Curved Yang-Mills gauge theories based on my preprint arXiv:2210.02924

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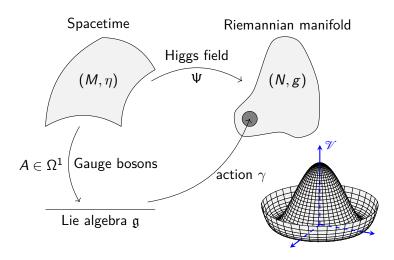
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22 February 2023

Conclusion

- Infinitesimal theory
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- Curved Yang-Mills gauge theory
 - Definition
 - Example
- Conclusion

Infinitesimal curved Yang-Mills-Higgs gauge theory



Guide: Infinitesimal curved Yang-Mills-Higgs gauge theory

Classical formalism	CYMH GT
Lie algebra $\mathfrak g$ as $M \times \mathfrak g$	Lie algebroid $E o N$
${\mathfrak g} ext{-action }\gamma$	Anchor ρ of <i>E</i> & <i>E</i> -connections
Canonical flat connection $ abla^0$ on $M imes \mathfrak{g}$	General connection ∇ on E

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Remarks (Why a "curved theory"?)

Usually, the field strength F is given by (abelian, for simplicity)

$$F := \mathrm{d}A = \mathrm{d}^{\nabla^0}A.$$

 \leadsto We will use a general connection ∇ instead of ∇^0 , and ∇ may not be flat.

Curved Yang-Mills gauge theory

Infinitesimal theory

We will only focus on Yang-Mills theories:

	Classical	Curved
	Lie algebra ${\mathfrak g}$	
Integrated	Lie group <i>G</i>	LGB ² 𝒯



¹LAB = Lie algebra bundle

²LGB = Lie group bundle

Principal bundles based on Lie group bundle actions

Definition (LGB actions, simplified)

$$\begin{array}{c} \mathscr{G} \\ \downarrow \\ \mathscr{P} \stackrel{\pi}{\longrightarrow} M \end{array}$$

 $\mathscr{P} \stackrel{\pi}{\to} M$ a fibre bundle. A **right-action of** \mathscr{G} **on** \mathscr{P} is a smooth map $\mathscr{P} * \mathscr{G} := \pi^* \mathscr{G} = \mathscr{P} \times_M \mathscr{G} \to \mathscr{P}$, $(p,g) \mapsto p \cdot g$, satisfying the following properties:

$$\pi(p \cdot g) = \pi(p),\tag{1}$$

$$(p \cdot g) \cdot h = p \cdot (gh), \tag{2}$$

$$p \cdot e_{\pi(p)} = p \tag{3}$$

for all $p \in \mathscr{P}$ and $g, h \in \mathscr{G}_{\pi(p)}$, where $e_{\pi(p)}$ is the neutral element of $\mathscr{G}_{\pi(p)}$.

Examples

Infinitesimal theory

Example

 \mathscr{G} acts canonically on itself:

$$\mathscr{G} * \mathscr{G} o \mathscr{G}, \ (q,h) \mapsto qh.$$

- Either by $M = \{*\}.$
- Or by $\mathscr{G} \cong M \times G$, then also $\mathscr{P} * \mathscr{G} \cong \mathscr{P} \times G$, and we can define

$$\mathscr{P} \times G \to \mathscr{P},$$

 $(p,g) \mapsto p \cdot g := p \cdot (\pi(p), g),$

which is equivalent to $\mathscr{P} * \mathscr{G} \to \mathscr{P}$.

Principal bundles based on Lie group bundle actions

Examples

Infinitesimal theory

Example

 \mathscr{G} acts canonically on itself:

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Example (Recovering Lie group action)

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Principal bundles based on Lie group bundle actions

Guide to recover the classical theory (Part 1 of 3)



Definition (Principal bundle)

Still a fibre bundle

$$G \longrightarrow \mathscr{P}$$

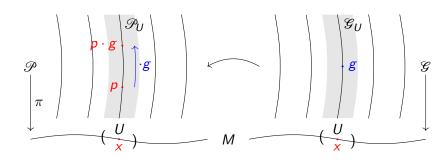
$$\downarrow^{\pi} M$$

but with \mathscr{G} -action

$$\mathscr{P} * \mathscr{G} \to \mathscr{P}$$
 $\mathscr{P} * \mathscr{G}$

simply transitive on fibres of \mathcal{P} , and "suitable" atlas.

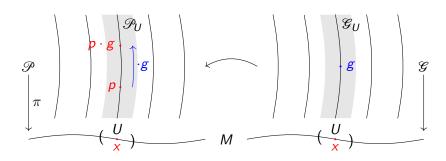
Connection on \mathcal{P} : Idea



But:

$$r_g: \mathscr{P}_X o \mathscr{P}_X$$
 $D_p r_g$ only defined on vertical structure

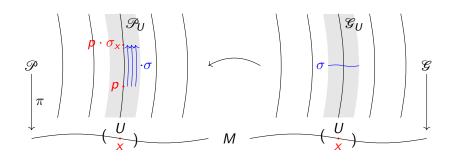
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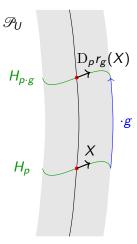
Connection on \mathcal{P} : Idea



Use
$$\sigma \in \Gamma(\mathcal{G})$$
: $r_{\sigma}(p) := p \cdot \sigma_{x}$

Connection on \mathcal{P} : Revisiting the classical setup

If \mathscr{P} a typical principal bundle (\mathscr{G} trivial, $\sigma \equiv g$ constant), and H a connection:

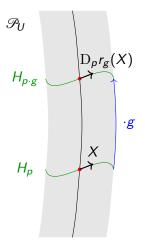


Connections as parallel transport

Infinitesimal theory

Connection on \mathcal{P} : Revisiting the classical setup

If \mathcal{P} a typical principal bundle (\mathcal{G} trivial, $\sigma \equiv g$ constant), and H a connection:



Remarks (Integrated case)

Parallel transport $PT^{\mathscr{P}}_{\alpha}$ in \mathscr{P} :

$$\mathsf{PT}_{\alpha}^{\mathscr{P}}(p \cdot g) = \mathsf{PT}_{\alpha}^{\mathscr{P}}(p) \cdot g$$

where $\alpha: I \to M$ is a base path

Connection on \mathcal{P} : General case

Remarks (Integrated case)

Ansatz: Introduce connection on \mathcal{G} ,

$$\mathsf{PT}^{\mathscr{P}}_{\alpha}(p \cdot g) = \mathsf{PT}^{\mathscr{P}}_{\alpha}(p) \cdot \mathsf{PT}^{\mathscr{G}}_{\alpha}(g).$$

- \bullet $\mathscr{G} \cong M \times G$
- 2 Equip \mathcal{G} with canonical flat connection

Connection on \mathcal{P} : General case

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$$\mathsf{PT}_{\alpha}^{\mathscr{P}}(p \cdot g) = \mathsf{PT}_{\alpha}^{\mathscr{P}}(p) \cdot \mathsf{PT}_{\alpha}^{\mathscr{G}}(g).$$

Guide to recover the classical theory (Part 2 of 3)

- $\mathfrak{G} \cong M \times G$
- 2 Equip \mathcal{G} with canonical flat connection

Basic notions

Infinitesimal theory

Classical situation: Differential of Lie group action

Remarks (Lie group G situation with Lie algebra \mathfrak{g})

In the case of a right G-action on \mathcal{P} , $\Phi: \mathcal{P} \times G \to \mathcal{P}$, we have

$$D_{(p,g)}\Phi(X,Y) = D_p r_g(X) + \overline{(\mu_G)_g(Y)}\Big|_{p,g}$$

for all $p \in \mathcal{P}$, $g \in G$, $X \in T_p \mathcal{P}$ and $Y \in T_g G$, where

- $\overline{\nu}$ denotes the fundamental vector field on \mathscr{P} of $\nu \in \mathfrak{g}$,
- μ_G is the Maurer-Cartan form of G.

Basic notions

Infinitesimal theory

Definition (Fundamental vector fields)

Fundamental vector fields defined by

$$\overline{\nu}_{p} := \frac{\mathrm{d}}{\mathrm{d}t}\Big|_{t=0} (p \cdot \mathrm{e}^{t \nu_{x}})$$

for all $\nu \in \Gamma(q)$ and $p \in \mathcal{P}_x$, where q is the LAB of \mathcal{G} .

Connection

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For $\sigma \in \Gamma(\mathscr{G})$ we define the **Darboux derivative** $\Delta \sigma \in \Omega^1(M; \mathscr{Q})$

$$\Delta \sigma = \sigma^! \mu_{\mathscr{C}},$$

where $\mu_{\mathscr{C}}$ is given by

$$(\mu_{\mathscr{G}})_{\mathsf{g}} := \mathrm{D}_{\mathsf{g}} \mathsf{L}_{\mathsf{g}^{-1}} \circ \pi^{\mathsf{v}},$$

 π^{ν} the projection onto the vertical bundle.

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Remarks

If \mathscr{G} a trivial LGB with canonical flat connection and if Lie group additionally a matrix group, then

$$\Delta \sigma = \sigma^{-1} d\sigma$$
.

Basic notions

Infinitesimal theory

Proposition (Differential of LGB action Φ , [S.-R. F.])

We have

$$\mathrm{D}_{(p,g)}\Phi(X,Y)=\mathrm{D}_p r_\sigma(X)-\left.\overline{(\pi^!\Delta\sigma)|_p(X)}\right|_{p\cdot g}+\left.\overline{(\mu_\mathscr{C})_g(Y)}\right|_{p\cdot g}$$

for all $(p,g) \in \mathscr{P}_X \times \mathscr{G}_X$, $(X,Y) \in T_{(p,g)}(\mathscr{P} * \mathscr{G})$, where σ is any section of \mathscr{G} with $\sigma_X = g$.

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for all $(p,g) \in \mathscr{P}_{\mathsf{X}} \times \mathscr{G}_{\mathsf{X}}$, $(X,Y) \in \mathrm{T}_{(p,g)}(\mathscr{P} * \mathscr{G})$, where σ is any section of \mathscr{G} with $\sigma_{\mathsf{x}} = \mathsf{g}$.

Definition (Modified right-pushforward, [S.-R. F.])

Define

$$r_{g*}(X) := \mathrm{D}_p r_{\sigma}(X) - \left. \overline{(\pi^! \Delta \sigma)|_p(X)} \right|_{p \cdot g}$$

Basic notions

Proposition (Well-defined isomorphism, [S.-R. F.])

We have that

$$T\mathscr{P}|_{\mathscr{P}_{\!\scriptscriptstyle X}} o T\mathscr{P}|_{\mathscr{P}_{\!\scriptscriptstyle X}}, \ X \mapsto r_{g*}(X),$$

is a well-defined automorphism over rg.

Definition (Ehresmann connection, [S.-R. F.])

H a horizontal distribution of $T\mathscr{P}$ with

$$\mathscr{V}_{g*}(H_p) = H_{p \cdot g}$$

Definitions

Infinitesimal theory

Definition (Ehresmann connection, [S.-R. F.])

H a horizontal distribution of $T\mathscr{P}$ with

$$\gamma_{g*}(H_p) = H_{p\cdot g}$$

Definition (Connection 1-form, [S.-R. F.])

 $A \in \Omega^1(\mathscr{P}; \pi^*_{\mathscr{Q}})$ with

$$A(\overline{\nu}) = \pi^* \nu,$$

 $\mathcal{P}_{\sigma}^! A = \mathrm{Ad}_{\sigma^{-1}} \circ A$

for all $\sigma \in \Gamma(\mathcal{G})$ and $\nu \in \Gamma(\mathcal{Q})$.

Remarks

$$\left(\mathscr{r}_{\sigma}^{!}A\right)_{p}(X)=A_{p\sigma_{X}}\left(\mathscr{r}_{\sigma*}(X)\right).$$

Theorem (Equivalence of both definitions, [S.-R. F.])

There is the usual 1:1 correspondence between both definitions:

Given H, define A by

$$A_p(\overline{\nu}_p + X) := (\pi^* \nu)_p,$$

where $X \in H_p$.

Given A, define H by

$$H_p := \operatorname{Ker}(A_p).$$

Theorem (Gauge transformation, [S.-R. F.])

Let s_i , s_i be two sections of \mathcal{P} over U_i and U_i , respectively, which are open subsets of M with $U_i \cap U_i \neq \emptyset$. Then over $U_i \cap U_i$

$$A_{s_i} = \operatorname{Ad}_{\sigma_{ji}^{-1}} \circ A_{s_j} + \Delta \sigma_{ji},$$

where $A_{s_i} := s_i! A$ and σ_{ii} a section of \mathscr{G} with $s_i = s_i \cdot \sigma_{ii}$.

Infinitesimal theory
OO

Compatibility conditions

Proposition (Connection on q, [S.-R. F.])

We have an induced vector bundle connection on g given by

$$\nabla^{\mathscr{G}}\nu := \left.\frac{\mathrm{d}}{\mathrm{d}t}\right|_{t=0} \Delta \mathrm{e}^{t\nu}.$$

Compatibility conditions

Infinitesimal theory

Remarks

Recall, \mathcal{G} a principal \mathcal{G} -bundle.

Definition (Compatibility conditions, [S.-R. F.])

 $\mu_{\mathscr{C}}$ a Yang-Mills connection (w.r.t. $\zeta \in \Omega^2(M; q)$) if it satisfies the **compatibility conditions**:

- **1** $\mu_{\mathscr{G}}$ a connection 1-form on $\mathscr{G} \stackrel{\pi_{\mathscr{G}}}{\to} M$:
- **2** $\mu_{\mathscr{C}}$ satisfies the **generalised Maurer-Cartan equation**

$$\left. \left(\mathrm{d}^{\pi_{\mathscr{G}}^* \nabla^{\mathscr{G}}} \mu_{\mathscr{G}} + \frac{1}{2} [\mu_{\mathscr{G}} \stackrel{\wedge}{,} \mu_{\mathscr{G}}]_{\pi_{\mathscr{G}}^* \mathscr{G}} \right) \right|_{\mathcal{G}} = \mathrm{Ad}_{g^{-1}} \circ \left. \pi_{\mathscr{G}}^! \zeta \right|_{\mathcal{G}} - \left. \pi_{\mathscr{G}}^! \zeta \right|_{\mathcal{G}}$$

Infinitesimal theory Compatibility conditions

Proposition ($\nabla^{\mathcal{G}}$ a Lie bracket derivation)

Let $\mu_{\mathscr{C}}$ be a connection 1-form on \mathscr{G} , then

$$\nabla^{\mathscr{G}}\left(\left[\mu,\nu\right]_{\mathscr{G}}\right) = \left[\nabla^{\mathscr{G}}\mu,\nu\right]_{\mathscr{G}} + \left[\mu,\nabla^{\mathscr{G}}\nu\right]_{\mathscr{G}}.$$

Compatibility conditions

Infinitesimal theory

Proposition ($\nabla^{\mathcal{G}}$ a Lie bracket derivation)

Let $\mu_{\mathscr{C}}$ be a connection 1-form on \mathscr{C} , then

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Theorem (Curvature of LAB connection exact, [S.-R. F.])

 $\mu_{\mathscr{C}}$ satisfies the generalized Maurer-Cartan equation w.r.t. ζ if and only if

$$R_{\nabla^{\mathscr{G}}} = \mathrm{ad} \circ \zeta.$$

Remarks

There is a simplicial differential δ on $\mathscr{G} \stackrel{\pi_{\mathscr{G}}}{\to} M$

$$\delta: \Omega^{\bullet}(\underbrace{\mathcal{G} * \ldots * \mathcal{G}}_{k \text{ times}}; \pi_{\mathscr{C}}^{*}g) \to \Omega^{\bullet}(\underbrace{\mathcal{G} * \ldots * \mathcal{G}}_{k+1 \text{ times}}; \pi_{\mathscr{C}}^{*}g)$$

such that the compatibility conditions are equivalent to

$$\delta\mu_{\mathscr{G}} = 0$$
 and $\mathrm{d}^{\pi_{\mathscr{G}}^* \nabla^{\mathscr{G}}} \mu_{\mathscr{G}} + \frac{1}{2} [\mu_{\mathscr{G}} \stackrel{\wedge}{,} \mu_{\mathscr{G}}]_{\pi_{\mathscr{G}}^* \mathscr{G}} = \delta\zeta.$

- \bigcirc $\mathcal{G}\cong M\times G$
- 2 Equip \mathcal{G} with canonical flat connection
- \bigcirc $\zeta \equiv 0$

Curved Yang-Mills gauge theory

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Guide to recover the classical theory (Part 3 of 3)

- $\mathfrak{G}\cong M\times G$
- 2 Equip \mathscr{G} with canonical flat connection
- $0 \ \zeta \equiv 0$

Given a Yang-Mills connection on \mathcal{G} :

Definition (Generalized curvature/field strength F of A, [S.-R. F.])

 π^H denotes the projection onto $H\subset T\mathscr{P}$, then we define

$$F := \mathrm{d}^{\pi^* \nabla^{\mathscr{G}}} A \circ \left(\pi^{\mathrm{H}}, \pi^{\mathrm{H}}\right) + \pi^! \zeta.$$

Theorem (Structure equation, [S.-R. F.])

$$F = \mathrm{d}^{\pi^* \nabla^{\mathcal{G}}} A + \frac{1}{2} [A \wedge A]_{\pi^* \mathcal{G}} + \pi^! \zeta.$$

Definition and properties

Proposition (Properties of F, [S.-R. F.])

- $F(X, \cdot) = 0$, if X vertical,
- $\bullet \ r_{\sigma}^! F = \mathrm{Ad}_{\sigma^{-1}} \circ F.$

Curved Yang-Mills gauge theory

Infinitesimal theory

Proposition (Properties of F, [S.-R. F.])

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$$F_{s_i} = \operatorname{Ad}_{\sigma_{ii}^{-1}} \circ F_{s_j},$$

where $F_{s_i} := s_i^! F$ and σ_{ii} a section of $\mathscr G$ with $s_i = s_i \cdot \sigma_{ii}$.

Theorem (Lagrangian, [S.-R. F.])

- κ be an Ad-invariant fibre metric on q,
- M a spacetime, and * its Hodge star operator,
- (*U_i*); open covering of M with subordinate gauges $s_i \in \Gamma(\mathscr{P}|_{II})$.

Then the Lagrangian $\mathfrak{L}_{CYM}[A]$, defined locally by

$$(\mathfrak{L}_{\mathrm{CYM}}[A])\big|_{U_i} := -\frac{1}{2}\kappa(F_{s_i} \stackrel{\wedge}{,} *F_{s_i}),$$

is well-defined, and

$$\mathfrak{L}_{\mathrm{CYM}}[L^!A] = \mathfrak{L}_{\mathrm{CYM}}[A]$$

for all principal bundle automorphisms L.

Example (Hopf fibration $\mathbb{S}^7 \to \mathbb{S}^4$, [S.-R. F.])

Let P be the Hopf bundle

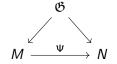
$$\mathrm{SU}(2)\cong \mathbb{S}^3\longrightarrow \mathbb{S}^7$$

Define $\mathscr{P} := \mathscr{G}$ as the inner group bundle of P,

$$\mathscr{G} := c_{\mathrm{SU}(2)}(P) := (P \times \mathrm{SU}(2)) / \mathrm{SU}(2).$$

This principal $c_{SU(2)}(P)$ -bundle admits the structure as curved Yang-Mills gauge theory; there is no description as flat theory.

Gauge theory	Structure
Yang-Mills	Lie group <i>G</i>
Curved Yang-Mills	Lie group bundle ${\mathscr G}$
Curved Yang-Mills-Higgs	Lie groupoid &?



Remarks

Infinitesimal theory

- Richer set of principal bundles, containing Lie groupoids equipped with "non-flat Maurer-Cartan forms".
- Principal bundle for the whole of Yang-Mills-Higgs theory
- Even if \mathscr{G} is trivial, what happens if its connection is not flat?
- May result into obstruction statements for curved Yang-Mills-Higgs gauge theories.

Thank you!