# Curved Yang-Mills gauge theories based on my preprint arXiv:2210.02924

Simon-Raphael Fischer

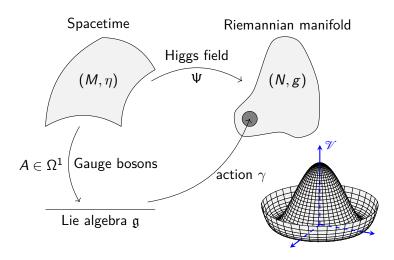
國家理論科學研究中心

30 December 2022

Conclusion

- Infinitesimal theory
- Integration: Ansatz
  - Principal bundles based on Lie group bundle actions
  - Connections as parallel transport
- Connection
  - Basic notions
  - Definitions
  - Gauge transformation
- 4 Curvature
  - Compatibility conditions
  - Definition and properties
- Curved Yang-Mills gauge theory
  - Definition
  - Example
- Conclusion

# **Infinitesimal** curved Yang-Mills-Higgs gauge theory



# Guide: Infinitesimal curved Yang-Mills-Higgs gauge theory

Classical formalism	CYMH GT
Lie algebra $\mathfrak g$ as $M \times \mathfrak g$	Lie algebroid $E o N$
${\mathfrak g} ext{-action }\gamma$	Anchor $\rho$ of <i>E</i> & <i>E</i> -connections
Canonical flat connection $ abla^0$ on $M  imes \mathfrak{g}$	General connection $\nabla$ on $E$

# Guide: Infinitesimal curved Yang-Mills-Higgs gauge theory

Classical formalism	CYMH GT
Lie algebra $\mathfrak g$ as $M \times \mathfrak g$	Lie algebroid $E o N$
${\mathfrak g} ext{-action }\gamma$	Anchor $\rho$ of <i>E</i> & <i>E</i> -connections
Canonical flat connection $ abla^0$ on $M  imes \mathfrak{g}$	General connection $\nabla$ on $E$

#### Remarks (Why a "curved theory"?)

Usually, the field strength F is given by (abelian, for simplicity)

$$F := \mathrm{d}A = \mathrm{d}^{\nabla^0}A.$$

 $\leadsto$  We will use a general connection  $\nabla$  instead of  $\nabla^0$ , and  $\nabla$  may not be flat.

Curved Yang-Mills gauge theory

Infinitesimal theory

#### We will only focus on Yang-Mills theories:

	Classical	Curved
	Lie algebra ${\mathfrak g}$	
Integrated	Lie group <i>G</i>	LGB <sup>2</sup> 𝒯



<sup>&</sup>lt;sup>1</sup>LAB = Lie algebra bundle

<sup>&</sup>lt;sup>2</sup>LGB = Lie group bundle

Principal bundles based on Lie group bundle actions

#### Definition (LGB actions, simplified)

$$\begin{array}{c} \mathscr{G} \\ \downarrow \\ \mathscr{P} \stackrel{\pi}{\longrightarrow} M \end{array}$$

 $\mathscr{P} \stackrel{\pi}{\to} M$  a fibre bundle. A **right-action of**  $\mathscr{G}$  **on**  $\mathscr{P}$  is a smooth map  $\mathscr{P} * \mathscr{G} := \pi^* \mathscr{G} = \mathscr{P} \times_M \mathscr{G} \to \mathscr{P}$ ,  $(p,g) \mapsto p \cdot g$ , satisfying the following properties:

$$\pi(p \cdot g) = \pi(p),\tag{1}$$

$$(p \cdot g) \cdot h = p \cdot (gh), \tag{2}$$

$$p \cdot e_{\pi(p)} = p \tag{3}$$

for all  $p \in \mathscr{P}$  and  $g, h \in \mathscr{G}_{\pi(p)}$ , where  $e_{\pi(p)}$  is the neutral element of  $\mathscr{G}_{\pi(p)}$ .

# Examples

Infinitesimal theory

#### Example

 $\mathscr{G}$  acts canonically on itself:

$$\mathscr{G} * \mathscr{G} o \mathscr{G}, \ (q,h) \mapsto qh.$$

- Either by  $M = \{*\}.$
- Or by  $\mathscr{G} \cong M \times G$ , then also  $\mathscr{P} * \mathscr{G} \cong \mathscr{P} \times G$ , and we can define

$$\mathscr{P} \times G \to \mathscr{P},$$
  
 $(p,g) \mapsto p \cdot g := p \cdot (\pi(p), g),$ 

which is equivalent to  $\mathscr{P} * \mathscr{G} \to \mathscr{P}$ .

Principal bundles based on Lie group bundle actions

# Examples

Infinitesimal theory

#### Example

 $\mathscr{G}$  acts canonically on itself:

$$\mathscr{G} * \mathscr{G} \to \mathscr{G},$$
  
 $(q,h) \mapsto qh.$ 

#### Example (Recovering Lie group action)

- Either by  $M = \{*\}$ .
- Or by  $\mathscr{G} \cong M \times G$ , then also  $\mathscr{P} * \mathscr{G} \cong \mathscr{P} \times G$ , and we can define

$$\mathscr{P} \times G \to \mathscr{P},$$
  
 $(p,g) \mapsto p \cdot g := p \cdot (\pi(p), g),$ 

which is equivalent to  $\mathscr{P} * \mathscr{G} \to \mathscr{P}$ .

Principal bundles based on Lie group bundle actions

Guide to recover the classical theory (Part 1 of 3)



### Definition (Principal bundle)

Still a fibre bundle

$$G \longrightarrow \mathscr{P}$$

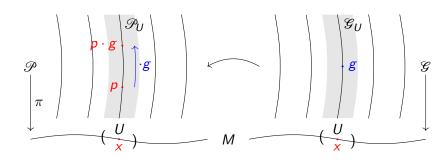
$$\downarrow^{\pi} M$$

but with  $\mathscr{G}$ -action

$$\mathscr{P} * \mathscr{G} \to \mathscr{P}$$
 $\mathscr{P} * \mathscr{G}$ 

simply transitive on fibres of  $\mathcal{P}$ , and "suitable" atlas.

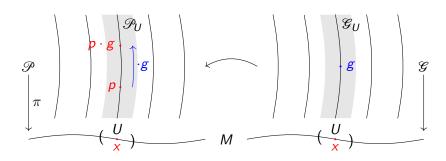
# Connection on $\mathcal{P}$ : Idea



But:

$$r_g: \mathscr{P}_X o \mathscr{P}_X$$
  $D_p r_g$  only defined on vertical structure

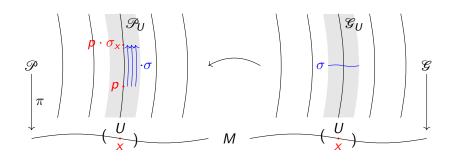
# Connection on $\mathcal{P}$ : Idea



But:

$$r_g:\mathscr{P}_{\mathsf{X}} o\mathscr{P}_{\mathsf{X}}$$
  $\mathrm{D}_p r_g$  only defined on vertical structure

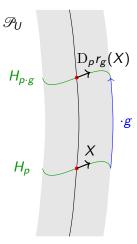
#### Connection on $\mathcal{P}$ : Idea



Use 
$$\sigma \in \Gamma(\mathcal{G})$$
:  $r_{\sigma}(p) := p \cdot \sigma_{x}$ 

# Connection on $\mathcal{P}$ : Revisiting the classical setup

If  $\mathscr{P}$  a typical principal bundle ( $\mathscr{G}$  trivial,  $\sigma \equiv g$  constant), and H a connection:

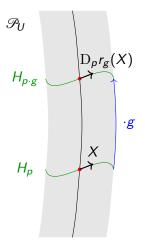


Connections as parallel transport

Infinitesimal theory

# Connection on $\mathcal{P}$ : Revisiting the classical setup

If  $\mathcal{P}$  a typical principal bundle ( $\mathcal{G}$  trivial,  $\sigma \equiv g$  constant), and H a connection:



#### Remarks (Integrated case)

Parallel transport  $PT^{\mathscr{P}}_{\alpha}$  in  $\mathscr{P}$ :

$$\mathsf{PT}_{\alpha}^{\mathscr{P}}(p \cdot g) = \mathsf{PT}_{\alpha}^{\mathscr{P}}(p) \cdot g$$

where  $\alpha: I \to M$  is a base path

#### Connection on $\mathscr{P}$ : General case

#### Remarks (Integrated case)

Ansatz: Introduce connection on  $\mathcal{G}$ ,

$$\mathsf{PT}^{\mathscr{P}}_{\alpha}(p \cdot g) = \mathsf{PT}^{\mathscr{P}}_{\alpha}(p) \cdot \mathsf{PT}^{\mathscr{G}}_{\alpha}(g).$$

- $\bullet$   $\mathscr{G} \cong M \times G$
- 2 Equip  $\mathcal{G}$  with canonical flat connection

#### Connection on $\mathcal{P}$ : General case

#### Remarks (Integrated case)

Ansatz: Introduce connection on  $\mathcal{G}$ ,

$$\mathsf{PT}_{\alpha}^{\mathscr{P}}(p \cdot g) = \mathsf{PT}_{\alpha}^{\mathscr{P}}(p) \cdot \mathsf{PT}_{\alpha}^{\mathscr{G}}(g).$$

Guide to recover the classical theory (Part 2 of 3)

- $\mathfrak{G} \cong M \times G$
- 2 Equip  $\mathcal{G}$  with canonical flat connection

Basic notions

Infinitesimal theory

# Classical situation: Differential of Lie group action

#### Remarks (Lie group G situation with Lie algebra $\mathfrak{g}$ )

In the case of a right G-action on  $\mathcal{P}$ ,  $\Phi: \mathcal{P} \times G \to \mathcal{P}$ , we have

$$D_{(p,g)}\Phi(X,Y) = D_p r_g(X) + \overline{(\mu_G)_g(Y)}\Big|_{p,g}$$

for all  $p \in \mathcal{P}$ ,  $g \in G$ ,  $X \in T_p \mathcal{P}$  and  $Y \in T_g G$ , where

- $\overline{\nu}$  denotes the fundamental vector field on  $\mathscr{P}$  of  $\nu \in \mathfrak{g}$ ,
- $\mu_G$  is the Maurer-Cartan form of G.

Basic notions

Infinitesimal theory

#### Definition (Fundamental vector fields)

Fundamental vector fields defined by

$$\overline{\nu}_{p} := \frac{\mathrm{d}}{\mathrm{d}t}\Big|_{t=0} (p \cdot \mathrm{e}^{t \nu_{x}})$$

for all  $\nu \in \Gamma(q)$  and  $p \in \mathcal{P}_x$ , where q is the LAB of  $\mathcal{G}$ .

Connection

00000000

For  $\sigma \in \Gamma(\mathscr{G})$  we define the **Darboux derivative**  $\Delta \sigma \in \Omega^1(M; \mathscr{Q})$ 

$$\Delta \sigma = \sigma^! \mu_{\mathscr{C}},$$

where  $\mu_{\mathscr{C}}$  is given by

$$(\mu_{\mathscr{G}})_{\mathsf{g}} := \mathrm{D}_{\mathsf{g}} \mathsf{L}_{\mathsf{g}^{-1}} \circ \pi^{\mathsf{v}},$$

 $\pi^{\nu}$  the projection onto the vertical bundle.

#### Definition (Darboux derivative)

For  $\sigma \in \Gamma(\mathcal{G})$  we define the **Darboux derivative**  $\Delta \sigma \in \Omega^1(M; q)$ 

$$\Delta \sigma = \sigma^! \mu_{\mathscr{C}},$$

where  $\mu_{\mathscr{C}}$  is given by

$$(\mu_{\mathscr{C}})_{\mathsf{g}} := \mathrm{D}_{\mathsf{g}} L_{\mathsf{g}^{-1}} \circ \pi^{\mathsf{v}},$$

 $\pi^{\nu}$  the projection onto the vertical bundle.

#### Remarks

If  $\mathscr{G}$  a trivial LGB with canonical flat connection and if Lie group additionally a matrix group, then

$$\Delta \sigma = \sigma^{-1} d\sigma$$
.

Basic notions

Infinitesimal theory

#### Proposition (Differential of LGB action $\Phi$ , [S.-R. F.])

We have

$$\mathrm{D}_{(p,g)}\Phi(X,Y)=\mathrm{D}_p r_\sigma(X)-\left.\overline{(\pi^!\Delta\sigma)|_p(X)}\right|_{p\cdot g}+\left.\overline{(\mu_\mathscr{C})_g(Y)}\right|_{p\cdot g}$$

for all  $(p,g) \in \mathscr{P}_X \times \mathscr{G}_X$ ,  $(X,Y) \in T_{(p,g)}(\mathscr{P} * \mathscr{G})$ , where  $\sigma$  is any section of  $\mathscr{G}$  with  $\sigma_X = g$ .

#### Proposition (Differential of LGB action $\Phi$ , [S.-R. F.])

We have

$$\mathrm{D}_{(p,g)}\Phi(X,Y)=\mathrm{D}_p r_\sigma(X)-\left.\overline{(\pi^!\Delta\sigma)|_p(X)}\right|_{p\cdot g}+\left.\overline{(\mu_\mathscr{E})_g(Y)}\right|_{p\cdot g}$$

for all  $(p,g) \in \mathscr{P}_{\mathsf{X}} \times \mathscr{G}_{\mathsf{X}}$ ,  $(X,Y) \in \mathrm{T}_{(p,g)}(\mathscr{P} * \mathscr{G})$ , where  $\sigma$  is any section of  $\mathscr{G}$  with  $\sigma_{\mathsf{x}} = \mathsf{g}$ .

#### Definition (Modified right-pushforward, [S.-R. F.])

Define

$$r_{g*}(X) := \mathrm{D}_p r_{\sigma}(X) - \left. \overline{(\pi^! \Delta \sigma)|_p(X)} \right|_{p \cdot g}$$

Basic notions

#### Proposition (Well-defined isomorphism, [S.-R. F.])

We have that

$$T\mathscr{P}|_{\mathscr{P}_{\!\scriptscriptstyle X}} o T\mathscr{P}|_{\mathscr{P}_{\!\scriptscriptstyle X}}, \ X \mapsto r_{g*}(X),$$

is a well-defined automorphism over rg.

# Definition (Ehresmann connection, [S.-R. F.])

H a horizontal distribution of  $T\mathscr{P}$  with

$$\mathscr{V}_{g*}(H_p) = H_{p \cdot g}$$

#### Definition (Ehresmann connection, [S.-R. F.])

H a horizontal distribution of  $T\mathscr{P}$  with

$$\mathscr{V}_{g*}(H_p) = H_{p\cdot g}$$

### Definition (Connection 1-form, [S.-R. F.])

 $A \in \Omega^1(\mathcal{P}; \pi^*_{\mathcal{Q}})$  with

$$A(\overline{\nu}) = \pi^* \nu,$$
  
 $\mathcal{P}_{\sigma}^! A = \mathrm{Ad}_{\sigma^{-1}} \circ A$ 

for all  $\sigma \in \Gamma(\mathcal{G})$  and  $\nu \in \Gamma(\mathcal{Q})$ .

#### Remarks

$$\left(r_{\sigma}^{!}A\right)_{\rho}(X)=A_{\rho\sigma_{X}}\left(r_{\sigma_{X}*}(X)\right).$$

#### Theorem (Equivalence of both definitions, [S.-R. F.])

There is the usual 1:1 correspondence between both definitions:

Given H, define A by

$$A_p(\overline{\nu}_p + X) := (\pi^* \nu)_p,$$

where  $X \in H_p$ .

Given A, define H by

$$H_p := \operatorname{Ker}(A_p).$$

# Theorem (Gauge transformation, [S.-R. F.])

Let  $s_i$ ,  $s_i$  be two sections of  $\mathcal{P}$  over  $U_i$  and  $U_i$ , respectively, which are open subsets of M with  $U_i \cap U_i \neq \emptyset$ . Then over  $U_i \cap U_i$ 

$$A_{s_i} = \operatorname{Ad}_{\sigma_{ji}^{-1}} \circ A_{s_j} + \Delta \sigma_{ji},$$

where  $A_{s_i} := s_i! A$  and  $\sigma_{ii}$  a section of  $\mathscr{G}$  with  $s_i = s_i \cdot \sigma_{ii}$ .

Infinitesimal theory
OO

Compatibility conditions

#### Proposition (Connection on q, [S.-R. F.])

We have an induced vector bundle connection on g given by

$$\nabla^{\mathscr{G}}\nu := \left.\frac{\mathrm{d}}{\mathrm{d}t}\right|_{t=0} \Delta \mathrm{e}^{t\nu}.$$

Compatibility conditions

Infinitesimal theory

#### Remarks

Recall,  $\mathcal{G}$  a principal  $\mathcal{G}$ -bundle.

#### Definition (Compatibility conditions, [S.-R. F.])

 $\mu_{\mathscr{C}}$  a Yang-Mills connection (w.r.t.  $\zeta \in \Omega^2(M; q)$ ) if it satisfies the **compatibility conditions**:

- **1**  $\mu_{\mathscr{G}}$  a connection 1-form on  $\mathscr{G} \stackrel{\pi_{\mathscr{G}}}{\to} M$ :
- **2**  $\mu_{\mathscr{C}}$  satisfies the **generalised Maurer-Cartan equation**

$$\left. \left( \mathrm{d}^{\pi_{\mathscr{G}}^* \nabla^{\mathscr{G}}} \mu_{\mathscr{G}} + \frac{1}{2} [\mu_{\mathscr{G}} \stackrel{\wedge}{,} \mu_{\mathscr{G}}]_{\pi_{\mathscr{G}}^* \mathscr{G}} \right) \right|_{\mathcal{G}} = \mathrm{Ad}_{g^{-1}} \circ \left. \pi_{\mathscr{G}}^! \zeta \right|_{\mathcal{G}} - \left. \pi_{\mathscr{G}}^! \zeta \right|_{\mathcal{G}}$$

Infinitesimal theory Compatibility conditions

#### Proposition ( $\nabla^{\mathcal{G}}$ a Lie bracket derivation)

Let  $\mu_{\mathscr{C}}$  be a connection 1-form on  $\mathscr{G}$ , then

$$\nabla^{\mathscr{G}}\left(\left[\mu,\nu\right]_{\mathscr{G}}\right) = \left[\nabla^{\mathscr{G}}\mu,\nu\right]_{\mathscr{G}} + \left[\mu,\nabla^{\mathscr{G}}\nu\right]_{\mathscr{G}}.$$

Compatibility conditions

Infinitesimal theory

#### Proposition ( $\nabla^{\mathcal{G}}$ a Lie bracket derivation)

Let  $\mu_{\mathscr{C}}$  be a connection 1-form on  $\mathscr{C}$ , then

$$\nabla^{\mathcal{G}}\left(\left[\mu,\nu\right]_{\mathcal{Q}}\right) = \left[\nabla^{\mathcal{G}}\mu,\nu\right]_{\mathcal{Q}} + \left[\mu,\nabla^{\mathcal{G}}\nu\right]_{\mathcal{Q}}.$$

#### Theorem (Curvature of LAB connection exact, [S.-R. F.])

 $\mu_{\mathscr{C}}$  satisfies the generalized Maurer-Cartan equation w.r.t.  $\zeta$  if and only if

$$R_{\nabla^{\mathscr{G}}} = \mathrm{ad} \circ \zeta.$$

#### Remarks

There is a simplicial differential  $\delta$  on  $\mathscr{G} \stackrel{\pi_{\mathscr{G}}}{\to} M$ 

$$\delta: \Omega^{\bullet}(\underbrace{\mathcal{G} * \ldots * \mathcal{G}}_{k \text{ times}}; \pi_{\mathscr{C}}^{*}g) \to \Omega^{\bullet}(\underbrace{\mathcal{G} * \ldots * \mathcal{G}}_{k+1 \text{ times}}; \pi_{\mathscr{C}}^{*}g)$$

such that the compatibility conditions are equivalent to

$$\delta\mu_{\mathscr{G}} = 0$$
 and  $\mathrm{d}^{\pi_{\mathscr{G}}^* \nabla^{\mathscr{G}}} \mu_{\mathscr{G}} + \frac{1}{2} [\mu_{\mathscr{G}} \stackrel{\wedge}{,} \mu_{\mathscr{G}}]_{\pi_{\mathscr{G}}^* \mathscr{G}} = \delta\zeta.$ 

- $\bigcirc$   $\mathcal{G}\cong M\times G$
- 2 Equip  $\mathcal{G}$  with canonical flat connection
- $\bigcirc$   $\zeta \equiv 0$

Curved Yang-Mills gauge theory

#### Remarks

There is a simplicial differential  $\delta$  on  $\mathscr{G} \stackrel{\pi_{\mathscr{G}}}{\to} M$ 

$$\delta: \Omega^{\bullet}(\underbrace{\mathcal{G} * \ldots * \mathcal{G}}_{k \text{ times}}; \pi_{\mathcal{C}}^{*}g) \to \Omega^{\bullet}(\underbrace{\mathcal{G} * \ldots * \mathcal{G}}_{k+1 \text{ times}}; \pi_{\mathcal{C}}^{*}g)$$

such that the compatibility conditions are equivalent to

$$\delta\mu_{\mathscr{G}}=0 \qquad \text{ and } \qquad \mathrm{d}^{\pi_{\mathscr{G}}^*\nabla^{\mathscr{G}}}\mu_{\mathscr{G}}+\frac{1}{2}[\mu_{\mathscr{G}}\stackrel{\wedge}{,}\mu_{\mathscr{G}}]_{\pi_{\mathscr{G}}^*\mathscr{Q}}=\delta\zeta.$$

#### Guide to recover the classical theory (Part 3 of 3)

- $\mathfrak{G}\cong M\times G$
- 2 Equip  $\mathscr{G}$  with canonical flat connection
- $0 \ \zeta \equiv 0$

#### Given a Yang-Mills connection on $\mathcal{G}$ :

# Definition (Generalized curvature/field strength F of A, [S.-R. F.])

 $\pi^H$  denotes the projection onto  $H\subset T\mathscr{P}$  , then we define

$$F := \mathrm{d}^{\pi^* \nabla^{\mathscr{G}}} A \circ \left(\pi^{\mathrm{H}}, \pi^{\mathrm{H}}\right) + \pi^! \zeta.$$

#### Theorem (Structure equation, [S.-R. F.])

$$F = \mathrm{d}^{\pi^* \nabla^{\mathcal{G}}} A + \frac{1}{2} [A \wedge A]_{\pi^* \mathcal{G}} + \pi^! \zeta.$$

Definition and properties

# Proposition (Properties of *F*, [S.-R. F.])

- $F(X, \cdot) = 0$ , if X vertical,
- $\bullet \ r_{\sigma}^{!}F = \mathrm{Ad}_{\sigma^{-1}} \circ F.$

Curved Yang-Mills gauge theory

Infinitesimal theory

#### Proposition (Properties of F, [S.-R. F.])

- $F(X, \cdot) = 0$ , if X vertical,
- $r_{\sigma}^! F = \operatorname{Ad}_{\sigma^{-1}} \circ F$ .

#### Theorem (Gauge transformation, [S.-R. F.])

Let  $s_i$ ,  $s_i$  be two sections of  $\mathcal{P}$  over  $U_i$  and  $U_i$ , respectively, which are open subsets of M with  $U_i \cap U_i \neq \emptyset$ . Then over  $U_i \cap U_i$ 

$$F_{s_i} = \operatorname{Ad}_{\sigma_{ii}^{-1}} \circ F_{s_j},$$

where  $F_{s_i} := s_i^! F$  and  $\sigma_{ii}$  a section of  $\mathscr G$  with  $s_i = s_i \cdot \sigma_{ii}$ .

### Theorem (Lagrangian, [S.-R. F.])

- $\kappa$  be an Ad-invariant fibre metric on q,
- M a spacetime, and \* its Hodge star operator,
- (*U<sub>i</sub>*); open covering of M with subordinate gauges  $s_i \in \Gamma(\mathscr{P}|_{II})$ .

Then the Lagrangian  $\mathfrak{L}_{CYM}[A]$ , defined locally by

$$(\mathfrak{L}_{\mathrm{CYM}}[A])\big|_{U_i} := -\frac{1}{2}\kappa(F_{s_i} \stackrel{\wedge}{,} *F_{s_i}),$$

is well-defined, and

$$\mathfrak{L}_{\mathrm{CYM}}[L^!A] = \mathfrak{L}_{\mathrm{CYM}}[A]$$

for all principal bundle automorphisms L.

# Example (Hopf fibration $\mathbb{S}^7 \to \mathbb{S}^4$ , [S.-R. F.])

Let P be the Hopf bundle

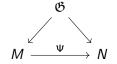
$$\mathrm{SU}(2)\cong \mathbb{S}^3\longrightarrow \mathbb{S}^7$$

Define  $\mathscr{P} := \mathscr{G}$  as the inner group bundle of P,

$$\mathscr{G} := c_{\mathrm{SU}(2)}(P) := (P \times \mathrm{SU}(2)) / \mathrm{SU}(2).$$

This principal  $c_{SU(2)}(P)$ -bundle admits the structure as curved Yang-Mills gauge theory; there is no description as flat theory.

Gauge theory	Structure
Yang-Mills	Lie group <i>G</i>
Curved Yang-Mills	Lie group bundle ${\mathscr G}$
Curved Yang-Mills-Higgs	Lie groupoid &?



#### Remarks

Infinitesimal theory

- Richer set of principal bundles, containing Lie groupoids equipped with "non-flat Maurer-Cartan forms".
- Principal bundle for the whole of Yang-Mills-Higgs theory
- Even if  $\mathscr{G}$  is trivial, what happens if its connection is not flat?
- May result into obstruction statements for curved Yang-Mills-Higgs gauge theories.

# Thank you!