$$



The time derivative of the position vector is the velocity

$$\overrightarrow{v_P} = \frac{d\overrightarrow{r_{OP}}}{dt} = \frac{dx(t)}{dt}i + \frac{dy(t)}{dt}j + \frac{dz(t)}{dt}k + \frac{di}{dt}x(t) + \frac{dj}{dt}y(t) + \frac{dk}{dt}z(t)$$

the rate of change in a frame in which i,j,k are fixed, gives $\frac{di}{dt} = \frac{dj}{dt} = \frac{dk}{dt} = 0$

$$\overrightarrow{v_P} = \frac{dx(t)}{dt}i + \frac{dy(t)}{dt}j + \frac{dz(t)}{dt}k$$

Similarly $\overrightarrow{a_P} = \ddot{x}(t)i + \ddot{y}(t)j + \ddot{z}(t)k$



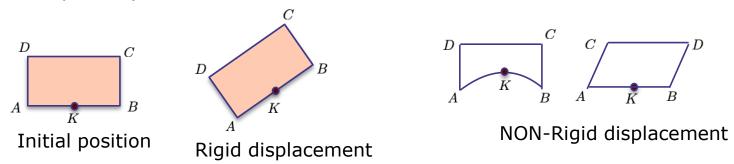
Kinematics of **RIGID BODIES**

Position: set of vectors defining the position of every point constituting the rigid body

Movement: is the description of the position of the rigid body (and of how all the vectors representing the position of every of its points) over time



Definition: the infinitesimal displacement is a displacement in which all the points constituting the body, change their position of an infinitesimal quantity.



Properties of a rigid body:

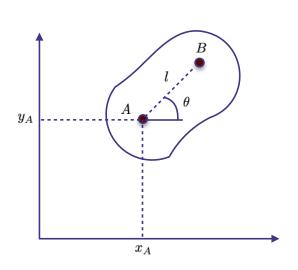
for any motion of the rigid body

- the distance between any arbitrary couple of points of the rigid body does not change.
- the angles between the segments connecting couples of points of the rigid body do not change



Kinematics of **RIGID BODIES** in the plane

We can reduce the number of DoF from infinite 2 (the coordinate of every point of the body) to 3: x,y,theta



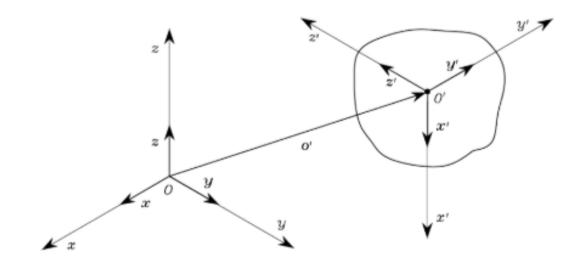
if we know the coordinates (x,y) of point A, we can represent the coordinates of point B by using 1 more coordinate, the retailing (provided that we know the distance between A and B)

given θ , every other couple of points will rotate of the same quantity. This means that the rotation is a property of the entire rigid body.



Rotations





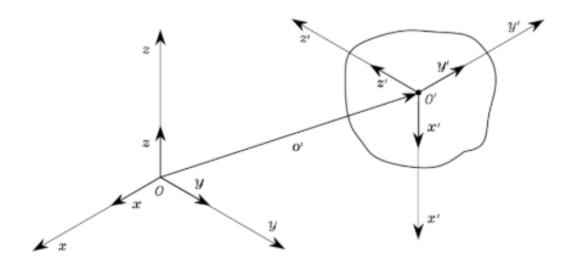
A body is represented in a base frame of reference by a vector that defines the position of the origin of a body fixed frame w.r.t. the base frame of reference

$$oldsymbol{o}' = egin{bmatrix} o_x' \ o_y' \ o_z' \end{bmatrix}$$

The orientation of the body w.r.t. the base frame is described by the projection of the unit vectors of the body-fixed frame into the base frame

$$egin{aligned} oldsymbol{x}' &= x_x' oldsymbol{x} + x_y' oldsymbol{y} + x_z' oldsymbol{z} \ oldsymbol{y}' &= y_x' oldsymbol{x} + y_y' oldsymbol{y} + y_z' oldsymbol{z} \ oldsymbol{z}' &= z_x' oldsymbol{x} + z_y' oldsymbol{y} + z_z' oldsymbol{z}. \end{aligned}$$





The rotation of a body can be described, in compact form, by a 3x3 matrix that combines the projection of the three unit vectors

$$oldsymbol{R} = egin{bmatrix} oldsymbol{x}' & oldsymbol{y}' & oldsymbol{z}' \ oldsymbol{x}' & oldsymbol{y}' & oldsymbol{z}' \ oldsymbol{x}'_x & oldsymbol{y}'_x & oldsymbol{z}'_x & oldsymbol{y}'_x & oldsymbol{z}'^Toldsymbol{x} & oldsymbol{z}'^Toldsymbol{z} & oldsymbol{$$

Matrix R is a rotation matrix.

Note that:

The column vectors of R are mutually orthogonal The column vectors of R have unit norm

 $\boldsymbol{x}'^T \boldsymbol{y}' = 0$ $\boldsymbol{y}'^T \boldsymbol{z}' = 0$ $\boldsymbol{z}'^T \boldsymbol{x}' = 0$

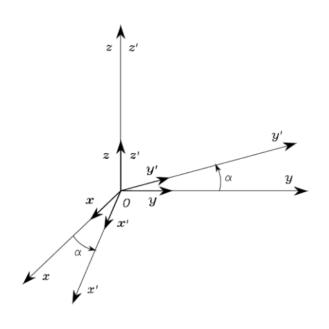
$$\boldsymbol{x'}^T\boldsymbol{x'}=1$$

$$oldsymbol{x}'^Toldsymbol{x}'=1 \qquad oldsymbol{y}'^Toldsymbol{y}'=1 \qquad oldsymbol{z}'^Toldsymbol{z}'=1$$

R is an orthogonal matrix

$$\boldsymbol{R}^T \boldsymbol{R} = \boldsymbol{I}_3$$





Rotation around axis z, by an angle α

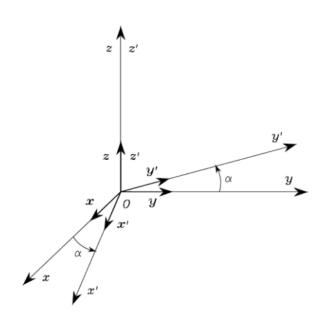
$$m{x}' = egin{bmatrix} \cos lpha \\ \sin lpha \\ 0 \end{bmatrix} \qquad m{y}' = egin{bmatrix} -\sin lpha \\ \cos lpha \\ 0 \end{bmatrix} \qquad m{z}' = egin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} \qquad m{R}_z(lpha) = egin{bmatrix} \cos lpha & -\sin lpha & 0 \\ \sin lpha & \cos lpha & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

Similarly, a rotation by β around axis y, and by γ around axis x are described by matrices

$$m{R}_y(m{eta}) = egin{bmatrix} \cosm{eta} & 0 & \sinm{eta} \ 0 & 1 & 0 \ -\sinm{eta} & 0 & \cosm{eta} \end{bmatrix} \qquad m{R}_x(\gamma) = egin{bmatrix} 1 & 0 & 0 \ 0 & \cos\gamma & -\sin\gamma \ 0 & \sin\gamma & \cos\gamma \end{bmatrix}$$

$$m{R}_x(\gamma) = egin{bmatrix} 1 & 0 & 0 \ 0 & \cos \gamma & -\sin \gamma \ 0 & \sin \gamma & \cos \gamma \end{bmatrix}$$



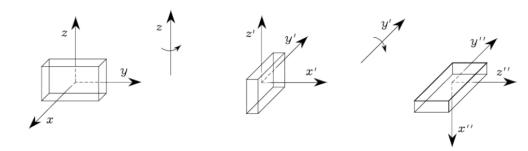


Composition of rotations:

An overall rotation can be expressed by a sequence of partial rotations, with each rotation defined w.r.t. the preceding one.

Composition of successive rotations is achieved by postmultiplication of the rotation matrices following the given order of rotations

$$egin{aligned} oldsymbol{p}^1 &= oldsymbol{R}_2^1 oldsymbol{p}^2 \ oldsymbol{p}^0 &= oldsymbol{R}_1^0 oldsymbol{p}^1 & oldsymbol{R}_2^0 &= oldsymbol{R}_1^0 oldsymbol{R}_2^1 \ oldsymbol{p}^0 &= oldsymbol{R}_2^0 oldsymbol{p}^2 \end{aligned}$$





Derivative of a Rotation matrix

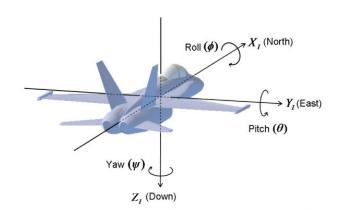
$$S(t) = \dot{R}(t)R^{T}(t)$$

$$R(t)R^{T}(t) = I \qquad \qquad \dot{R}(t)R^{T}(t) + R(t)\dot{R}^{T}(t) = 0 \qquad \qquad \dot{R}(t) = S(t)R(t)$$

S(t) Is a Skew symmetric matrix, as $S(t) + S^{T}(t) = 0$

$$S(t) = \begin{bmatrix} 0 & -\omega_z & -\omega_y \\ -\omega_z & 0 & -\omega_x \\ -\omega_y & -\omega_x & 0 \end{bmatrix}$$





Minimum representation of orientation

A rotation can be represented with a minimum of 3 parameters

9 elements 6 orthogonality constraints

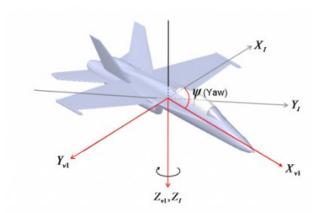


3 parameters,

Euler Angles

RPY angles

A set of Euler angles following the order ZYX (Roll-Pitch-Yaw angles)

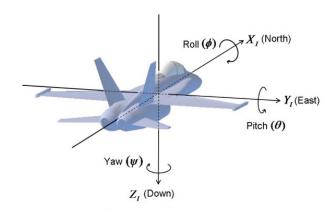


$$R_I^{v1}(\psi) = egin{pmatrix} \cos(\psi) & \sin(\psi) & 0 \ -\sin(\psi) & \cos(\psi) & 0 \ 0 & 0 & 1 \end{pmatrix}$$

Figure 2 - Yaw rotation into the Vehicle-1 Frame

Singularity problems in inverse relation





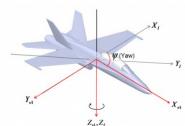


Figure 2 - Yaw rotation into the Vehicle-1 Frame

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Minimum representation of orientation

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9 elements



6 orthogonality constraints



3 parameters, Euler Angles

RPY angles

A set of Euler angles following the order ZYX (Roll-Pitch-Yaw angles)

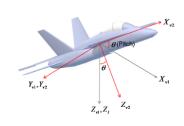
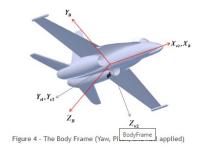


Figure 3 - The Vehicle-3 Frame (Yaw and Pitch Rotation Applied)

$$R_{v1}^{v2}(heta) = egin{pmatrix} \cos(heta) & 0 & -\sin(heta) \ 0 & 1 & 0 \ \sin(heta) & 0 & \cos(heta) \end{pmatrix}$$



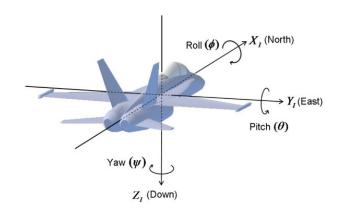
$$R^B_{v2}(\phi) = egin{pmatrix} 1 & 0 & 0 \ 0 & \cos(\phi) & \sin(\phi) \ 0 & -\sin(\phi) & \cos(\phi) \end{pmatrix}$$

B represents the body-frame I represents the Inertial frame

$$R_{I}^{B}(\phi, \theta, \psi) = R_{v2}^{B}(\phi)R_{v1}^{v2}(\theta)R_{I}^{v1}(\psi)$$

$$R_B^I(\phi, heta,\psi) = R_I^{v1}(-\psi)R_{v1}^{v2}(- heta)R_{v2}^B(-\phi)$$





How to obtain inertial frame accelerometer data from an accelerometer placed on the body?

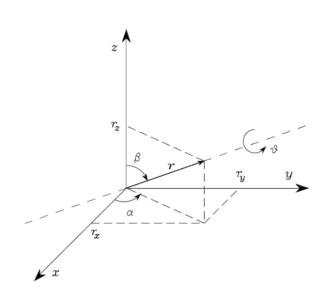
$$\mathbf{v}_I = R_B^I(\phi, \theta, \psi)\mathbf{v}_B$$

How to convert rate gyro data (rotational velocity of the body, or body angular rates) to derivative of Euler angles

$$D(\phi, \theta, \psi) = egin{pmatrix} 1 & \sin(\phi) \tan(\theta) & \cos(\phi) \tan(\theta) \\ 0 & \cos(\phi) & -\sin(\phi) \\ 0 & \sin(\phi)/\cos(\theta) & \cos(\phi)/\cos(\theta) \end{pmatrix} \qquad egin{pmatrix} \dot{\phi} \\ \dot{\dot{\theta}} \\ \dot{\psi} \end{pmatrix} = egin{pmatrix} p + q\sin(\phi) \tan(\theta) + r\cos(\phi) \tan(\theta) \\ q\cos(\phi) - r\sin(\phi) \\ q\sin(\phi)/\cos(\theta) + r\cos(\phi)/\cos(\theta) \end{pmatrix}$$

http://www.chrobotics.com/library/understanding-euler-angles





Non-minimal representation of orientation: axis-angle

A rotation can be represented with of 4 parameters, defining an axis of rotation (3 parameters) and the rotation around such axis (1 parameter)

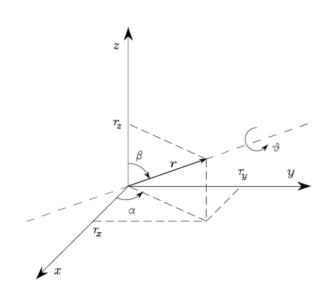
Procedure to compute R, given the axis and the angle

- Align r with z, which is obtained as the sequence of a rotation by $-\alpha$ about z and a rotation by $-\beta$ about y.
- Rotate by ϑ about z.
- Realign with the initial direction of r, which is obtained as the sequence of a rotation by β about y and a rotation by α about z.

$$R(\vartheta, r) = R_z(\alpha)R_y(\beta)R_z(\vartheta)R_y(-\beta)R_z(-\alpha)$$

being
$$\sin\alpha = \frac{r_y}{\sqrt{r_x^2 + r_y^2}} \qquad \cos\alpha = \frac{r_x}{\sqrt{r_x^2 + r_y^2}} \qquad \sin\beta = \sqrt{r_x^2 + r_y^2} \qquad \cos\beta = r_z$$





Non-minimal representation of orientation: axis-angle

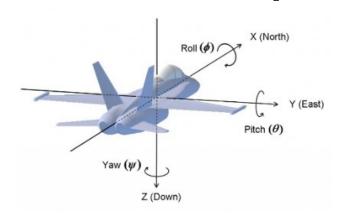
A rotation can be represented with of 4 parameters, defining an axis of rotation (3 parameters) and the rotation around such axis (1 parameter)

Inverse problem

$$\vartheta = \cos^{-1}\left(\frac{r_{11} + r_{22} + r_{33} - 1}{2}\right) \qquad r = \frac{1}{2\sin\vartheta} \begin{bmatrix} r_{32} - r_{23} \\ r_{13} - r_{31} \\ r_{21} - r_{12} \end{bmatrix}$$

Singularity when $\vartheta = 0$





Non-minimal representation of orientation: unit quaternion

A rotation can be represented with of 4 parameters, consisting of 1 real and 3 complex numbers.

How can we use quaternions to encode rotations from inertial frame to body-fixed frame?

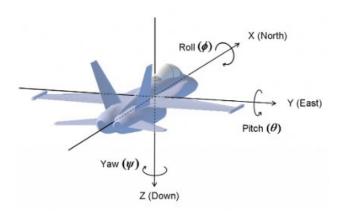
Let the vector q_i^b be defined as the unit-vector quaternion encoding rotation from the inertial frame to the body frame of the sensor

$$\mathbf{q}_i^b = \left(egin{array}{cccc} a & b & c & d \end{array}
ight)^T$$

The elements **b**, **c**, and **d** are the "vector part" of the quaternion, and can be thought of as a vector about which rotation should be performed.

The element **a** is the "scalar part" that specifies the amount of rotation that should be performed about the vector part.





Non-minimal representation of orientation: unit quaternion

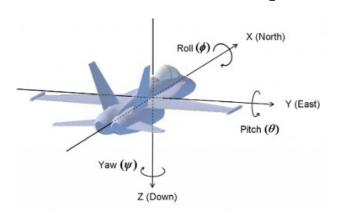
A rotation can be represented with of 4 parameters, consisting of 1 real and 3 complex numbers.

How can we use quaternions to encode rotations from inertial frame to body-fixed frame?

if θ is the angle of rotation and the vector $(v_x v_y v_z)^T$ is a unit vector representing the axis of rotation, then the quaternion elements are defined as

$$egin{pmatrix} a \ b \ c \ d \end{pmatrix} = egin{pmatrix} \cos(0.5 heta) \ v_x \sin(0.5 heta) \ v_y \sin(0.5 heta) \ v_z \sin(0.5 heta) \end{pmatrix}$$





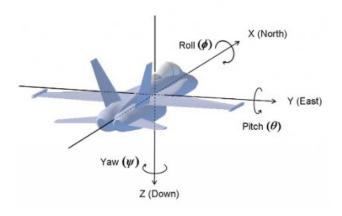
Quaternions

The attitude quaternion q_i^b can be used to rotate an arbitrary 3-element vector from the inertial frame to the body frame using the operation

$$\mathbf{v}_B = \mathbf{q}_i^b \left(egin{array}{c} 0 \ \mathbf{v}_I \end{array}
ight) (\mathbf{q}_i^b)^{-1}.$$

- a vector can rotated by treating it like a quaternion with zero real-part and multiplying it by the attitude quaternion and its inverse.
- The inverse of a quaternion is equivalent to its conjugate, which means that all the vector elements (the last three elements in the vector) are negated





Quaternions

The attitude quaternion q_i^b can be used to rotate an arbitrary 3-element vector from the inertial frame to the body frame using the operation

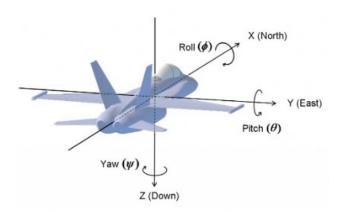
$$\mathbf{v}_B = \mathbf{q}_i^b \left(egin{array}{c} 0 \ \mathbf{v}_I \end{array}
ight) (\mathbf{q}_i^b)^{-1}.$$

Quaternion multiplications

Given two quaternions

Given two quaternions
$$\mathbf{q}_1 = \begin{pmatrix} a_1 & b_1 & c_1 & d_1 \end{pmatrix}^T$$
 $\mathbf{q}_2 = \begin{pmatrix} a_1a_2 - b_1b_2 - c_1c_2 - d_1d_2 \\ a_1b_2 + b_1a_2 + c_1d_2 - d_1c_2 \\ a_1c_2 - b_1d_2 + c_1a_2 + d_1b_2 \\ a_1d_2 + b_1c_2 - c_1b_2 + d_1a_2 \end{pmatrix}$





Quaternions

We can reconstruct a rotation matrix from an attitude quaternion.

The rotation matrix from the inertial frame to the body-fixed frame is given by:

$$R_i^b(\mathbf{q}_i^b) = egin{pmatrix} a^2 + b^2 - c^2 - d^2 & 2bc - 2ad & 2bd + 2ac \ 2bc + 2ad & a^2 - b^2 + c^2 - d^2 & 2cd - 2ab \ 2bd - 2ac & 2cd + 2ab & a^2 - b^2 - c^2 + d^2 \end{pmatrix}$$

We can therefore represent vectors in the inertial frame, as vector in the body-fixed frame as:

$$\mathbf{v}_B = R_i^b(\mathbf{q}_i^b)\mathbf{v}_I$$

http://www.chrobotics.com/library/understanding-quaternions



Derivative of a Rotation matrix

Consider a constant vector p' and the vector p(t) = R(t)p'

The derivative of p(t) is $\dot{p}(t) = \dot{R}(t)p'$

Which can be written as

$$\dot{p}(t) = S(t)R(t)p' = \omega(t) \times R(t)p'$$

Where $\omega = [\omega_x \quad \omega_y \quad \omega_z]^T$ is the vector of the angular velocity of frame R(t) with respect to the reference frame at time t



If we represent the the rotation of a frame with Euler angles (e.g. RPY -> $\vec{\theta} = [\varphi \quad \vartheta \quad \psi]^T$), is the derivative of the Euler angles the same as ω ?

NO!

How are they related?

It depends on the angle representation that we chose

$$ec{\omega} = egin{bmatrix} 1 & 0 & -s_{ heta} \ 0 & c_{\phi} & c_{ heta} s_{\phi} \ 0 & -s_{\phi} & c_{ heta} egin{bmatrix} \dot{ heta} \ 0 & -s_{\phi} & c_{ heta} egin{bmatrix} \dot{ heta} \ \end{pmatrix} \dot{ heta} \end{pmatrix} \qquad D(\phi, heta, \psi) = egin{bmatrix} 1 & \sin(\phi) \tan(\theta) & \cos(\phi) \tan(\theta) \ 0 & \cos(\phi) & -\sin(\phi) \ 0 & \sin(\phi)/\cos(\theta) & \cos(\phi)/\cos(\theta) \end{pmatrix}$$



Dynamics

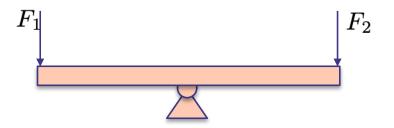


Discipline of mechanics that studies the equilibrium of a body

In other words, in a static system, all forces and moments are ballanced

... what does it mean?

Example: find the relations and constraint forces that keep the system in equilibrium

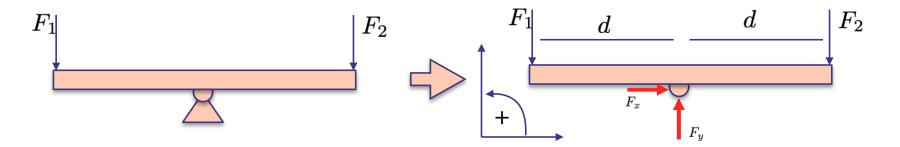


We need to make a ballance of forces in x and y direction, in order to define the constraint forces, and write one or more ballance of moments to make sure we have rotational equilibrium

How does a ballance of forces looks like?



How does a ballance of forces looks like?



$$F_y - F_1 - F_2 = 0$$

Ballance of forces in y-direction

$$F_x = 0$$

Ballance of forces in x-direction

$$F_1d - F_2d = 0$$

Ballance of moments around the pivot



Basics notions:

Laws of Mechanics

1) An object at rest stays at rest and an object in motion stays in motion with the same speed and in the same direction unless acted upon by an unbalanced force (Law of Inertia)

There exists a frame, with reference to which, an isolated system that is fixed, stays fixed. This frame of reference is called *Inertial frame*.





Basics notions:

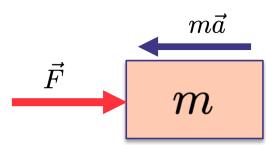
Laws of Mechanics

2) The acceleration of an object as produced by a net force is directly proportional to the magnitude of the net force, in the same direction as the net force, and inversely proportional to the mass of the object.

A material point s.t. a force \vec{F} moves with an acceleration \vec{a} that is proportional to the force \vec{F} , following:

$$\vec{F} = m\vec{a}$$

m is the inertial mass of the object



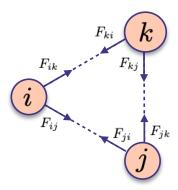


Basics notions:

Laws of Mechanics

3) For every action, there is an equal and opposite reaction.

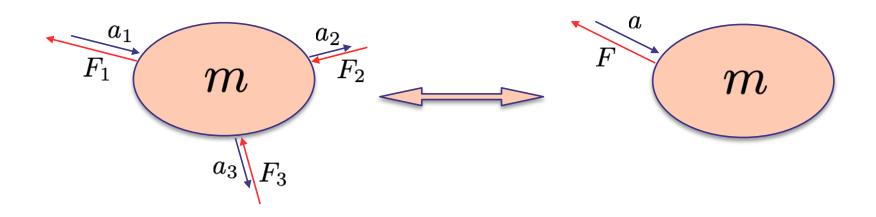
Given a system of material points that interact among each other, the forces exchanged among those points are equal and opposite in couples, and directed on the same line.





Basics notions:

If a force F_1 applied to a point produces an acceleration a_1 , and if a force F_2 applied to the same point produces an acceleration a_2 , then the two forces, when applied simultaneously, will produce an acceleration $a=a_1+a_2$ that correspond to the effect of the overall force $F=F_1+F_2$





Types of Forces: some definitions

A force \vec{F} acting on a point, generally depends on:

- \vec{r} the position of the point \vec{v} the velocity of point P
- time

$$\vec{F} = \vec{F}(\vec{r}, \ \vec{v}, \ t)$$

We define:

ACTIVE: those forces, known a priori, that depend on the position and velocity of a point

REACTIVE: those forces that are a reaction to the motion (e.g. constraint forces)

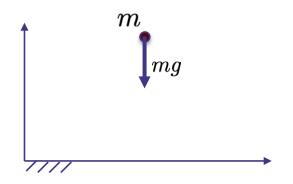


Types of Forces: some definitions

Constant Force:

A force that does not change over time in both modulus and direction

An example is the weight of an object



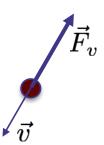
$$ec{w} = -mgec{j}$$



Types of Forces: some definitions

Viscous Force:

it is not a positional force!



$$\vec{F}_v = -\gamma \vec{v}$$

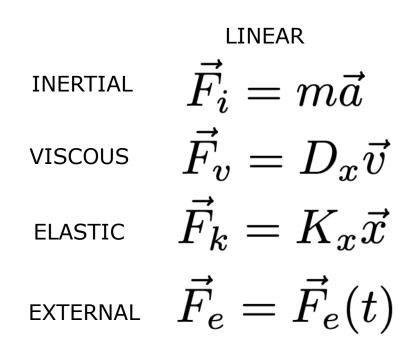
A viscous force is opposed to the motion of the object.

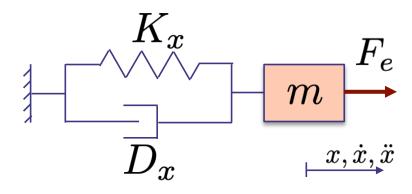
Examples of viscous force are:

- air resistance
- fluid resistance (e.g. water)



Some exercises (dynamics):

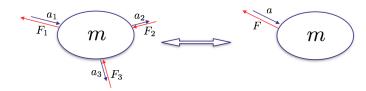




DYNAMICS

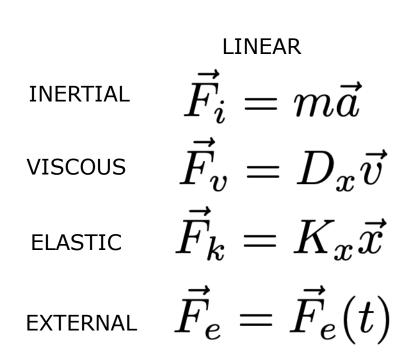
Basics notions:

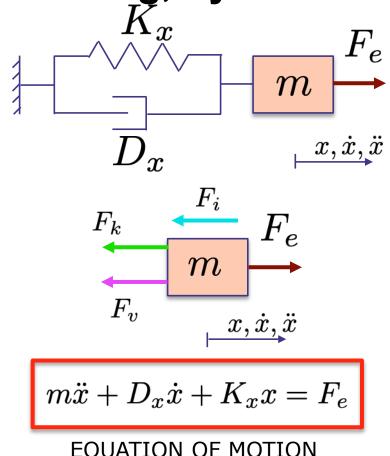
If a force F_1 applied to a point produces an acceleration a_1 , and if a force F_2 applied to the same point produces an acceleration a_2 , then the two forces, when applied simultaneously, will produce an acceleration $a=a_1+a_2$ that correspond to the effect of the overall force $F=F_1+F_2$





Some exercises (dynamics):

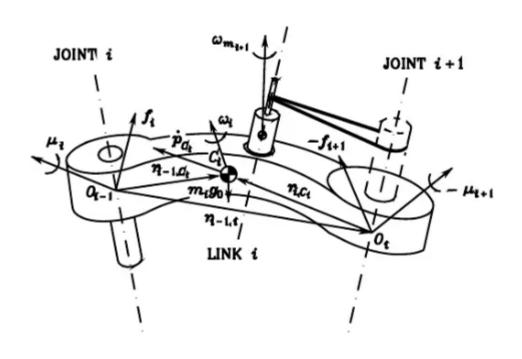




EQUATION OF MOTION



Recursive Newton-Euler: recap from robotic course



parameters

 m_i mass of augmented link,

 \bar{I}_i inertia tensor of augmented link,

 I_{m_i} moment of inertia of rotor,

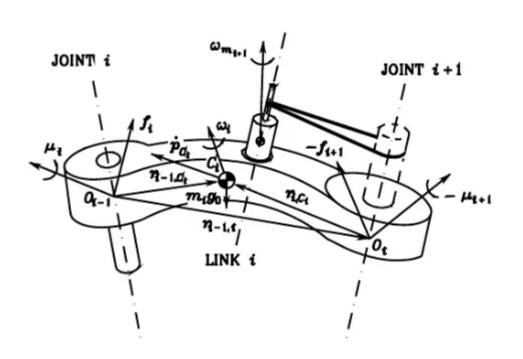
 r_{i-1,C_i} vector from origin of Frame (i-1) to centre of mass C_i ,

 r_{i,C_i} vector from origin of Frame i to centre of mass C_i ,

 $r_{i-1,i}$ vector from origin of Frame (i-1) to origin of Frame i.



Recursive Newton-Euler: recap from robotic course

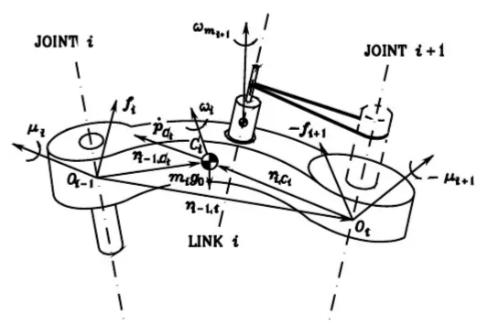


Kinematic variables (velocities and accelerations to be considered)

 \dot{p}_{C_i} linear velocity of centre of mass C_i , \dot{p}_i linear velocity of origin of Frame i, ω_i angular velocity of link, ω_{m_i} angular velocity of rotor, \ddot{p}_{C_i} linear acceleration of centre of mass C_i , \ddot{p}_i linear acceleration of origin of Frame i, $\dot{\omega}_i$ angular acceleration of link, $\dot{\omega}_{m_i}$ angular acceleration of rotor, g_0 gravity acceleration.



Recursive Newton-Euler: recap from robotic course



Dynamic variables (forces and moments to be considered)

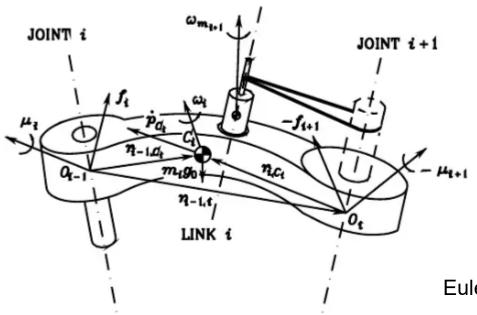
 f_i force exerted by Link i-1 on Link i, $-f_{i+1}$ force exerted by Link i+1 on Link i,

 μ_i moment exerted by Link i-1 on Link i with respect to origin of Frame i-1,

 $-\mu_{i+1}$ moment exerted by Link i+1 on Link i with respect to origin of Frame i.



Recursive Newton-Euler: recap from robotic course



Quantities expressed w.r.t. base frame

Newton Equations (translational motion of CoM)

$$\boldsymbol{f}_i - \boldsymbol{f}_{i+1} + m_i \boldsymbol{g}_0 = m_i \ddot{\boldsymbol{p}}_{C_i}$$

Euler Equations (rotational motion of CoM)

$$\mu_i + f_i \times r_{i-1,C_i} - \mu_{i+1} - f_{i+1} \times r_{i,C_i} = \frac{d}{dt} (\bar{I}_i \omega_i + k_{r,i+1} \dot{q}_{i+1} I_{m_{i+1}} z_{m_{i+1}})$$



Euler Equations (rotational motion of CoM)

$$\mu_i + f_i \times r_{i-1,C_i} - \mu_{i+1} - f_{i+1} \times r_{i,C_i} = \frac{d}{dt} (\bar{I}_i \omega_i + k_{r,i+1} \dot{q}_{i+1} I_{m_{i+1}} z_{m_{i+1}})$$

It is convenient to express the inertia tensor in the current frame, to have a constant tensor.

$$\begin{split} \frac{d}{dt}(\bar{\boldsymbol{I}}_{i}\boldsymbol{\omega}_{i}) &= \dot{\boldsymbol{R}}_{i}\bar{\boldsymbol{I}}_{i}^{i}\boldsymbol{R}_{i}^{T}\boldsymbol{\omega}_{i} + \boldsymbol{R}_{i}\bar{\boldsymbol{I}}_{i}^{i}\dot{\boldsymbol{R}}_{i}^{T}\boldsymbol{\omega}_{i} + \boldsymbol{R}_{i}\bar{\boldsymbol{I}}_{i}^{i}\boldsymbol{R}_{i}^{T}\dot{\boldsymbol{\omega}}_{i} \\ &= \boldsymbol{S}(\boldsymbol{\omega}_{i})\boldsymbol{R}_{i}\bar{\boldsymbol{I}}_{i}^{i}\boldsymbol{R}_{i}^{T}\boldsymbol{\omega}_{i} + \boldsymbol{R}_{i}\bar{\boldsymbol{I}}_{i}^{i}\boldsymbol{R}_{i}^{T}\boldsymbol{S}^{T}(\boldsymbol{\omega}_{i})\boldsymbol{\omega}_{i} + \boldsymbol{R}_{i}\bar{\boldsymbol{I}}_{i}^{i}\boldsymbol{R}_{i}^{T}\dot{\boldsymbol{\omega}}_{i} \\ &= \bar{\boldsymbol{I}}_{i}\dot{\boldsymbol{\omega}}_{i} + \boldsymbol{\omega}_{i}\times(\bar{\boldsymbol{I}}_{i}\boldsymbol{\omega}_{i}) \end{split}$$



$$\begin{aligned} \boldsymbol{\mu}_{i} + \boldsymbol{f}_{i} \times \boldsymbol{r}_{i-1,C_{i}} & -\boldsymbol{\mu}_{i+1} - \boldsymbol{f}_{i+1} \times \boldsymbol{r}_{i,C_{i}} = \bar{\boldsymbol{I}}_{i} \dot{\boldsymbol{\omega}}_{i} + \boldsymbol{\omega}_{i} \times (\bar{\boldsymbol{I}}_{i} \boldsymbol{\omega}_{i}) \\ & + k_{r,i+1} \ddot{q}_{i+1} I_{m_{i+1}} \boldsymbol{z}_{m_{i+1}} + k_{r,i+1} \dot{q}_{i+1} I_{m_{i+1}} \boldsymbol{\omega}_{i} \times \boldsymbol{z}_{m_{i+1}} \end{aligned}$$



For studying the mathematical model of the Quad-rotor, we need to make some assumptions for system dynamics. Some of these assumptions are as below.

- Quad-rotor's structure is symmetric around.
- Quad-rotor UAV consists of a rigid frame equipped with four rotors.
- No mass changes during motion (time invariant).
- COG (Center of Gravity) is fixed at the origin of the Quad-rotor.
- Thrust and drag constants are proportional to the square value of the motor's speed.



Kinematics

Linear position and velocity

$$\mathbf{x} = (x, y, z)^{\top}$$

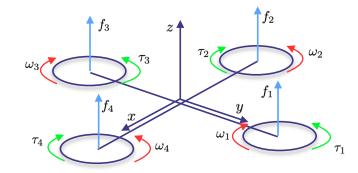
$$\dot{\mathbf{x}} = (\dot{x}, \dot{y}, \dot{z})^{ op}$$

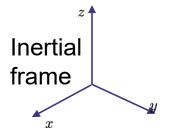
Angular position and velocity

$$\theta = (\phi, \theta, \psi)^\top$$

$$\dot{ heta} = (\dot{\phi},\dot{ heta},\dot{\psi})^{ op}$$

Body-fixed frame







Kinematics

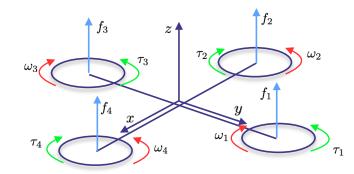
$$egin{aligned} heta &= (\phi, heta, \psi)^{ op} \ \dot{ heta} &= (\dot{\phi}, \dot{ heta}, \dot{\psi})^{ op} \end{aligned}$$

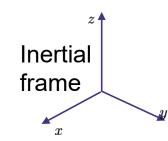
We need to know the angular velocity ω of the body-fixed frame

$$\mathbf{R}(\gamma) = \mathbf{R}_z(\psi)\mathbf{R}_y(heta)\mathbf{R}_x(\phi) = egin{pmatrix} c_\psi c_ heta & c_\psi s_ heta s_\phi - s_\psi c_\phi & c_\psi s_ heta c_\phi + s_\psi s_\phi \ s_\psi c_ heta & s_\psi s_ heta s_\phi + c_\psi c_\phi & s_\psi s_ heta c_\phi - c_\psi s_\phi \ -s_ heta & c_ heta s_\phi & c_ heta c_\phi \end{pmatrix}$$

$$ec{\omega} = egin{bmatrix} 1 & 0 & -s_{ heta} \ 0 & c_{\phi} & c_{ heta}s_{\phi} \ 0 & -s_{\phi} & c_{ heta}egin{bmatrix} ec{ heta} \end{bmatrix} \dot{ec{ heta}} \end{pmatrix}$$

Body-fixed frame







Forces

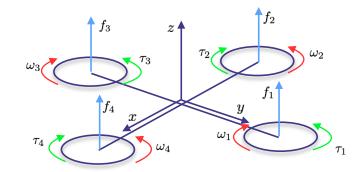
Propeller total trust

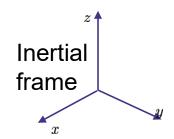
$$F_B = \sum_{i=1}^4 f_i = k egin{bmatrix} 0 \ 0 \ \sum_{i=1}^4 \omega_i^2 \end{bmatrix}$$

Viscous damping

$$F_D = egin{bmatrix} -k_d \dot{x} \ -k_d \dot{y} \ -k_d \dot{z} \end{bmatrix}$$

Body-fixed frame







Torques

Propellers' drag torques

$$\tau_D = \frac{1}{2} R \rho C_D A v^2 = \frac{1}{2} R \rho C_D A (\omega R)^2 = b\omega^2$$

total torque for one propeller around its z-axis:

$$\tau_z = b\omega^2 + I_M \dot{\omega}$$

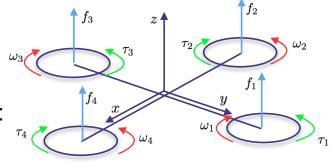
Assumption: $\dot{\omega} \approx 0$

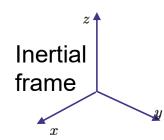
total torque for one propeller around z-axis of the body:

$$\tau_z = (-1)^{i+1} b\omega_i^2$$

$$\tau_z = (-1)^{i+1} b\omega_i^2$$
 $\tau_\psi = b(\omega_1^2 - \omega_2^2 + \omega_3^2 - \omega_4^2)$

Body-fixed frame







Body-fixed frame

Torques

Torque on x-axis of the rigid body

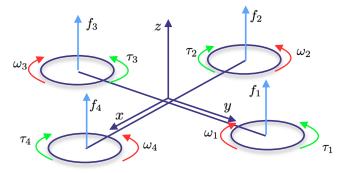
$$\tau_{\phi} = \sum r \times T = L(k\omega_1^2 - k\omega_3^2) = Lk(\omega_1^2 - \omega_3^2)$$

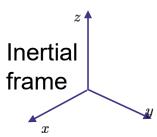
Torque on x-axis of the rigid body

$$\tau_{\theta} = \sum r \times T = L(k\omega_2^2 - k\omega_4^2) = Lk(\omega_2^2 - \omega_4^2)$$

Total body torque

$$au_{B} = egin{bmatrix} Lk(\omega_{1}^{2} - \omega_{3}^{2}) \ Lk(\omega_{2}^{2} - \omega_{4}^{2}) \ b(\omega_{1}^{2} - \omega_{2}^{2} + \omega_{3}^{2} - \omega_{4}^{2}) \end{bmatrix}$$







Body-fixed frame

Equations of motion

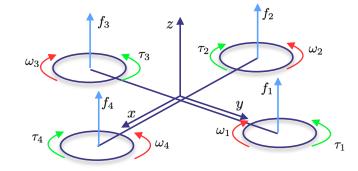
Linear dynamics (in the inertial frame):

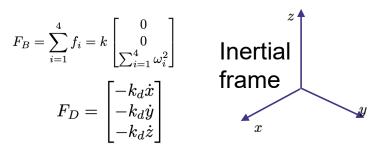
$$\mathbf{m}\ddot{\mathbf{x}} = \begin{bmatrix} 0 \\ 0 \\ -mg \end{bmatrix} + RF_B + F_D$$

Rotational Dynamics (in body-fixed frame)

$$I\dot{\omega} + \omega \times (I\omega) = \tau$$
 $(\tau = \tau_B)$

$$\dot{\omega} = egin{bmatrix} \dot{\omega}_x \ \dot{\omega}_y \ \dot{\omega}_z \end{bmatrix} = I^{-1}(au_B - \omega imes (I\omega))$$





$$\mathbf{m} = egin{bmatrix} m & 0 & 0 \ 0 & m & 0 \ 0 & 0 & m \end{bmatrix} \quad I = egin{bmatrix} I_{xx} & 0 & 0 \ 0 & I_{yy} & 0 \ 0 & 0 & I_{zz} \end{bmatrix}$$



Equations of motion

Rotational Dynamics (in body-fixed frame)

$$\dot{\omega} = egin{bmatrix} au_{\phi} I_{xx}^{-1} \ au_{ au} I_{yy}^{-1} \ au_{ au} I_{zz}^{-1} \end{bmatrix} - egin{bmatrix} rac{I_{yy} - I_{zz}}{I_{xx}} \omega_y \omega_z \ rac{I_{zz} - I_{xx}}{I_{yy}} \omega_x \omega_z \ rac{I_{xx} - I_{yy}}{I_{zz}} \omega_x \omega_y \end{bmatrix}$$

Body-fixed frame

