

VE230 RC slides Week 3

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Overview

Content

- Electrostatics in Free Space
- Coulomb's Law
- Gauss's Law and Application
- Electric Potential

Electrostatics in Free Space

Static electric charges (source) in free space \rightarrow electric field

Electric field intensity

$$\mathbf{E} = \lim_{q \rightarrow 0} \frac{\mathbf{F}}{q} \quad (\text{V/m})$$

Fundamental Postulates of Electrostatics

- Differential form:

$$\vec{\nabla} \cdot \mathbf{E} = \frac{\rho}{\epsilon_0} \quad (\text{divergence})$$

$$\vec{\nabla} \times \mathbf{E} = 0 \quad (\text{curl})$$

where ρ is the volume charge density of free charges (C/m^3), ϵ_0 is the permittivity of free space, a universal constant.

- Integral form:

$$\oint_S \mathbf{E} \cdot d\mathbf{s} = \frac{Q}{\epsilon_0}$$

$$\oint_C \mathbf{E} \cdot d\boldsymbol{\ell} = 0$$

where Q is the total charge contained in volume V bounded by surface S . Also, the scalar line integral of the static electric field intensity around any closed path vanishes.

\mathbf{E} is **not solenoidal** (unless $\rho = 0$), but **irrotational**

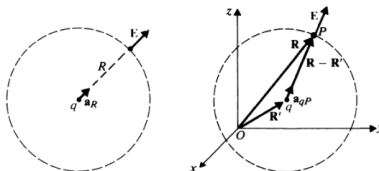
Electric Field due to a System of Discrete Charges

- a single point charge (charge on the origin):

$$\mathbf{E} = \mathbf{a}_R E_R = \mathbf{a}_R \frac{q}{4\pi\epsilon_0 R^2} \quad (\text{V/m})$$

- a single point charge (charge is not on the origin):

$$\mathbf{E}_p = \frac{q (\mathbf{R} - \mathbf{R}')}{4\pi\epsilon_0 |\mathbf{R} - \mathbf{R}'|^3} \quad (\text{V/m})$$



(a) Point charge at the origin.

(b) Point charge not at the origin.

FIGURE 3-2

Electric field due to a point charge.

Electric Field due to a System of Discrete Charges

- **a single point charge (charge is not on the origin):** When a point charge q_2 is placed in the field of another point charge q_1 at the origin, a force F_{12} is experienced by q_2 due to the electric field intensity E_{12} of q_1 at q_2 . Then we have:

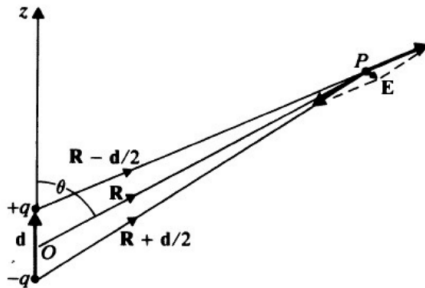
$$\vec{F}_{12} = q_2 \vec{E}_{12} = a_R \frac{q_1 q_2}{4\pi\epsilon_0 R^2}$$

- **several point charges:**

$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \sum_{k=1}^n \frac{q_k (\mathbf{R} - \mathbf{R}'_k)}{|\mathbf{R} - \mathbf{R}'_k|^3}$$

Electric Dipole

■ Electric Field



Electric Dipole

■ Electric Field

general expression:

$$\mathbf{E} = \frac{q}{4\pi\epsilon_0} \left\{ \frac{\mathbf{R} - \frac{\mathbf{d}}{2}}{|\mathbf{R} - \frac{\mathbf{d}}{2}|^3} - \frac{\mathbf{R} + \frac{\mathbf{d}}{2}}{|\mathbf{R} + \frac{\mathbf{d}}{2}|^3} \right\}$$

if $d \ll R$:

$$\mathbf{E} \cong \frac{q}{4\pi\epsilon_0 R^3} \left[3 \frac{\mathbf{R} \cdot \mathbf{d}}{R^2} \mathbf{R} - \mathbf{d} \right]$$

■ Electric Dipole Moment

Definition:

$$\mathbf{p} = q\mathbf{d}$$

,where q is the charge, vector \mathbf{d} goes from $-q$ to $+q$.

$$\mathbf{p} = \mathbf{a}_z p = p (\mathbf{a}_R \cos \theta - \mathbf{a}_\theta \sin \theta)$$

$$\mathbf{R} \cdot \mathbf{p} = R p \cos \theta$$

Electric Dipole

- **Electric Field:** (spherical coordinate)

$$\mathbf{E} = \frac{p}{4\pi\epsilon_0 R^3} (\mathbf{a}_R 2 \cos \theta + \mathbf{a}_\theta \sin \theta) \quad (\text{V/m})$$

Electric Field due to a Continuous Distribution of Charge

■ General Differential Element:

$$d\mathbf{E} = \mathbf{a}_R \frac{\rho dv'}{4\pi\epsilon_0 R^2},$$

where dv' is the differential volume element.

■ Line Charge:

$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \int_{L'} \mathbf{a}_R \frac{\rho_\ell}{R^2} d\ell' \quad (\text{V/m})$$

■ Surface Charge:

$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \int_{S'} \mathbf{a}_R \frac{\rho_s}{R^2} ds' \quad (\text{V/m})$$

■ Volume Charge:

$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \int_{V'} \mathbf{a}_R \frac{\rho}{R^2} dv' \quad (\text{V/m})$$

Gauss's Law and Application

Definition

The total outward flux of the E-field over any closed surface in free space is equal to **the total charge enclosed in the surface** divided by ϵ_0 . (Note that we can choose arbitrary surface S for our convenience.)

$$\oint_S \mathbf{E} \cdot d\mathbf{s} = \frac{Q}{\epsilon_0}$$

Application

■ Conditions for Maxwell's Integral Equations:

There is **a high degree of symmetry** in the charge distribution or in the electrical field (i.e., spherically symmetric, planar, line charge, etc).

Gauss's Law Application

Example:

Determine the electric field intensity of an infinitely long, straight, line charge of a uniform density ρ_s in air.

Gauss's Law Application

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Determine the electric field intensity of an infinitely long, straight, line charge of a uniform density ρ_s in air.

Answer

$$E \cdot 2\pi \cdot r \cdot l = \frac{\rho \cdot l}{\epsilon_0}$$

Several Useful Models

Note: The charge distribution should be **uniform**.

different models	$E(\text{magnitude})$
infinitely long, line charge	$E = \frac{\rho_\ell}{2\pi r \epsilon_0}$
infinite planar charge	$E = \frac{\rho_s}{2\epsilon_0}$
uniform spherical surface charge with radius R	$\begin{cases} E = 0 (r < R) \\ E = \frac{Q}{4\pi r^2 \epsilon_0} (r > R) \end{cases}$
uniform sphere charge with radius R	$\begin{cases} E = \frac{Qr}{4\pi R^3} (r < R) \\ E = \frac{Q}{4\pi r^2 \epsilon_0} (r > R) \end{cases}$
infinitely long, cylindrical charge with radius R	$\begin{cases} E = \frac{\rho_v r}{2\epsilon_0} (r < R) \\ E = \frac{\rho_v R^2}{2r \epsilon_0} (r > R) \end{cases}$

Electric Potential

- **Expression:**

$$\mathbf{E} = -\nabla V$$

the reason for the negative sign: consistent with the convention that in going against the \mathbf{E} field, the electric potential V increases.

- **Electric Potential Difference:**

$$V_2 - V_1 = - \int_{P_1}^{P_2} \mathbf{E} \cdot d\mathbf{l}$$

- **Electric Potential due to a Charge Distribution**

$$V = \frac{q}{4\pi\epsilon_0 R}$$

Electric Potential

■ Line Charge:

$$V = \frac{1}{4\pi\epsilon_0} \int_{L'} \frac{\rho_l}{R} dl' \quad (V)$$

■ Surface Charge:

$$V = \frac{1}{4\pi\epsilon_0} \int_{S'} \frac{\rho_s}{R} ds' \quad (V)$$

■ Volume Charge:

$$V = \frac{1}{4\pi\epsilon_0} \int_{V'} \frac{\rho}{R} dv' \quad (V)$$

Assignment

P.3-22

The polarization in a dielectric cube of side L centered at the origin is given by $\mathbf{P} = P_0(\mathbf{a}_x x + \mathbf{a}_y y + \mathbf{a}_z z)$.

- a) Determine the surface and volume bound-charge densities
- b) Show that the total bound charge is zero.

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- Determine the surface and volume bound-charge densities
- Show that the total bound charge is zero.

$$a) \rho_{ps} = \bar{\mathbf{P}} \cdot \bar{\mathbf{a}}_n = P_0 \frac{L}{2} \text{ on all six faces of the cube.}$$

$$\rho_p = -\bar{\nabla} \cdot \bar{\mathbf{P}} = -3P_0.$$

$$b) Q_s = 6L^2 \rho_{ps} = 3P_0 L^3, \quad Q_v = L^3 \rho_p = -3P_0 L^3.$$

$$\text{Total bound charge} = Q_s + Q_v = 0.$$

Assignment

P.3-23

Determine the electric field intensity at the center of a small spherical cavity cut out of a large block of dielectric in which a polarization \mathbf{P} exists.

Assignment

P.3-23

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1-19 Assume $\bar{\rho} = \bar{a}_z P$. Surface charge density $\rho_{ps} = \bar{\rho} \cdot \bar{a}_n$
 $= (\bar{a}_z P) \cdot (\bar{a}_n)$
 $= P \cos \theta.$



The z-component
 of the electric field
 intensity due to a ring of
 ρ_{ps} contained in width $R d\theta$ at θ is

$$dE_z = \frac{\rho_{ps} \cos \theta}{4\pi\epsilon_0 R^2} (2\pi R \sin \theta) (R d\theta) \cos \theta$$

$$= \frac{P}{2\epsilon_0} \cos^2 \theta \sin \theta d\theta.$$

At the center
 of the cavity : $\bar{E} = \bar{a}_z E_z = \bar{a}_z \frac{P}{2\epsilon_0} \int_0^\pi \cos^2 \theta \sin \theta d\theta = \frac{\bar{P}}{3\epsilon_0}.$

Assignment

P.3-25

Assume that the $z=0$ plane separates two lossless dielectric regions with $\epsilon_{r1} = 2$ and $\epsilon_{r2} = 3$. If we know that \mathbf{E}_1 in region 1 is $\mathbf{a}_x 2y - \mathbf{a}_y 3x + \mathbf{a}_z (5 + z)$, what do we also know about \mathbf{E}_2 and \mathbf{D}_2 in region 2? Can we determine \mathbf{E}_2 and \mathbf{D}_2 at any point in region 2? Explain.

Assignment

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P.3-21 At the $z=0$ plane : $\mathbf{E}_1 = \mathbf{a}_x 2y - \mathbf{a}_y 3x + \mathbf{a}_z 5$.

$$\bar{E}_{1t}(z=0) = \bar{E}_{2t}(z=0) = \mathbf{a}_x 2y - \mathbf{a}_y 3x.$$

$$\bar{D}_{1n}(z=0) = \bar{D}_{2n}(z=0) \rightarrow 2\bar{E}_{1n}(z=0) = 3\bar{E}_{2n}(z=0)$$

$$\rightarrow \bar{E}_{2n}(z=0) = \frac{2}{3}(\mathbf{a}_z 5) = \mathbf{a}_z \frac{10}{3}$$

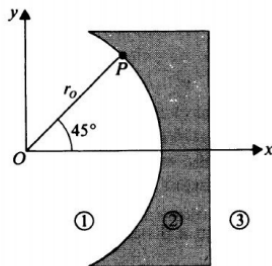
$$\therefore \bar{E}_2(z=0) = \mathbf{a}_x 2y - \mathbf{a}_y 3x + \mathbf{a}_z \frac{10}{3},$$

$$\bar{D}_2(z=0) = (\mathbf{a}_x 6y - \mathbf{a}_y 9x + \mathbf{a}_z 10)\epsilon_0.$$

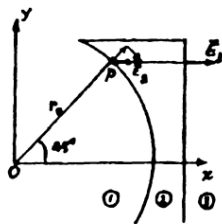
Assignment

P.3-28

Dielectric lenses can be used to collimate electromagnetic fields. In Fig.1 the left surface of the lens is that of a circular cylinder, and the right surface is a plane. If \mathbf{E}_1 at point $P(r_0, 45^\circ, z)$ in region 1 is $a_r 5 - a_\phi 3$, what must be the dielectric constant of the lens in order that \mathbf{E}_3 in region 3 is parallel to the x-axis?



Assignment

P. 3-24

Assume $\vec{E}_2 = \bar{a}_r E_{2r} + \bar{a}_\phi E_{2\phi}$

$$\text{B.C.: } \bar{a}_n \times \vec{E}_1 = \bar{a}_n \times \vec{E}_2 \rightarrow E_{2\phi} = -$$

For \vec{E}_3 , and hence \vec{E}_2 , to be parallel to the x -axis,

$$E_{2\phi} = -E_{2r} \rightarrow E_{2r} = 3.$$

$$\text{B.C.: } \bar{a}_n \cdot \vec{D}_1 = \bar{a}_n \cdot \vec{D}_2 \rightarrow s = 3$$

$$\therefore \epsilon_{r2} = 5/3.$$

Assignment

P.3-32

The radius of the core and the inner radius of the outer conductor of a very long coaxial transmission line are r_i and r_o , respectively. The space between the conductors is filled with two coaxial layers of dielectrics. The dielectric constants of the dielectrics are ϵ_{r1} for $r_i < r < b$ and ϵ_{r2} for $b < r < r_o$. Determine its capacitance per unit length.

Assignment

P.3-32

The radius of the core and the inner radius of the outer conductor of a very long coaxial transmission line are r_i and r_o , respectively. The space between the conductors is filled with two coaxial layers of dielectrics. The dielectric constants of the dielectrics are ϵ_{r1} for $r_i < r < b$ and ϵ_{r2} for $b < r < r_o$. Determine its capacitance per unit length.

$$\begin{aligned}
 \underline{P.3-30} \quad D &= \bar{a}_r \frac{\rho_L}{2\pi r} \cdot \quad E_1 = \bar{a}_r \frac{\rho_L}{2\pi\epsilon_0\epsilon_{r1}r}, r_i < r < b; \quad E_2 = \bar{a}_r \frac{\rho_L}{2\pi\epsilon_0\epsilon_{r2}r}, b < r < r_o \\
 V &= -\int_{r_o}^{r_i} \vec{E} \cdot d\vec{r} = \frac{\rho_L}{2\pi\epsilon_0} \left[\frac{1}{\epsilon_{r1}} \ln\left(\frac{b}{r_i}\right) + \frac{1}{\epsilon_{r2}} \ln\left(\frac{r_o}{b}\right) \right], \\
 C &= \frac{\rho_L}{V} = \frac{2\pi\epsilon_0}{\frac{1}{\epsilon_{r1}} \ln\left(\frac{b}{r_i}\right) + \frac{1}{\epsilon_{r2}} \ln\left(\frac{r_o}{b}\right)} \quad (\text{F/m}).
 \end{aligned}$$