1. [60] RWA too! In class we found the time evolution operator under the Jaynes-Cummings Hamiltonian to be:

$$\widehat{U} = \begin{pmatrix} \cos\left(\frac{\Omega't}{2}\right) - i\frac{\delta}{\Omega'}\sin\left(\frac{\Omega't}{2}\right) & -i\frac{\Omega}{\Omega'}\sin\left(\frac{\Omega't}{2}\right) \\ -i\frac{\Omega}{\Omega'}\sin\left(\frac{\Omega't}{2}\right) & \cos\left(\frac{\Omega't}{2}\right) + i\frac{\delta}{\Omega'}\sin\left(\frac{\Omega't}{2}\right) \end{pmatrix}$$

on the basis of  $|\psi\rangle=a|1,n+1\rangle+b|0,n\rangle$ , where the first label in the ket is the qubit state and the second the harmonic oscillator state. Here  $\Omega = \frac{g}{2}\sqrt{n+1}$ ,  $\delta = \omega - \omega_o$ , and  $\Omega' = \sqrt{\Omega^2 + \delta^2}$ . When deriving this time-evolution operator, we originally started with the Hamiltonian

$$\frac{H}{\hbar} = \frac{\omega_o}{2} \hat{\sigma}_z + \omega \left( \hat{a}^\dagger \, \hat{a} + \frac{1}{2} \right) + \frac{g}{2} \, \hat{\sigma}_x (\hat{a} + \hat{a}^\dagger)$$

before using the RWA to drop the counter rotating terms. Use a numerical solver (e.g QuTIP) to find the evolution under this original Hamiltonian and compare it to that predicted by our time evolution operator for the cases:

a. [15] 
$$|\psi(t=0)\rangle = |0,0\rangle$$
,  $\omega = \omega_0 = 2\pi$  and  $g = \omega/100$ 

b. [15] 
$$|\psi(t=0)\rangle = |0,0\rangle$$
,  $\omega = \omega_0 = 2\pi$  and  $g = \frac{\omega}{2}$ 

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$$|\psi(t=0)\rangle = |0,0\rangle$$
,  $\omega = \omega_o = 2\pi$  and  $g = \frac{\omega}{2}$   
c. [15]  $|\psi(t=0)\rangle = |0,10\rangle$ ,  $\omega = \omega_o = 2\pi$  and  $g = \frac{\omega}{2}$ 

d. [15] 
$$|\psi(t=0)\rangle=|00\rangle$$
,  $\omega=\omega_o=2\pi$  and  $g=\omega$ 

2. [60] Jayne says too! Suppose that two qubits coupled to the same harmonic oscillator can described by the Hamiltonian:

$$\frac{H}{\hbar} = \frac{\omega_o}{2} \hat{\sigma}_z^{(1)} + \frac{\omega_o}{2} \hat{\sigma}_z^{(2)} + \omega \left( \hat{a}^\dagger \, \hat{a} + \frac{1}{2} \right) + \frac{g}{2} \left( (\hat{\sigma}_+^{(1)} + \hat{\sigma}_+^{(2)}) \hat{a} + (\hat{\sigma}_-^{(1)} + \hat{\sigma}_-^{(2)}) \hat{a}^\dagger \right)$$

Implement this Hamiltonian in QuTip and use it to find the time dynamics of the following situations. For simplicity assume  $\omega_o = 2\pi$ , and  $g = \omega_o/100$ . Since you cannot use an infinite dimensioned Hilbert space in QuTip you'll need to truncate the basis at some maximum  $|n\rangle$ . Make sure and choose that maximum n large enough to not affect your answer. For  $\omega = 2\omega_0$  find the time evolution of the population in the initial state for the following initial states over one period of oscillation:

- i.  $[10] | \psi(t=0) \rangle = |110 \rangle$  -- for this problem our ordering in the ket is  $|qubit 1, qubit 2, QHO\rangle$ .
- ii.  $[10] | \psi(t=0) \rangle = |010 \rangle$
- iii.  $[10] | \psi(t=0) \rangle = |100 \rangle$
- iv. [10]  $|\psi(t=0)\rangle = |000\rangle$
- v. [10]  $|\psi(t=0)\rangle = \frac{1}{\sqrt{2}}(|010\rangle + |100\rangle)$
- vi.  $[10] |\psi(t=0)\rangle = |019\rangle$
- vii. [10] Compare your answer in part (vi) to the analytical expression we derived in class. Comment on similarities and differences.

3. [50] Quantum computing runs on effective Hamiltonians. As we saw in class, two qubits can be coupled by a harmonic oscillator to produce an effective Hamiltonian that reproduces their evolution but does not contain the harmonic oscillator. For example, we saw that two qubits coupled by a harmonic oscillator on resonance with the qubit transitions produced an effective Hamiltonian of the form

$$\frac{\hat{H}_{eff}}{\hbar} = \frac{g}{4} \Big( \hat{\sigma}_{x}^{(1)} \hat{\sigma}_{x}^{(2)} + \hat{\sigma}_{y}^{(1)} \hat{\sigma}_{y}^{(2)} \Big).$$

It turns out that using very similar schemes these system can be engineered to produce other effective Hamiltonians. In this problem we will explore the time-evolution due to these effective Hamiltonians. Suppose an effective Hamiltonian is created between two qubits of the form

$$\frac{\widehat{H}_{eff}}{\hbar} = \frac{g}{2}\,\widehat{\sigma}_x^{(1)}\,\widehat{\sigma}_x^{(2)}.$$

- a. [20] Calculate the time evolution operator of this Hamiltonian.
- b. [30] A controlled-NOT gate can be constructed from an XX gate by sandwiching it between some single qubit rotations. The paper:
  - D. Maslov, *New J. Phys.* **19** 023035 (2017) shows one way of doing this in Figure 1 incidentally, this is how lonQ implements a CNOT gate on their hardware. Show that this prescription indeed produces a CNOT gate up to a global phase. A few tips:
    - 1. Remember that to apply the unitaries from left to right means that you operate first with the leftmost operator.
    - 2. Choose  $\frac{gt}{2}$  to reproduce the needed angle  $\chi$  for the XX gate  $\chi$  is defined in the text above Fig. 1.
    - 3. Remember that to apply e.g. a Y rotation to just qubit 1 means you are also applying the identity to qubit 2 and therefore the appropriate operator is the tensor produce of the Y rotation and the identity operator.