

Detecting Quantum Entanglement In Multiqubit Systems Using Covariance Matrix

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Why talk about characterization using covariance matrix?

① Requirement of CMC:

Devise techniques to detect entangled states present in the system and how they respond to certain operations. For eg.- To assess whether an operation preserves or destroys entanglement.

② How is CMC a better criteria ?

By measuring the quadratures in gaussian states¹ and their variances and covariances, we can build the experimental covariance matrix, encoding information about the means and fluctuations (variances) of the quadratures..

③ Combining with State Tomography:

Quantum state tomography involves reconstructing the entire density matrix of the state, for better characterization.

¹However, we confine our discussion to the discrete case.

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Covariance Matrix

- Let ρ be a pure or mixed quantum state, described by a (positive) density operator in a d -dimensional Hilbert space \mathcal{H} and let $\{M_k : k = 1, \dots, N\}$ a suitable set of observables.

Covariance Matrix

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 - These observables are orthonormal observables with respect to the Hilbert-Schmidt scalar product between observables, i.e., they can be defined using $Tr(M_i M_j)$.

Covariance Matrix

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 - These observables are orthonormal observables with respect to the Hilbert-Schmidt scalar product between observables, i.e., they can be defined using $Tr(M_i M_j)$.
 - As an example for such a set of observables for the case of a single qubit, one can consider the (appropriately normalized) Pauli matrices, $M_1 = \frac{\mathbf{I}}{\sqrt{2}}$, $M_2 = \frac{\sigma_x}{\sqrt{2}}$, $M_3 = \frac{\sigma_y}{\sqrt{2}}$, $M_4 = \frac{\sigma_z}{\sqrt{2}}$.

Bipartite case

- The $d^2 \times d^2$ covariance matrix $\gamma = \gamma(\rho, \{M_k\})$ and the $d^2 \times d^2$ symmetrized covariance matrix (CM) $\gamma^S = \gamma^S(\rho, \{M_k\})$ are defined by their matrix entries as -

$$\gamma_{i,j} = \langle M_i M_j \rangle - \langle M_i \rangle \langle M_j \rangle, \quad \gamma_{i,j}^S = \frac{\langle M_i M_j \rangle + \langle M_j M_i \rangle}{2} - \langle M_i \rangle \langle M_j \rangle. \quad (1)$$

- γ is a complex Hermitian matrix. The matrix γ^S in turn is real and symmetric. Both γ and γ^S are positive semidefinite, i.e., $\gamma, \gamma^S \geq 0$

Block structure

Construction

Let ρ be a state of a bipartite system, and let

$M_k = \{A_k \otimes \mathbb{I}, \mathbb{I} \otimes B_k\}$ be a set of observables and $\{M_k\}$ is a pairwise complete set. Then, the block covariance matrix $\gamma(\rho, \{M_k\})$ has the entries $\gamma_{i,j} = \langle M_i M_j \rangle - \langle M_i \rangle \langle M_j \rangle$ and consequently a block structure:

$$\gamma = \begin{pmatrix} A & C \\ C^T & B \end{pmatrix},$$

where

$$A = \gamma(\rho_A, \{A_k\}) \text{ and } B = \gamma(\rho_B, \{B_k\})$$

are CMs of the reduced states of systems A and B , and

$$C_{i,j} = \langle A_i \otimes B_j \rangle - \langle A_i \rangle \langle B_j \rangle.$$

²Similarly, $\gamma^S\{M_k\}$.

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Main Idea

- We want to exploit the structure of the covariance matrix and see whether if we notice a pattern.
- As we see later, a diagonal block structure for separable states is identified by using the constraints on each of these blocks dependent of the dimension of the quantum state.

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The CM criterion:

Let $\kappa_a = \sum_k p_k |\phi_a\rangle\langle\phi_a|$ and $\kappa_b = \sum_k p_k |\psi_b\rangle\langle\psi_b|$ be spread over the Hilbert spaces H_A and H_B , then it is observed that-

$$\gamma(\rho, \{M_k\}) \geq \kappa_a \oplus \kappa_b, \quad (2)$$

$$\gamma(\rho, \{M_k\}) - t(\kappa_a \oplus \kappa_b) \geq 0, \quad (3)$$

for all covariance matrices over a separable ρ . CMC works regardless of our choice of observables.

$$M_k \rightarrow \tilde{M}_k \equiv U^\dagger M_k U = \sum_I O_{k,I} M_I \quad (4)$$

Some Important Facts about the criterion

- The transformation of observables leave the *eigenvalues invariant.*

$$\gamma(\rho, \{\tilde{M}_k\}) = \mu \gamma(\rho, \{M_k\}) \mu^T \quad (5)$$

- If ρ is a d -dimensional pure state and the $\{M_k\}$ are orthogonal, then $\gamma(\rho_A) = \frac{P}{2}$, where $P^2 = P$ is a projector onto a $2(d - 1)$ dimensional subspace of a total d^2 -dim space. It follows that $\text{Tr}(\kappa_A) = d_A - 1$ and $\text{Tr}(\kappa_B) = d_B - 1$.
- So for the 2- qubit case (i.e. with $\text{dim}(\mathcal{H}_A) = \text{dim}(\mathcal{H}_B) = 2$, we will have $\text{Tr}(\kappa_A) = \text{Tr}(\kappa_B) = 1$. This is a particularly useful result for two qubits when $\{A_k\} = \{B_k\} = \{I, \sigma_x, \sigma_y, \sigma_z\}/\sqrt{2}$, an 8x8 matrix is reduced to a 6x6 matrix since the column and rows corresponding to identity are zero, leading to a $\text{dim} = 3$ value of $\kappa_{A/B}$.

Designing Covariance Matrix for 2 Qubits

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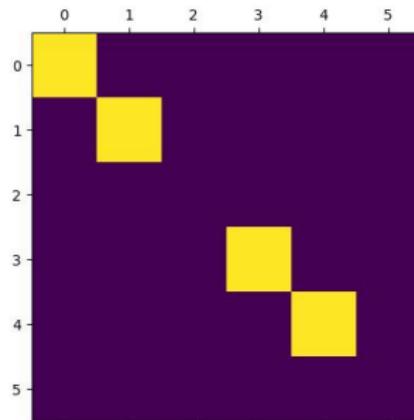
2 Qubit Separable State

- $|\psi\rangle = |01\rangle = \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix}$

Co-variance Matrix :

$$\begin{bmatrix} 0.5 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0.5 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0.5 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0.5 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix}$$

Pictorial Representation of Blocks :



$t = 1$

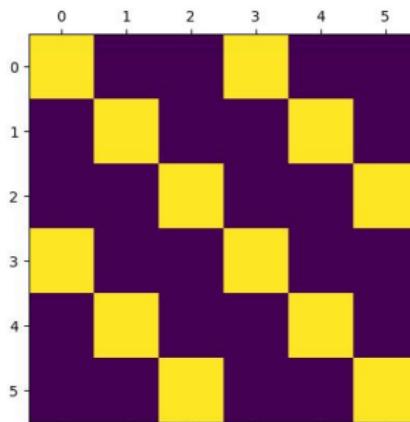
2 Qubit Bell (Werner) State

$$\bullet \quad |\psi\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) = \begin{bmatrix} 0.707 \\ 0 \\ 0 \\ 0.707 \end{bmatrix}$$

Co-variance Matrix :

$$\begin{bmatrix} 0.5 & 0 & 0 & 0.5 & 0 & 0 \\ 0 & 0.5 & 0 & 0 & -0.5 & 0 \\ 0 & 0 & 0.5 & 0 & 0 & 0.5 \\ 0.5 & 0 & 0 & 0.5 & 0 & 0 \\ 0 & -0.5 & 0 & 0 & 0.5 & 0 \\ 0 & 0 & 0.5 & 0 & 0 & 0.5 \end{bmatrix}$$

Pictorial Representation of Blocks :



$$t < 1$$

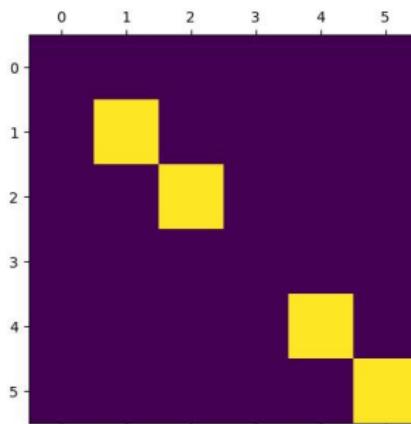
2 Qubit Separable Werner State

- $|\psi\rangle = \frac{1}{2}(|00\rangle + |01\rangle + |10\rangle + |11\rangle) = \begin{bmatrix} 0.5 \\ 0.5 \\ 0.5 \\ 0.5 \end{bmatrix}$

Co-variance Matrix :

$$\begin{bmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0.5 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0.5 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0.5 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0.5 \end{bmatrix}$$

Pictorial Representation of Blocks :



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Three qubit extension

Block structure for 3 qubits

For three qubit, the block structure with $\{A_k\}=\{B_k\}=\{C_k\}=\\ \{\mathbb{I}, \sigma_x, \sigma_y, \sigma_z\}/\sqrt{2}$

$$\gamma(\rho_{ABC}, \{M_k\}) = \begin{pmatrix} A & D & E \\ D^T & B & F \\ E^T & F^T & C \end{pmatrix} \quad (6)$$

The CMC is then given by-

$$\gamma - (\kappa_A \oplus \kappa_B \oplus \kappa_C) \geq 0 \quad (7)$$

- The above equation is modified to $\gamma - t(X_a \oplus X_b \oplus X_c) \geq 0$ and is treated as a feasibility problem.
Hence, $t < 1 \implies$ Entangled state.

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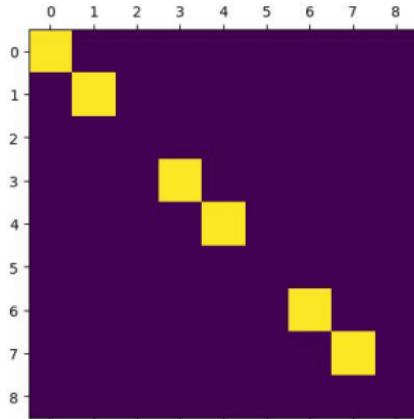
3 Qubit Separable State

- $|\psi\rangle = |011\rangle = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \end{bmatrix}$

Co-variance Matrix :

$$\begin{bmatrix} 0.5 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0.5 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0.5 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0.5 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0.5 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0.5 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix}$$

Pictorial Representation of Blocks :



$t = 1$

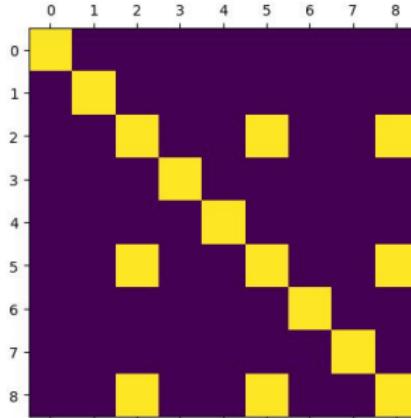
3 Qubit GHZ State

- $|\psi\rangle = \frac{1}{\sqrt{2}}(|000\rangle + |111\rangle) = \begin{bmatrix} 0.707 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0.707 \end{bmatrix}$

Co-variance Matrix :

$$\begin{bmatrix} 0.5 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0.5 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0.5 & 0 & 0 & 0.5 & 0 & 0 & 0.5 \\ 0 & 0 & 0 & 0.5 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0.5 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0.5 & 0 & 0 & 0.5 \\ 0 & 0 & 0.5 & 0 & 0 & 0.5 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0.5 & 0 & 0 \\ 0 & 0 & 0 & 0.5 & 0 & 0 & 0 & 0.5 & 0 \end{bmatrix}$$

Pictorial Representation of Blocks :



$$t < 1$$

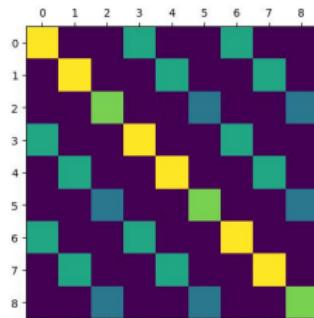
3 Qubit W State

- $|\psi\rangle = \frac{1}{\sqrt{3}}(|100\rangle + |010\rangle + |001\rangle) = \begin{bmatrix} 0 \\ 0.577 \\ 0.577 \\ 0 \\ 0.577 \\ 0 \\ 0 \\ 0 \end{bmatrix}$

Co-variance Matrix :

$$\begin{bmatrix} 0.5 & 0 & 0 & 0.3 & 0 & 0 & 0.3 & 0 & 0 \\ 0 & 0.5 & 0 & 0 & 0.3 & 0 & 0 & 0.3 & 0 \\ 0 & 0 & 0.4 & 0 & 0 & -0.2 & 0 & 0 & -0.2 \\ 0.3 & 0 & 0 & 0.5 & 0 & 0 & 0.3 & 0 & 0 \\ 0 & 0.3 & 0 & 0 & 0.5 & 0 & 0 & 0.3 & 0 \\ 0 & 0 & -0.2 & 0 & 0 & 0.4 & 0 & 0 & -0.2 \\ 0.3 & 0 & 0 & 0.3 & 0 & 0 & 0.5 & 0 & 0 \\ 0 & 0.3 & 0 & 0 & 0.3 & 0 & 0 & 0.5 & 0 \\ 0 & 0 & -0.2 & 0 & 0 & -0.2 & 0 & 0 & 0.4 \end{bmatrix}$$

Pictorial Representation of Blocks :



$t < 1$

3 Qubit Bi separable State

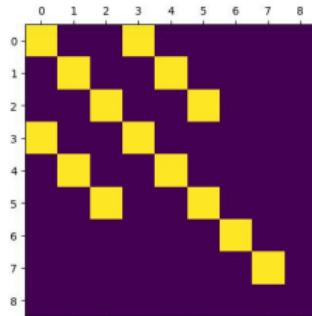
- $|\psi\rangle = \frac{1}{\sqrt{2}}(|000\rangle + |110\rangle) = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) \otimes |0\rangle =$

$$\begin{bmatrix} 0.707 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0.707 \\ 0 \end{bmatrix}$$

Co-variance Matrix :

$$\begin{bmatrix} 0.5 & 0 & 0 & 0.5 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0.5 & 0 & 0 & -0.5 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0.5 & 0 & 0 & 0.5 & 0 & 0 & 0 \\ 0.5 & 0 & 0 & 0.5 & 0 & 0 & 0 & 0 & 0 \\ 0 & -0.5 & 0 & 0 & 0.5 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0.5 & 0 & 0 & 0.5 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0.5 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0.5 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix}$$

Pictorial Representation of Blocks :



$$t < 1$$

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CM for 4-qubits

Consider the Hilbert space $H_A \otimes H_B \otimes H_C \otimes H_D$ where the observable

$M_k = \{A_k \otimes I \otimes I \otimes I, I \otimes B_k \otimes I \otimes I, \dots, I \otimes I \otimes I \otimes D_k\}$, where M_k is a pairwise distinct set and γ is defined as usual . The CM would now be seen as follows-

$$\gamma(\rho_{ABCD}, \{M_k\}) = \begin{pmatrix} A & E & F & G \\ E^T & B & H & I \\ F^T & H^T & C & J \\ G^T & I^T & J^T & D \end{pmatrix}$$

where A,B,C and D are the reduced covariance matrices, with $E_{kl} = \langle A_k B_l \rangle - \langle A_k \rangle \langle B_l \rangle$ and so on for the rest of the blocks.

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CMC for 4-qubits

Formalism for 4-qubits

Let $\kappa_a = \sum_k p_k |\phi_a\rangle\langle\phi_a|$, $\kappa_b = \sum_k p_k |\psi_b\rangle\langle\psi_b|$,
 $\kappa_c = \sum_k p_k |\phi_c\rangle\langle\phi_c|$, $\kappa_d = \sum_k p_k |\psi_d\rangle\langle\psi_d|$ be spread over the Hilbert spaces \mathcal{H}_A , \mathcal{H}_B , \mathcal{H}_C and \mathcal{H}_d , Then it is observed that-

$$\gamma(\rho_{ABCD}, \{M_k\}) \geq \kappa_a \oplus \kappa_b \oplus \kappa_c \oplus \kappa_d$$

- $\gamma(\rho, \{\tilde{M}_k\}) = \mu \gamma(\rho, \{M_k\}) \mu^T$
- $M_k \rightarrow \tilde{M}_k \equiv U^\dagger M_k U = \sum_I O_{k,I} M_I$
- We know that the reduced covariance matrices form a Projector ($P/2$), where P is a projector onto the $2(d-1)$ dim subspace out of total d^2 dimensions.

More about the criterion

For a pure state with the state dimension two (qubit case) we readily see that-

$$\gamma(|a\rangle\langle a|, \{A_k\}) = \frac{P}{2} = \frac{I_3 - |\phi_a\rangle\langle\phi_a|}{2} = \kappa_a$$

Thus allowing us to comment that $\text{Tr}(\kappa_{a/b/c/d}) = 1$.

We restate the problem again as a feasibility one, i.e.,

$$\gamma - t(X_a \oplus X_b \oplus X_c \oplus X_d) \geq 0.$$

The objective of this would be to try and propose a measure of quantification for entanglement.

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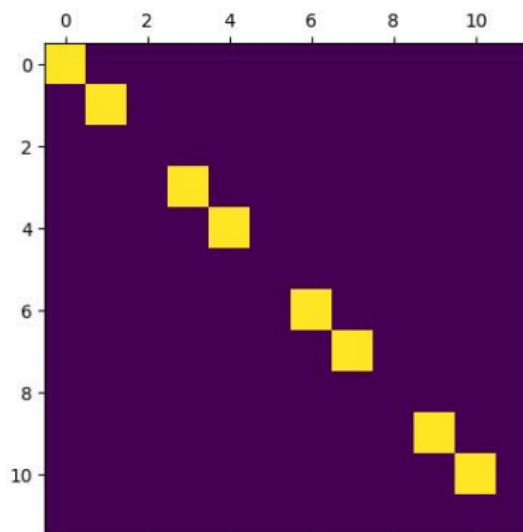
④ References

4 Qubit Separable State

•

$$|\psi\rangle = |0101\rangle = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 0 \\ 1 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \end{bmatrix}$$

Pictorial Representation of Blocks :



$$t = 1$$

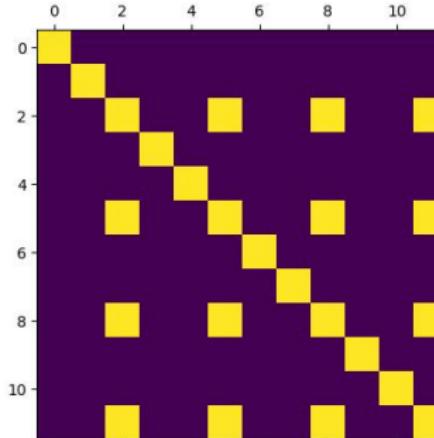
4 Qubit GHZ like State

- 1

$$|\psi\rangle = \frac{1}{\sqrt{2}}(|0000\rangle + |1111\rangle)$$

$$= \boxed{0.707}$$

Pictorial Representation of Blocks :



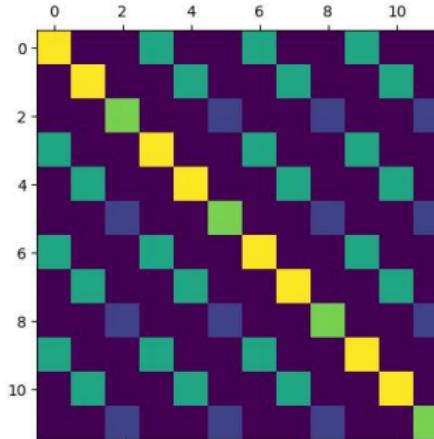
$$t < 1$$

4 Qubit W like State

- $|\psi\rangle = \frac{1}{2}(|1000\rangle + |0100\rangle + |0010\rangle + |0001\rangle)$

$$= \begin{bmatrix} 0 \\ 0.5 \\ 0.5 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \end{bmatrix}$$

Pictorial Representation of Blocks :



$$t < 1$$

4 Qubit W like State

Co-variance Matrix :

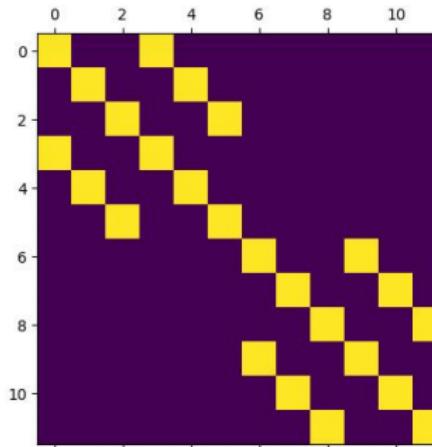
$$\begin{bmatrix} 0.5 & 0 & 0 & 0.3 & 0 & 0 & 0.3 & 0 & 0 & 0.3 & 0 & 0 & 0 \\ 0 & 0.5 & 0 & 0 & 0.3 & 0 & 0 & 0.3 & 0 & 0 & 0.3 & 0 & 0 \\ 0 & 0 & 0.4 & 0 & 0 & -0.1 & 0 & 0 & -0.1 & 0 & 0 & -0.1 & 0 \\ 0.3 & 0 & 0 & 0.5 & 0 & 0 & 0.3 & 0 & 0 & 0.3 & 0 & 0 & 0 \\ 0 & 0.3 & 0 & 0 & 0.5 & 0 & 0 & 0.3 & 0 & 0 & 0.3 & 0 & 0 \\ 0 & 0 & -0.1 & 0 & 0 & 0.4 & 0 & 0 & -0.1 & 0 & 0 & -0.1 & 0 \\ 0.3 & 0 & 0 & 0.3 & 0 & 0 & 0.5 & 0 & 0 & 0.3 & 0 & 0 & 0 \\ 0 & 0.3 & 0 & 0 & 0.3 & 0 & 0 & 0.5 & 0 & 0 & 0.3 & 0 & 0 \\ 0 & 0 & -0.1 & 0 & 0 & -0.1 & 0 & 0 & 0.4 & 0 & 0 & -0.1 & 0 \\ 0.3 & 0 & 0 & 0.3 & 0 & 0 & 0.3 & 0 & 0 & 0.5 & 0 & 0 & 0 \\ 0 & 0.3 & 0 & 0 & 0.3 & 0 & 0 & 0.3 & 0 & 0 & 0.5 & 0 & 0 \\ 0 & 0 & -0.1 & 0 & 0 & -0.1 & 0 & 0 & -0.1 & 0 & 0 & 0.4 & 0 \end{bmatrix}$$

4 Qubit Bi Separable State

- $|\psi\rangle = \frac{1}{2}(|0000\rangle + |0011\rangle + |1100\rangle + |1111\rangle) =$

$$\frac{(|00\rangle + |11\rangle) \otimes (|00\rangle + |11\rangle)}{2} = \begin{bmatrix} 0.5 \\ 0 \\ 0 \\ 0 \\ 0.5 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0.5 \\ 0 \\ 0 \\ 0.5 \end{bmatrix}$$

Pictorial Representation of Blocks :



$$t < 1$$

4 Qubit Bi Separable like State

Co-variance Matrix :

$$\begin{bmatrix} 0.5 & 0 & 0 & 0.5 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0.5 & 0 & 0 & -0.5 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0.5 & 0 & 0 & 0.5 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0.5 & 0 & 0 & 0.5 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & -0.5 & 0 & 0 & 0.5 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0.5 & 0 & 0 & 0.5 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0.5 & 0 & 0 & 0.5 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0.5 & 0 & 0 & -0.5 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0.5 & 0 & 0 & 0.5 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0.5 & 0 & 0 & 0.5 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -0.5 & 0 & 0 & 0.5 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0.5 & 0 & 0 & 0 & 0.5 \end{bmatrix}$$

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Let's look at the feasibility formulation

Remember

$$\gamma - t(\kappa_a \oplus \kappa_b \oplus \kappa_c \oplus \kappa_d) \geq 0$$

Entanglement Parameter

Let ρ be a 4 party quantum state with CM $\gamma(\rho)$. We define a function $V(\rho)$ as

$$V(\rho) = \max_{t, \kappa_A, \kappa_B} \{t \leq 1 : \gamma(\rho) - t\kappa_A \oplus \kappa_B \oplus \kappa_C \oplus \kappa_D \geq 0\}.$$

The entanglement parameter $E(\rho)$ is then defined as

$$E(\rho) = 1 - V(\rho).$$

The entanglement parameter $E(\rho)$ is invariant under local unitary transformations and is convex in the state, that is for

$$\rho = p\rho_1 + (1-p)\rho_2 \text{ we have that } E(\rho) \leq pE(\rho_1) + (1-p)E(\rho_2).$$

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Thanks For Your Attention!