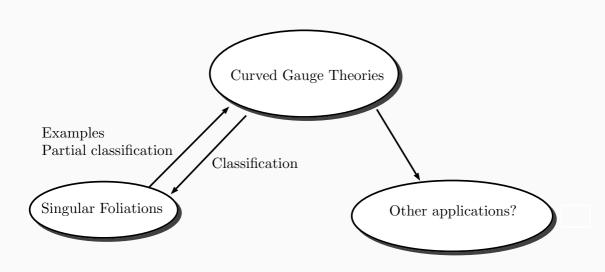
Curved Yang-Mills-Higgs theories

Simon-Raphael Fischer, based on joint works with Camille Laurent-Gengoux, and with Mehran Jalali Farahani, Hyungrok Kim (金炯錄), Christian Sämann





Motivation

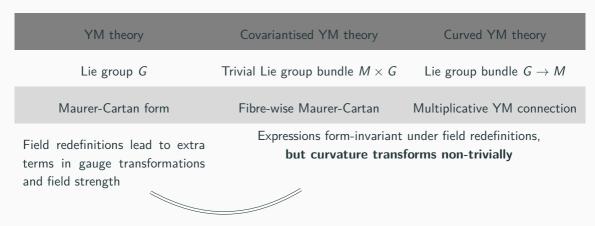
Motivation: Covariantisation of Yang-Mills(-Higgs) theory

Covariantization

Classical theory	Covariantised flat theory	Curved Theory
Vector space V	Trivial vector bundle $M \times V$	Vector bundle $V o M$
$rac{\partial}{\partial x^i}$	Canonical flat connection $ abla^0$	Vector bundle connection $ abla$
Coordinate changes may lead to extra terms	Coordinate expressions form-invariant under coordinate changes	

Curved Yang-Mills gauge theory (curved YM theory)

Covariantization



S.-R. Fischer. *Integrating curved Yang–Mills gauge theories*, arXiv: 2210.02924, 2022.

S.-R. Fischer. Geometry of curved Yang–Mills–Higgs gauge theories, Ph.D. thesis, Institut Camille Jordan [Villeurbanne], France, U. Geneva, Switzerland, 2021; doi: 10.13097/archive-ouverte/unige:152555

Curved Yang-Mills-Higgs theory (curved YMH theory)

Covariantization

YMH theory	Covariantised YMH theory	Curved YMH theory
Lie group G with right-action on N	Action groupoid $N \times G$	Lie groupoid $G ightrightarrows N$
Maurer-Cartan form	Fibre-wise Maurer-Cartan	Covariant adjustments
Field redefinitions lead to extra terms in gauge transformations and field strength	Expressions form-invariant under field redefinitions, but curvature transforms non-trivially	

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Lie groupoids



$$\mathcal{G}$$
 $\downarrow \downarrow s$
 N

Definition (Lie groupoids)

 \mathscr{G} a **Lie groupoid** if there are surjective submersions $s,t\colon \mathscr{G}\to N$, source and target, respectively, and a smooth multiplication map $\mathscr{G}_s\times_t\mathscr{G}\to\mathscr{G}$ such that

$$s(g'g) = s(g),$$
 $t(g'g) = t(g')$

for all $(g',g) \in \mathcal{G}_{s} \times_{t} \mathcal{G}$ (i.e. s(g') = t(g)), satisfying the typical expected properties, that is,

Associativity:
$$(g''g')g = g''(g'g),$$
 Units:
$$ge_{s(g)} = g, \qquad e_{t(g)}g = g,$$
 Inverse:
$$g^{-1}g = e_{s(g)}, \qquad gg^{-1} = e_{t(g)}$$

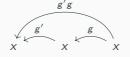
for all $(g'', g', g) \in \mathcal{G}_s \times_t \mathcal{G}_s \times_t \mathcal{G}$, where one requires the existence of the *unit* e as a global section of both, s and t, and the *inverse* $g^{-1} \in \mathcal{G}$ of g.



Example (Lie groups)

Lie groups G

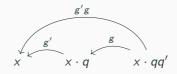
$$G$$
 $t \downarrow \downarrow s$
 $\{*\}$



Example (Lie group bundles (LGBs))

LGB $\pi_G \colon G \to M$

$$G$$
 $\pi_G \bigcup_{M} \pi_G$



Example (Action groupoid (trivial))

Lie group G with action $\Psi \colon \mathcal{N} \times G \to \mathcal{N}$, $(p,q) \mapsto p \cdot q$, on \mathcal{N} .

$$\begin{array}{c}
N \times G \\
\operatorname{pr}_1 \downarrow \downarrow \Psi \\
N
\end{array}$$

$$(x,q) (x \cdot q, q') = (x, qq'),$$

 $e_x = (x, e),$
 $(x,q)^{-1} = (x \cdot q, q^{-1})$



Definition (Groupoid right-action)

A **right-action** is a smooth map $P {}_{\Phi} \times_t \mathscr{G} \to P$ such that

$$\Phi(p \cdot g') = s(g'),$$

$$(p \cdot g') \cdot g = p \cdot (g'g),$$

$$p \cdot e_{\Phi(p)} = p$$

for all $(p, g', g) \in P_{\Phi} \times_t \mathcal{G}_t \times_s \mathcal{G}$.

Principal bundles and their

Ehresmann connections



Definition (Principal groupoid-bundles)

 $\pi: P \to M$ surjective submersion is a **principal** \mathscr{G} -bundle if

$$\pi(p\cdot g)=\pi(p)$$

for all $(p,g) \in P_{\Phi} \times_t \mathcal{G}$, and if

$$P {}_{\Phi} \times_t \mathscr{G} \to P {}_{\pi} \times_{\pi} P,$$

 $(p,g) \mapsto (p,p \cdot g)$

is a diffeomorphism.

Ehresmann connection



Remarks

Infinitesimal action:

$$\mathsf{T}P_{\mathsf{D}\Phi} \times_{\mathsf{D}t} \mathsf{T}\mathscr{G} \to \mathsf{T}P,$$

 $(X,Y) \mapsto X \cdot Y.$

For $r_g(p) := p \cdot g$ with $\Phi(p) = t(g)$, its infinitesimal version corresponds to

$$Dr_g(X) = X \cdot 0$$

for all X with $D\Phi(X) = 0$.

Ehresmann connection



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Ehresmann connection



Infinitesimal \mathcal{G} -action on P:

$$\mathsf{T}P_{\mathsf{D}\Phi} \times_{\mathsf{D}t} \mathsf{T}\mathscr{G} \to \mathsf{T}P,$$

 $(X, Y) \mapsto X \cdot Y.$

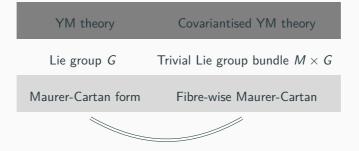
Idea

Horizontal distribution HP (w.r.t. π) with

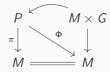
$$\mathsf{D}\Phi(\mathsf{H}P)=0.$$

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- D. Signori and M. Stiénon. On nonlinear gauge theories, J. Geom. Phys. 59 1063, 2009.

We lose covariantization



Me would lose the equivalence to the covariantised theory where we set $\Phi=\pi$



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New idea!



Idea (Ehresmann connection on P)

Equip $\mathscr G$ with a horizontal distribution $H\mathscr G$ (w.r.t. t). An Ehresmann connection on P is a horizontal distribution HP (w.r.t. π) so that the infinitesimal $\mathscr G$ -action on P restricts

$$\mathsf{H}P \,_{\mathsf{D}\Phi} \times_{\mathsf{D}t} \mathsf{H}\mathscr{G} \to \mathsf{H}P.$$

Invariance then via the modified right-pushforward \mathcal{P}_{g*}

$$r_{g*}(X) := X \cdot Y$$

for all $X \in T_p P$, where $Y \in H_g \mathcal{G}$ is the unique lift of $D_p \Phi(X)$.

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Idea

HP Ehresmann, then

$$\mathsf{H}P_{\mathsf{D}\Phi} \times_{\mathsf{D}t} \mathsf{H}\mathscr{G} \to \mathsf{H}P.$$

Theorem

Given such an HP, then HG is Cartan,

$$H\mathscr{G}_{Ds} \times_{Dt} H\mathscr{G} \to H\mathscr{G}$$
.

Sketch as a proof for LGBs.

W.r.t. parallel transports we have

$$PT^{P}(p \cdot g) = PT^{P}(p) \cdot PT^{\mathscr{G}}(g).$$

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Given such an HP, then H \mathscr{G} is **Cartan**,

$$\mathsf{H}\mathscr{G}_{\mathsf{D}s}\times_{\mathsf{D}t}\mathsf{H}\mathscr{G}\to\mathsf{H}\mathscr{G}.$$

Sketch as a proof for LGBs.

W.r.t. parallel transports we have

$$\mathsf{PT}^P(p \cdot g) = \mathsf{PT}^P(p) \cdot \mathsf{PT}^{\mathscr{G}}(g),$$

then

$$PT^{P}(p \cdot (gq)) = PT^{P}(p) \cdot PT^{\mathscr{G}}(gq)$$

$$\parallel$$

$$PT^{P}((p \cdot g) \cdot q) = PT^{P}(p) \cdot (PT^{\mathscr{G}}(g) PT^{\mathscr{G}}(q)).$$

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Remarks

Vice versa, if $H\mathscr{G}$ is Cartan, then such HP exist.

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What else?

Idea

Open a textbook about ordinary gauge theory, and replace:

- Right-pushforward → Modified right-pushforward,
- Adjoint representation → Adjoint representation via Cartan connections (?)
- Generalising the concept of field strength accounting for a possible curvature coming from ${\mathscr G}$

Too much to introduce! (Lie algebroids, Lie algebroid connections, basic curvature, basic connections, and so on...)

Thus, for today: LGBs!

Curved Yang-Mills gauge theory

Remarks

The groupoid is now an LGB $\pi_{\mathscr{G}} : \mathscr{G} \to M$ with structural Lie group G.

Gauge invariance requires us to study special Cartan connections:

Definition (Multiplicative Yang-Mills connections)

A connection 1-form $\mu_{\mathscr{G}}$ on the LGB \mathscr{G} is a multiplicative Yang-Mills connection (w.r.t. a $\zeta \in \Omega^2(M; \mathscr{Q})$) if it is a Cartan connection satisfying the generalised Maurer-Cartan equation:

$$\left. \left(\mathrm{d}^{\pi_{\mathscr{G}}^* \nabla^{\mathscr{G}}} \mu_{\mathscr{G}} + \frac{1}{2} [\mu_{\mathscr{G}} \stackrel{\wedge}{,} \mu_{\mathscr{G}}]_{\pi_{\mathscr{G}}^* \mathscr{G}} \right) \right|_{g} = \mathrm{Ad}_{g^{-1}} \circ \pi_{\mathscr{G}}^! \zeta \big|_{g} - \pi_{\mathscr{G}}^! \zeta \big|_{g}$$

for all $g \in \mathcal{G}$, where g is the Lie algebra bundle (LAB) of \mathcal{G} , and $\nabla^{\mathcal{G}}$ is the naturally induced LAB connection.

Remarks (Closed connection, exact curvature)

$$R_{\nabla^{\mathcal{G}}} = \operatorname{ad} \circ \zeta.$$

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Definition

The field strength F of an Ehresmann connection $A \in \Omega^1(P; \pi^*g)$ on $\pi: P \to M$ is defined by

$$F := \mathrm{d}^{\pi^* \nabla^{\mathcal{G}}} A + \frac{1}{2} [A \stackrel{\wedge}{,} A]_{\pi^* \mathscr{Q}} + \pi^! \zeta.$$

Theorem

By contracting F with an ad-invariant scalar product we achieve a gauge-invariant Lagrangian $\mathcal{L}[A]$, that is,

$$\mathscr{L}[L^!A] = \mathscr{L}[A$$

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Examples

Idea

Let us assume that the structural Lie group G is semisimple, then:

Theorem

There is a unique (up to isomorphism) principal G-bundle Q such that

$$\mathscr{G} = (Q \times G)/G,$$

multiplicative Yang-Mills connections on & are precisely associated connections, thus in 1:1 correspondence to Ehresmann connections on Q, and ζ is the corresponding curvature on Q.

Corollary

There are only curved multiplicative Yang-Mills connections on G if and only if Q is not flat.

Example

Hopf fibration
$$Q := \mathbb{S}^7 \to \mathbb{S}^4$$
,

$$\mathscr{G} = (Q \times SU(2))/SU(2).$$

- This can be relaxed to S with a trivial centre, using a nice geometric interpretation via singular

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Example

Hopf fibration $Q := \mathbb{S}^7 \to \mathbb{S}^4$.

$$\mathscr{G} = (Q \times SU(2))/SU(2).$$

Remarks

Some technical details:

- This is stable under field redefinitions.
- This can be relaxed to S with a trivial centre, using a nice geometric interpretation via singular foliations. However, dynamics of gauge theory require Ad-invariance...

See Camille's talk on Friday ©

About Lie groupoids...

Idea

Assume that G is semisimple **and** acts faithfully on \mathbb{R}^d fixing 0. Then \mathscr{G} acts faithfully on

$$N := \left(Q \times \mathbb{R}^d\right) / G$$
,

giving rise to an action groupoid structure on $\phi^*\mathscr{G}$ $(\phi\colon N o M)$

$$\phi^* \mathcal{G}$$

$$t \iiint_{S} s$$

$$N$$

Theorem

The pullback of the multiplicative Yang-Mills connection on $\phi^*\mathcal{G}$ is flat up to field redefinitions if and only if Q is flat.

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Other technical details:

- Also stable under field redefinitions!
- Here semi-simplicity is rather important.

Example

Again Q the Hopf fibration $\mathbb{S}^7 \to \mathbb{S}^4$, then naturally choose $\mathbb{R}^d = \mathbb{C}^2$.

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A quick funny example, again by \mathbb{S}^7

Example

The Lie groupoid $\mathscr G$ given as the pair groupoid of $\mathbb S^7$ leads to descriptions of curved Yang-Mills-Higgs theories, which cannot be described in an ordinary way due to the lack of non-associativity on $\mathbb S^7$.

Thank you!