

# Quantum Chemistry on Quantum Computers

## #1 General Introduction

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# Self introduction

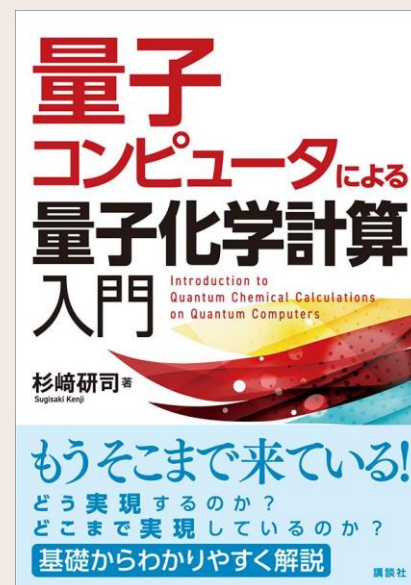
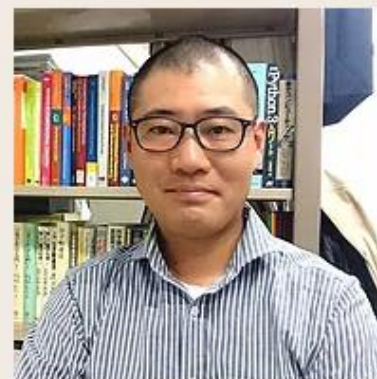
## Kenji Sugisaki

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&  
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- JST PRESTO (2019 Oct.–2023 Mar.)
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- Special Appointment Assistant Professor, Osaka City University (2013–2016)
- Visiting Postdoctoral Researcher, Harvard University (2011 Feb.–2011 Apr.)
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- Research Associate, Osaka City University (2006–2008)
- DSci, Osaka City University (2006)
- MSc, Osaka City University (2004)
- BSc, Osaka City University (2003)



## Schedule of this lecture series

#	Date Time in IST	Theme
1	10/28 (Thu) 10:00	<b>General introduction of quantum computers</b> History of quantum computers, qubits, quantum gates, quantum circuits, universal gate sets, NISQ devices, quantum error correction, etc.
2	11/11 (Thu) 10:00	<b>Quantum algorithms</b> Deutsch–Jozsa algorithm, quantum teleportation, quantum fingerprinting (SWAP test), etc.
3	11/17 (Wed) 10:00	<b>Quantum phase estimation algorithm</b> Hadamard test (1-qubit QPE), iterative QPE, Bayesian QPE, relationship between evolution time length and phase precisions, resource estimations, etc.
4	11/25 (Thu) 10:00	<b>Fermion–qubit transformation and quantum circuit constructions</b> Jordan–Wigner transformation, parity basis, Bravyi–Kitaev transformation, Trotter decomposition, etc.
5	12/2 (Thu) 10:00	<b>Variational quantum algorithms</b> Chemistry-inspired ansatzes, hardware-efficient ansatzes, error mitigations, excited states calculations, resource estimations, etc.
6	12/9 (Thu) 10:00	<b>Adiabatic quantum algorithms</b> Adiabatic theorem, quantum annealing, adiabatic state preparation, etc.
7	12/16 (Thu) 10:00	<b>Computational complexities</b> Computational complexity classes (in classical computing; P, NP, etc., and on quantum computing; BPP, BQP, QMA, etc)
8	12/23 (Thu) 10:00	<b>Techniques for resource (qubits/quantum gates) and error reductions</b> Symmetry adaptations, quantum circuit optimizations, Trotter decomposition error reductions, etc.

# Quantum computers

R. P. Feynman, Simulating Physics with Computers. *Int. J. Theor. Phys.* **1982**, 21, 467–488.



Photo taken from wikipedia

$\psi(x_1, x_2, \dots, x_R, t)$  with large  $R$  cannot be simulated with a normal computer, because it has too many variables.

The problem is, how can we simulate the quantum mechanics?

***Let the computer itself be built of quantum mechanical elements which obey quantum mechanical laws.***

# Quantum Turing machine

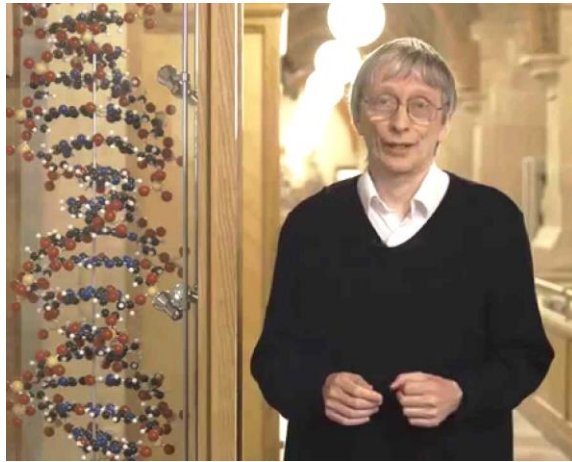


Photo taken from wikipedia

*Proc. R. Soc. Lond. A* **400**, 97–117 (1985)  
*Printed in Great Britain*

## Quantum theory, the Church–Turing principle and the universal quantum computer

BY D. DEUTSCH

*Department of Astrophysics, South Parks Road, Oxford OX1 3RQ, U.K.*

*(Communicated by R. Penrose, F.R.S. – Received 13 July 1984)*

ANNALS OF SCIENCE MAY 2, 2011 ISSUE

## DREAM MACHINE

*The mind-expanding world of quantum computing.*

By Rivka Galchen  
 May 2, 2011

<https://www.newyorker.com/magazine/2011/05/02/dream-machine>

According to Deutsch, the insight for that paper came from a conversation in the early eighties with the physicist Charles Bennett, of I.B.M., about computational-complexity theory, at the time a sexy new field that investigated the difficulty of a computational task. Deutsch questioned whether computational complexity was a fundamental or a relative property. Mass, for instance, is a fundamental property, because it remains the same in any setting; weight is a relative property, because an object's weight depends on the strength of gravity acting on it. Identical baseballs on Earth and on the moon have equivalent masses, but different weights. If computational complexity was like mass—if it was a fundamental property—then complexity was quite profound; if not, then not.



# Quantum computers

- ✓ Computers that operate according to **quantum mechanical principles**
- ✓ Use **quantum bits (qubits)** as the minimal unit of information
- ✓ Some problems those computational cost increase exponentially against the problem size on classical computer can be solved with polynomial costs

Article

*Nature* **2019**, 574, 505–510.

## Quantum supremacy using a programmable superconducting processor

Supercomputer “Summit”



Photo taken from website:  
<https://phys.org/news/2018-06-ornl-summit-supercomputer.html>

53 qubit quantum computer  
“Sycamore”

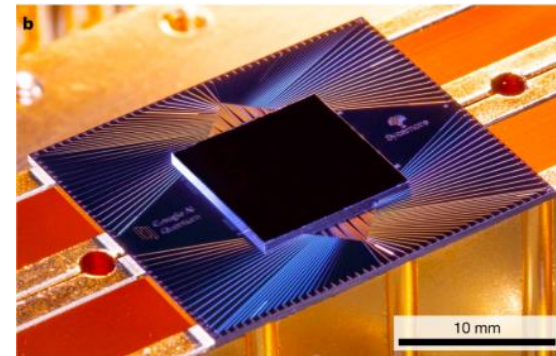


Photo taken from *Nature* **2019**, 574, 505–510.

**10,000 years**



**200 seconds**

# Misconceptions in quantum computing

*Quantum computers can solve any problems faster than the classical computers.*

Not any problems can be accelerated to solve by using quantum computers

*Quantum computers can solve particular problems faster than the classical computers.*

Still not exact.

Exponential speedup is guaranteed for some quantum algorithms (e.g., prime factorization and hidden subgroup problems). However, this states only for the **scaling of computational cost against the problem size**.

$n$ : size of problem

Classical computer

Quantum computer

$$O(a^n)$$

*Exponential scaling*

*Exponential speedup*

$$O(n^c)$$

*Polynomial scaling*

*Exponential improvement of computational cost scaling against the problem size  $n$*

# Misconceptions in quantum computing

*Exponential improvement of the computational cost scaling against the system size (number of electrons/basis sets) is possible by using quantum computers.*

Depends on the algorithm being used.

**Quantum phase estimation (QPE):** Obtain eigenfunction/eigenvalue with **polynomial cost**. On classical computer it needs exponential cost.

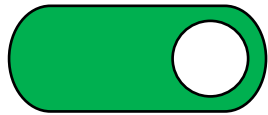
**Variational quantum eigensolver (VQE):** Obtain expectation value of approximated wave function with **(probably) polynomial cost**. On classical computer it can be solved polynomial cost, too.

	Classical computer	Quantum computer
Full-CI calculation	Exponential	Polynomial using QPE
Approximated methods (e.g., CCSD)	Polynomial	Polynomial using VQE



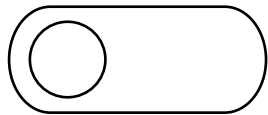
# Quantum bits (Qubits)

**Bit:** Either 1 or 0



ON

**1**



OFF

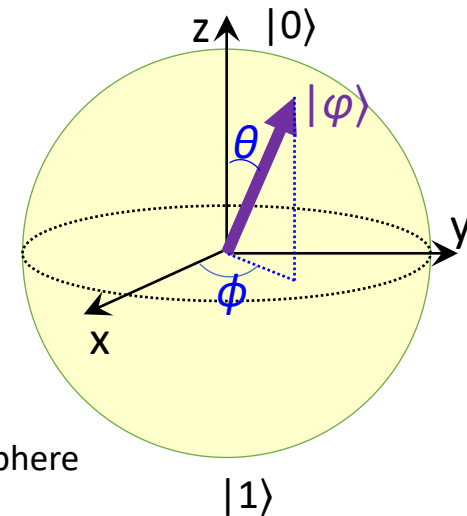
**0**

**Qubit:** Arbitrary superposition of  $|0\rangle$  and  $|1\rangle$  states

$$|\varphi\rangle = c_0|0\rangle + c_1|1\rangle$$

$c_0$  and  $c_1$  are complex numbers satisfying  $|c_0|^2 + |c_1|^2 = 1$  (normalization condition)

$$|\varphi\rangle = \cos\frac{\theta}{2}|0\rangle + e^{i\phi}\sin\frac{\theta}{2}|1\rangle$$



Bloch sphere

## Quantum bits (Qubits)

$$|\varphi\rangle = 0.6|0\rangle + 0.8|1\rangle \xrightarrow{\text{Measurement}} \begin{cases} |\varphi\rangle = |0\rangle & (36\% \text{ of probability}) \\ |\varphi\rangle = |1\rangle & (64\% \text{ of probability}) \end{cases}$$

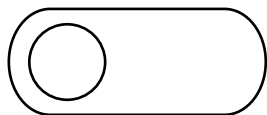
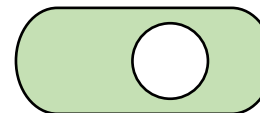
- When we measure qubit to readout the results, the quantum superposition is destroyed and projected to the eigenstate of the observable
- The probability to obtain  $|0\rangle$  and  $|1\rangle$  states in the measurement is proportional to the square of the coefficient

**Bit:** Either 1 or 0



ON

**1**



OFF

**0**

$$0.64 \neq 0.6|0\rangle + 0.8|1\rangle$$

## Quantum bits (Qubits)

Quantum states can be entangled, in which measurement of one qubit affects the quantum state of other qubits

$$\frac{1}{2}(|00\rangle + |01\rangle + |10\rangle + |11\rangle) \quad (\text{Unentangled state})$$

Measurement outcome of the 1 <sup>st</sup> qubit	Measurement outcome of the 2 <sup>nd</sup> qubit
$ 0\rangle$	$ 0\rangle$ or $ 1\rangle$ with 50:50 probability
$ 1\rangle$	$ 0\rangle$ or $ 1\rangle$ with 50:50 probability

$$\frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) \quad (\text{ENTANGLED state})$$

Measurement outcome of the 1 <sup>st</sup> qubit	Measurement outcome of the 2 <sup>nd</sup> qubit
$ 0\rangle$	$ 0\rangle$ with 100% probability
$ 1\rangle$	$ 1\rangle$ with 100% probability

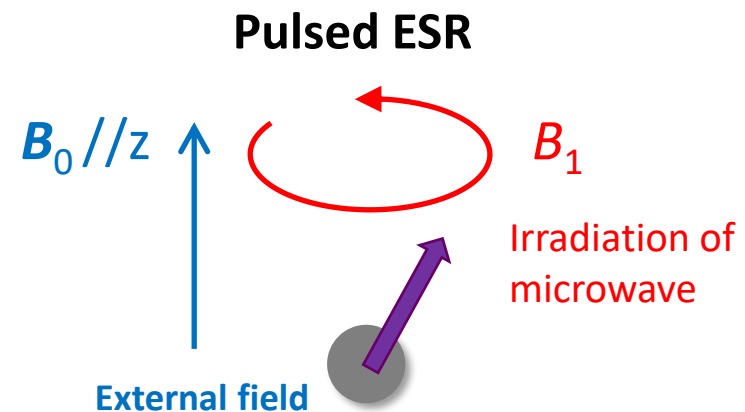
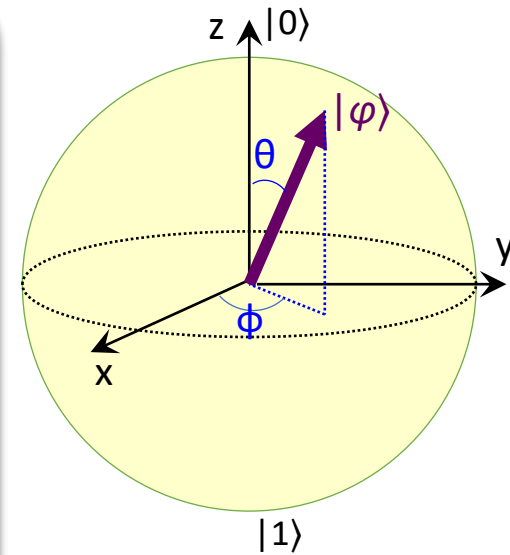
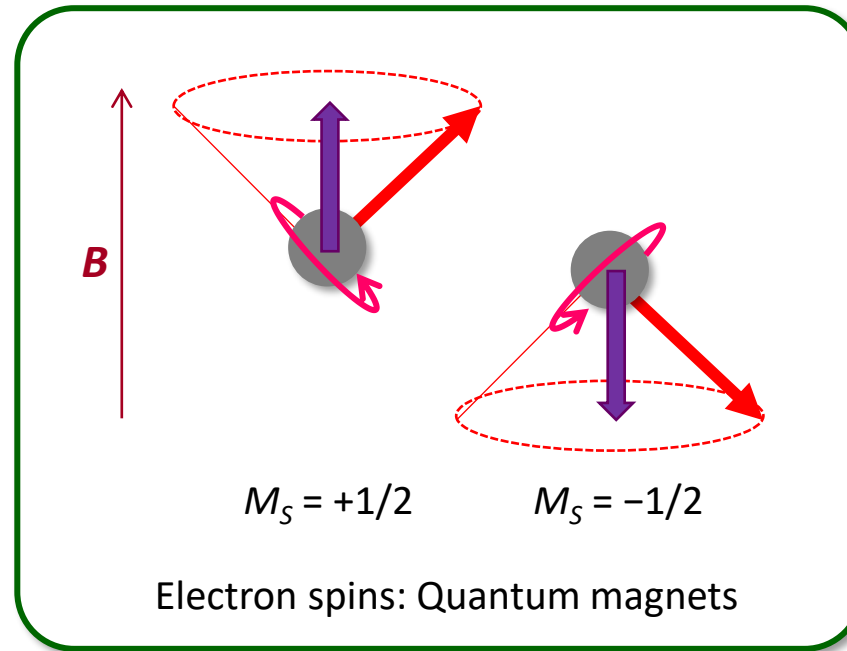
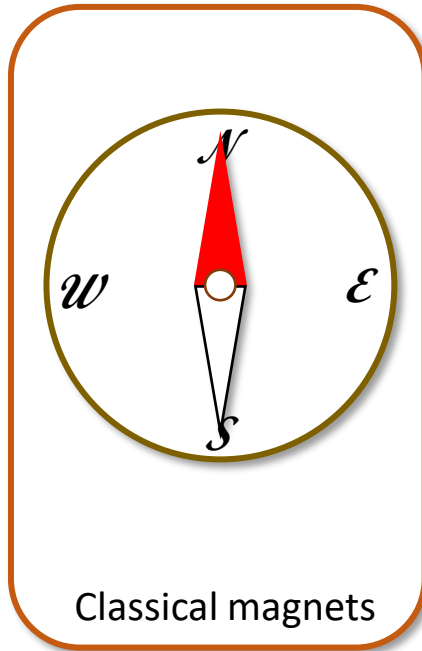
## Examples of qubits



Figure taken from "The Economist", 20th June, 2015.

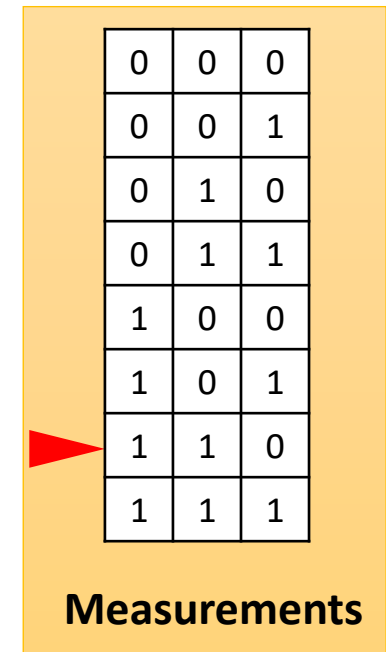
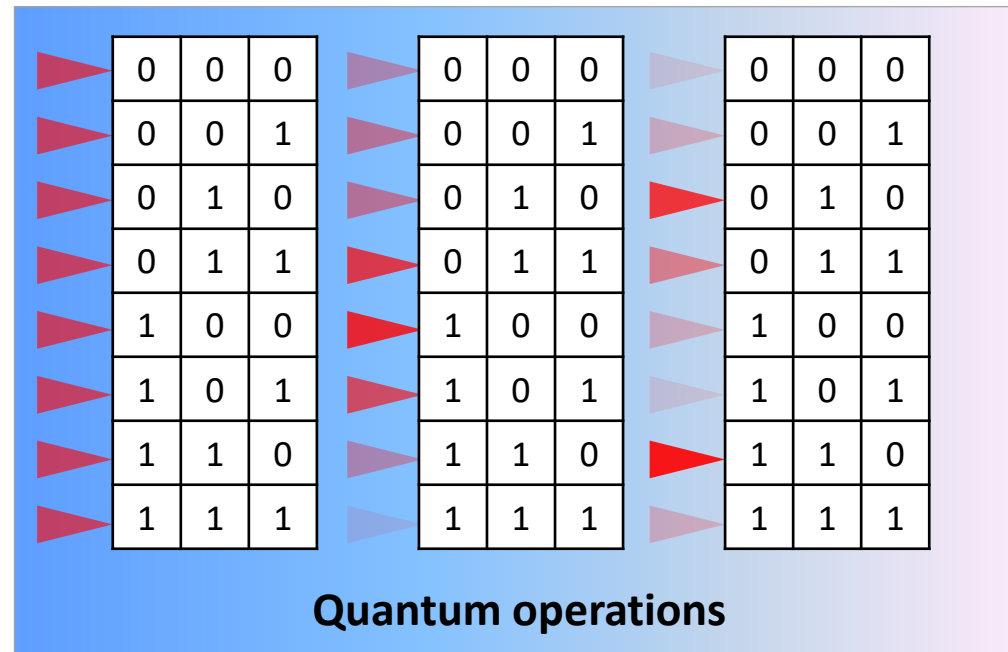
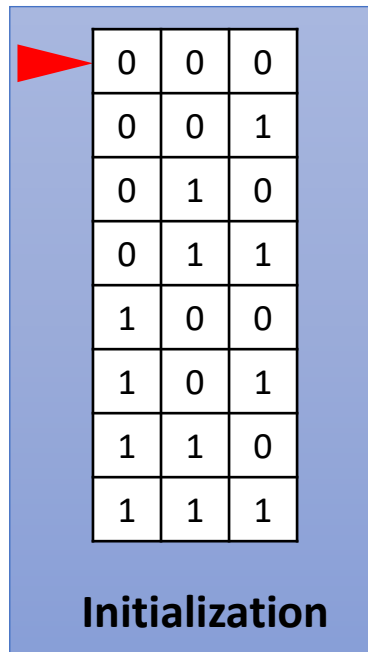
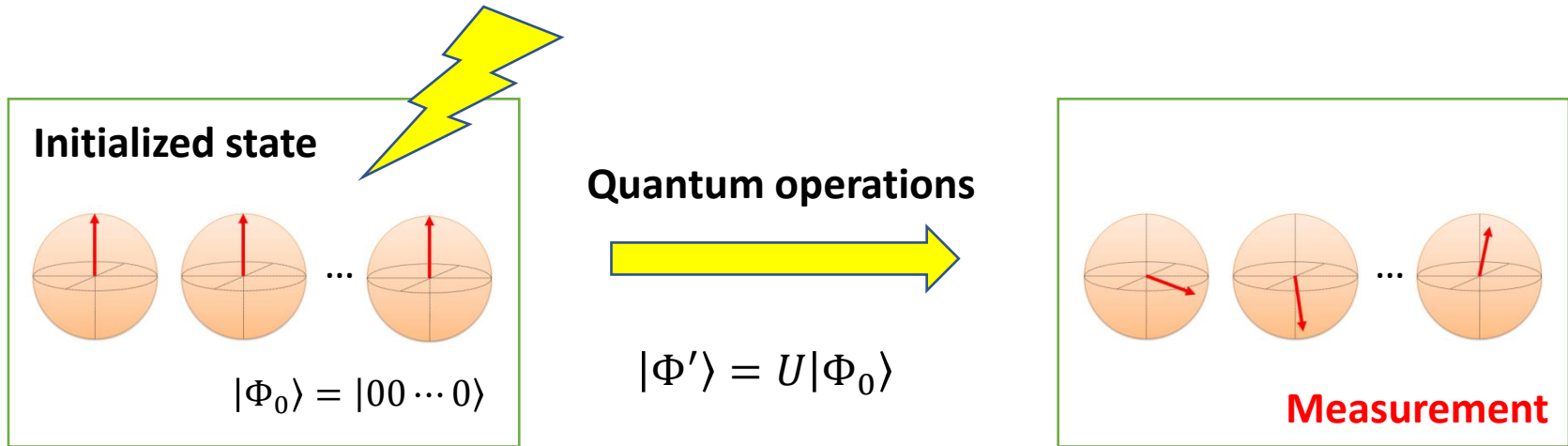
- Superconducting circuits
- Photons (polarized light)
- Trapped ions
- Nuclear and electron spins
- Semiconductors
- Majorana fermions

# An example: Electron spin qubits

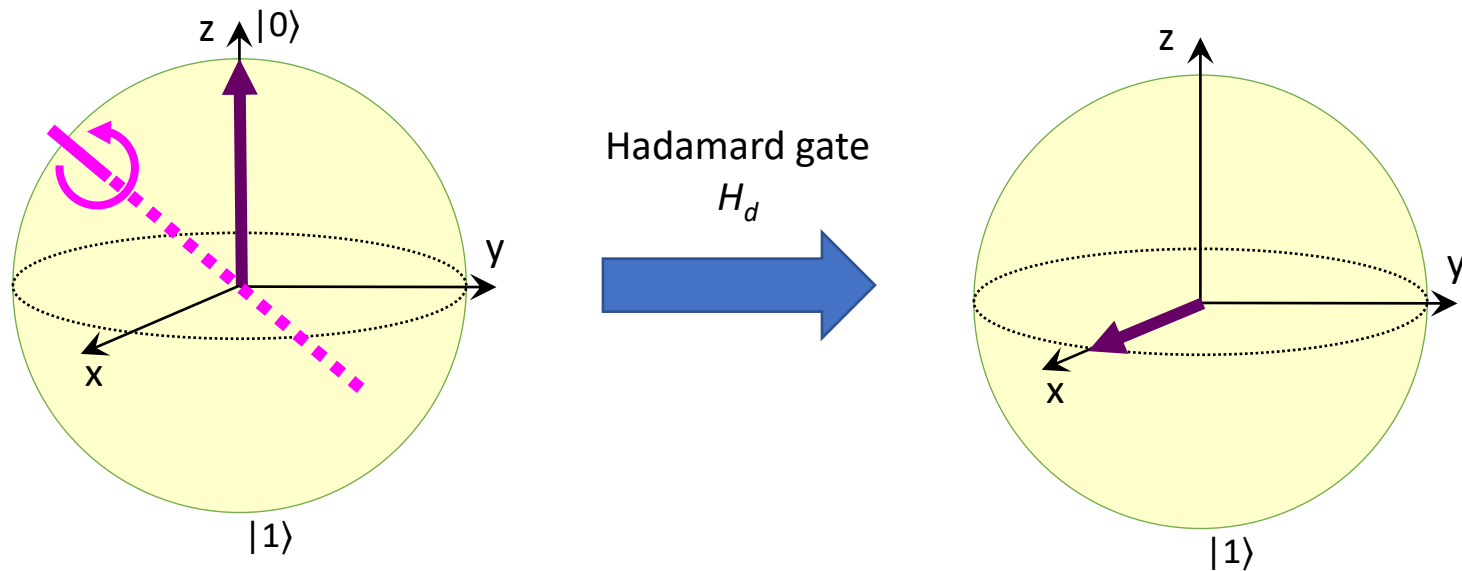




# Quantum computing



# An example of quantum operations







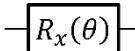
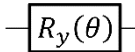
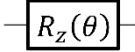
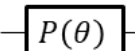


180° rotation along the axis  
that bisects z and x axes

We can calculate the quantum state after quantum operations by matrix algebra

$$|\varphi\rangle = c_0|0\rangle + c_1|1\rangle = \begin{pmatrix} c_0 \\ c_1 \end{pmatrix} \quad H_d = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$

$$H_d|\varphi\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} c_0 \\ c_1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} c_0 + c_1 \\ c_0 - c_1 \end{pmatrix}$$

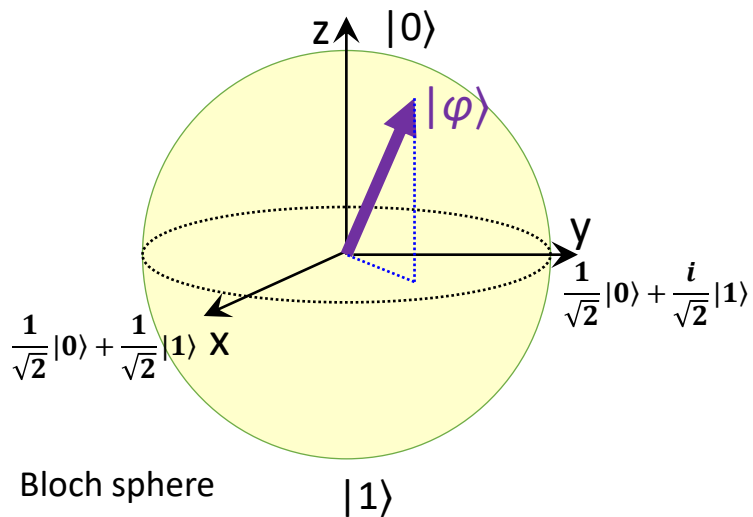
## Quantum gates

Name	Circuit symbol	Matrix representation
Hadamard ( $H_d$ )		$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$
Pauli-X		$\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$
Pauli-Y		$\begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$
Pauli-Z		$\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$
$R_x(\theta)$		$\begin{pmatrix} \cos \frac{\theta}{2} & -i \sin \frac{\theta}{2} \\ i \sin \frac{\theta}{2} & \cos \frac{\theta}{2} \end{pmatrix}$
$R_y(\theta)$		$\begin{pmatrix} \cos \frac{\theta}{2} & -\sin \frac{\theta}{2} \\ \sin \frac{\theta}{2} & \cos \frac{\theta}{2} \end{pmatrix}$
$R_z(\theta)$		$\begin{pmatrix} e^{-i\theta/2} & 0 \\ 0 & e^{i\theta/2} \end{pmatrix}$
$P(\theta)$		$\begin{pmatrix} 1 & 0 \\ 0 & e^{i\theta} \end{pmatrix}$
CNOT		$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$
SWAP		$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$

## Pauli gates

Name	Circuit symbol	Matrix representation
Pauli-X	$\boxed{X}$	$\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$
Pauli-Y	$\boxed{Y}$	$\begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$
Pauli-Z	$\boxed{Z}$	$\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

The quantum states on the x, y, and z axes are eigenfunction of the Pauli-X, Y, and Z operators, with the eigenvalues either 1 or -1.



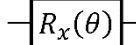
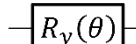
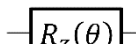
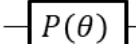
$$X|+\rangle = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1/\sqrt{2} \\ 1/\sqrt{2} \end{pmatrix} = \begin{pmatrix} 1/\sqrt{2} \\ 1/\sqrt{2} \end{pmatrix} = |+\rangle$$

$$X|-\rangle = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1/\sqrt{2} \\ -1/\sqrt{2} \end{pmatrix} = \begin{pmatrix} -1/\sqrt{2} \\ 1/\sqrt{2} \end{pmatrix} = -|-\rangle$$

$$Y|+_y\rangle = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 1/\sqrt{2} \\ i/\sqrt{2} \end{pmatrix} = \begin{pmatrix} -i^2/\sqrt{2} \\ i/\sqrt{2} \end{pmatrix} = |+_y\rangle$$

$$Y|-_y\rangle = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 1/\sqrt{2} \\ -i/\sqrt{2} \end{pmatrix} = \begin{pmatrix} i^2/\sqrt{2} \\ i/\sqrt{2} \end{pmatrix} = -|-_y\rangle$$

## Rotational gates and phase shift gates

Name	Circuit symbol	Matrix representation
$R_x(\theta)$		$\begin{pmatrix} \cos \frac{\theta}{2} & -i \sin \frac{\theta}{2} \\ -i \sin \frac{\theta}{2} & \cos \frac{\theta}{2} \end{pmatrix}$
$R_y(\theta)$		$\begin{pmatrix} \cos \frac{\theta}{2} & -\sin \frac{\theta}{2} \\ \sin \frac{\theta}{2} & \cos \frac{\theta}{2} \end{pmatrix}$
$R_z(\theta)$		$\begin{pmatrix} e^{-i\theta/2} & 0 \\ 0 & e^{i\theta/2} \end{pmatrix}$
$P(\theta)$		$\begin{pmatrix} 1 & 0 \\ 0 & e^{i\theta} \end{pmatrix}$

$R_x$ ,  $R_y$ , and  $R_z$  gates are rotation gates along x, y, and z axes.

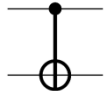
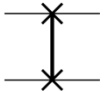
$P(\theta)$  is a phase shift gate.

The phase difference between the  $|0\rangle$  and  $|1\rangle$  states after the application of  $R_z(\theta)$  and  $P(\theta)$  gates are the same. Only global phase is different.

$$P(\pi) = Z, P(\pi/2) = S, P(\pi/4) = T$$



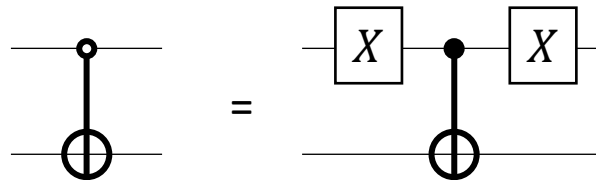
## Two qubit gates

Name	Circuit symbol	Matrix representation
CNOT		$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$
SWAP		$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$

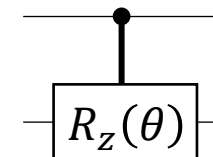
CNOT (controlled-NOT): Apply an NOT operation to the target qubit **if and only if the control qubit is the  $|1\rangle$  state**.

SWAP exchanges quantum states of two qubits:  $\text{SWAP}[j, k]|q_j, q_k\rangle = |q_k, q_j\rangle$

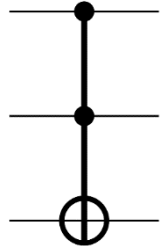
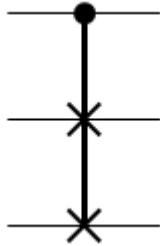
Open circle notation is also used when apply the quantum gate to the target qubit if and only if the control is the  $|0\rangle$  state.



One can define a two qubit gate by adding a control qubit, like controlled- $R_z(\theta)$  gate,

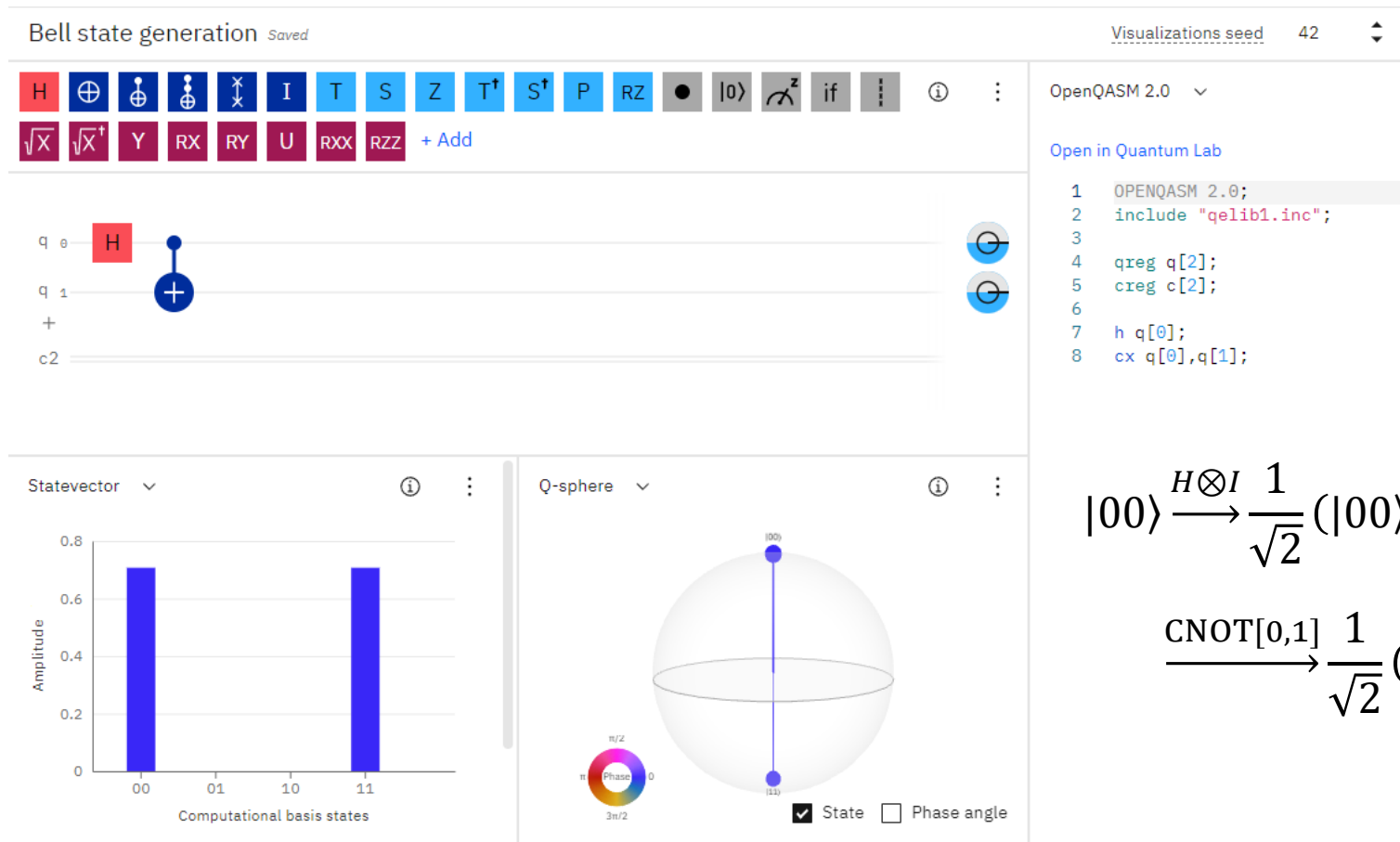


# Three qubit gates

Name	Circuit symbol	Matrix representation
CCNOT (Toffoli)		$\begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \end{pmatrix}$
C-Swap (Fredkin)		$\begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \end{pmatrix}$
Deutsch's gate		$\begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & i \cos \theta & \sin \theta \\ 0 & 0 & 0 & 0 & 0 & 0 & \sin \theta & i \cos \theta \end{pmatrix}$

# Bell states generations

Maximally entangled states in two-qubit system, known as Bell states, can be generated by combining Hadamard, Pauli-X, and CNOT gates.

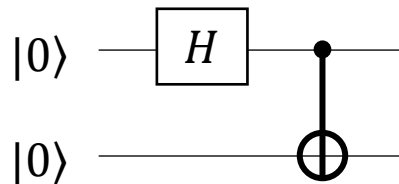


$$|00\rangle \xrightarrow{H \otimes I} \frac{1}{\sqrt{2}} (|00\rangle + |10\rangle)$$

$$\xrightarrow{\text{CNOT}[0,1]} \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$$

## Bell states generations

Maximally entangled states in two-qubit system, known as Bell states, can be generated by combining Hadamard, Pauli-X, and CNOT gates.



$$|00\rangle \xrightarrow{H \otimes I} \frac{1}{\sqrt{2}} (|00\rangle + |10\rangle) \xrightarrow{\text{CNOT}[0,1]} \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$$

### [Question]

**How to construct quantum circuits to generate other three Bell states**

$$\frac{1}{\sqrt{2}} (|01\rangle + |10\rangle), \frac{1}{\sqrt{2}} (|01\rangle - |10\rangle), \text{ and } \frac{1}{\sqrt{2}} (|00\rangle - |11\rangle)$$

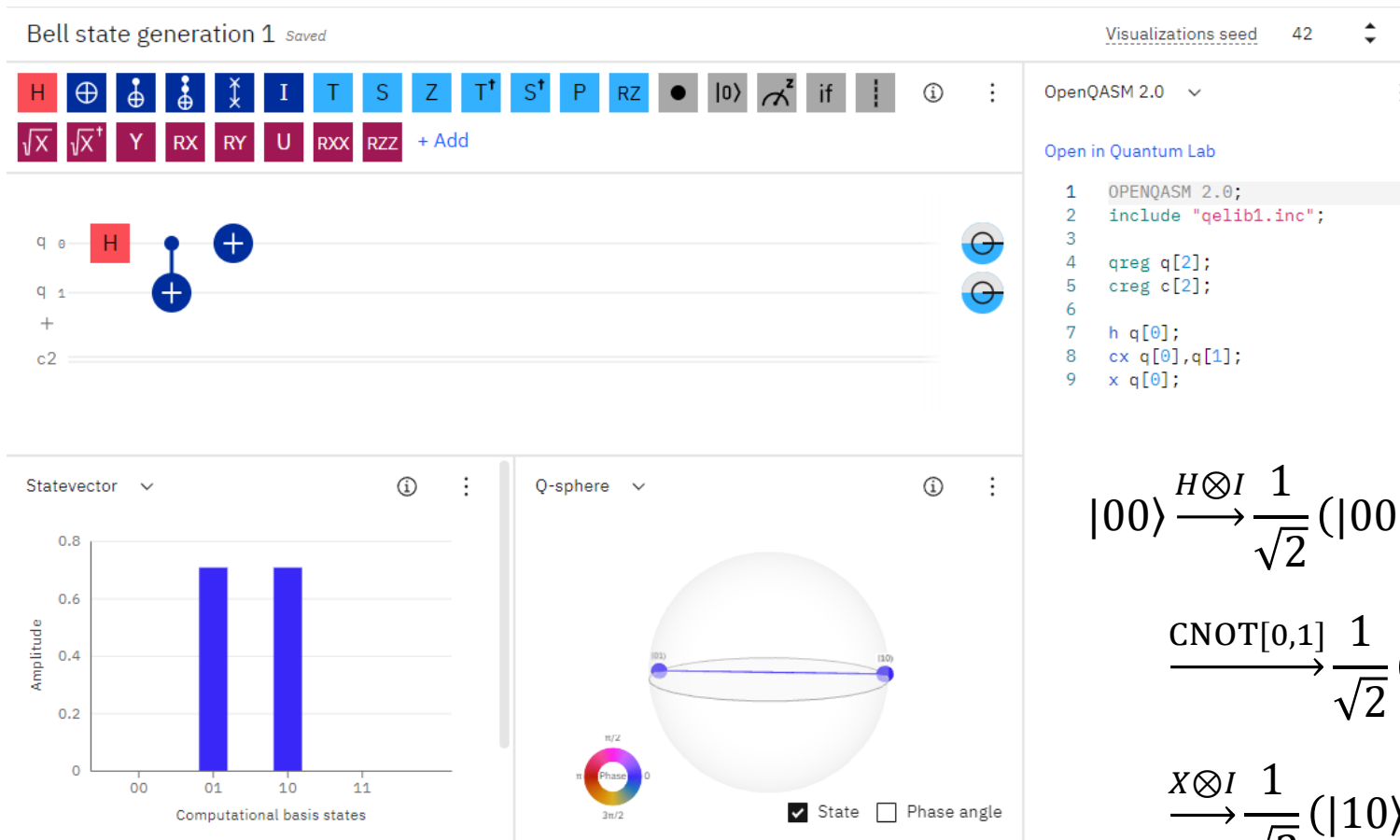
**from the quantum state initialized to the  $|00\rangle$  state using Hadamard, Pauli-X, and CNOT gates?**

$$\left\{ \begin{array}{l} H|0\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \\ H|1\rangle = \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle) \end{array} \right. \quad \left\{ \begin{array}{l} X|0\rangle = |1\rangle \\ X|1\rangle = |0\rangle \end{array} \right. \quad \left\{ \begin{array}{l} \text{CNOT}|00\rangle = |00\rangle \\ \text{CNOT}|01\rangle = |01\rangle \\ \text{CNOT}|10\rangle = |11\rangle \\ \text{CNOT}|11\rangle = |10\rangle \end{array} \right.$$

# Bell states generations

[Answer]

Generation of  $\frac{1}{\sqrt{2}}(|01\rangle + |10\rangle)$



$$|00\rangle \xrightarrow{H \otimes I} \frac{1}{\sqrt{2}}(|00\rangle + |10\rangle)$$

$$\xrightarrow{\text{CNOT}[0,1]} \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$$

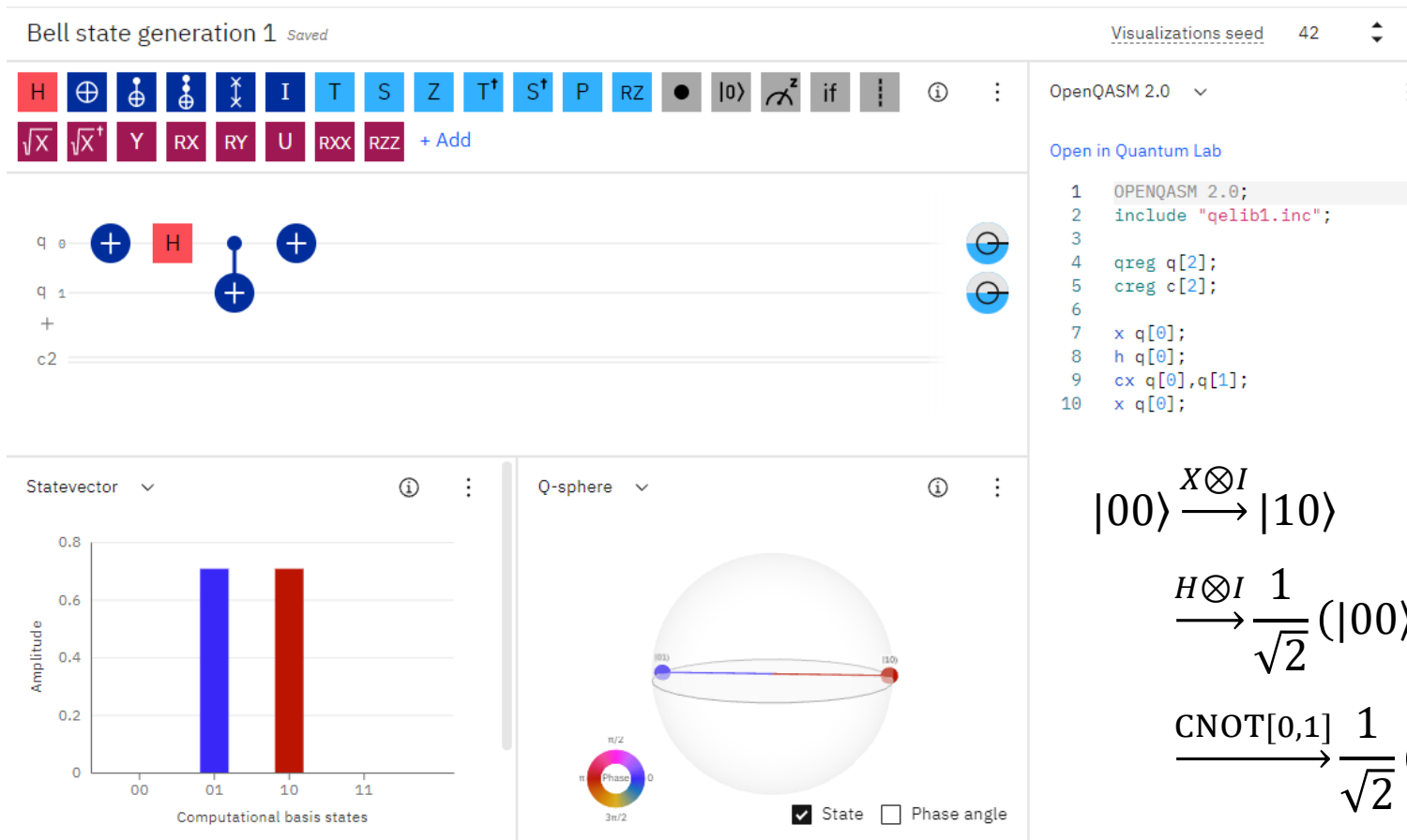
$$\xrightarrow{X \otimes I} \frac{1}{\sqrt{2}}(|10\rangle + |01\rangle)$$



# Bell states generations

[Answer]

Generation of  $\frac{1}{\sqrt{2}} (|01\rangle - |10\rangle)$



$$|00\rangle \xrightarrow{X \otimes I} |10\rangle$$

$$\xrightarrow{H \otimes I} \frac{1}{\sqrt{2}} (|00\rangle - |10\rangle)$$

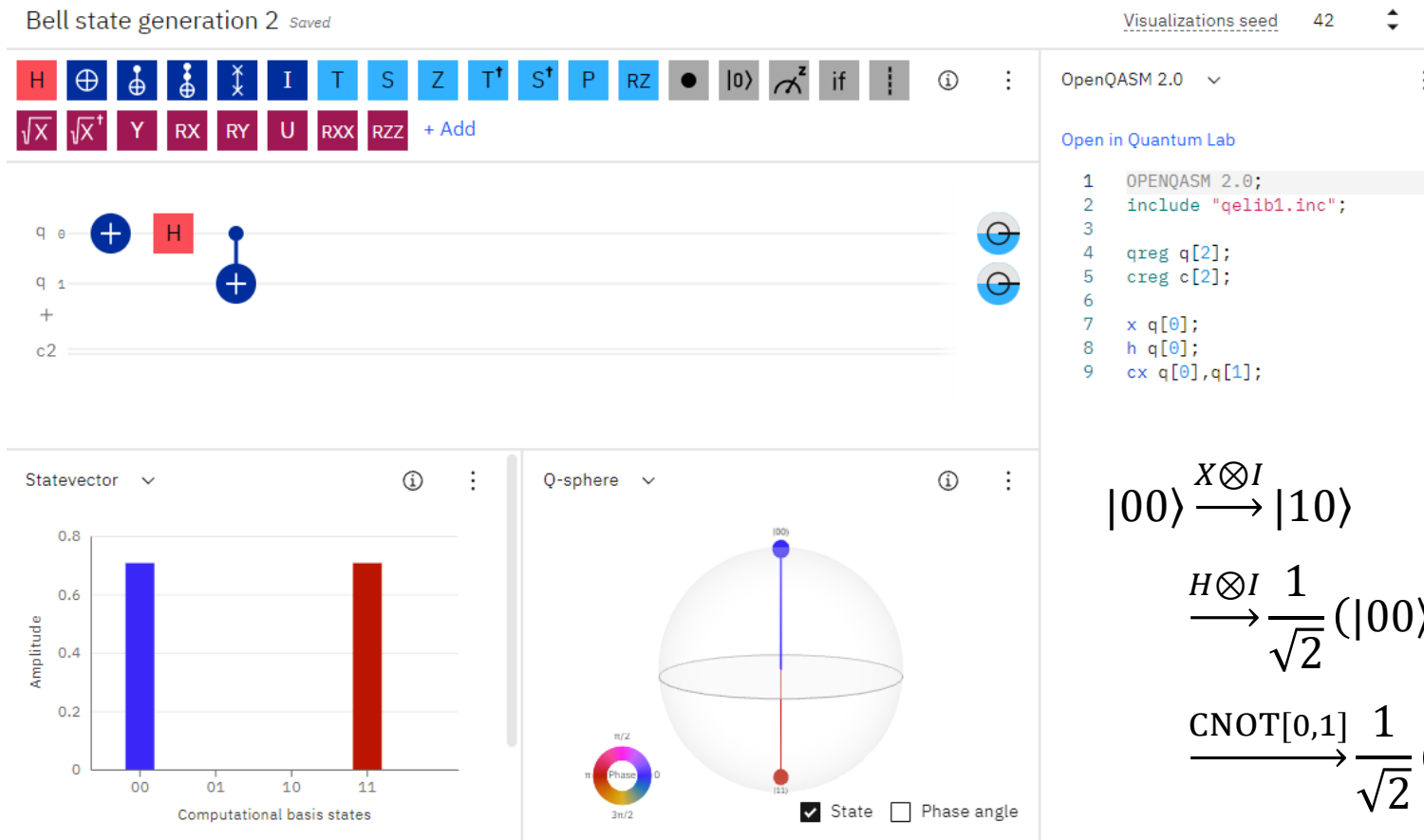
$$\xrightarrow{\text{CNOT}[0,1]} \frac{1}{\sqrt{2}} (|00\rangle - |11\rangle)$$

$$\xrightarrow{X \otimes I} \frac{1}{\sqrt{2}} (|10\rangle - |01\rangle)$$

# Bell states generations

[Answer]

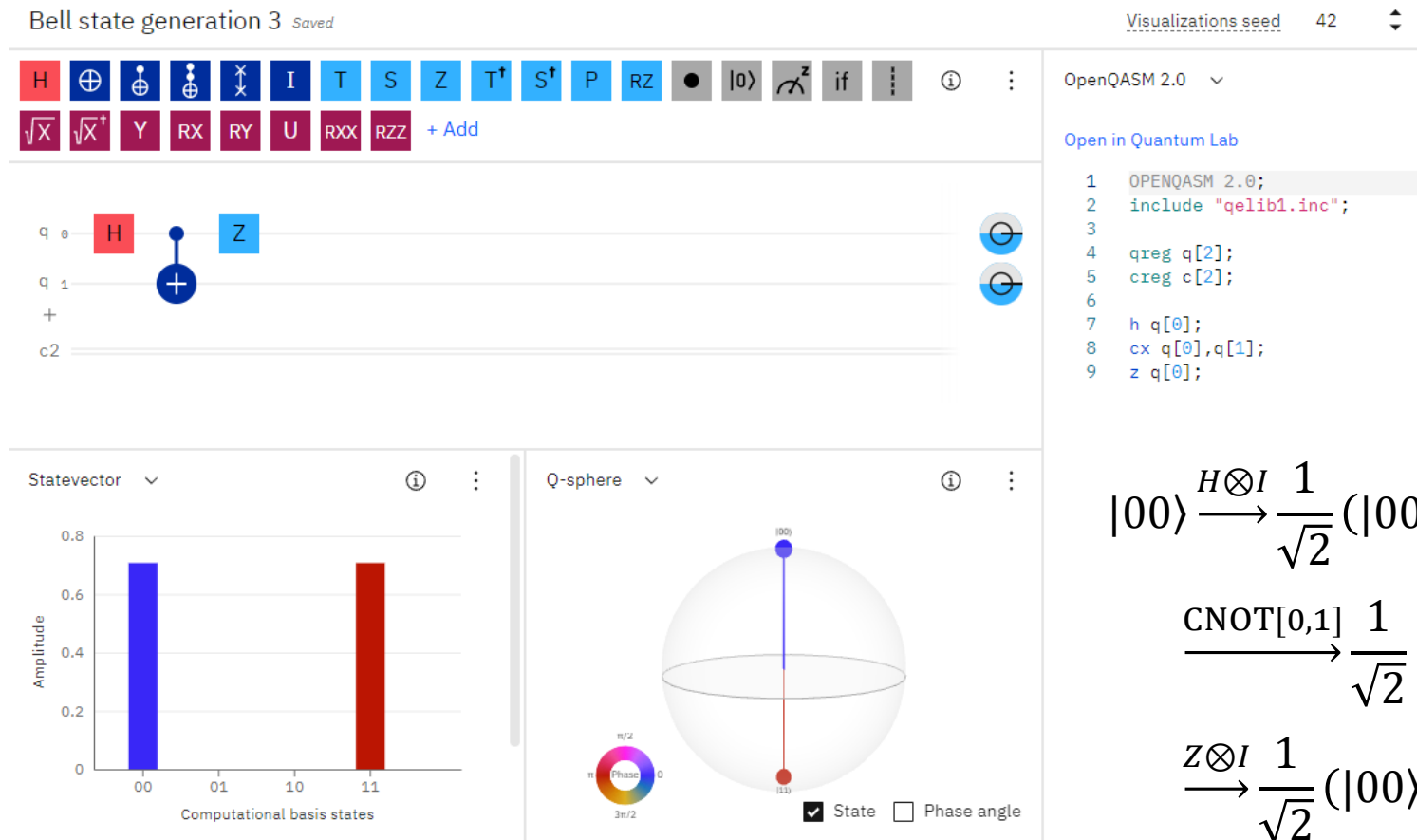
Generation of  $\frac{1}{\sqrt{2}} (|00\rangle - |11\rangle)$



\* There are many other quantum circuits to generate the same Bell state.

# Bell states generations

Generation of  $\frac{1}{\sqrt{2}}(|00\rangle - |11\rangle)$



$$\begin{aligned}
 |00\rangle &\xrightarrow{H \otimes I} \frac{1}{\sqrt{2}}(|00\rangle + |10\rangle) \\
 &\xrightarrow{\text{CNOT}[0,1]} \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) \\
 &\xrightarrow{Z \otimes I} \frac{1}{\sqrt{2}}(|00\rangle - |11\rangle)
 \end{aligned}$$

# Universal quantum gate sets

Any unitary operation in the Hilbert space of  $n$ -qubits can be decomposed by using universal quantum gate sets.

- 1) Arbitrary one qubit rotation gates and CNOT gate
- 2) Hadamard + phase ( $S$ ) + CNOT +  $\pi/8$  ( $T$ ) gates
- 3) Deutsch's gate

## ☆ Gottesman–Knill theorem

Stabilizer circuits can be efficiently simulated on classical computers.

Preparation of qubits in computational basis states.

Quantum gates from the Clifford group ( $H_d$ ,  $S$ ,  $S^\dagger$ , CNOT, and Pauli- $X$ ,  $Y$ , and  $Z$ )

Measurement in computational basis

# Quantum circuits

H

$\oplus$

$\otimes$

$\otimes$

$\otimes$

I

T

S

Z

$T^\dagger$

$S^\dagger$

P

RZ

●

$|0\rangle$

$\otimes^z$

if

|

$\sqrt{X}$

$\sqrt{X}^\dagger$

Y

ⓘ

⋮

RX

RY

U

RXX

RZZ

+ Add

—

—

—

—

OpenQASM 2.0 ⌵

[Open in Quantum Lab](#)

```

1 OPENQASM 2.0;
2 include "qelib1.inc";
3
4 qreg q[4];
5 creg c[4];
6
7 h q[0];
8 h q[1];
9 cx q[0],q[3];
10 cx q[1],q[3];
11 ccx q[0],q[1],q[2];
12 measure q[0] -> c[0];
13 measure q[1] -> c[1];
14 measure q[2] -> c[2];
15 measure q[3] -> c[3];
        
```

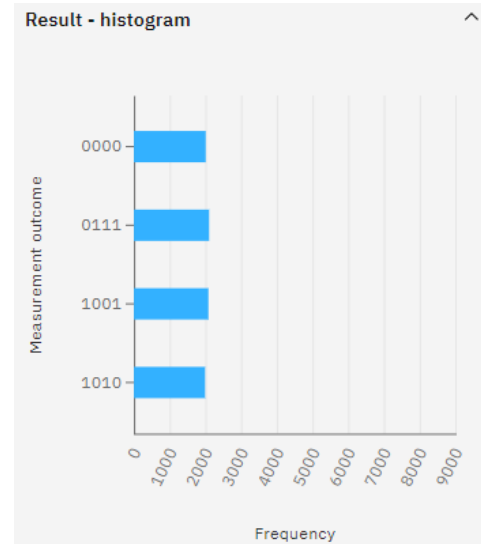
$$|0000\rangle \xrightarrow{H \otimes I \otimes I \otimes I} \frac{1}{\sqrt{2}} (|0000\rangle + |1000\rangle)$$

$$\xrightarrow{I \otimes H \otimes I \otimes I} \frac{1}{2} (|0000\rangle + |0100\rangle + |1000\rangle + |1100\rangle)$$

$$\xrightarrow{\text{CNOT}[0,3]} \frac{1}{2} (|0000\rangle + |0100\rangle + |1001\rangle + |1101\rangle)$$

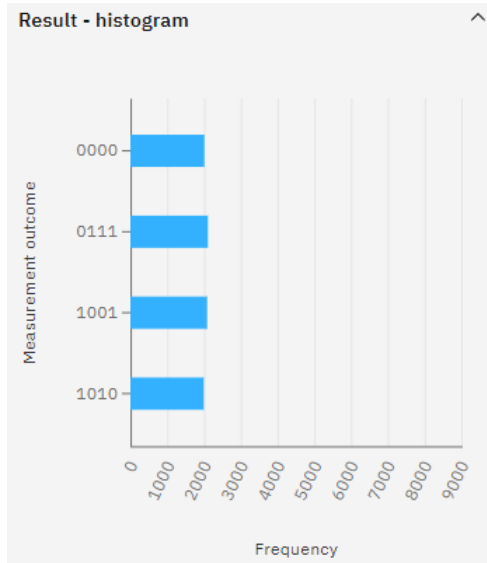
$$\xrightarrow{\text{CNOT}[1,3]} \frac{1}{2} (|0000\rangle + |0101\rangle + |1001\rangle + |1100\rangle)$$

$$\xrightarrow{\text{CCNOT}[0,1,2]} \frac{1}{2} (|0000\rangle + |0101\rangle + |1001\rangle + |1110\rangle)$$



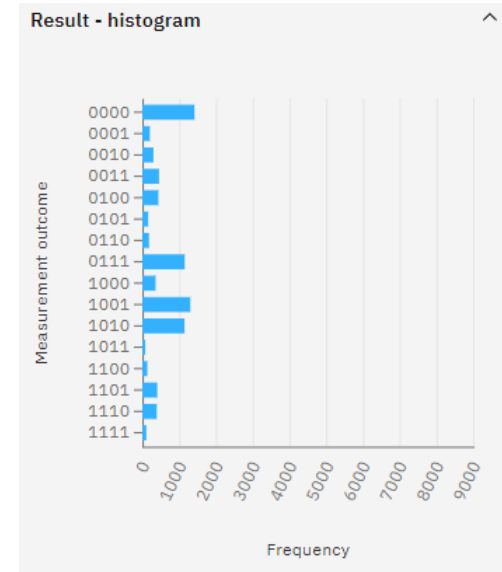


# Quantum circuits

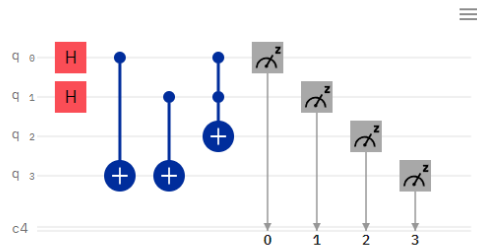


Ideal results

Results from real  
quantum computer  
(ibmq\_manila)

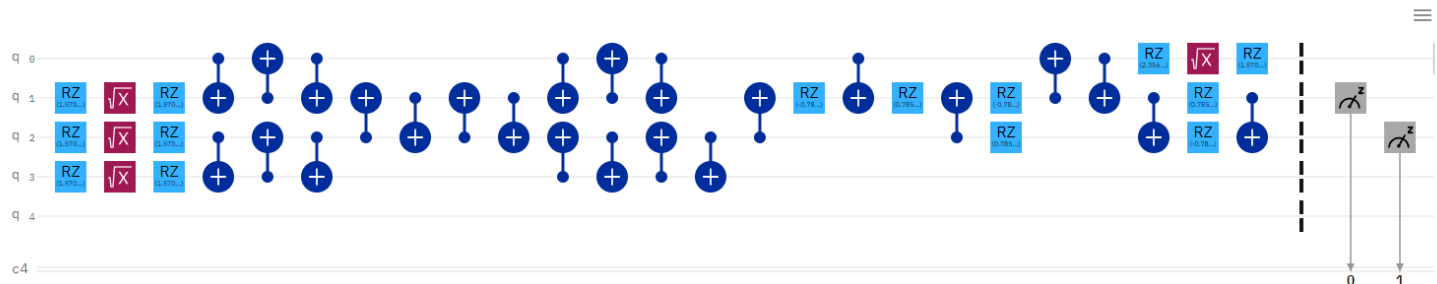


Original circuit



5 quantum gates –(transpilation)→ 42 quantum gates

Transpiled circuit



# Quantum devices

ibmq\_manila



## Details



5	Status:	● Online	Avg. CNOT Error:	8.525e-3
Qubits	Total pending jobs:	73 jobs	Avg. Readout Error:	2.218e-2
32	Processor type ⓘ:	Falcon r5.11L	Avg. T1:	135.26 us
Quantum Volume	Version:	1.0.10	Avg. T2:	71.21 us
	Basis gates:	CX, ID, RZ, SX, X	Providers with access:	<a href="#">2 Providers</a> ↓
	Your usage:	66 jobs	Supports Qiskit Runtime:	Yes

## Your upcoming reservations 0



## Calibration data

Last calibrated: 25 minutes ago



Map view

Graph view

Table view

Qubit:

Frequency (GHz)



Avg 4.971



Connection:

CNOT error



Avg 8.525e-3

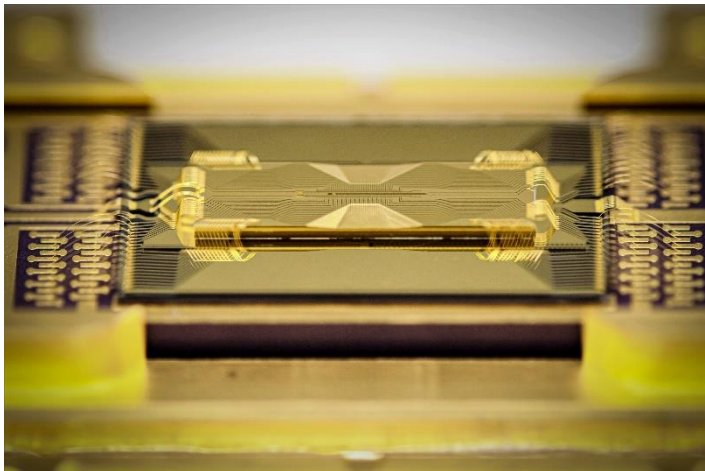


# NISQ devices

## Noisy Intermediate-Scale Quantum devices

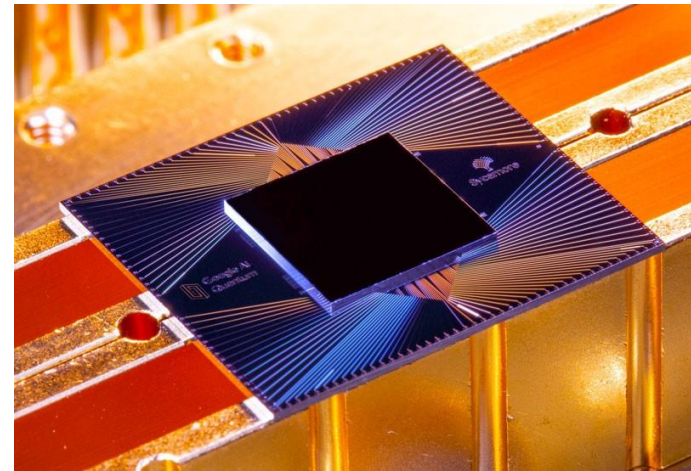
J. Preskill, *Quantum* **2018**, 2, 79.

- About 50–100 qubits
- Too small to implement quantum error correction
- The number of executable coherent gates is at most 1,000



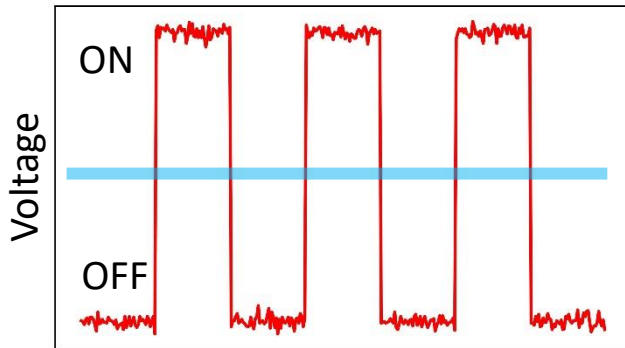
IonQ's EGT Series Ion Trap Chip  
Photo taken from IonQ's website

<https://ionq.com/news/august-25-2021-reconfigurable-multicore-quantum-architecture>



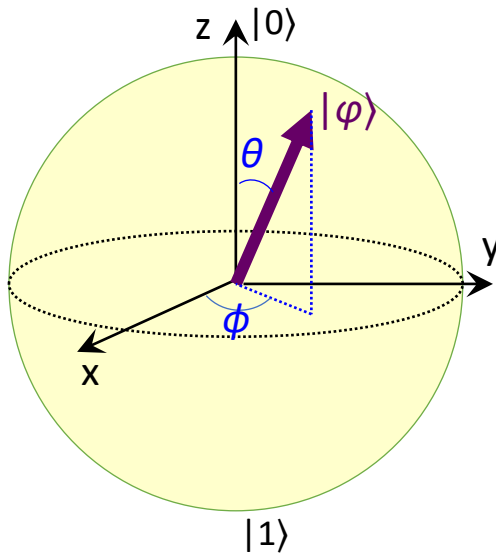
Google's 54-qubit quantum chip "Sycamore"  
Photo taken from *Nature* **2019**, 574, 461–462.

# Quantum error correction



## Classical bit is robust against noise

- “On” if the voltage is higher than the threshold, otherwise “Off”.
- Noises and errors will reset every time the operation proceeds.



## Quantum bit is very fragile against noise

- Different  $\theta$  and  $\phi$  means different quantum states
- The error accumulates by proceeding operations.
- Quantum superposition can be broken by system-environment interactions (decoherence)



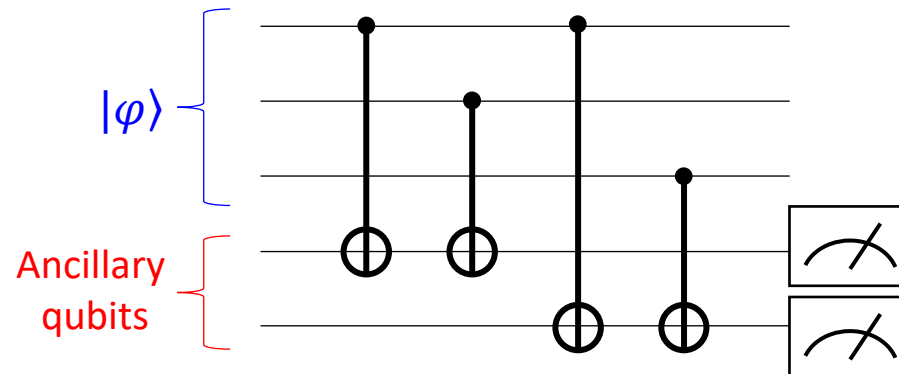
**Use multiple physical qubits to retain information of one logical qubit!**

# The three-qubit bit-flip code

- Use **three physical qubits** for one logical qubits.
- Need **two ancillary qubits** to detect/correct single bit-flip error.

$$|0\rangle \rightarrow |0_L\rangle \equiv |000\rangle \quad |1\rangle \rightarrow |1_L\rangle \equiv |111\rangle$$

$$|\varphi\rangle = c_0|0_L\rangle + c_1|1_L\rangle \equiv c_0|000\rangle + c_1|111\rangle$$



Situation	Quantum state before error correction	Quantum state after measurement	Measurement outcome
No bit-flip error	$(c_0 000\rangle + c_1 111\rangle) \otimes  00\rangle$	$(c_0 000\rangle + c_1 111\rangle) \otimes  00\rangle$	00
1st qubit flip	$(c_0 100\rangle + c_1 011\rangle) \otimes  00\rangle$	$(c_0 100\rangle + c_1 011\rangle) \otimes  11\rangle$	11
2nd qubit flip	$(c_0 010\rangle + c_1 101\rangle) \otimes  00\rangle$	$(c_0 010\rangle + c_1 101\rangle) \otimes  10\rangle$	10
3rd qubit flip	$(c_0 001\rangle + c_1 110\rangle) \otimes  00\rangle$	$(c_0 001\rangle + c_1 110\rangle) \otimes  01\rangle$	01