Curved Yang-Mills gauge theories and their recent applications

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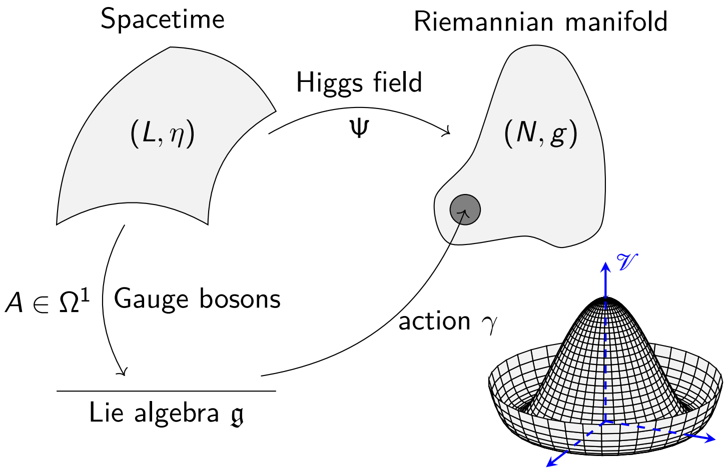
Curved Yang-Mills gauge theory

Infinitesimal curved Yang-Mills-Higgs gauge theory



Sombrero hat potential made by Mark J.D. Hamilton.

Motivation 1 by Thomas Strobl and Alexei Kotov

Classical formalism	CYMH GT
Lie algebra $\mathfrak{g}$ as $L \times \mathfrak{g}$	Lie algebroid $E \to N$
$\mathfrak{g}$-action γ	Anchor ρ of E & E-connections
Canonical flat connection ∇^0 on $L \times \mathfrak{g}$	General connection ∇ on E

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Remarks (Why a "curved theory"?)

Usually, the field strength F is given by (abelian, for simplicity)

$$F := \mathrm{d}A = \mathrm{d}^{\nabla^0} A.$$

$\rightsquigarrow$ We will use a general connection ∇ instead of ∇^0, and ∇ may not be flat.

Motivation 2 by S.-R. F.

Consider a semisimple Lie group G and a principal G-bundle $P \to L$:

$$(P \times \mathfrak{g})/G \hookrightarrow \mathrm{T}P/G \twoheadrightarrow \mathrm{T}L$$

Gedankenexperiment

1. Adjoint connection $\leftrightarrow$ Ehresmann connection on P.
2. As parallel transport along a curve γ:

$$\mathrm{PT}_{\gamma}^{\mathrm{Ad}(P)}([p, v]) = \left[\mathrm{PT}_{\gamma}^{P}(p), v\right]$$

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Lie algebra $\mathfrak{g}$ as trivial bundle w/ canonical flat connection

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1. Adjoint connection $\leftrightarrow$ Ehresmann connection on P.
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$$\mathrm{PT}_\gamma^{\mathrm{Ad}(P)}([p, v]) = \left[\mathrm{PT}_\gamma^P(p) \cdot \kappa_\gamma, \kappa_\gamma^{-1} \cdot \mathrm{PT}_\gamma^0(v)\right],$$

Lie algebra $\mathfrak{g}$ as trivial bundle w/ canonical flat connection, κ_γ values in G & "suitable"

Theorem (Field Redefinitions S.-R. F.)

This leads to an equivalence relation of gauge theories, preserving dynamics and kinematics.
But: In the curved sense! Curvature terms appear.

Motivation (S.-R. F.)

1. How to formulate gauge theory such that it is invariant under field redefinitions?
2. Are there curved theories which are **not** equivalent to classical ones?

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We will only focus on Yang-Mills theories:

	Classical	Curved
Infinitesimal	Lie algebra $\mathfrak{g}$	LAB[1] $\mathscr{g}$
Integrated	Lie group G	LGB[2] $\mathscr{G}$

$$\begin{tikzcd} G \arrow[r] & \mathscr{G} \arrow[d] \\ & M \end{tikzcd}$$

[1]LAB = Lie algebra bundle
[2]LGB = Lie group bundle

Definition (LGB actions, simplified)

$$\begin{array}{ccc} & & \mathscr{G} \\ & & \downarrow \\ \mathscr{P} & \xrightarrow{\pi} & L \end{array}$$

$\mathscr{P} \xrightarrow{\pi} L$ a fibre bundle. A **right-action of $\mathscr{G}$ on $\mathscr{P}$** is a smooth map $\mathscr{P} * \mathscr{G} := \pi^*\mathscr{G} = \mathscr{P} \times_L \mathscr{G} \to \mathscr{P}$, $(p, g) \mapsto p \cdot g$, satisfying the following properties:

$$\pi(p \cdot g) = \pi(p), \tag{1}$$

$$(p \cdot g) \cdot h = p \cdot (gh), \tag{2}$$

$$p \cdot e_{\pi(p)} = p \tag{3}$$

for all $p \in \mathscr{P}$ and $g, h \in \mathscr{G}_{\pi(p)}$, where $e_{\pi(p)}$ is the neutral element of $\mathscr{G}_{\pi(p)}$.

K. Mackenzie. General Theory of Lie Groupoids and Algebroids. *London Mathematical Society Lecture Note Series*, 213, 2005.

Definition (Principal bundle)

Still a fibre bundle

$$\begin{tikzcd} G \arrow[r] & \mathscr{P} \arrow[d, "\pi"] \\ & L \end{tikzcd}$$

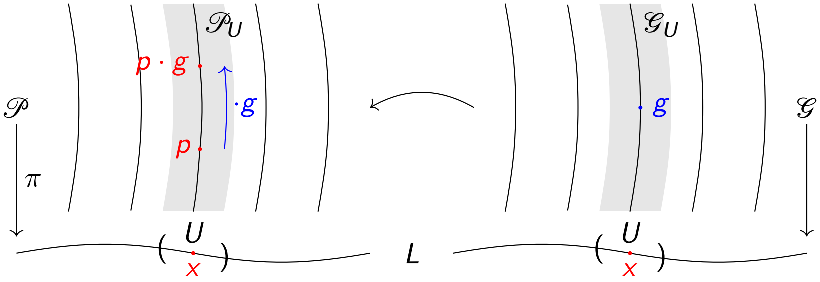
but with $\mathscr{G}$-action

$$\underset{\mathscr{P} * \mathscr{G}}{\cancel{\mathscr{P} \times G}} \to \mathscr{P}$$

simply transitive on fibres of $\mathscr{P}$, and "suitable" atlas.

Ieke Moerdijk, Janez Mrcun. Introduction to Foliations and Lie Groupoids.
Cambridge Studies in Advanced Mathematics 91, Cambridge University Press, Cambridge, 2003

Connection on $\mathscr{P}$: Idea



But:

$$r_g : \mathscr{P}_x \to \mathscr{P}_x$$

$\Rightarrow$ $\mathrm{D}_p r_g$ only defined on vertical structure

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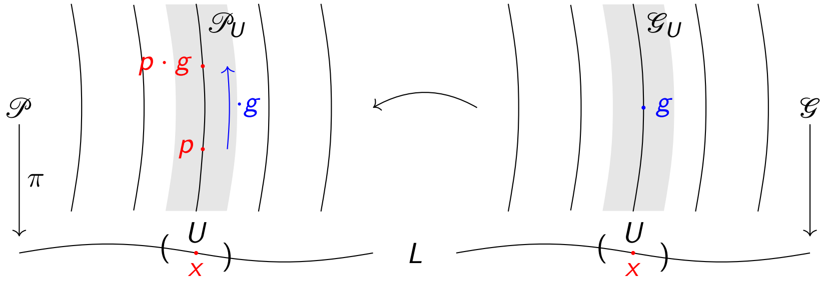

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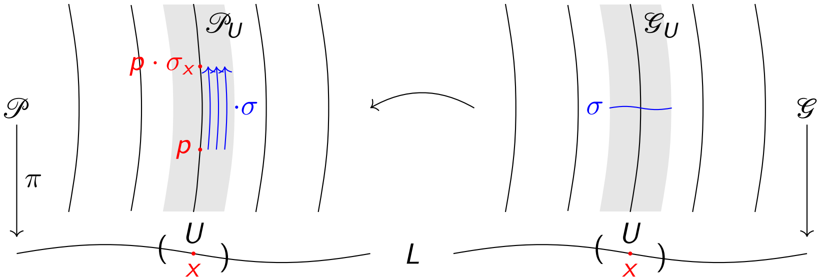
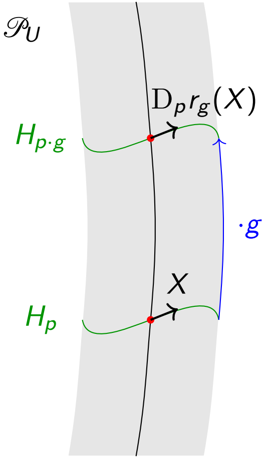


Use $\sigma \in \Gamma(\mathscr{G}) : r_\sigma(p) := p \cdot \sigma_x$

Connection on $\mathscr{P}$: Revisiting the classical setup

If $\mathscr{P}$ a typical principal bundle ($\mathscr{G}$ trivial, $\sigma \equiv g$ constant), and H a connection:



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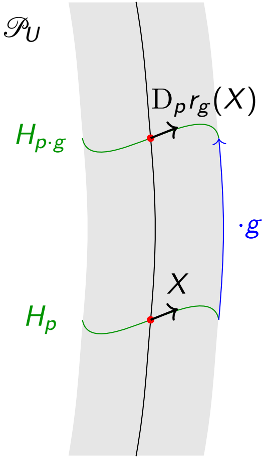

Remarks (Integrated case)

Parallel transport $\mathrm{PT}_\gamma^{\mathscr{P}}$ in $\mathscr{P}$:

$$\mathrm{PT}_\gamma^{\mathscr{P}}(p \cdot g) = \mathrm{PT}_\gamma^{\mathscr{P}}(p) \cdot g$$

where $\gamma : I \to L$ is a base path

Connection on $\mathscr{P}$: General case

Remarks (Integrated case)

Ansatz: Introduce connection on $\mathscr{G}$,

$$\mathrm{PT}_\gamma^{\mathscr{P}}(p \cdot g) = \mathrm{PT}_\gamma^{\mathscr{P}}(p) \cdot \mathrm{PT}_\gamma^{\mathscr{G}}(g).$$

Back to the roots

1. $\mathscr{G} \cong L \times G$
2. Equip $\mathscr{G}$ with canonical flat connection

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Definition (Ehresmann/Yang-Mills connection, [C. L.-G., S.-R. F.])

A surjective submersion $\pi_{\mathscr{T}} : \mathscr{T} \to L$ so that one has a commuting diagram

$$\begin{tikzcd} & \mathscr{G} \arrow[dl, bend right] \arrow[d, "\pi_{\mathscr{G}}"] \\ \mathscr{T} \arrow[r, "\pi_{\mathscr{T}}"] & L \end{tikzcd}$$

1. **Ehresmann connection:**

$$\mathrm{PT}^{\mathscr{T}}_{\gamma}(t \cdot g) = \mathrm{PT}^{\mathscr{T}}_{\gamma}(t) \cdot \mathrm{PT}^{\mathscr{G}}_{\gamma}(g)$$

2. **Yang-Mills connection:** Additionally

$$\mathrm{PT}^{\mathscr{T}}_{\gamma_0}(t) = t \cdot g_{\gamma_0}$$

for some $g_{\gamma_0} \in \mathscr{G}_{\pi_{\mathscr{T}}(t)}$, where γ_0 is a contractible loop.

Definition (Multiplicative YM connection, [S.-R. F.])

On $\mathscr{G}$ there is also the notion of **multiplicative Yang-Mills connections**, that is,

$$\mathrm{PT}^{\mathscr{G}}_{\gamma}(q \cdot g) = \mathrm{PT}^{\mathscr{G}}_{\gamma}(q) \cdot \mathrm{PT}^{\mathscr{G}}_{\gamma}(g),$$
$$\mathrm{PT}^{\mathscr{G}}_{\gamma_0}(q) = g_{\gamma_0} \cdot q \cdot g_{\gamma_0}^{-1}$$

Remarks

There is a simplicial differential δ on $\mathcal{G} \overset{\pi_{\mathcal{G}}}{\to} L$ with Lie algebra bundle $\mathscr{g}$

$$\delta : \Omega^\bullet(\underbrace{\mathcal{G} * \ldots * \mathcal{G}}_{k \text{ times}}; \pi_{\mathcal{G}}^* \mathscr{g}) \to \Omega^\bullet(\underbrace{\mathcal{G} * \ldots * \mathcal{G}}_{k+1 \text{ times}}; \pi_{\mathcal{G}}^* \mathscr{g})$$

such that the definition of the multiplicative Yang-Mills connection is equivalent to the **compatibility conditions**

- Connection closed
- Curvature exact ([S.-R. F.])

 For differential: Marius Crainic, Maria Amelia Salazar, and Ivan Struchiner. Multiplicative forms and Spencer operators. *Mathematische Zeitschrift*, 279(3):939–979, 2015.

Remarks

On the Lie algebra bundle $\mathcal{g}$ we have a connection ∇ with

$$\nabla\left([\mu,\nu]_{\mathcal{g}}\right) = [\nabla\mu,\nu]_{\mathcal{g}} + [\mu,\nabla\nu]_{\mathcal{g}},$$
$$R_\nabla = \mathrm{ad}\circ\zeta.$$

Example

Given a short exact sequence of algebroids

$$\mathcal{g} \hookrightarrow E \twoheadrightarrow \mathrm{T}L$$

with splitting $\chi\colon \mathrm{T}L \to E$, then

$$\nabla_X\nu = [\chi(X),\nu]_E,$$
$$\zeta(X,X') = [\chi(X),\chi(X')]_E - \chi([X,X']).$$

For the first equation: Camille Laurent-Gengoux, Mathieu Stiénon, and Ping Xu. Non-abelian differentiable gerbes. *Advances in Mathematics*, 220(5):1357–1427, 2009.

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Integrating Alexei Kotov's and Thomas Strobl's idea

Definition (Principal bundle connection, [S.-R. F.])

- On $\mathscr{G}$: Multiplicative Yang-Mills connection
- On $\mathscr{P}$: Ehresmann connection

Definition (Generalized curvature/field strength F of A, [S.-R. F.])

We define

$$F := \mathrm{d}^{\pi^*\nabla} A + \frac{1}{2}[A \stackrel{\wedge}{,} A]_{\pi^*\mathcal{g}} + \pi^! \zeta.$$

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Theorem (Lagrangian, [S.-R. F.])

- *κ be an Ad-invariant fibre metric on $\mathscr{g}$,*
- *L a spacetime, and $*$ its Hodge star operator,*
- *$(U_i)_i$ open covering of L with subordinate gauges $s_i \in \Gamma(\mathscr{P}|_{U_i})$.*

Then the Lagrangian $\mathfrak{L}_{\mathrm{CYM}}[A]$, defined locally by

$$\left(\mathfrak{L}_{\mathrm{CYM}}[A]\right)\big|_{U_i} := -\frac{1}{2}\kappa(F_{s_i} \overset{\wedge}{,} *F_{s_i}),$$

is well-defined, and

$$\mathfrak{L}_{\mathrm{CYM}}\left[K^! A\right] = \mathfrak{L}_{\mathrm{CYM}}[A]$$

for all principal bundle automorphisms K.

Back to the roots

1. $\mathscr{G} \cong L \times G$
2. Equip $\mathscr{G}$ with canonical flat connection
3. $\zeta \equiv 0$

Example (Hopf fibration $\mathbb{S}^7 \to \mathbb{S}^4$, [S.-R. F.])

Let P be the Hopf bundle

$$\begin{array}{ccc} \mathrm{SU}(2) \cong \mathbb{S}^3 & \longrightarrow & \mathbb{S}^7 \\ & & \downarrow \\ & & \mathbb{S}^4 \end{array}$$

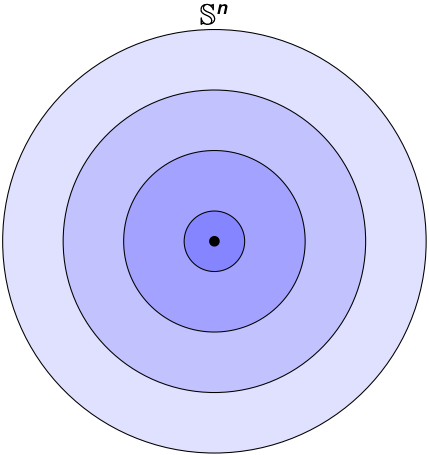
Define $\mathscr{P} := \mathscr{G}$ as the inner group bundle of P,

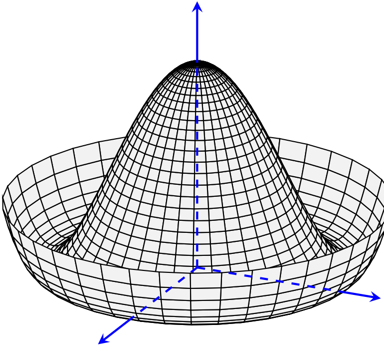
$$\mathscr{G} := c_{\mathrm{SU}(2)}(P) := \left(P \times \mathrm{SU}(2)\right)\Big/\mathrm{SU}(2).$$

This principal $c_{\mathrm{SU}(2)}(P)$-bundle admits the structure as curved Yang-Mills gauge theory; there is no description as classical gauge theory.

Applications: Singular foliations (joint work w/ Camille Laurent-Gengoux)

Why foliations?





Right diagram made by Mark J.D. Hamilton.

Singular Foliations:

- Gauge Theory
  (Ex.: Singular foliation $\leftrightarrow$ Symmetry breaking $\rightarrow$ Higgs mechanism)
- Poisson Geometry
  (Singular foliation of symplectic leaves)
- Lie groupoids and algebroids
- Dirac structures
- Generalised complex manifolds
- Non-commutative geometry
- ...

Definition (Smooth singular foliation)

A **smooth singular foliation** $\mathscr{F}$ on a smooth manifold is a subspace of $\mathfrak{X}_c(M)$ so that

- it is **involutive**,
- it is **stable under** $C^\infty(M)$**-multiplication**,
- it is **locally finitely generated**.

Peter Stefan, Accessible sets, orbits, and foliations with singularities. *Proc. London Math. Soc.*, 29, 1974.
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- it is **locally finitely generated**, *i.e. around each* $p \in M$ *there is an open neighbourhood* U *and a finite family* $(X^i)_i^r$ $(X^i \in \mathscr{F})$ *such that for all* $X \in \mathscr{F}$ *there are* $f_i \in C^\infty(M)$ *satisfying on* U.

$$X = \sum_i f_i X^i.$$

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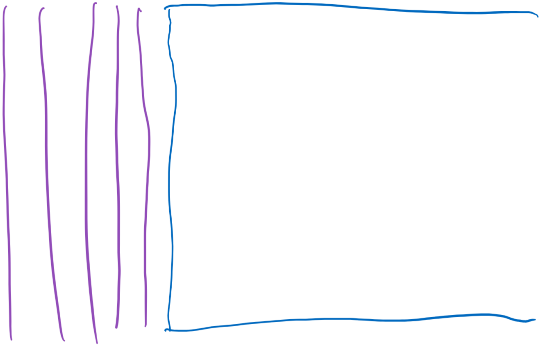
Remarks (Leaves)

Following the flows in $\mathscr{F}$, this gives rise to a partition of connected immersed submanifolds in M.



Peter Stefan, Accessible sets, orbits, and foliations with singularities. *Proc. London Math. Soc.*, 29, 1974.
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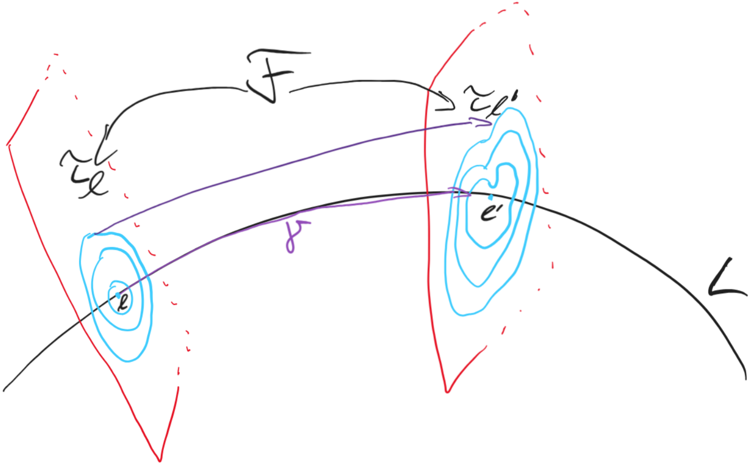
Idea: Relation to gauge theory



Camille Laurent-Gengoux and Leonid Ryvkin, The holonomy of a singular leaf, *Selecta Mathematica 28*, no. 2, 45, 2022.

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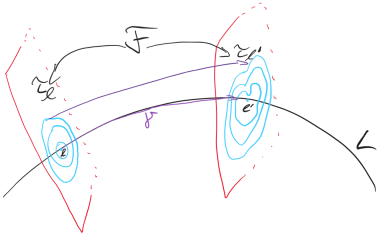

Theorem ($\mathscr{F}$-connections)

There is a connection on the normal bundle of a leaf L:

- *Horizontal vector fields are in $\mathscr{F}$.*
- *Parallel transport PT_γ has values in $\mathrm{Sym}(\tau_l, \tau_{l'})$.*
- *For a contractible loop γ_0 at l: PT_{γ_0} values in $\mathrm{Inner}(\tau_l)$.*

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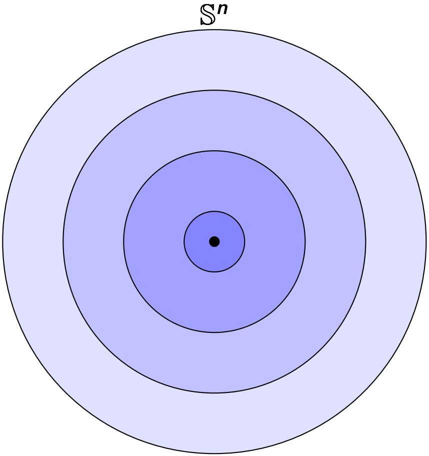
Example of a transverse foliation τ:

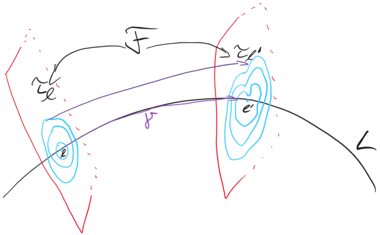


Remarks

- Inner(τ_I) maps each circle to itself
- Sym(τ_I) allows to exchange circles
- Both preserve τ_I and fix the origin

Camille Laurent-Gengoux and Leonid Ryvkin, The holonomy of a singular leaf, *Selecta Mathematica 28*, no. 2, 45, 2022.

Idea



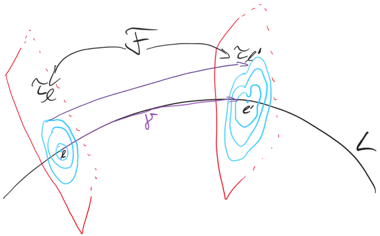
Idea

Generators of $\mathscr{F}$ given by $\mathscr{F}_{\text{projectable}}$:

$$\mathbb{H}(X) + \overline{\nu},$$

where $X \in \mathfrak{X}(L)$, $\mathbb{H}(X)$ its projectable horizontal lift, $\nu \in \Gamma(\mathrm{inner}(\tau))$ and $\overline{\nu}$ its fundamental vector field.

Idea: Relation to gauge theory



Idea

Fix l and given τ_l: Reconstruct $\mathscr{F}$.

$$\begin{aligned}
[\mathbb{H}(X)+\overline{\nu},\mathbb{H}(X')+\overline{\mu}] &= \mathbb{H}([X,X'])+\overline{\ldots} \\
&= \underbrace{[\mathbb{H}(X),\mathbb{H}(X')]}_{\rightsquigarrow\ \text{curvature}} \\
&\qquad + \underbrace{[\mathbb{H}(X),\overline{\mu}]-[\mathbb{H}(X'),\overline{\nu}]}_{\rightsquigarrow\ \text{connection}} + \overline{[\nu,\mu]}
\end{aligned}$$

Theorem ([C. L.-G., S.-R. F.])

Given a multiplicative Yang-Mills connection on $\mathscr{G}$ and a Yang-Mills connection $\mathbb{H}$ on $\mathscr{T}$, then there is a natural foliation on $\mathscr{T}$ generated by

$$\mathbb{H}(X) + \overline{\nu},$$

where $X \in \mathfrak{X}(L)$ and $\nu \in \Gamma(\mathscr{g})$.

Proof.

We have

$$[\mathbb{H}(X), \overline{\nu}] = \overline{\nabla_X \nu},$$
$$[\mathbb{H}(X), \mathbb{H}(X')] = \mathbb{H}([X, X']) + \overline{\zeta(X, X')},$$

where $\zeta \in \Omega^2(L; \mathscr{g})$.

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Idea (Leaf L simply connected)

Fix a point $l \in L$ with transverse model $\left(\mathbb{R}^d, \tau_l\right)$:

1. $G = \mathrm{Inn}(\tau_l)$
2. P a principal G-bundle, equipped with an ordinary connection

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4. $\mathscr{T} := \left(P \times \mathbb{R}^d\right)\big/G$, the **normal bundle**

Remarks

- Think of the induced connection on $\mathscr{T}$ as the $\mathscr{F}$-connection.
- $\mathscr{G}$ acts on $\mathscr{T}$ (canonically from the left).

Idea (Leaf L simply connected)

Fix a point $l \in L$ with transverse model $\left(\mathbb{R}^d, \tau_l\right)$:

1. $G = \mathrm{Inn}(\tau_l)$
2. P a principal G-bundle, equipped with an ordinary connection
3. $\mathscr{G} := (P \times G)\big/G$, the **inner group bundle**
4. $\mathscr{T} := \left(P \times \mathbb{R}^d\right)\big/G$, the **normal bundle**

Remarks

- Think of the induced connection on $\mathscr{T}$ as the $\mathscr{F}$-connection.
- $\mathscr{G}$ acts on $\mathscr{T}$ (canonically from the left).

Proposition ([C. L.-G., S.-R. F.])

The associated connection on $\mathscr{G}$ is a multiplicative Yang-Mills connection and the one on $\mathscr{T}$ is a corresponding Yang-Mills connection.

Remarks

Thus, we have a singular foliation on $\mathscr{T}$, which, by construction, admits L as a leaf and τ_l as transverse data.

Proposition ([C. L.-G., S.-R. F.])

The reconstructed foliation is independent of the choice of connection on P.

$$
\begin{tikzcd}
\Gamma(\mathrm{Ad}(P)) \arrow[rrrr, "\nu \mapsto \overline{\nu}"] \arrow[dr, hook] & & & & \mathcal{T} \arrow[dl, hook'] \\
& \Gamma(\mathrm{At}(P)) \arrow[rr, "\xi \mapsto \overline{\xi}"] \arrow[dr, two heads] & & \mathscr{F}_{\text{projectable}} \arrow[dl, two heads] & \\
& & \mathfrak{X}(L) \arrow[ul, bend left, "A"] \arrow[ur, bend right, "\mathbb{H}"'] & &
\end{tikzcd}
$$

Proposition ([C. L.-G., S.-R. F.])

The reconstructed foliation is independent of the choice of connection on P.

$$\begin{tikzcd}
\Gamma(\mathrm{Ad}(P)) \arrow[rrrr, "\nu \mapsto \overline{\nu}"] \arrow[dr, hook] & & & & \tau \arrow[dl, hook'] \\
& \Gamma(\mathrm{At}(P)) \arrow[rr, "\xi \mapsto \overline{\xi}"] \arrow[dr, two heads] & & \mathscr{F}_{\text{projectable}} \arrow[dl, two heads] & \\
& & \mathfrak{X}(L) \arrow[ul, bend left, "A"] \arrow[ur, bend right, "\mathbb{H}"'] & &
\end{tikzcd}$$

Summary

Remarks ([C. L.-G., S.-R. F.])

In the simply connected case, the following are equivalent:

- Singular foliations with leaf L and transverse model $\left(\mathbb{R}^d, \tau_I\right)$
- Principal $\mathrm{Inner}(\tau_I)$-bundles over L

Remarks (Classification of curved Yang-Mills gauge theories)

If $\mathscr{G}$ acts faithfully on $\mathscr{T}$, preserving L, then a curved Yang-Mills gauge theory can be flattened if and only if P is flat.

Curved YM Gauge Theory	Singular Foliations $\mathscr{F}$
Multiplicative Yang-Mills connection	$\mathscr{F}$-connection
Flat gauge theory	Flat singular foliation
Field redefinition of connection on $\mathscr{G}$	Different choice of $\mathscr{F}$-connection

Thank you!