

# Geometry of curved Yang-Mills-Higgs gauge theories

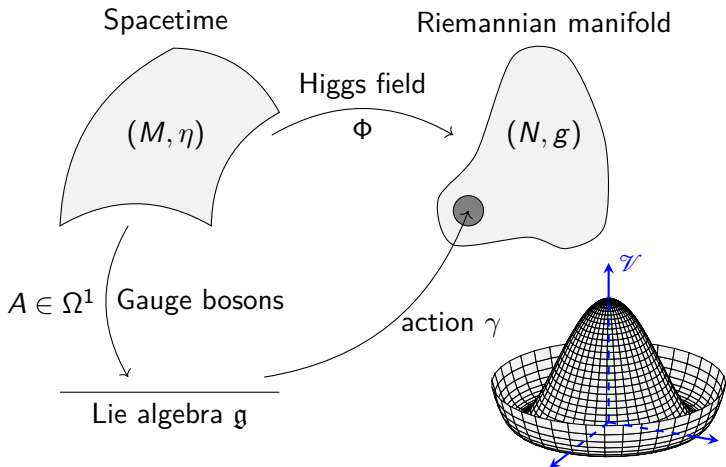
Simon-Raphael Fischer

30 June 2021

# Table of contents

- 1 Curved Yang-Mills-Higgs gauge theory
  - Motivation
  - Mathematical basics
  - Curved Yang-Mills-Higgs gauge theory
- 2 Field Redefinition
  - General definition
  - Lie algebra bundles
- 3 Integration (WIP)
- 4 Conclusion

# Infinitesimal gauge theory



# Guide: Curved Yang-Mills-Higgs gauge theory

Classical formalism	CYMH GT
Lie algebra $\mathfrak{g}$ as $M \times \mathfrak{g}$	Lie algebroid $E \rightarrow N$
$\mathfrak{g}$ -action $\gamma$	Anchor $\rho$ of $E$ & $E$ -connections
Canonical flat connection $\nabla^0$ on $M \times \mathfrak{g}$	General connection $\nabla$ on $E$

# Guide: Curved Yang-Mills-Higgs gauge theory

Classical formalism	CYMH GT
Lie algebra $\mathfrak{g}$ as $M \times \mathfrak{g}$	Lie algebroid $E \rightarrow N$
$\mathfrak{g}$ -action $\gamma$	Anchor $\rho$ of $E$ & $E$ -connections
Canonical flat connection $\nabla^0$ on $M \times \mathfrak{g}$	General connection $\nabla$ on $E$

## Remarks (Why a "curved theory"?)

Usually, the field strength  $F$  is given by (abelian, for simplicity)

$$F := dA = d^{\nabla^0} A.$$

$\leadsto$  We will use a general connection  $\nabla$  instead of  $\nabla^0$ , and  $\nabla$  may not be flat.

## Definition (Lie algebroids)

Let  $E \rightarrow N$  be a vector bundle. Then  $E$  is a Lie algebroid, if there is a bundle map  $\rho : E \rightarrow TN$ , called the **anchor**, and a Lie algebra structure on  $\Gamma(E)$  with Lie bracket  $[\cdot, \cdot]_E$  satisfying

$$[\mu, f\nu]_E = f[\mu, \nu]_E + \mathcal{L}_{\rho(\mu)}(f) \nu \quad (1)$$

for all  $f \in C^\infty(N)$  and  $\mu, \nu \in \Gamma(E)$ .

## Example

- $E = TN, \rho = \mathbb{1}_{TN}$

## Definition (Lie algebroids)

Let  $E \rightarrow N$  be a vector bundle. Then  $E$  is a Lie algebroid, if there is a bundle map  $\rho : E \rightarrow TN$ , called the **anchor**, and a Lie algebra structure on  $\Gamma(E)$  with Lie bracket  $[\cdot, \cdot]_E$  satisfying

$$[\mu, f\nu]_E = f[\mu, \nu]_E + \mathcal{L}_{\rho(\mu)}(f) \nu \quad (1)$$

for all  $f \in C^\infty(N)$  and  $\mu, \nu \in \Gamma(E)$ .

## Example

- $E = TN$ ,  $\rho = \mathbb{1}_{TN}$
- $E$  a bundle of Lie algebras,  $\rho = 0$

## Definition (Lie algebroids)

Let  $E \rightarrow N$  be a vector bundle. Then  $E$  is a Lie algebroid, if there is a bundle map  $\rho : E \rightarrow TN$ , called the **anchor**, and a Lie algebra structure on  $\Gamma(E)$  with Lie bracket  $[\cdot, \cdot]_E$  satisfying

$$[\mu, f\nu]_E = f[\mu, \nu]_E + \mathcal{L}_{\rho(\mu)}(f) \nu \quad (1)$$

for all  $f \in C^\infty(N)$  and  $\mu, \nu \in \Gamma(E)$ .

## Example

- $E = TN$ ,  $\rho = \mathbb{1}_{TN}$
- $E$  a bundle of Lie algebras,  $\rho = 0$



As for an action  $\gamma : \mathfrak{g} \rightarrow \mathfrak{X}(N)$  we have:

**Proposition (Anchor as homomorphism)**

$\rho : \Gamma(E) \rightarrow \mathfrak{X}(N)$  is a homomorphism of Lie algebras.

The classical formalism will correspond to:

### Proposition (Action Lie algebroids)

Let  $(\mathfrak{g}, [\cdot, \cdot]_{\mathfrak{g}})$  be a Lie algebra with a  $\mathfrak{g}$ -action  $\gamma$  on  $N$ . Then there is a **unique** Lie algebroid structure on  $E := N \times \mathfrak{g}$  as a vector bundle over  $N$  such that

$$\rho(\nu) = \gamma(\nu), \quad (2)$$

$$[\mu, \nu]_E = [\mu, \nu]_{\mathfrak{g}} \quad (3)$$

for all constant sections  $\mu, \nu \in \Gamma(E)$ .

Classical formalism	CYMH GT
Lie algebra $\mathfrak{g}$ as $M \times \mathfrak{g}$	Lie algebroid $E \rightarrow N$
$\mathfrak{g}$ -action $\gamma$	Anchor $\rho$ of $E$ & <b><math>E</math>-connections</b> as lift of $\rho$
Canonical flat connection $\nabla^0$ on $M \times \mathfrak{g}$	General connection $\nabla$ on $E$

$\rightsquigarrow$   $E$ -connections similar to vector bundle connections but Leibniz rule along  $\rho$ , i.e. only lifting vector fields in the image of  $\rho$

Classical formalism	CYMH GT
Lie algebra $\mathfrak{g}$ as $M \times \mathfrak{g}$	Lie algebroid $E \rightarrow N$
$\mathfrak{g}$ -action $\gamma$	Anchor $\rho$ of $E$ & <b><math>E</math>-connections</b> as lift of $\rho$
Canonical flat connection $\nabla^0$ on $M \times \mathfrak{g}$	General connection $\nabla$ on $E$

$\rightsquigarrow$   $E$ -connections similar to vector bundle connections but Leibniz rule along  $\rho$ , i.e. only lifting vector fields in the image of  $\rho$

### Example ( $E$ -connection ${}^E\nabla$ on $V$ )

$\nabla'$  a vector bundle connection on  $V \rightarrow N$ , then

$${}^E\nabla_\nu v := \nabla'_{\rho(\nu)} v \quad (4)$$

for all  $\nu \in \Gamma(E)$  and  $v \in \Gamma(V)$ . In short denoted by  $\nabla'_\rho$ .

## Example

For  $\nabla$  a connection on  $E$  we have the **basic connection** given as a pair of  $E$ -connections on  $E$  and on  $\mathbb{T}N$  by

$$\nabla_{\nu}^{\text{bas}} \mu := [\nu, \mu]_E + \nabla_{\rho(\mu)} \nu, \quad (5)$$

$$\nabla_{\nu}^{\text{bas}} X := [\rho(\nu), X] + \rho(\nabla_X \nu) \quad (6)$$

for all  $X \in \mathfrak{X}(N)$  and  $\nu, \mu \in \Gamma(E)$ .

## Remarks (Encoding of Lie algebra representations)

Test this with trivial bundles and canonical flat connection  $\nabla^0$ ,  
i.e.  $E = N \times \mathfrak{g}$  and  $\nabla^0 \nu = 0$  for constant sections  $\nu$ .

## Example

For  $\nabla$  a connection on  $E$  we have the **basic connection** given as a pair of  $E$ -connections on  $E$  and on  $TN$  by

$$\nabla_{\nu}^{\text{bas}} \mu := [\nu, \mu]_E + \nabla_{\rho(\mu)} \nu, \quad (5)$$

$$\nabla_{\nu}^{\text{bas}} X := [\rho(\nu), X] + \rho(\nabla_X \nu) \quad (6)$$

for all  $X \in \mathfrak{X}(N)$  and  $\nu, \mu \in \Gamma(E)$ .

## Remarks (Encoding of Lie algebra representations)

Test this with trivial bundles and canonical flat connection  $\nabla^0$ , i.e.  $E = N \times \mathfrak{g}$  and  $\nabla^0 \nu = 0$  for constant sections  $\nu$ .

## Definition (Basic curvature)

Let  $\nabla$  be a connection on  $E$ . The **basic curvature**  $R_{\nabla}^{\text{bas}}$  is defined as an element of  $\Gamma(\wedge^2 E^* \otimes T^*N \otimes E)$  by

$$R_{\nabla}^{\text{bas}}(\mu, \nu)X := \nabla_X([\mu, \nu]_E) - [\nabla_X \mu, \nu]_E - [\mu, \nabla_X \nu]_E - \nabla_{\nabla_{\nu}^{\text{bas}} X} \mu + \nabla_{\nabla_{\mu}^{\text{bas}} X} \nu, \quad (7)$$

where  $\mu, \nu \in \Gamma(E)$  and  $X \in \mathfrak{X}(N)$ .

## Proposition

*We recover the curvature of the basic connection:*

$$R_{\nabla^{\text{bas}}} = \begin{cases} -R_{\nabla}^{\text{bas}}(\cdot, \cdot) \circ \rho, & \text{on } E, \\ -\rho \circ R_{\nabla}^{\text{bas}}, & \text{on } TN. \end{cases} \quad (8)$$



## Definition (Basic curvature)

Let  $\nabla$  be a connection on  $E$ . The **basic curvature**  $R_{\nabla}^{\text{bas}}$  is defined as an element of  $\Gamma(\wedge^2 E^* \otimes T^*N \otimes E)$  by

$$R_{\nabla}^{\text{bas}}(\mu, \nu)X := \nabla_X([\mu, \nu]_E) - [\nabla_X \mu, \nu]_E - [\mu, \nabla_X \nu]_E - \nabla_{\nabla_{\nu}^{\text{bas}} X} \mu + \nabla_{\nabla_{\mu}^{\text{bas}} X} \nu, \quad (7)$$

where  $\mu, \nu \in \Gamma(E)$  and  $X \in \mathfrak{X}(N)$ .

## Proposition

*We recover the curvature of the basic connection:*

$$R_{\nabla^{\text{bas}}} = \begin{cases} -R_{\nabla}^{\text{bas}}(\cdot, \cdot) \circ \rho, & \text{on } E, \\ -\rho \circ R_{\nabla}^{\text{bas}}, & \text{on } TN. \end{cases} \quad (8)$$

## Definition (Space of fields)

Fields are a pair consisting of:

- Higgs field  $\Phi \in C^\infty(M; N)$
- Field of gauge bosons  $A \in \Omega^1(M; \Phi^* E)$

## Definition (Space of fields)

Fields are a pair consisting of:

- Higgs field  $\Phi \in C^\infty(M; N)$
- Field of gauge bosons  $A \in \Omega^1(M; \Phi^* E)$

## Definition (Space of fields)

Fields are a pair consisting of:

- Higgs field  $\Phi \in C^\infty(M; N)$
- Field of gauge bosons  $A \in \Omega^1(M; \Phi^* E)$

## Definition (Minimal coupling)

**Minimal coupling**  $\mathfrak{D}$ ,  $(\Phi, A) \mapsto \mathfrak{D}(\Phi, A) \in \Omega^1(M; \Phi^* TN)$ , by

$$\mathfrak{D}(\Phi, A) := \mathfrak{D}^A \Phi := D\Phi - (\Phi^* \rho)(A), \quad (9)$$

where  $D\Phi : TM \rightarrow TN$  is the tangent map.

## Definition (Space of fields)

Fields are a pair consisting of:

- Higgs field  $\Phi \in C^\infty(M; N)$
- Field of gauge bosons  $A \in \Omega^1(M; \Phi^* E)$

## Definition (Minimal coupling)

**Minimal coupling**  $\mathfrak{D}$ ,  $(\Phi, A) \mapsto \mathfrak{D}(\Phi, A) \in \Omega^1(M; \Phi^* TN)$ , by

$$\mathfrak{D}(\Phi, A) := \mathfrak{D}^A \Phi := D\Phi - (\Phi^* \rho)(A), \quad (9)$$

where  $D\Phi : TM \rightarrow TN$  is the tangent map.

## Definition (Field strength)

Let  $\nabla$  be a connection on  $E$ . We define the **field strength**  $F$ ,  $(\Phi, A) \mapsto F(\Phi, A) \in \Omega^2(M; \Phi^*E)$ , by

$$F(\Phi, A) := d^{\Phi^*\nabla} A + \frac{1}{2}(\Phi^* t_{\nabla^{\text{bas}}})(A \wedge A), \quad (10)$$

where  $t_{\nabla^{\text{bas}}}$  is the torsion of  $\nabla^{\text{bas}}$  on  $E$  and  $d^{\Phi^*\nabla}$  the exterior covariant derivative of  $\Phi^*\nabla$ .

## Definition (Field strength)

Let  $\nabla$  be a connection on  $E$ . We define the **field strength**  $F$ ,  $(\Phi, A) \mapsto F(\Phi, A) \in \Omega^2(M; \Phi^*E)$ , by

$$F(\Phi, A) := d^{\Phi^*\nabla} A + \frac{1}{2}(\Phi^* t_{\nabla^{\text{bas}}})(A \wedge A), \quad (10)$$

where  $t_{\nabla^{\text{bas}}}$  is the torsion of  $\nabla^{\text{bas}}$  on  $E$  and  $d^{\Phi^*\nabla}$  the exterior covariant derivative of  $\Phi^*\nabla$ .

## Definition (Generalised field strength)

Let  $\zeta$  be an element of  $\Omega^2(N; E)$ , then we define the **generalised field strength**  $G$  by

$$G(\Phi, A) := F(\Phi, A) + \frac{1}{2}(\Phi^*\zeta)(\mathfrak{D}^A\Phi \wedge \mathfrak{D}^A\Phi). \quad (11)$$

## Definition (Curved Yang-Mills-Higgs (CYMH) Lagrangian)

Let  $\kappa$  be a fibre metric on  $E$ , then the **curved Yang-Mills-Higgs Lagrangian**  $\mathfrak{L}_{\text{CYMH}}$ ,  $(\Phi, A) \mapsto \mathfrak{L}_{\text{CYMH}}(\Phi, A) \in \Omega^{\dim(M)}(M)$ , is defined by

$$\begin{aligned} \mathfrak{L}_{\text{CYMH}}(\Phi, A) := & -\frac{1}{2}(\Phi^*\kappa)(G(\Phi, A) \wedge *G(\Phi, A)) \\ & + (\Phi^*g)\left(\mathfrak{D}^A\Phi \wedge *\mathfrak{D}^A\Phi\right) - *(\Phi^*\mathcal{V}), \quad (12) \end{aligned}$$

where  $*$  is the Hodge star operator related to the spacetime metric  $\eta$ .



## Definition (CYMH GT)

Assume we have additionally the **compatibility conditions**

$$R_{\nabla} + d^{\nabla^{\text{bas}}} \zeta = 0, \quad (13)$$

$$R_{\nabla}^{\text{bas}} = 0, \quad (14)$$

$$\nabla^{\text{bas}} \kappa = 0, \quad (15)$$

$$\nabla^{\text{bas}} g = 0, \quad (16)$$

$$\mathcal{L}_{\rho} \mathcal{V} = 0, \quad (17)$$

then we say that we have a **curved Yang-Mills-Higgs gauge theory**.

We say that we have a **pre-classical gauge theory**, if  $\nabla$  is flat.

## Definition (CYMH GT)

Assume we have additionally the **compatibility conditions**

$$R_{\nabla} + d^{\nabla^{\text{bas}}} \zeta = 0, \quad (13)$$

$$R_{\nabla}^{\text{bas}} = 0, \quad (14)$$

$$\nabla^{\text{bas}} \kappa = 0, \quad (15)$$

$$\nabla^{\text{bas}} g = 0, \quad (16)$$

$$\mathcal{L}_{\rho} \mathcal{V} = 0, \quad (17)$$

then we say that we have a **curved Yang-Mills-Higgs gauge theory**.

We say that we have a **pre-classical gauge theory**, if  $\nabla$  is flat.

If we have additionally  $\zeta \equiv 0$ , then we say that we have a **classical gauge theory**.

## Definition (CYMH GT)

Assume we have additionally the **compatibility conditions**

$$R_{\nabla} + d^{\nabla^{\text{bas}}} \zeta = 0, \quad (13)$$

$$R_{\nabla}^{\text{bas}} = 0, \quad (14)$$

$$\nabla^{\text{bas}} \kappa = 0, \quad (15)$$

$$\nabla^{\text{bas}} g = 0, \quad (16)$$

$$\mathcal{L}_{\rho} \mathcal{V} = 0, \quad (17)$$

then we say that we have a **curved Yang-Mills-Higgs gauge theory**.

We say that we have a **pre-classical gauge theory**, if  $\nabla$  is flat.

If we have additionally  $\zeta \equiv 0$ , then we say that we have a **classical gauge theory**.

# Summary

$$\begin{array}{ccccc}
 \Phi^* E & & (E, [\cdot, \cdot]_E, \kappa, \nabla) & \xleftarrow{\zeta \in \Omega^2} & \\
 \downarrow & & \downarrow & \searrow \rho & \\
 M & \xrightarrow{\Phi} & (N, g) & \xleftarrow{\quad} & TN
 \end{array}$$

$\rightsquigarrow$  Together with the compatibility conditions we will have gauge invariance, that is,

$$\delta_\varepsilon \mathfrak{L}_{\text{CYMH}} = 0. \quad (18)$$

## Motivation

- Are there CYMH GTs which are neither pre-classical nor classical?
- Difficulty: There is an equivalence relation of CYMH GTs keeping the same Lagrangian and preserving the physics, possibly turning theories into (pre-)classical ones.

This was provided by Edward Witten in a private communication with Thomas Strobl about a specific example of a CYMH GT.

## Motivation

- Are there CYMH GTs which are neither pre-classical nor classical?
- Difficulty: There is an equivalence relation of CYMH GTs keeping the same Lagrangian and preserving the physics, possibly turning theories into (pre-)classical ones.

This was provided by Edward Witten in a private communication with Thomas Strobl about a specific example of a CYMH GT.

## Definition (Field redefinition, [S.-R. F.])

Let  $\lambda \in \Omega^1(N; E)$  such that  $\Lambda := \mathbb{1}_E - \lambda \circ \rho$  is an automorphism of  $E$ . We then define the **field redefinitions** by

$$\tilde{A}^\lambda := (\Phi^* \Lambda)(A) + \Phi^! \lambda, \quad (19)$$

$$\tilde{\nabla}^\lambda := \nabla + \left( \Lambda \circ d^{\nabla^{\text{bas}}} \circ \Lambda^{-1} \right) \lambda, \quad (20)$$

$$\tilde{\kappa}^\lambda := \kappa \circ \left( \Lambda^{-1}, \Lambda^{-1} \right), \quad (21)$$

$$\tilde{g}^\lambda := g \circ \left( \hat{\Lambda}^{-1}, \hat{\Lambda}^{-1} \right), \quad (22)$$

where  $\hat{\Lambda} := \mathbb{1}_{TN} - \rho \circ \lambda$ , and for all  $X, Y \in \mathfrak{X}(N)$  we also define

$$\begin{aligned} & \tilde{\zeta}^\lambda \left( \hat{\Lambda}(X), \hat{\Lambda}(Y) \right) \\ & := \Lambda(\zeta(X, Y)) - \left( d^{\tilde{\nabla}^\lambda} \lambda \right) (X, Y) + t_{\tilde{\nabla}^\lambda \rho}(\lambda(X), \lambda(Y)). \end{aligned} \quad (23)$$

## Proposition ([S.-R. F.]

- *Field redefinitions define an equivalence relation of CYMH gauge theories*
- $\tilde{\mathcal{L}}_{\text{CYMH}}^\lambda = \mathcal{L}_{\text{CYMH}}$

Let us now apply a field redefinition in order to study whether  $\nabla$  and  $\zeta$  can become flat and zero, respectively.



## Proposition ([S.-R. F.]

- *Field redefinitions define an equivalence relation of CYMH gauge theories*
- $\tilde{\mathcal{L}}_{\text{CYMH}}^\lambda = \mathcal{L}_{\text{CYMH}}$

Let us now apply a field redefinition in order to study whether  $\nabla$  and  $\zeta$  can become flat and zero, respectively.

# What happens in the case of Lie algebra bundles?

## Example (Lie algebra bundles (LABs))

- $E = \mathfrak{g}$  an LAB ( $\rho \equiv 0$ ) with a field of Lie brackets  $[\cdot, \cdot]_{\mathfrak{g}} \in \Gamma(\wedge^2 \mathfrak{g}^* \otimes \mathfrak{g})$  which restricts to the bracket of a given Lie algebra  $\mathfrak{g}$

# What happens in the case of Lie algebra bundles?

## Example (Lie algebra bundles (LABs))

- $E = \mathfrak{g}$  an LAB ( $\rho \equiv 0$ ) with a field of Lie brackets  $[\cdot, \cdot]_{\mathfrak{g}} \in \Gamma(\wedge^2 \mathfrak{g}^* \otimes \mathfrak{g})$  which restricts to the bracket of a given Lie algebra  $\mathfrak{g}$

Compatibilities:

- $\kappa$  needs to be ad-invariant

# What happens in the case of Lie algebra bundles?

## Example (Lie algebra bundles (LABs))

- $E = \mathcal{g}$  an LAB ( $\rho \equiv 0$ ) with a field of Lie brackets  $[\cdot, \cdot]_{\mathcal{g}} \in \Gamma(\wedge^2 \mathcal{g}^* \otimes \mathcal{g})$  which restricts to the bracket of a given Lie algebra  $\mathfrak{g}$

Compatibilities:

- $\kappa$  needs to be ad-invariant
- We need

$$\nabla_Y([\mu, \nu]_{\mathcal{g}}) = [\nabla_Y \mu, \nu]_{\mathcal{g}} + [\mu, \nabla_Y \nu]_{\mathcal{g}}, \quad (24)$$

$$R_{\nabla}(Y, Z)\mu = [\zeta(Y, Z), \mu]_{\mathcal{g}} \quad (25)$$

for all  $Y, Z \in \mathfrak{X}(N)$  and  $\mu, \nu \in \Gamma(\mathcal{g})$ .

# What happens in the case of Lie algebra bundles?

## Example (Lie algebra bundles (LABs))

- $E = \mathcal{G}$  an LAB ( $\rho \equiv 0$ ) with a field of Lie brackets  $[\cdot, \cdot]_{\mathcal{G}} \in \Gamma(\wedge^2 \mathcal{G}^* \otimes \mathcal{G})$  which restricts to the bracket of a given Lie algebra  $\mathfrak{g}$

Compatibilities:

- $\kappa$  needs to be ad-invariant
- We need

$$\nabla_Y([\mu, \nu]_{\mathcal{G}}) = [\nabla_Y \mu, \nu]_{\mathcal{G}} + [\mu, \nabla_Y \nu]_{\mathcal{G}}, \quad (24)$$

$$R_{\nabla}(Y, Z)\mu = [\zeta(Y, Z), \mu]_{\mathcal{G}} \quad (25)$$

for all  $Y, Z \in \mathfrak{X}(N)$  and  $\mu, \nu \in \Gamma(\mathcal{G})$ .

### Theorem (Invariant for LABs, [S.-R. F.])

*We have*

$$d^{\tilde{\nabla}^\lambda} \tilde{\zeta}^\lambda = d^\nabla \zeta, \quad (26)$$

*and  $d^\nabla \zeta$  has values in the centre of  $\mathfrak{g}$ .*

# Behaviour of the field redefinition of $\zeta$

Theorem (Existence of non-classical theories, [S.-R. F.])

*If  $d^\nabla \zeta \neq 0$ , then there is no field redefinition such that  $\tilde{\zeta}^\lambda = 0$ .*

## Remarks

Starting with a classical theory:

If  $\dim(N) \geq 3$  and if Lie algebra  $\mathfrak{g}$  has a non-zero centre, then we can always construct a pre-classical CYMH GT which is not a classical one by adding a  $\zeta$  with  $d^\nabla \zeta \neq 0$ .

# Behaviour of the field redefinition of $\zeta$

Theorem (Existence of non-classical theories, [S.-R. F.])

*If  $d^\nabla \zeta \neq 0$ , then there is no field redefinition such that  $\tilde{\zeta}^\lambda = 0$ .*

## Remarks

Starting with a classical theory:

If  $\dim(N) \geq 3$  and if Lie algebra  $\mathfrak{g}$  has a non-zero centre, then we can always construct a pre-classical CYMH GT which is not a classical one by adding a  $\zeta$  with  $d^\nabla \zeta \neq 0$ .

However, by  $R_\nabla = \text{ad}_g \circ \zeta$  it may still be that  $\nabla$  becomes flat.



# Behaviour of the field redefinition of $\zeta$

Theorem (Existence of non-classical theories, [S.-R. F.])

*If  $d^\nabla \zeta \neq 0$ , then there is no field redefinition such that  $\tilde{\zeta}^\lambda = 0$ .*

## Remarks

Starting with a classical theory:

If  $\dim(N) \geq 3$  and if Lie algebra  $\mathfrak{g}$  has a non-zero centre, then we can always construct a pre-classical CYMH GT which is not a classical one by adding a  $\zeta$  with  $d^\nabla \zeta \neq 0$ .

However, by  $R_\nabla = \text{ad}_g \circ \zeta$  it may still be that  $\nabla$  becomes flat.

# Turning to the field redefinition of $\nabla$ :

Theorem (Differential on centre-valued forms, [S.-R. F.])

$\nabla$  restricts to the centre of  $\mathfrak{g}$  and induces a differential  $d^\Xi$  on centre-valued forms. Moreover,  $d^\Xi$  is independent of the field redefinitions.

Sketch of proof.

Recall

$$\begin{aligned}\nabla_Y([\mu, \nu]_{\mathfrak{g}}) &= [\nabla_Y \mu, \nu]_{\mathfrak{g}} + [\mu, \nabla_Y \nu]_{\mathfrak{g}}, \\ R_\nabla(Y, Z)\mu &= [\zeta(Y, Z), \mu]_{\mathfrak{g}}, \\ \tilde{\nabla}_Y^\lambda \mu &= \nabla_Y \mu - [\lambda(Y), \mu]_{\mathfrak{g}},\end{aligned}$$

for all  $Y, Z \in \mathfrak{X}(N)$  and  $\mu, \nu \in \Gamma(\mathfrak{g})$ . Then insert  $\mu$  with values in the centre.

# Turning to the field redefinition of $\nabla$ :

Theorem (Differential on centre-valued forms, [S.-R. F.])

$\nabla$  restricts to the centre of  $\mathfrak{g}$  and induces a differential  $d^\Xi$  on centre-valued forms. Moreover,  $d^\Xi$  is independent of the field redefinitions.

Sketch of proof.

Recall

$$\nabla_Y([\mu, \nu]_{\mathfrak{g}}) = [\nabla_Y \mu, \nu]_{\mathfrak{g}} + [\mu, \nabla_Y \nu]_{\mathfrak{g}},$$

$$R_\nabla(Y, Z)\mu = [\zeta(Y, Z), \mu]_{\mathfrak{g}},$$

$$\tilde{\nabla}_Y^\lambda \mu = \nabla_Y \mu - [\lambda(Y), \mu]_{\mathfrak{g}},$$

for all  $Y, Z \in \mathfrak{X}(N)$  and  $\mu, \nu \in \Gamma(\mathfrak{g})$ . Then insert  $\mu$  with values in the centre.

## Theorem (Closedness of $d^\nabla \zeta$ , [S.-R. F.])

*We have*

$$d^\Xi d^\nabla \zeta = 0. \quad (27)$$

## Definition (Obstruction class, [S.-R. F.])

We define the **obstruction class** by

$$\text{Obs}(\Xi) := [d^\nabla \zeta]_{d^\Xi}. \quad (28)$$

### Theorem (Closedness of $d^\nabla \zeta$ , [S.-R. F.])

We have

$$d^\Xi d^\nabla \zeta = 0. \quad (27)$$

### Definition (Obstruction class, [S.-R. F.])

We define the **obstruction class** by

$$\text{Obs}(\Xi) := [d^\nabla \zeta]_{d^\Xi}. \quad (28)$$

### Proposition ([S.-R. F.])

- *An invariant of the field redefinitions.*

Theorem (Closedness of  $d^\nabla \zeta$ , [S.-R. F.])

We have

$$d^\Xi d^\nabla \zeta = 0. \quad (27)$$

## Definition (Obstruction class, [S.-R. F.])

We define the **obstruction class** by

$$\text{Obs}(\Xi) := [d^\nabla \zeta]_{d^\Xi}. \quad (28)$$

## Proposition ([S.-R. F.])

- *An invariant of the field redefinitions.*
- *If  $\nabla$  flat, then  $\text{Obs}(\Xi) = 0$ .*

Theorem (Closedness of  $d^\nabla \zeta$ , [S.-R. F.])

We have

$$d^\Xi d^\nabla \zeta = 0. \quad (27)$$

## Definition (Obstruction class, [S.-R. F.])

We define the **obstruction class** by

$$\text{Obs}(\Xi) := [d^\nabla \zeta]_{d^\Xi}. \quad (28)$$

## Proposition ([S.-R. F.])

- *An invariant of the field redefinitions.*
- *If  $\nabla$  flat, then  $\text{Obs}(\Xi) = 0$ .*

# Theorem (Obstruction for non-pre-classical theories, [S.-R. F.]

*If  $\text{Obs}(\Xi) \neq 0$ , then there is no field redefinition such that  $\tilde{\nabla}^\lambda$  is flat.*

# Theorem (Locally always pre-classical)

*If  $N$  is contractible, then there is a field redefinition such that  $\tilde{\nabla}^\lambda$  is flat.*



### Theorem (Obstruction for non-pre-classical theories, [S.-R. F.] )

*If  $\text{Obs}(\Xi) \neq 0$ , then there is no field redefinition such that  $\tilde{\nabla}^\lambda$  is flat.*

### Theorem (Locally always pre-classical)

*If  $N$  is contractible, then there is a field redefinition such that  $\tilde{\nabla}^\lambda$  is flat.*

### Remarks

Second theorem follows as a result by K. Mackenzie (General Theory of Lie Groupoids and Algebroids. *London Mathematical Society Lecture Note Series*, 213, 2005). Mackenzie derived  $\text{Obs}(\Xi)$  in the context of extending Lie algebroids by LABs.

### Theorem (Obstruction for non-pre-classical theories, [S.-R. F.] )

*If  $\text{Obs}(\Xi) \neq 0$ , then there is no field redefinition such that  $\tilde{\nabla}^\lambda$  is flat.*

### Theorem (Locally always pre-classical)

*If  $N$  is contractible, then there is a field redefinition such that  $\tilde{\nabla}^\lambda$  is flat.*

### Remarks

Second theorem follows as a result by K. Mackenzie (General Theory of Lie Groupoids and Algebroids. *London Mathematical Society Lecture Note Series*, 213, 2005). Mackenzie derived  $\text{Obs}(\Xi)$  in the context of extending Lie algebroids by LABs.

## Example (Zero obstruction class not necessarily pre-classical)

Let  $P$  be the Hopf fibration

$$\begin{array}{ccc} \mathrm{SU}(2) & \longrightarrow & \mathbb{S}^7 \\ & & \downarrow \\ & & \mathbb{S}^4 \end{array}$$

Then for the adjoint bundle

$$\mathcal{Q} := P \times_{\mathrm{SU}(2)} \mathfrak{su}(2) := \left( \mathbb{S}^7 \times \mathfrak{su}(2) \right) / \mathrm{SU}(2)$$

we have a non-flat  $\nabla$  satisfying the compatibility conditions such that all of its field redefinitions are not flat either, but  $\mathrm{Obs}(\Xi) = 0$ .

# Summary

## Remarks

Locally, LABs are always pre-classical but not necessarily classical.  
In general,  $\text{Obs}(\Xi) = 0$  does not imply a flat connection.

So, what actually happens in the adjoint bundle of  $S^7 \rightarrow S^4$ ?  
 $\rightsquigarrow$  Integration

# Summary

## Remarks

Locally, LABs are always pre-classical but not necessarily classical.  
In general,  $\text{Obs}(\Xi) = 0$  does not imply a flat connection.

So, what actually happens in the adjoint bundle of  $S^7 \rightarrow S^4$ ?  
 $\rightsquigarrow$  Integration

# Idea

## Remarks

Observe that the Higgs field  $\Phi$  doesn't really matter at all in the case of LABs

$\rightsquigarrow$  Think in terms of bundles over  $M$  directly

	Classical	Curved
Infinitesimal	Lie algebra $\mathfrak{g}$	LAB $\mathcal{G}$
Integrated	Lie group $G$	LGB <sup>1</sup> $\mathcal{G}$ ?

$$\begin{array}{ccc} G & \longrightarrow & \mathcal{G} \\ & & \downarrow \pi \\ & & M \end{array}$$

---

<sup>1</sup>LGB = Lie group bundle

# Idea

## Remarks

Observe that the Higgs field  $\Phi$  doesn't really matter at all in the case of LABs

$\rightsquigarrow$  Think in terms of bundles over  $M$  directly

	Classical	Curved
Infinitesimal	Lie algebra $\mathfrak{g}$	LAB $\mathcal{G}$
Integrated	Lie group $G$	LGB <sup>1</sup> $\mathcal{G}?$

$$\begin{array}{ccc} G & \longrightarrow & \mathcal{G} \\ & & \downarrow \pi \\ & & M \end{array}$$

---

<sup>1</sup>LGB = Lie group bundle

# LGB actions

## Definition

$$\begin{array}{ccc} & \mathcal{G} & \\ & \downarrow \pi & \\ N & \xrightarrow{f} & M \end{array}$$

A **right-action of  $\mathcal{G}$  on  $N$**  is a smooth map  $N * \mathcal{G} := f^* \mathcal{G} \rightarrow N$ ,  $(p, g) \mapsto p \cdot g$ , satisfying the following properties:

$$f(p \cdot g) = \pi(g), \quad (29)$$

$$(p \cdot g) \cdot h = p \cdot (gh), \quad (30)$$

$$p \cdot e_{f(p)} = p \quad (31)$$

for all  $p \in N$  and  $g, h \in \mathcal{G}_{f(p)}$ , where  $e_{f(p)}$  is the neutral element of  $\mathcal{G}_{f(p)}$ .



# Examples

## Example

$\mathcal{G}$  acts canonically on itself:

$$\begin{aligned}\mathcal{G} * \mathcal{G} &\rightarrow \mathcal{G}, \\ (q, h) &\mapsto qh.\end{aligned}$$

## Example (Recovering Lie group action)

- Either by  $M = \{*\}$ .
- Or by  $\mathcal{G} \cong M \times G$ , then also  $N * \mathcal{G} \cong N \times G$ , and we can define

$$\begin{aligned}N \times G &\rightarrow N, \\ (p, g) &\mapsto p \cdot g := p \cdot (f(p), g),\end{aligned}$$

which is equivalent to  $N * \mathcal{G} \rightarrow N$ .

# Examples

## Example

$\mathcal{G}$  acts canonically on itself:

$$\begin{aligned}\mathcal{G} * \mathcal{G} &\rightarrow \mathcal{G}, \\ (q, h) &\mapsto qh.\end{aligned}$$

## Example (Recovering Lie group action)

- Either by  $M = \{*\}$ .
- Or by  $\mathcal{G} \cong M \times G$ , then also  $N * \mathcal{G} \cong N \times G$ , and we can define

$$\begin{aligned}N \times G &\rightarrow N, \\ (p, g) &\mapsto p \cdot g := p \cdot (f(p), g),\end{aligned}$$

which is equivalent to  $N * \mathcal{G} \rightarrow N$ .

# Principal bundle

Still a fibre bundle

$$\begin{array}{ccc} G & \longrightarrow & \mathcal{P} \\ & & \downarrow \\ & & M \end{array}$$

but with  $\mathcal{G}$ -action

$$\begin{array}{ccc} \cancel{\mathcal{P}} \times \cancel{G} & \longrightarrow & \mathcal{P} \\ \mathcal{P} * \mathcal{G} & & \end{array}$$

# Summary: When are theories flat?

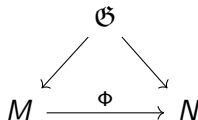
	Locally	Globally
LABs	Pre-classical	$\text{ad}(\mathbb{S}^7 \rightarrow \mathbb{S}^4)$ curved
Tangent bundles	Pre-classical	$T\mathbb{S}^7$ curved

## Remarks

What about general Lie algebroids? Find an invariant like  $d^\nabla \zeta$   
 $\rightsquigarrow d^\nabla \zeta$  measures the failure of the Bianchi identity of  $G$   
 $\rightsquigarrow$  Generalised Bianchi identity?

# Hope: Structural Lie groupoids

Gauge theory	Structure
Yang-Mills	Lie group $G$
Curved Yang-Mills	Lie group bundle $\mathcal{G}$
Curved Yang-Mills-Higgs	Lie groupoid $\mathcal{G}?$



**Thank you!**