NBODY6++

Manual for the Computer Code

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eqmass10k_03.ps		

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Table of contents 3

Table of contents

1	Introdu	ction	4
2	Code v	ersions	6
3	Getting	started	7
	3.1	Installation	7
	3.2	Start a simulation	7
	3.3	Output	8
	3.4	Code stop	8
	3.5	Restart a simulation	9
	3.6	Code structure	9
4	Input v	ariables	0
5	Thresh	olds for the variables	22
6	How to	read the diagnostics	23
7	Runs o	n parallel machines	28
8	The He	ermite integration method	29
9	Individ	ual and block time steps	31
10	The Al	nmad–Cohen scheme	3
11	KS–Re	gularization	86
12	Nbody	-units	8
13	Output		39
Refe	rences .		51

4 1 Introduction

1 Introduction

Gravity is an ever–present force in the Universe and is involved into the dynamics of all kinds of bodies, from the tiny atom to the clusters of galaxies. At small spatial scales, its influence is covered by other strong forces (e.g. magnetic, pressure, radiation induced), while on the very large scale it becomes the most dominant power. In astrophysics, it governs the dynamical evolution of many self–gravitating systems. Here, we concentrate on such systems that are dominated by mutual gravitation between particles.

The numerical star-by-star simulation of a simple cluster containing some more than hundred thousand members still places heavy demands on the available hard- and software. A balance has to be found between two constraints: On one hand the *realism*, i.e. the input of profound physics, inclusion of all astrophysical effects as well as the maintenance of the accuracy of calculations; and on the other hand, the *efficiency*, i.e. the limitations given by the computational possibilities and suitable codes to be finished in a reasonable time. Many different kinds of approaches have been undertaken to suffice both:

- codes based on the direct force integration [2], [5], [6], see also: http://www.sverre.com/,
- statistical models, which themselves divide into several subgroups (Fokker–Planck approximation by [10]; Monte–Carlo method by [13]; Gas models by [27]),
- usage of high-performance parallel computers [28], [11],
- or the construction of special hardware devoted for these purposes (GRAPE [19], see also: http://www.astrogrape.org/ and http://www.cs.rit.edu/~grapecluster/.

The code NBODY6++ described in this manual is designed for an accurate integration of many bodies (e.g. in a star cluster, planetary system, galactic nucleus) based on the direct integration of the Newtonian equations of motion. It is optimal for collisional systems, where long times of integration and high accuracy or both are required, in order to follow with high precision the secular evolution of the objects.

NBODY6++ is a descendant of the family of NBODY codes initiated by Sverre Aarseth [4], which has been extended to be suitable for parallel computers [28]. The basic features of the code increasing the efficiency may be considered under four separate headings: fourth order prediction—correction method (Hermite scheme), individual and block time—steps, regularization of close encounters and few-body subsystems, and a neighbour scheme (Ahmad—Cohen scheme). We briefly describe these ideas in this booklet, while a detailed description can be found in [3] as well as his book [6].

While NBODY6++ is not that different from NBODY6 to justify a completely new name, the user should, however, be aware that in order to make a parallelization of regular and irregular force computations possible at all, some significant changes in the order of operations became necessary. As a consequence, trajectories of the same initial system, simulated by NBODY6 and NBODY6++ will diverge from each other, due to the inherent exponential instability and deterministic chaos in *N*-body systems. Still one should always expect that the *global* properties are well behaved in both cases (e.g. energy conservation). While much effort is taken to keep NBODY6 and NBODY6++ as close as possible this is never 100% the case, and the interested should always contact Sverre Aarseth or Rainer Spurzem if in doubt about these matters.

This manual should serve as a practical starter kit for new students working with NBODY6++. It is not meant as a complete reference or scientific paper; for that see the references and in particular the excellent compendium of Aarseth's book on Gravitational *N*-Body Simulations [6].

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6 2 Code versions

2 Code versions

The development of the NBODY code has begun in the 1960s [1], though there exist some earlier precursors [29], [30]. It has set a quasi-standard for the precise direct integration of gravitating many-body systems. There exist several code groups (NBODY0–7, and a number of special implementations) for different usage, some of which are rather of historical interest.

The current NBODY6++ code is available publicly under Subversion or Github. You can download the beta version by using: svn co http://silkroad.bao.ac.cn/repos/betanb6 git clone https://github.com/lwang-astro/betanb6pp.git

The stable version will be avaiable under svn co http://silkroad.bao.ac.cn/repos/releasenb6 The documents and input samples are included.

The original N-body codes can be accessed publicly via Sverre Aarseth's ftp and web sites at ftp://ftp.ast.cam.ac.uk/pub/sverre/ and http://www.sverre.com/.

A brief comparison of the code versions:

ITS: Individual time-steps

ACS: Neighbour scheme (Ahmad-Cohen scheme) with block time-steps

KS: KS-regularization of few-body subsystems

HITS: Hermite scheme integration method combined with hierarchical block time steps

PN: Post-Newtonian terms AR: Algorithmic regularization

	ITS	ACS	KS	HITS	PN	AR
NBODY1	√					
NBODY2		✓		✓		
NBODY3	✓		✓			
NBODY4			✓	✓		
NBODY5	✓	✓	✓			
NBODY6		✓	✓	✓		
NBODY7		✓	✓	✓	√	✓

3 Getting started

3.1 Installation

After downloading NBODY6++GPU from Github, a directory will be created containing all the source files (routines and functions), documents and input samples. By default the directory is called nbody6ppGPU. The README file provides a short guideline for using the code, including some tips for common issues.

To compile the code, first the configure should be executed. More details of configure options can be found by using:

```
./configure -help
```

The simple way to use configure is just type: ./configure Then the configure script will check your system environments to find avaiable Fortran/c++/CUDA compilers. If all checking pass, there will be a summary showing the name of excutable file (nbody6++.**), the supported features, the installation path and the maximum array sizes (NMAX, KMAX, LMAX,MMAX). Here NMAX is the maximum number of particles, KMAX is the maximum number of KS pairs, LMAX is the maximum neighbor number and MMAX is the maximum merger number (≥ 3 bodies stable hierarchical system).

The default installation path is "/user/local". If you want to change it, use:

./configure -prefix=Installpath Then the code will be installed in "Installpath".

After configure, type

make

to compile the code and use

make install

for installation.

The most important options of configure you need is shown in Table 3.1.

Figure 3.0: Major options of configure script

Option	Description
-prefix=path	Installation path
-disable-gpu	Disable GPU acceleration (In the case you don't have
	Nvidia GPU with cuda support)
-enable-simd=avx/sse/no	Switch the features of SIMD parallel method (AVX / SSE
	/ NONE)
-disable-mpi	Disable MPI parallelization
-disable-openmp	Disable OpenMP support
-with-par=size	Choose the maximum array sizes (NMAX, KMAX,
	LMAX, MMAX), see detail by "./configure -help"

The manual file is stored in "Installpath/share/doc". The input samples are in "samples" of the source code directory.

3.2 Start a simulation

Depending on the user's individual research, the Nbody code opens a wide field of application possibilities. The user has to define the initial model by a number of input control variables, e.g. number of stars, the size of the cluster, a mass function and a density profile. An input file with

8 3 Getting started

many parameters is required to start a simulation. The detailed explanation of all input parameters is given in Chapter 4.

In the "samples" directory, three examples are shown. The "N16k.input" file generates an initial model with a Plummer sphere and no primordial binaries. The initial virial radius is 1 pc. The Kroupa (2001) type IMF is applied with mass range from 0.08 to $150\,M_\odot$. The solar neighbor tidal field is switched on. The mass, positions and velocities of stars are generated by the code during the initialization. A user-defined initial particle data, dat .10, can also be used by setting the option KZ(22). The file contains mass, positions and velocities of individual stars (one line per star).

For no MPI mode, the simulation can be started by type in a linux shell:

```
./nbody6++.** < input > output &
```

where "input" is the input file and output is the file will be generated by the code. With MPI parallelization, the standard way to start the simulation is:

```
mpirun -np [number of processes] ./nbody6++.** < input > output &
```

However, depending on the configuration of computing clusters, the way to start a MPI job can be very different. Users should contact administrator of the clusters to help for starting simulations.

3.3 Output

After the simulation start, the output file generated provides snapshots of the state of the system for a brief overview of some key data of the simulation to judge about the quality and performance of the run.

There are several files also created. The description of all these files are shown in Chapter 13. Most important files are comm.1/2[_*] which contain dumps of the complete common blocks. This files are for a restart and checkpoint purposes. single/binary/merger.40_* contain the snapshots of data (mass, position, velocity and stellar evolution parameters). global.30 and status.36 contain the important physical parameters every output time. coal.24 and coll.13 contains the information of mergers.

3.4 Code stop

A run will be finished when one of 4 conditions becomes true:

- the specified CPU-time on the computer is exceeded (variable TCOMP in the input file), or
- the maximum Nbody-time (see Ch. 4) is reached (variable TCRIT), or
- the physical cluster time in Myr is reached (variable TCRITp), or
- the number of cluster stars has fallen below a minimum (variable NCRIT).

However, in many cases, the simulation may abort abnormally. The most common error shown for such stop is "Calculation halted". This error message suggests that the relative energy error is too large ($>5 \times QE$). QE is defined in the input file provided by users.

There are many reasons for large energy errors. One major reason is an inaccurate integration of a few-body interaction. This can related to the switching criterion of regularization method. To solve such an issue, a first try is to modify *RMIN* and *DTMIN* and restart the simulation. *RMIN* and *DTMIN* are two major criterion for triggering KS regularization for a binary or close encounter. If this does not help, the second choice is to reduce *ETAI*, *ETAR* and *SMAX* which

control the time steps of Hermite integration. Usually if the large energy error is due to the integration error of dynamics, these two methods can help to avoid large energy error. If the same error appear no matter how these parameters are modified, it indicates that either a bug may exist for a special condition or the stellar evolution of a special type of binary is not properly treated. In such case, a detail debugging may needed and you may want to contact with the code developers for help (e.g. push an issue in GitHub). Or a simple solution is to start a new simulation.

3.5 Restart a simulation

The program dump all necessary data of simulations in a COMMON file comm. 1 or comm. 2[_*]. These files can be used to restart a simulation with a restart input file.

Before a restart, it is recommendable to copy or rename the files, otherwise they may be overwritten. The code only read comm.1 to restart a run. Thus firstly the comm.2[_*] should be renamed to comm.1

To restart a run, a different input file is needed. Only the first line is similar to the standard input file, The first parameter, *KSTART*, has to be changed to 2 or higher number. A simple example is:

```
2 1000000.0 1.E6 40 40 0
```

This simple input file will restart the simulation from the dumped COMMON file comm. 1. If the user wants to change some parameters for the restart, KSTART=3,5 can be used. For example:

```
3 1000000.0 1.E6 40 40 640
2.0 2.0 0.0 0.0 0.0 16 0
```

This restart file will change DTADJ and DELTAT to 2.0. The KZ(16) is changed to 0. All other parameters that are set to 0.0 (TADJ, TNEXT, TCRIT, QE) keep same as before. The detail description of restart parameters are shown in Chapter 4.

3.6 Code structure

The code NBODY6++ is written in Fortran 77 and consists of about 300 files. Their functionality was improved as well as new routines included all the way through the decades along with the technological achievements of the hardware. The starting (main) routine is called nbody6.F.

Most of the files have the suffix .f, .F, .cu or .h. All .f files are directly read by a Fortran compiler. The .F files will pass preprocessor first, which selects code lines separated by preprocessor options, e.g. between #ifdef PARALLEL and #endif, for they activate the parallel code on different multiprocessor machines. By this, some portability between different hardware is ensured at least, and a single processor version of the code can easily be compiled as well. The .h are header files and declare the variables and their blocks. .cu are CUDA codes for GPU parallelization. .cpp are C++ codes for SIMD parallelization.

4 Input variables

The input control file of NBODY6++ (see below), contains a minimum of 90 parameters which guide one simulation run for its technical and physical properties (it is very similar but not identical to the one used for NBODY6). As for the technical aspect, the file supervises the run e.g. for its duration, intervals of the output, or error check; the physical parameters concern the size of a cluster, initial conditions, or a number of optional features related to the numerical problem to be studied. The handling of this input file appears rather entangled at first sight, for it has grown rather historically and "ready-for-use" than custom-oriented. Thus, the input variables are read by different routines (functions) in the code, and the nature of the parameters are woven with each other in some cases. Also, some parameters require additional input, such that the total number of lines and parameters may vary.

In the following, we explain the main input file and give an example of typical values for a simulation of an isolated globular cluster. Then, we proceed to the thresholds.

Input	with	all	options:
-------	------	-----	----------

nbody6.F	KSTART	TCOMP	TCRITp	isernb	iserreg	iserks				
input.F	N	NFIX	NCRIT	NRAND	NNBOPT	NRUN	NCOMM			
	ETAI	ETAR	RS0	DTADJ	DELTAT	TCRIT	QE	RBAR	ZMBAR	
	KZ(1)	KZ(2)	KZ(3)	KZ(4)	KZ(5)	KZ(6)	KZ(7)	KZ(8)	KZ(9)	KZ(10)
	KZ(11)	KZ(12)	KZ(13)	KZ(14)	KZ(15)	KZ(16)	KZ(17)	KZ(18)	KZ(19)	KZ(20)
	KZ(21)	KZ(22)	KZ(23)	KZ(24)	KZ(25)	KZ(26)	KZ(27)	KZ(28)	KZ(29)	KZ(30)
	KZ(31)	KZ(32)	KZ(33)	KZ(34)	KZ(35)	KZ(36)	KZ(37)	KZ(38)	KZ(39)	KZ(40)
	KZ(41)	KZ(42)	KZ(43)	KZ(44)	KZ(45)	KZ(46)	KZ(47)	KZ(48)	KZ(49)	KZ(50)
	DTMIN	RMIN	ETAU	ECLOSE	GMIN	GMAX	SMAX			
data.F	ALPHA	BODY1	BODYN	NBINO	NHIO	ZMET	EPOCH0	DTPLOT		
setup.F	APO	ECC	N2	SCALE						(KZ(5)=2)
	APO	ECC	SCALE							(KZ(5)=3)
	APO	ECC	SCALE							(KZ(5)=3)
	SEMI	ECC	M1	M2						(KZ(5)=4)
	ZMH	RCUT							(KZ(5) =	6&&KZ(24)<0)
scale.F	Q	VXROT	VZROT	RTIDE						
xtrn10.F	GMG	RG0								(KZ(14)=2)
	GMG	DISK	A	В	VCIRC	RCIRC	GMB	AR	G AM	
	RG[1:3]	VG[1:3]								(KZ(14)=3)
	MP	AP	MPDOT	TDELAY					(KZ(14)	=3 KZ(14)=4
	RG[1:3]	VG[1:3]								(KZ(14)=5)
binpop.F	SEMI	ECC	RATIO	RANGE	NSKIP	IDORM			(KZ(8)=	1 KZ(8)>4)
hipop.F	SEMI	ECC	RATIO	RANGE					(KZ(8)>	-0&&KZ(18)>1)
imbhinit.F	MMBH	XBH(1)	XBH(2)	XBH(3)	VBH(1)	VBH(2)	VBH(3)	DTBH		(KZ(24)=1)
cloud0.F	NCL	RB2	VCL	SIGMA	CLM	RCL2				(KZ(13)>0)

nbody6.F:

Run control index
=1: new run (construct new model or read from dat.10)
=2: restart/continuation of a run, needs fort.1
=3: restart + changes of DTADJ, DELTAT, TADJ, TNEXT, TCRIT, QE, J, KZ(J)
=4: restart + changes of ETAI, ETAR, ETAU, DTMIN, RMIN, NCRIT, NNBOPT,
SMAX
=5: restart containing the combination of the control index 3 and 4
Maximum wall-clock time in seconds (parallel runs: wall clock)
Termination time in Myr
For MPI parallel runs: only irregular block sizes larger than this value are executed
in parallel mode (dummy variable for single CPU)

For MPI parallel runs: only regular block sizes larger than this value are executed in iserreg parallel mode (dummy variable for single CPU) For MPI parallel runs: only ks block sizes larger than this value are executed in iserks parallel mode (dummy variable for single CPU) input.F: Total number of particles (single + c.m.s. of binaries; singles + $3 \times$ c.m.s. of binaries < NMAX - 2)NFIX Multiplicator for output interval of data on conf. 3 and of data for binary stars (output each DELTAT×NFIX time steps; compare KZ(3) and KZ(6)) NCRIT Minimum particle number (alternative termination criterion) NRAND Random number seed; any positive integer Desired optimal neighbour number (< LMAX-5) NNBOPT **NRUN** Run identification index NCOMM Frequency to store the dumping data (mydump) **ETAI** Time-step factor for irregular force polynomial Time-step factor for regular force polynomial **ETAR** RS0 Initial guess for all radii of neighbour spheres (*N*–body units) DTADJ Time interval for parameter adjustment and energy check (*N*–body units) DELTAT Time interval for writing output data and diagnostics, multiplied by NFIX (N-body units) **TCRIT** Termination time (*N*–body units) Energy tolerance: QE – immediate termination if DE/E > $5*QE \& KZ(2) \le 1$; – restart if DE/E > 5*QE & KZ(2) > 1 and termination after second restart attempt. RBAR Scaling unit in pc for distance (*N*–body units) **ZMBAR** Scaling unit for average particle mass in solar masses (in scale-free simulations RBAR and ZMBAR can be set to zero; depends on KZ(20)) KZ(1)Save COMMON to file comm. 1 = 1: at end of run or when dummy file STOP is created = 2: every 100*NMAX steps KZ(2)Save COMMON to file comm. 2 and comm. 2 [TIME [NB]] = 1: save at output time and every time interval of NCOMM (NB time unit). Notice after restart, the adjust/output time (TADJ) should be integer times of NCOMM, otherwise the COMMON data would not be dumped. = 2: save at output time and every time interval of NCOMM and restart simulation if energy error DE/E > 5*QESave basic data to file conf.3 at output time (unformatted) KZ(3)KZ(4)(Suppressed) Binary diagnostics on bdat.4 (# = threshold levels <10) KZ(5)Initial conditions of the particle distribution, needs KZ(22)=0 = 0: uniform & isotropic sphere = 1: Plummer random generation

= 2: two Plummer models in orbit (extra input)

- = 3: massive perturber and planetesimal disk (each pariticle has circular orbit, constant separation along radial direction between each neighbor and random phase) (extra input)
- = 4: massive initial binary (extra input)
- = 5: Jaffe model (extra input)
- > 6: Zhao BH cusp model (extra input if KZ(24)<0)
- KZ(6) Output of significant and regularized binaries at main output (bodies.f)
 - = 1: output regularized and significant binaries (IEI>0.1 ECLOSE)
 - = 2: output regularized binaries only
 - = 3: output significant binaries at output time and regularized binaries with time interval DELTAT
 - = 4: output of regularized binaries only at output time
- KZ(7) Determine Lagrangian radii and average mass, particle counters, average velocity, velocity dispersion, rotational velocity within Lagrangian radii (lagr.f)
 - = 1: Get actual value of half mass radius RSCALE by using current total mass
 - \geq 2: Output data at main output and lagr. 7
 - \geq 6: Output Lagrangian radii for two mass groups at lagr.31 and lagr.32 (lagr2.f; based on KZ(5)=1,2; cost is $O(N^2)$)
 - --- methods:
 - = 2,4: Lagrangian radii calculated by initial total mass
 - $= 3, \ge 5$: Lagrangian radii calculated by current total mass (The single/K.S-binary Lagrangian radii are still calculated by initial single/binary total mass)
 - = 2,3: All parameters are averaged within the shell between two Lagrangian radii neighbors
 - > 4: All parameters are averaged from center to each Lagrangian radius
- KZ(8) Primordial binaries initialization and output (binpop.f)
 - —- Initialization:
 - = 0: No primordial binaries
 - $=1, \ge 3$: generate primordial binaries based on KZ(41) and KZ(42) (binpop.F)
 - = 2: Input primordial binaries from first $2 \times NBIN0$ lines of dat.10
 - —- Output:
 - > 0: Save information of primordial binary that change member in pbin.18; binary diagnostics at main output (binout.f)
 - > 2: Output KS binary in bdat.9, soft binary in bwdat.19 at output time
- KZ(9) Binary diagnostics
 - = 1,3: Output diagnostics for the hardest binary below ECLOSE in hbin.39 (adjust.f)
 - ≥ 2: Output binary evolution stages in binev.17 (binev.f)
 - \geq 3: Output binary with degenerate stars in degen. 4 (degen. f)
- KZ(10) K.S. regularization diagnostics at main output
 - > 0: Output new K.S. information
 - > 1: Output end K.S. information
 - \geq 3: Output each integrating step information
- KZ(11) (Suppressed)
- KZ(12) > 0: HR diagnostics of evolving stars with output time interval DTPLOT in sse.83 (single star) and bse.82 (K.S. binary)
 - =-1: used if KZ(19)= 0 (see details in KZ(19) description)
- KZ(13) Interstellar clouds

```
= 1: constant velocity for new cloud
            > 2: Gaussian velocity for new cloud
KZ(14)
           External tidal force
           = 1: standard solar neighbor tidal field
           = 2: point-mass galaxy with circular orbit (extra input)
           = 3: point-mass + disk + halo + Plummer (extra input)
           = 4: Plummer model (extra input)
           = 5: Milky-Way potential (extra input) based on Python Galpy package (Bovy, 2015)
KZ(15)
           Triple, quad, chain and merger search
            \geq 1: Switch on triple, quad, chain (KZ(30)>0) and merger search (impact.f)
            \geq 2: Diagnostics at main output at begin and end of triple, quad
            ≥ 3: Save first five outer orbits every half period of wide quadruple before merger
           and stable quadruples accepted for merger in quastab.89
KZ(16)
           Auto-adjustment of regularization parameters
            ≥ 1: Adjust RMIN, DTMIN & ECLOSE every DTADJ time
           \geq 3: modify RMIN for GPERT > 0.05 or < 0.002 in chain; output diagnostics at
           kscrit.77
KZ(17)
           Auto-adjustment of ETAI, ETAR and ETAU by tolerance QE every DTADJ time
           (check.f)
           > 1: Adjust ETAI, ETAR
            > 2: Adjust ETAU
KZ(18)
           Hierarchical systems
           = 1,3: diagnostics (hiarch.f)
           ≥ 2: Initialize primordial stable triples, number is NHIO (hipop.F)
            ≥ 4: Data bank of stable triple, quad in hidat.87 (hidat.f)
KZ(19)
           Stellar evolution mass loss
           = 0: if KZ(12)= -1, the output data will keep the input data unit if KZ(22)= 2-4
           or N-body units if KZ(22) = 6 - 10
           = 1,2: supernova scheme
            \geq 3: Eggleton, Tout & Hurley
            > 5: extra diagnostics (mdot.F)
           = 2,4: Input stellar parameters from fort.21 (instar.f)
                N lines of (MI, KW, M0, EPOCH1, OSPIN)
                MI: Current mass
                KW: Kstar type
                M0: Initial mass
                EPOCH1: evolved age of star (Age = TIME[Myr] - EPOCH1)
                OSPIN: angular velocity of star
KZ(20)
           Initial mass functions, need KZ(22)=0 or 9:
           = 0: self-defined power-law mass function using ALPHAS (data.F)
           = 1: Miller-Scalo-(1979) IMF (imf.f)
           = 2,4: KTG (1993) IMF (imf2.f)
           = 3,5: Eggleton-IMF (imf2.f)
           = 6,7: Kroupa(2001) (imf2.f), extended to Brown Dwarf regime (imfbd.f)
           — Primordial binary mass
           = 2,6: random pairing (imf2.f)
           = 3,4,5,7: binary mass ratio corrected by (m_1/m_2)' = (m_1/m_2)^{0.4} + \text{constant} (Eggle-
           ton, imf2.f)
```

- = 8: binary mass ratio $q = m_1/m_2$ ($m_2 \le m_1$) use distribution $0.6q^{-0.4}$ (Kouwenhoven)
- KZ(21) Extra diagnostics information at main output every DELTAT interval (output.F)
 - \geq 1: output NRUN, MODEL, TCOMP, TRC, DMIN, AMIN, RMAX, RSMIN, NEFF
 - \geq 2: Number of escapers NESC at main output will be counted by Jacobi escape criterion (cost is $O(N^2)$, jacobi.f)
- KZ(22) Initialization of basic particle data mass, position and velocity (data.F)
 - Initialization with internal method
 - = 0,1: Initial position, velocity based on KZ(5), initial mass based on KZ(20)
 - = 1: write initial conditions in dat.10 (scale.F)
 - --- Initialization by reading data from dat . 10
 - = 2: input through NBODY-format (7 parameters each line: mass, position(1:3), velocity(1:3))
 - = 3: input through Tree-format (data.F)
 - = 4: input through Starlab-format
 - = 6: input through NBODY-format and do scaling
 - = 7: input through Tree-Format and do scaling
 - = 8: input through Starlab-format and do Scaling
 - = 9: input through NBODY-format but ignore mass (first column) and use IMF based on KZ(20), then do scaling
 - = 10: input through NBODY-format and all units are astrophysical units (mass: M_{\odot} ; position: pc; velocity: km/s)
- KZ(23) Removal of escapers (escape.F)
 - \geq 1: remove escapers and ghost particles generated by two star coalescence (collision)
 - =2,4: write escaper diagnostics in esc 11
 - \geq 3: initialization & integration of tidal tail
- KZ(24) Initial conditions for subsystems
 - < 0: ZMH & RCUT (N-body units) Zhao model (Need KZ(5)≥6, setup.F)
 - =1: Add one massive black hole (extra input: mass, position, velocity and output frequency), will output black hole data in mbh.45 and its neighbor data in mbhnb.46
- KZ(25) Velocity kicks for white dwarfs (kick.F)
 - = 1: Type 10 Helium white dwarf & 11 Carbon-Oxygen white dwarf
 - = 2: All WDs (type 10, 11 and type 12 Oxygen-Neon white dwarf)
- KZ(26) Slow-down of two-body motion, increase the regularization integration efficiency
 - \geq 1: Apply to KS binary
 - \geq 2: Apply to chain
 - = 3: Rectify to get better energy conservation
- KZ(27) Two-body tidal circularization (Mardling & Aarseth, 2001; Portegies Zwart et al. 1997)
 - (Please suppress in KS parallel version)
 - = 1: sequential
 - = 2: chaos
 - =-1: Only detect collision and suppress coalescence
- KZ(28) Magnetic braking and gravitational radiation for NS or BH binaries (Need KZ(19)=3)
 - ≥ 1: GR coalescence for NS & BH (brake.f, brake3.f)
 - ≥ 2: Diagnostics at main output (brake.f)

- = 3: Input of ZMH = 1/SQRT(2*N) (Need KZ(5) \geq 6) (setup.F)
- = 4: Set every star as type 13 Neutron star (instar.f)
- KZ(29) (Suppressed) Boundary reflection for hot system
- KZ(30) Hierarchical system regularization
 - = -1: Use chain only
 - = 0: No triple, quad and chain regularization, only merger
 - = 1: Use triple, quad and chain (impact.f)
 - ≥ 2: Diagnostics at begin/end of chain at main output
 - \geq 3: Diagnostics at each step of chain at main output
- KZ(31) Centre of mass correction after energy check (cmcorr.f)
- KZ(32) Adjustment (increase) of adjust interval DTADJ, output interval DELTAT and energy error criterion QE based on binding energy of cluster (check.f)
- KZ(33) Block-step statistics at main output (diagnostics)
 - \geq 1: Output irregular block step; and K.S. binary step if KZ(8)>0
 - \geq 2: Output regular block step
- KZ(34) Roche-lobe overflow
 - = 1: Roche & Spin synchronisation on binary with circular orbit (synch.f)
 - = 2: Roche & Tidal synchronisation on binary with circular orbit by BSE method (bsetid.f)
- KZ(35) TIME reset to zero every 100 time units, total time is TTOT = TIME + TOFF (offset.f)
- KZ(36) (Suppressed) Step reduction for hierarchical systems
- KZ(37) Neighbour list additions (checkl.F)
 - ≥ 1: Add high-velocity particles into neighbor list
 - \geq 2: Add small time step particle (like close encounter particles near neighbor radius) into neighbor list
- KZ(38) Force polynomial corrections during regular block step calculation
 - = 0: no corrections
 - = 1: all gains & losses included
 - = 2: small regular force change skipped
 - = 3: fast neighbour loss only
- KZ(39) Neighbor radius adjustment method
 - = 0: The system has unique density centre and smooth density profile
 - $=1,\geq 3$: The system has no unique density centre or smooth density profile skip velocity modification of RS(I) (regint.f, regcor_gpu.f)
 - do not reduce neighbor radius if particle is outside half mass radius
 - reduce RS(I) by multiply 0.9 instead of estimation of RS(I) based on NNBOPT/NNB when neighbor list overflow happens (fpoly0.F, util_gpu.F)
 - = 2,3: Consider sqrt(particle mass / average mass) as the factor to determine the particle's neighbor membership. (fpoly0.F, util_gpu.F)
- KZ(40) = 0: For the initialization of particle time steps, use only force and its first derivative to estimate. This is very efficient.
 - > 0: Use Fploy2 (second and third order force derivatives calculation) to estimate the initial time steps. This method provide more accurate time steps and avoid incorrent time steps for some special cases like initially cold systems, but the computing cost is much higher $(O(N^2))$
- KZ(41) proto-star evolution of eccentricity and period for primordial binaries initialization (proto_star_evol, binpop.F)

KZ(42) Initial binary distribution

= 0: RANGE>0: uniform distribution in log(semi) between SEMI0 and SEMI0/RANGE

RANGE<0: uniform distribution in semi between SEMI0 and -1*RANGE.

- = 1: linearly increasing distribution function f = 0.03438 * logP
- = 2: f = 3.5logP/[100 + (logP) **2]
- = 3: f = 2.3(logP 1)/[45 + (logP 1)**2]; This is a "3rd" iteration when pre-ms evolution is taken into account with KZ(41)=1
- = 4: f = 2.5(logP 1)/[45 + (logP 1)**2]; This is a "34th" iteration when pre-ms evolution is taken into account with KZ(41)=1 and RBAR<1.5
- = 5: Duquennoy & Mayor 1991, Gaussian distribution with mean $\log P = 4.8$, SDEV in $\log P = 2.3$. Use Num.Recipes routine gasdev.f to obtain random deviates given "idum1"
- KZ(43) (Unused)
- KZ(44) (Unused)
- KZ(45) (Unused)
- KZ(46) HDF5/BINARY/ANSI format output and global parameter output (main output, see chapter 13 for details)
 - = 1,3: HDF5(if HDF5 is compiled)/BINARY format
 - = 2,4: ANSI format
 - = 1,2: Only output active stars with time interval defined by KZ(47)
 - = 3,4: Output full particle list with time interval defined by KZ(47)
- KZ(47) Frequency for KZ(46) output

Output data with time interval $0.5^{KZ(47)} \times SMAX$

- KZ(48) (Unused)
- KZ(49) Computation of Moments of Inertia (with Chr. Theis) in fort .60 (ellan.f)
- KZ(50) For unperted KS binary. The neighbor list is searched for finding next KS step. It is safer to get correct step but not efficient when unperted binary number is large. To suppress this, set to 1

DTMIN	Time–step criterion for regularization search
RMIN	Distance criterion for regularization search
ETAU	Regularized time-step parameter (6.28/ETAU steps/orbit)
ECLOSE	Binding energy per unit mass for hard binary (positive)
GMIN	Relative two-body perturbation for unperturbed motion
GMAX	Secondary termination parameter for soft KS binaries
SMAX	Maximum time-step (factor of 2 commensurate with 1.0)

data.F:

ALPHA	Power-law index for initial mass function, routine data.F
BODY1	Maximum particle mass before scaling (based on KZ(20); solar mass unit)
BODYN	Minimum particle mass before scaling
NBIN0	Number of primordial binaries (need KZ(8)>0)
	by routine imf2.F using a binary IMF (KZ(20)≥2)
	by routine binpop.F splitting single stars (KZ(8)>0)
	- by reading subsystems from dat . 10 (KZ(22)≥2)
NHI0	Number of primordial hierarchical systems (need $KZ(18) \ge 2$)

ZMET EPOCH0 DTPLOT	Metal abundance (in range 0.03 - 0.0001) Evolutionary epoch (in 10^6 yrs) Plotting interval for stellar evolution HRDIAG (N-body units; \geq DELTAT)
setup.F:	if (kz(5)=2)
APO	Separation of two Plummer models in N -body units (SEMI = APO/(1 + ECC). (No-
AIO	tice SEMI will be limited between 2.0 and 50.0)
ECC	Eccentricity of two-body orbit (ECC ≥ 0 and ECC < 0.999)
N2	Membership of second Plummer model (N2 <= N)
SCALE	Scale factor for the second Plummer model, second cluster will be generated by first
	Plummer model with $X \times SCALE$ and $V \times \sqrt{SCALE} (\ge 0.2 \text{ for limiting minimum})$
	size)
	if (kz(5)=3)
APO	Separation between the perturber and Sun in <i>N</i> –body units
ECC	Eccentricity of orbit (=1 for parabolic encounter)
SCALE	Perturber mass scale factor, perturber mass = Center star mass \times SCALE (=1 for
	Msun)
	if $(kz(5)=4)$
SEMI	Semi-major axis (slightly modified; ignore if ECC > 1)
ECC	Eccentricity (ECC > 1: NAME = 1 & 2 free-floating)
M1	Mass of first member (in units of mean mass)
<u>M2</u>	Mass of second member (rescaled total mass = 1)
77) ([]	if $(kz(5)\geq 6)$ and $(kz(24)<0)$
ZMH	Mass of single BH (in N-body units)
RCUT	Radial cutoff in Zhao cusp distribution (MNRAS, 278, 488)
scale.F:	
Q	Virial ratio (routine scale.F; Q=0.5 for equilibrium)
VXROT	XY-velocity scaling factor (> 0 for solid-body rotation)
VZROT	Z-velocity scaling factor (not used if $VXROT = 0$)
RTIDE	Unscaled tidal radius for $KZ(14)=2$ and $KZ(22)\geq 2$. If not zero, RBAR = RT/RTIDE
	where RT[pc] is tidal radius calculated from input GMG and RG0
xtrn10.F:	: if (kz(14)=2)
GMG	Point-mass galaxy (solar masses, linearized tidel field in circular orbit)
RG0	Central distance (in kpc)
	if (kz(14)=3)
GMG	Point-mass galaxy (solar masses)
DISK	Mass of Miyamoto disk (solar masses)
A	Softening length in Miyamoto potential (in kpc)
В	Vertical softening length (kpc)
VCIRC	Galactic circular velocity (km/sec) at RCIRC (=0: no halo)
RCIRC	Central distance for VCIRC with logarithmic potential (kpc)
GMB	Dehnen model budge mass (solar masses)
AR	Dehnen model budge scaling radius (kpc)
GAM	Dehnen model budge profile power index gamma

RG	Initial position; DISK+VCIRC=0, VG(3)=0: A(1+E)=RG(1), E=RG(2)
VG	Initial cluster velocity vector (km/sec)
	if (kz(14)=3,4)
MP	Total mass of Plummer sphere (in scaled units)
AP	Plummer scale factor (N-body units; square saved in AP2)
MPDOT	Decay time for gas expulsion (MP = $MPO/(1 + MPDOT*(T-TD))$
TDELAY	Delay time for starting gas expulsion (T > TDELAY)
	if (kz(14)=5)
RG	Initial cluster c.m. position vector [kpc]
VG	Initial cluster c.m. velocity vector [km/sec]
binpop.F:	if (kz(8)=1 or kz(8)>2)
SEMI	Initial semi-major axis limit
ECC	Initial eccentricity
	< 0: thermal distribution, $f(e) = 2e$
	≥ 0 and ≤ 1 : fixed value of eccentricity
	= 20: uniform distribution
	= 30: distribution with $f(e) = 0.1765/(e^2)$
	= 40: general $f(e) = a * e^b$, $e^0 <= e <= 1$ with $a = (1+b)/(1-e^0(1+b))$, current
	values: $e0 = 0$ and $b = 1$ (thermal distribution)
RATIO	$KZ(42) \le 1$: Binary mass ratio $M1/(M1+M2)$
	$KZ(42) = 1.0$: $M1 = M2 = \langle M \rangle$
RANGE	KZ(42)= 0: semi-major axis range for uniform logarithmic distribution;
Tu ii (OL	not used for other $KZ(42)$
NSKIP	Binary frequency of mass spectrum (starting from body #1)
IDORM	Indicator for dormant binaries (> 0: merged components)
hipop.F:	if (kz(8)>0 and kz(18)>1)
SEMI	Max semi-major axis in model units (all equal if RANGE = 0)
ECC	Initial eccentricity (< 0 for thermal distribution)
RATIO	Mass ratio (= 1.0: $M1 = M2$; random in [0.5 \sim 0.9])
RANGE	Range in SEMI for uniform logarithmic distribution (>0)
imbhinit.	F:if (kz(24)=1)
MMBH	Mass of massive black hole in solar mass unit
XBH(1:3)	3 dimensional position of massive black hole in pc
VBH(1:3)	3 dimensional velocity of massive black hole in km/s
DTBH	Output interval for massive black hole data in mbh. 45 and mbhnb. 46 (N-body unit)
cloud0.F:	if (kz(13)>0)
NCL	Number of interstellar clouds
RB2	Radius of cloud boundary in pc (square is saved)
VCL	Mean cloud velocity in km/sec
SIGMA	Velocity dispersion (KZ(13)>1: Gaussian)
CLM	The state of the s

Individual cloud masses in solar masses (maximum MCL)

CLM

RCL2 Half-mass radii of clouds in pc (square is saved)

A typical input file can look like as follows. It defines a new simulation running for 1,000,000 CPU-minutes with N=16,000 particles distributed from a Plummer profile (KZ(5)=1). The run may alternatively terminate when TCRIT=1000.0 N-body units. or if the final particle number of NCRIT=10 has been reached. The output and adjustment time interval DELTAT/DTADJ are 1.0 N-body unit. The initial mass function follows Kroupa, (2001) with mass ranging from $m_{\text{max}}=20.0M_{\odot}$ to $m_{\text{min}}=0.08M_{\odot}$ (BODY1 and BODYN). The initial virial ratio is 0.5 (equilibrium). The stellar evolution is switched on (KZ(19)=3) and initial metallicity is 0.001. Multiples and chain regularization are switched on (KZ(15)=2 and KZ(30)=2). It uses solar neighbor tidal field (KZ(14)=1).

```
1 1000000.0 1.E6 40 40 640

16000 1 10 43532 100 1

0.02 0.02 0.1 1.0 1.0 1000.0 2.0E-05 1.0 0.7

0 1 1 0 1 0 4 0 0 2

0 1 0 1 2 1 0 0 3 6

1 0 2 0 0 2 0 0 0 0 2

1 0 2 1 1 0 1 1 0 0

0 0 0 0 0 0 0 0 0 0

1.0E-06 1E-4 0.2 1.0 1.0E-06 0.01 1.0

2.35 20.0 0.08 0 0 0.001 0 1.0

0.5 0.0 0.0 0.0
```

Input variables for primordial Binaries

Many star clusters contain initial hard binaries with binding energies much larger than the thermal energy (the threshold ECLOSE is a suitable division between hard and soft binaries). There are two ways to initialise primordial binaries:

The first one always starts from some initial mass function (IMF) provided by the routines imf.f or imf2.f. The option KZ(8)=1 or ≥ 3 invokes the routine binpop.F, which reads the last line of the input file containing NBIN and the parameters of their distribution (see above). In this case, binaries are created either by random pairing of single stars obtained from the IMF or by splitting them, depending on the value of KZ(20) (see above).

The second way assumes that particle data, including the binaries, are provided via the input data on file dat.10 (as e.g. in the Kyoto–II collaborative experiment). In such a case KZ(8)=2 and NBINO should be set to the expected number of primordial binaries from the file. The code will first create NBINO centers of masses, and then use those for scaling, before regularizing the pairs and the calculation begins.

A typical input file with primordial binaries looks as follows. Here, we use binary random pairing from imf2.f and binpop.F (KZ(20)=6 and KZ(8)=3, respectively) for 1000 initial binaries. The semi-major axes of binaries use uniform distribution in log(semi) with a range from 41.3 AU to 0.00413 AU. The eccentricity of binaries use thermal distribution. It was created from this input file running for 1000 time units. Stellar evolution was also switched on in this file (KZ(19)=3). In the package of the code, the file N10k_B1k.input is included.

```
1 1000000.0 1.E6 40 40 640
10000 1 10 43532 100 1
0.02 0.02 0.17 1.0 1.0 800.0 5.0E-05 1.0 0.7
0 1 1 0 1 0 4 3 0 2
0 1 0 1 2 1 0 0 3 6
1 0 2 0 0 2 0 0 0 2
1 0 2 1 1 0 1 1 0 0
0 0 0 0 0 0 0 0 0
5.0E-06 3E-4 0.2 1.0 1.0E-06 0.01 0.5
2.35 100.0 0.08 1000 0 0.001 0 1.0
0.5 0.0 0.0 0.0
2E-4 -1.0 1.0 1E4 5 0
```

Stellar Evolution

Stellar evolution is invoked by KZ(19)=1,2 or $KZ(19)\geq 3$, offering two different schemes. The simpler one is KZ(19)=1, while the more complex one, $K(19)\geq 3$, is based on the Cambridge stellar evolution package (Hurley, Pols, Tout 2000). The common envelope, roche transfering binaries are also considered. The main effects are changing stellar masses, radii, and luminosities, which give rise to cluster mass loss. The mass is assumed to escape from the cluster immediately and possible collisions depend on stellar radii.

With the additional option KZ(12)>0, information on binaries and single stars is written on two files (unit 82, file bev.82 and unit 83, file sev.83) in regular time intervals determined by TPLOT (See details in Section Output).

Restart

Users may want to restart the simulation from a stopped run. The input file for restart is different from an input for a new simulation. The first line is similar, only the first parameter KSTART is needed to change from 1 to 2-5. The second/third lines are based on the value of KSTART:

KSTART	Function
1	new run, start from initial values given in data. F
2	continuation of a run without changes
3	restart of a run with changes of the following parameters given in
	the second line of a newly created input file:
	DTADJ, DELTAT, TADJ, TNEXT, TCRIT, QE, J, K
	where the options KZ can be changed via KZ(J)=K
4	restart of a run with following parameters changed in the second
	line: ETAI, ETAR, ETAU, DTMIN, RMIN, NCRIT, NNBOPT,
	SMAX
5	restart of a run with all parameter changes in the run control index
	3 and 4. The changes must succeed the first line.

[&]quot;0" values in the fields are interpreted as: Do not change the value of this parameter.

Notice the COMMON file comm.1 is required together with input file of parameters for successfully restarting simulations. Based on KZ(1) and KZ(2), comm.1/2[_*] are generated during simulations. The NCOMM control the frequency of the comm.2[_*]. For example, if NCOMM=10, every 10 NB time unit, a comm.2[_*] file is generated. All these files can be used to restart a simulation from the time shown in the suffix [_*]. Before using comm.2[_*] for restarting, the filename should be firstly renamed to comm.1.

5 Thresholds for the variables

Before the compilation of the code (Chapter 3), the parameter file (params . h) should be consulted to check whether some vector dimensions are in the desired range. Most important are

- the maximum particle number NMAX,
- the maximum number of regularised KS pairs KMAX, and
- the maximum number of neighbours per particle LMAX.

The particles are saved in various lists which serve to distinguish between their funcionality. The table below explains their nomenclature. "KS-pairs" are particles that approach each other in a hyperbolic encounter; they are given a special treatment by the code (see Chapter 11). If NPAIRS is the amount of KS-pairs, then IFIRST = 2*NPAIRS + 1 is the first single particle (not member of a KS pair), and N the last one. NTOT = N + NPAIRS is the total number of particles plus c.m.'s. Therefore NMAX, the dimension of all vectors containing particle data should be at least of size N + KMAX, where N is the number of particles and KMAX the maximum number of expected KS pairs. If one starts with single particles, KMAX = 10 or 20 should usually be enough, but in clusters with a large number of primordial binaries, KMAX must be large.

N: Total number of particles

NBINO: number of primordial binaries (physical bound stars)

NBIN: ???

NPAIRS: Number of binaries (KS-pairs, see Chapter 11), transient unbound pairs as well as

persistent binaries

NTOT: = N + NPAIRS;

Number of single particles plus centres of masses of regularized (KS) pairs

KMAX: threshold for the amount of allowed KS pairs

NMAX: = N + KMAX; threshold for the total number of particles and the centre of masses

Hier gibt's noch ein Bildchen!

6 How to read the diagnostics

The diagnostics is the ASCII readable text printed on unit 6 *stdout* ("out1000" in Chapter 3) that gives a brief overview of the global status and progress of the cluster simulation. Different routines write into that file, depending on the options chosen as the input variables. The following lines occur:

N NFIX NCRIT NRAND NNBOPT NRUN	
1000 5 10 1006 50 1	
ETAI ETAR RSO DTADJ DELTAT TCRIT _P TCRIT QE RBAR ZMBAR	
1.0E-02 2.0E-02 3.0E-01 1.0E+01 1.0E+01 1.0E+06 2.0E+01 2.0E-05 1.0E+00 7.0E-01	
OPTIONS	written by the routine:
1 2 3 4 5 6 7 8 9 10 11 12 13 14 15 16 17 18 19 20 21 22 23 24 25 26 27 28 29 30 31 32 33 34 35 36 37 38 39 40	input.F
1 1 1 0 1 0 4 0 0 2 1 0 0 0 1 1 1 0 0 0 1 0 1 0 1 0 0 2 0 0 2 0 0 2 0 1 0 1	Input.i
OPTIONS BK:	Usage: Repetition of the input variables
1 2 3 4 5 6 7 8 9 10	input variables
0 0 0 0 0 0 0 0	
DTMIN RMIN ETAU ECLOSE GMIN GMAX	
1.0E-04 1.0E-02 1.0E-01 1.0E+00 1.0E-06 1.0E-02	
****** NOTE: new random number seed initialisation! ****** AND new ran2 from new ed. of Press et al.	
STANDARD IMF ALPHA = 2.35 BODY1 = 20.0 BODYN = 0.10 ZMASS = 3.36752E+02 NBINO= 0 ZMET = 0.00 EPOCHO = 0.00	data.F, (if KZ(20)=0 &
DIMADN CTAR INC. ND = 400 DANGE = 2.075404 0.447 04 7MD = 2.745400 4MD = 0.087 04	BODY1≠BODYN)
BINARY STAR IMF: NB = 400 RANGE = 3.27E+01 2.14E-01 ZMB = 3.74E+02 <mb> = 9.36E-01</mb>	,
SINGLE STAR IMF: NS = 1200 RANGE = 7.76E+00 1.01E-01 ZMS = 5.25E+02 <ms> = 4.38E-01</ms>	or
MF power law index, max. mass, min. mass, total mass, # of primordial bin., metallicity, evolution. epoch [Myrs].	$imf2.F$, if $KZ(20) \ge 2$
umber of objects, mass range, average mass before scaling.	Information about initial mass function (IMF).

```
SCALING: SX = 1.00421D+00 E = -2.49E-01 M(1) = 5.94E-02 M(N) = 2.97E-04 <M> = 1.00E-03

TIME SCALES: TRH = 2.8E+01 TCR = 2.8E+00 2<R>/<V> = 2.8E+00

PHYSICAL SCALING: R* = 1.0000E+00 M* = 7.0000E+02 V* = 1.7348E+00 T* = 5.6466E-01 <M> = 7.0000E-01

SU = 4.4335E+07 AU = 2.0627E+05 YRS = 3.5408E+06
```

scaling factor for energy, total energy, max. mass, min. mass, average mass after scaling;

Spitzer's half-mass relaxation time, crossing time obtained from total energy and mass, crossing time obtained from virial radius (see 12);

information about physical scaling: values of one N-body unit in length (pc), mass (solar masses), velocity (km/s), time (million years), average mass of particles (solar masses), astronomical units (one N-body unit) and years (one N-body unit).

```
fpoly1 time= 0.1200000035762785
fpoly2 time= 0.2100000062584875
```

CPU (wall clock in parallel execution) time for initialising the force and its time derivative (fpoly1, fpoly_mpi.f) and the second and third time derivative of the force (fpoly2, fpoly2_mpi.f). The mpi-versions are called for initialisation in case of parallel runs.

start.F

scale.F, units.f

```
lagr.F
TIME M/MT: 1.00D-02 2.00D-02 5.00D-02 1.00D-01 2.00D-01 3.00D-01 4.00D-01 5.00D-01 7.00D-01 9.00D-01 1.00D+00 <RC
 0.0 RLAGR: 1.52D-01 1.81D-01 1.91D-01 2.83D-01 4.33D-01 5.22D-01 6.17D-01 7.52D-01 1.16D+00 1.96D+00 5.86D+00 2.97D-01
 0.0 AVMASS: 6.26D-04 4.26D-03 6.13D-04 9.45D-04 7.82D-04 1.11D-03 1.09D-03 8.35D-04 9.19D-04 1.25D-03 9.01D-04 1.23D-03
                                        53
                                                 130
                                                                    91
                                                                            122
                17
                          9
                                  2
                                                           90
                                                                                  216
                                                                                             163
                                                                                                      107
 0.0 SIGR2: 2.19D-01 1.70D-01 7.37D-01 2.44D-01 1.89D-01 2.33D-01 2.16D-01 2.43D-01 1.72D-01 8.36D-02 5.09D-02 2.09D-01
 0.0 SIGT2: 2.01D-01 8.25D-02 4.46D-02 2.12D-01 3.39D-01 2.42D-01 1.71D-01 1.83D-01 1.51D-01 1.03D-01 6.16D-02 1.63D-01
 0.0 VROT: -8.34D-02 4.41D-01 -6.28D-01 3.14D-02 -1.54D-01 -9.44D-02 -6.45D-02 1.17D-02 5.35D-02 1.98D-02 -3.46D-02 4.46D-01
```

Time, specification of the Lagrangian radii, core radius

Time, Lagrangian radii, core radius (if primordial binaries: separately for singles and binaries, not shown above)

Time, average mass between Lagrangian radii, avmass in the core

Time, number of particles within the shell, in the core

Time, radial velocity dispersion within the shell, in the core

Time, tangential vel. dispersion within the shell, in the core

Time, rotational vel. within the shell, in the core (not shown above)

rank, "ADJUST:", total time in NB units, physical time, virial ratio, relative energy error, total energy, total energy of regularized pairs, energy of mergers

```
RMIN = 1.1E-03 DTMIN = 3.5E-05 RHOM = 3.5E+02 RSCALE = 9.5E-01 RSMIN = 2.2E-01 ECLOSE = 1.05 TC = 3

PE N ttot tree tirr tyredtot tint tinit tks ttoomm tadj tmov tyrednb tsub tsub2 xtsub1 xtsub2

0 1000 41.46000 29.54 7.23 0.63 40.39 0.99 0.07 0.00 0.59 0.00 1.50 0.00 0.00 0.0000D+00 0.0000D+00
```

close encounter distance and minimum time step (for regularization search, updated from input parameters if KZ(16)=1), maximum density, virial radius, minimum neighbour sphere, hard binary threshold energy, total run time in units of initial crossing times

number of processors, number of particles, total processing time, total regular processing, total irregular processing, processing of prediction, time spent in intgrt.F, for initialisation, for KS integration, for communication, for adjust and energy check, for overhead of moving data in parallel runs, for neighbour predictions, for MPI communication after irregular (tsub) and regular (tsub2) blocks, number of bytes transferred respectively. From xtsub1/tsub and xtsub2/tsub2 the sustained bandwidth of MPI communication can be read off. Note, that the determination of these quantities involves a certain overhead by many calls of cputim.F per block, so for critically large production runs one may want to comment these out (most of them in intgrt.F).

adjust.F

```
O NM = O MM = O NS = 1000 NSTEPS =
                                                                               273 321696
                                                                                                                             -0.250004
        M# = 1 CPH = 6 91000E-01 TRC = 0 0 DMIN = 6 6E-05 6 6E-05 1 0E+02 1 0E+02 AMIN = 1 0E+02 RMAX = 0 0E+00 RSMIN = 0 22 NEFF = 128
                        NC MC RHOD RHOM CMAX <Cn> Ir/R
#1 0 95 9 5 0 21 0 08
                       5 0.073 159. 350. 5. 37.0 0.13
                                                           0 0 0.000 0.0000 0.006197 0.000 0.000 2.83
            NBCORR NBFULL NBVOID NRCONV NICONV NBSMIN NBDIS NBDIS2 NCMDER NBDER NFAST NBFAST
                                         9227 1576
                                                       0
                                                              0
                                                                    33
                        NKSPER NPRECT NKSREF NKSMOD NTTRY NTRIP NQUAD NCHAIN NMERG NSTEPT NSTEPQ NSTEPC
                                0 0
                                              0
                                                   0
                                                         0
                                                               0
```

output.F

time, actual particle number, average neighbour number, number of KS pairs, number of merged KS pairs, number of hierarchical subsystems, number of single stars, step numbers (irregular, irr. c.m., regular, KS), relative energy error since last output, total energy

several more lines uncommented here....

```
STEP I 0 3 63 91 154 220 160 133 109 44 19 4
STEP R 0 4 77 133 249 310 179 45 3
Max Speedup Irr: 4 3.76D+00 8 6.89D+00 16 1.12D+01 32 1.66D+01 64 2.15D+01 128 2.49D+01 256 2.66D+01 512 2.74D+01 1024 2.77D+01
Max Speedup Reg: 4 3.71D+00 8 6.69D+00 16 1.12D+01 32 1.67D+01 64 2.22D+01 128 2.49D+01 256 2.91D+01 512 3.04D+01 1024 3.11D+01
```

histogram of distribution of irregular (STEP I), regular (STEP R)

If there are p step distribution (not appearing here, STEP U, in physical time), statistics of parallel work for irr. and reg. steps, figures given are theoretical speedups for infinitely fast communication (limit of large block sizes)

```
END RUN TIME[Myr] = 11.29 TOFF/TIME/TTOT= 0.00000000 20.00000000 CPUTOT = 1.6 ERRTOT =-5.15000D-05 DETOT =-1.28197D-05

0 INTEGRATION INTERVAL = 20.00 NIRR= 3237662 NIRRB= 1245 NREG= 779010 NKS= 4625

PER TIME UNIT: NIRR= 1.61883D+05 NIRRB= 6.22500D+01 NREG= 3.89505D+04 NKS= 2.31250D+02

Total CPU= 97.11000289410342
```

This is the regular end of a run giving: the integration time, total cumulative absolute and relative errors, cumulative number of regular, irregular, KS steps, the step numbers per time unit and the total CPU (wall clock for parallel) time in minutes.

adjust.F

levels.f

To check a regular stop of the run, look at the end of the diagnostics first. If there are failures, the line "CALCULATION HALTED" appears and means that the energy conservation could not be guaranteed. A restart with smaller steps (ETAI, ETAR) and larger neighbour number NNBOPT may cure the problem, but not always; persistent problems should be reported to Rainer Spurzem.

The unix command on the output file, e.g.

homedir> grep ADJUST out1000

produces an overview of the accuracy (energy error at every DTADJ interval). It may show where problems originated; a restart from the last ADJUST before the error with smaller output intervals is one way to look after it. Watch out, because sometimes errors are not reproducible, because changes in ADJUST intervals change frequencies of prediction and small differences can build up. A quick possibility to see the real evolution of the system is to grep for the lines with Lagrangian radii and other quantities (see above), which can directly be plotted, e.g. with gnuplot, because the first column is always the time.

7 Runs on parallel machines

For parallel runs, the file mpif.h is very important, and system specialists should be consulted in addition to us what to use. Again, for some standard systems templates are provided (e.g. mpif.t3e.h or mpif.mpich.h). The routine providing CPU-time measurements, cputim.F, and the use of the function flush.f may need special attention depending on the hardware.

8 The Hermite integration method

Each particle is completely specified by its mass m, position \mathbf{r}_0 , and velocity \mathbf{v}_0 , where the subscript 0 denotes an initial value at a time t_0 . The equation of motion for a particle i is given by its momentary acceleration $\mathbf{a}_{0,i}$ due to all other particles and its time derivative $\dot{\mathbf{a}}_{0,i}$ as

$$\mathbf{a}_{0,i} = -\sum_{i \neq j} Gm_j \frac{\mathbf{R}}{R^3}, \tag{1}$$

$$\dot{\mathbf{a}}_{0,i} = -\sum_{i \neq j} Gm_j \left[\frac{\mathbf{V}}{R^3} + \frac{3\mathbf{R}(\mathbf{V} \cdot \mathbf{R})}{R^5} \right], \tag{2}$$

where *G* is the gravitational constant; $\mathbf{R} = \mathbf{r}_{0,i} - \mathbf{r}_{0,j}$ is the relative coordinate; $R = |\mathbf{r}_{0,i} - \mathbf{r}_{0,j}|$ the modulus; and $\mathbf{V} = \mathbf{v}_{0,i} - \mathbf{v}_{0,j}$ the relative space velocity to the particle *j*.

The Hermite scheme employed in NBODY6++ follows the trajectory of the particle by firstly "predicting" a new position and new velocity for the next time step t. A Taylor series for $\mathbf{r}_i(t)$ and $\mathbf{v}_i(t)$ is formed:

$$\mathbf{r}_{p,i}(t) = \mathbf{r}_0 + \mathbf{v}_0(t - t_0) + \mathbf{a}_{0,i} \frac{(t - t_0)^2}{2} + \dot{\mathbf{a}}_{0,i} \frac{(t - t_0)^3}{6},$$
(3)

$$\mathbf{v}_{p,i}(t) = \mathbf{v}_0 + \mathbf{a}_{0,i}(t - t_0) + \dot{\mathbf{a}}_{0,i}\frac{(t - t_0)^2}{2}.$$
 (4)

The predicted values of \mathbf{r}_p and \mathbf{v}_p , which result from this simple Taylor series evaluation, using the force and its time derivative at t_0 , do not fulfil the requirements for an accurate high-order integrator; they just give a first approximation to \mathbf{r}_1 and \mathbf{v}_1 at the upcoming time t_1 . Even if the time step, $t_1 - t_0$, is chosen impracticably small, a considerable error will quickly occur, let alone the inadequate computational effort. Therefore, an improvement is made by the Hermite interpolation which approximates the higher accelerating terms by another Taylor series:

$$\mathbf{a}_{i}(t) = \mathbf{a}_{0,i} + \dot{\mathbf{a}}_{0,i} \cdot (t - t_{0}) + \frac{1}{2} \mathbf{a}_{0,i}^{(2)} \cdot (t - t_{0})^{2} + \frac{1}{6} \mathbf{a}_{0,i}^{(3)} \cdot (t - t_{0})^{3}, \tag{5}$$

$$\dot{\mathbf{a}}_{i}(t) = \dot{\mathbf{a}}_{0,i} + \mathbf{a}_{0,i}^{(2)} \cdot (t - t_{0}) + \frac{1}{2} \mathbf{a}_{0,i}^{(3)} \cdot (t - t_{0})^{2}. \tag{6}$$

Here, the values of $\mathbf{a}_{0,i}$ and $\dot{\mathbf{a}}_{0,i}$ are already known, but a further derivation of equation (2) for the two missing orders on the right hand side turns out to be quite cumbersome. Instead, one determines the additional acceleration terms from the predicted ("provisional") \mathbf{r}_p and \mathbf{v}_p ; we calculate their acceleration and time derivative according to the equations (1) and (2) anew and call these new terms $\mathbf{a}_{p,i}$ and $\dot{\mathbf{a}}_{p,i}$, respectively. Because these values ought to be generated by the former high–order terms also (which we avoided), we put them into the left–hand sides of (5) and (6). Solving equation (6) for $\mathbf{a}_{0,i}^{(2)}$, then substituting it into (5) and simplifying yields the third derivative:

$$\mathbf{a}_{0,i}^{(3)} = 12 \frac{\mathbf{a}_{0,i} - \mathbf{a}_{p,i}}{(t - t_0)^3} + 6 \frac{\dot{\mathbf{a}}_{0,i} + \dot{\mathbf{a}}_{p,i}}{(t - t_0)^2}. \tag{7}$$

Similarly, substituting (7) into (5) gives the second derivative:

$$\mathbf{a}_{0,i}^{(2)} = -6 \frac{\mathbf{a}_{0,i} - \mathbf{a}_{p,i}}{(t - t_0)^2} - 2 \frac{2\dot{\mathbf{a}}_{0,i} + \dot{\mathbf{a}}_{p,i}}{t - t_0}.$$
 (8)

Note, that the desired high–order accelerations are found just from the combination of the low–order terms for \mathbf{r}_0 and \mathbf{r}_p . We never derived higher than the first derivative, but achieved the higher orders easily through (1) and (2). This is called the Hermite scheme.

Previously, a four–step Adams–Bashforth–Moulton integrator was used (especially in NBODY5, [2]), however, the new Hermite scheme allows twice as large timesteps for the same accuracy. Also its storage requirements are less [16], [17], [4], [5].

Finally, we extend the Taylor series for $\mathbf{r}_i(t)$ and $\mathbf{v}_i(t)$, eqs. (3) and (4), by two more orders, and find the "corrected" position $\mathbf{r}_{1,i}$ and velocity $\mathbf{v}_{1,i}$ of the particle i at the computation time t_1 as

$$\mathbf{r}_{1,i}(t) = \mathbf{r}_{p,i}(t) + \mathbf{a}_{0,i}^{(2)} \frac{(t-t_0)^4}{24} + \mathbf{a}_{0,i}^{(3)} \frac{(t-t_0)^5}{120},$$
 (9)

$$\mathbf{v}_{1,i}(t) = \mathbf{v}_{p,i}(t) + \mathbf{a}_{0,i}^{(2)} \frac{(t-t_0)^3}{6} + \mathbf{a}_{0,i}^{(3)} \frac{(t-t_0)^4}{24}. \tag{10}$$

The integration cycle for other upcoming steps may now be repeated from the beginning, eqs. (1) and (2). The local error in \mathbf{r} and \mathbf{v} within the two time steps $\Delta t = t_1 - t_0$ is expected to be of order $\mathcal{O}(\Delta t^5)$, the global error for a fixed physical integration time scales with $\mathcal{O}(\Delta t^4)$ [15].

9 Individual and block time steps

Stellar systems are characterized by a huge dynamical range in radial and temporal scales. The time scale varies e.g. in a star cluster from orbital periods of binaries of some days up to the relaxation of a few hundred million years, or even billions of years. Even if we put for a moment the very close binaries aside, which are treated differently (by regularization methods), there typically is a large dynamic range in the average local stellar density from its centre to the very outskirts, where it dissolves into the galactic tidal field. In a classical picture, the two closest bodies would determine the time–step of force calculation for the whole rest of the system. However, for bodies in regions where the changes of the force are relatively small, a permanent re–computing of the terms appears time consuming. So, in order to economize the calculation, these objects shall be allowed to move a longer distance before a recomputation becomes necessary. In between there is always the possibility to acquire particle positions and velocities via a Taylor series prediction, as described in Chapter 8. This is the idea of a vital method for assigning different time–steps, $\Delta t = t_1 - t_0$, between the force computations, the so–called "individual time–step scheme" [1], which was later advanced to the hierarchical block steps.

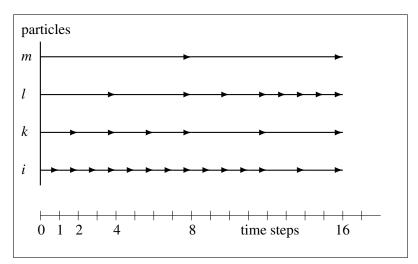


Figure 9.1: Block time steps exemplary for four particles.

Each particle is assigned its own Δt_i which is first illustrated for the case of "block time–steps" in Figure 9.1. The particle named i has the smallest time step at the beginning, so its phase space coordinates are determined at each time step. The time step of k is twice as large as i's, and its coordinates are just extrapolated ("predicted") at the odd time steps, while a full force calculation is due at the dotted times. The step width may be altered or not after the end of the integration cycle for the special particle, as demonstrated for k and l beyond the label "8". The time steps have to stay commensurable with both, each other as well as the total time, such that a hierarchy is guaranteed. This is the block step scheme.

As a first estimate, the rate of change of the acceleration seems to be a reasonable quantity for the choice of the time step: $\Delta t_i \propto \sqrt{\mathbf{a}_i/\dot{\mathbf{a}}_i}$. But it turns out that for special situations in a many-body system, it provides some undesired numerical errors. After some experimentation, the following formula was adopted [2]:

$$\Delta t_i = \sqrt{\eta \frac{|\mathbf{a}_{1,i}||\mathbf{a}_{1,i}^{(2)}| + |\dot{\mathbf{a}}_{1,i}|^2}{|\dot{\mathbf{a}}_{1,i}||\mathbf{a}_{1,i}^{(3)}| + |\mathbf{a}_{1,i}^{(2)}|^2}},$$
(11)

where η is a dimensionless accuracy parameter which controls the error. In most applications it is taken to be $\eta \approx 0.01$ to 0.02, see also next chapter.

For the block–time steps, the synchronization is made by taking the next–lowest integer of Δt_i ; the time steps are quantized to powers of 2 [15]. Then, there will be a group (block) of several particles which are due to movement at each time step. If one keeps the exact Δt_i 's evaluated from (11) for each particle, the commensurability is destroyed, and we arrive at the so–called "individual time steps"; in this case, there exists one sole particle being due. The latter concept is realized in the earlier codes NBODY1, NBODY3, NBODY5, where a neighbour scheme is renounced. NBODY4, NBODY6, and NBODY6++ use a block step scheme.

Subsystems like star binaries, triples or a similar subgroups (they are termed KS pairs, chains, hierarchies) enter the time–step scheme with their respective centre's of masses only. Their internal motion is treated in a different way by a regularized integration (Chapter 11).

10 The Ahmad–Cohen scheme

The computation of the full force for each particle in the system makes simulations very time—consuming for large memberships. Therefore, it is desirable to construct a method in order to speed up the calculations while retaining the collisional approach. One way to achieve this is to employ a "neighbour scheme", suggested by [9].

The basic idea is to split the force polynomial (5) on a given particle i into two parts, an irregular and a regular component:

$$\mathbf{a}_i = \mathbf{a}_{i,\text{irr}} + \mathbf{a}_{i,\text{reg}}.\tag{12}$$

The irregular acceleration $\mathbf{a}_{i,\text{irr}}$ results from particles in a certain neighbourhood of i (in the code, FI and FIDOT are the irregular force and its time derivative at the last irregular step; internally some routines use FIRR and FD as a local variable). They give rise to a stronger fluctuating gravitational force, so it is determined more frequently than the regular one of the more distant particles that do not change their relative distance to i so quickly (in the code, FR and FRDOT are the regular force and its time derivative at the last regular step; some routines use as a local variable FREG and FDR). We can replace the full summation in eq. (1) by a sum over the $N_{\rm nb}$ nearest particles for $\mathbf{a}_{i,\text{irr}}$ and add a distant contribution from all the others. This contribution is updated using another Taylor series up to the order FRDOT, the time derivative of FR at the last regular force computation¹.

Wether a particle is a neighbour or not is determined by its distance; all members inside a specified sphere ("neighbour sphere" with radius r_s) are held in a list, which is modified at the end of each "regular time–step" when a total force summation is carried out. In addition, approaching particles within a surrounding shell satisfying $\mathbf{R} \cdot \mathbf{V} < 0$ are included. This "buffer zone" serves to identify fast approaching particles before they penetrate too far inside the neighbour sphere. The neighbour criterion should be improved according to relative forces rather than distances, in particular, if there are very strong mass differences between particles (black holes!) — such kind of work is under progress.

Figures 10.1 and 10.2 show how the Ahmad–Cohen scheme works for one particle [17]. At the beginning of the force calculation, a list of neighbour objects around the particle i is created first (filled dots). From this neighbour list the irregular component $\mathbf{a}_{i,\text{irr}}$ is calculated, and then the summation is continued to the distant particles obtaining $\mathbf{a}_{i,\text{reg}}$. At the same time we also calculate the first time derivative. From the equations (5) and (6) the position and velocity of the particle i are predicted. At time $t_{1,\text{irr}}$ we apply the "corrector" only for $\mathbf{a}_{i,\text{irr}}$ from the neighbours; the regular component we do not correct, but obtain by extrapolating $\mathbf{a}_{i,\text{reg}}$. At the next step, $t_{2,\text{irr}}$, the same predictor–corrector method proceeds for the neighbour particles, while the correction of the distant acceleration term is still neglected. When t_1 is reached, the total force is calculated on the basis of the full application of the Hermite predictor–corrector method. Also, a new neighbour list is constructed using the positions at time t_1 . Thus, we calculate at certain times only the forces from neighbours (irregular time–step, t_{irr}), while at other times we calculate both the forces from neighbours and distant particles (regular time–step, t_{reg}).

For a neighbour list of size $N_{\rm nb} \ll N$, this procedure can lead to a significant gain in efficiency, provided the respective time scales for $\mathbf{a}_{i,\rm irr}$ and $\mathbf{a}_{i,\rm reg}$ are well separated.

¹Note, that the code also keeps the variables F and FDOT, which contain one half (!) of the *total* force, and one sixth (!) of the *total* time derivative of the force; this just a handy assignment for the frequent predictions of equation 3.

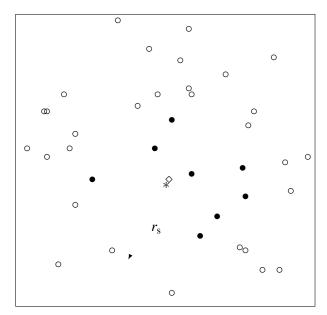


Figure 10.1: Illustration of the neighbour scheme for particle i marked as the asterisk (after [2]).

The actual size of neighbour spheres in NBODY6++ is controlled iteratively by a requirement in order to keep a certain optimal number of neighbours. This variable, NNBOPT, can be adjusted according to performance requirements. Its typical values are between 50 and 200 for a very wide range of total particle numbers N. Outside of the half-mass radius, the requirement of having NNBOPT neighbours is relaxed due to low local densities. Insisting on NNBOPT neighbours could result in undesired large amplitude fluctuations of the neighbour radii.

While [18] claim that the optimal neighbour number should grow as $N^{3/4}$ (which would be unsuitable for the performance on parallel computers), this is still an unsettled question. [2] advocates the coupling of the neighbour radius to the local density contrast, but NBODY6++ does *not* use that, since it makes average neighbour numbers much less predictable, which is bad for the performance and profiling issues on supercomputers, again.

Resuming, the method of the two particle groups is squeezed into the hierarchical time-step scheme making the overall view quite complex. Each particle is moved due to its time-step order and the time-steps, because the force calculation is divided: In eq. (11) a further subscript is needed which distinguishes the regular and irregular time step. The accuracy can be tuned by

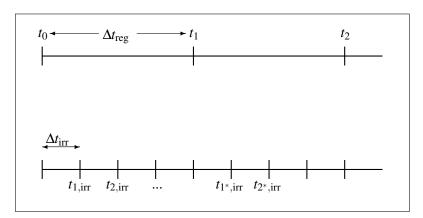


Figure 10.2: Regular and irregular time steps (after [17]).

 $\eta_{
m irr} pprox 0.01$ and $\eta_{
m reg} pprox 0.02$, again.

Both, the neighbour scheme and the hierarchical time-step scheme have in common that they are centered on one particle *i*, and they distinguish between nearby and remote stars, and they save computational time. One may ask: What is the intriguing difference between them? — The neighbour scheme is a *spatial* hierarchy, which avoids a frequent force calculation of the remote particles, because their totality provides a smooth potential which does not vary so much with respect to the particle *i*; that potential is rather superposed by some fluctuating peaks of close-by stars which will be "worked in" by the more often force determination. The time step scheme, in contrast, exhibits the *temporal* behaviour of the intervals for re-calculation of the full force in order to maintain the exactness of the trajectory; time steps chosen too small slow down the advancing calculation losing the computer's efficiency.

11 KS-Regularization

The fourth main feature of the codes since NBODY3 is a special treatment of close binaries. A close encounter is characterised by an impact parameter that is smaller than the parameter for a 90 degree deflection

$$p_{90} = 2G(m_1 + m_2)/v_{\infty}^2 \tag{13}$$

where G, m_1 , m_2 , v_∞ are the gravitational constant, the masses of the two particles and their relative velocity at infinity. In the cluster centre, it is very likely that two (or even more) stars come very close together in a hyperbolic encounter. As the relative distance of the two bodies becomes small $(R \to 0)$, their timesteps are reduced to prohibitively small values, and truncation errors grow due to the singularity in the gravitational potential, eqs. (1) and (2). In the NBODY code, the parameter RMIN is used to define a close encounter, and it is kept to the value of equation 13 (if KZ(16) > 0 is chosen in the control parameters). The corresponding time step DTMIN can be estimated from

$$dt_{\min} = \kappa \left[\frac{\eta}{0.03} \right] \left(\frac{r_{\min}^3}{\langle m \rangle} \right)^{1/2} \tag{14}$$

where κ is a free numerical factor, η the general time step factor, and $\langle m \rangle$ the average stellar mass [2]. If two particles are getting closer to each other than RMIN, and their time steps getting smaller than DTMIN, then they are candidates for "regularization".

Regularization is an elegant trick in order to deal with such particles which are as close as the diamond in the Figure 10.1. The idea is to take both stars out of the main integration cycle, replace them by their centre of mass (c.m.) and advance the usual integration with this composite particle instead of resolving the two components. The two members of the regularized pair (henceforth KS pair) will be relocated to the beginning of all vectors containing particle data, while at the end one additional c.m. particle is created (see below). One of the purposes of the code variable NAME(I) is to identify particles after such a reshuffling of data.

To be actually regularized, the two particles have to fulfil two more sufficient criteria: that they are approaching each other, and that their mutual force is dominant. In the equations in routine search.f, these sufficient criteria are defined as

$$\mathbf{R} \cdot \mathbf{V} > 0.1 \sqrt{(G(m_1 + m_2)R)}$$

 $\gamma := \frac{|\mathbf{a}_{pert}| \cdot R^2}{G(m_1 + m_2)} < 0.25$

Here, \mathbf{a}_{pert} is the vectorial differential force exerted by other perturbing particles onto the two candidates, R, \mathbf{R} , \mathbf{V} are scalar and vectorial distance and relative velocity vector between the two candidate, respectively. The factor 0.1 in the upper equation allows nearly circular orbits to be regularized; $\gamma < 0.25$ demands that the relative strength of the perturbing forces to the pairwise force is one quarter of the maximum. These conditions describe quantitatively that a two-body subsystem is dynamically separated from the rest of the system, but not unperturbed.

The internal motion of a KS pair will be determined by switching to a different (regularized) coordinate system. This transformation can be traced back to the square in quaternion space, where — by sacrificing some commutativity rules — it is guaranteed that the real-space motion does not leave the three-dimensional Cartesian space. It involves a set of four regular spatial coordinates and a fictitious time s(t), obtained in its simplest variant by the transformation dt = Rds. Any unperturbed two-body orbit in real space is mapped onto a harmonic oscillator in KS-space with double the frequency. Since the harmonic potential is regular, numerical integration with high accuracy can proceed with much better efficiency, and there is no danger of truncation errors for

arbitrarily small separations. The internal time–step of such a KS–regularized pair is independent of the eccentricity and, depending on the parameter ETAU, of the order of some 50–100 steps per orbit. The method of regularization goes back to [14] and makes an accurate calculation of a perturbed two–body motion possible. A modern theoretical approach to this subject can be found in [25]; the Hamiltonian formalism of the underlying transformations is nicely explained in [20].

While regularization can be used for any analytical two–body solution across a mathematical collision, it is practically applied to perturbed pairs only. Once the perturbation γ falls below a critical value (input parameter GMIN $\approx 10^{-6}$), a KS–pair is considered unperturbed, and the analytical solution for the Keplerian orbit is used instead of doing numerical integration. A little bit misleading is that such unperturbed KS–pairs are denoted in the code as "mergers", e.g. in the number or merges (NM) and the energy of the mergers (EMERGE). Merged pairs can be resolved at any time if the perturbation changes. The two–body KS regularization occurs in the code either for short-lived hyperbolic encounters or for persistent binaries.

In the code, the KS-pair appears as a new particle at the postion of the centre of mass. The variable NTOT, that contains the total number of particles *N* plus the c.m.'s, is increased by 1. When the pair is disrupted, NTOT is decreased again. The maximum number of possible KS-pairs is saved in the variable KMAX, which sets a threshold for the extension of the vector NTOT (see Chapter 5).

Close encounters between single particles and binary stars are also a central feature of cluster dynamics. Such temporary triple systems often reveal irregular motions, ranging from just a perturbed encounter to a very complex interaction, in which disruption of binaries, exchange of components and ejection of one star may occur. Although not analytically solvable, the general three–body problem has received much attention. So, the KS–regularization was expanded to the isolated 3– and 4–body problem, and later on to the perturbed 3–, 4–, and finally to the *N*–body problem. The routines are called

- triple.f (unperturbed 3-body subsystems, [8]),
- quad.f (unperturbed 4-body subsystems), and
- chain.f with different stages of implementation (slow-down, Stumpff functions, see for consecutive references Mikkola & Aarseth 1990, 1993, 1996, 1998, and [20]).

While occurrences of "triple" and "quad" will be rare in a simulation, the chain regularization is invoked if a KS-pair has a close encounter with another single star or another pair. Especially, if systems start with a large number of primordial binaries, such encounters may lead to stable (or quasi-stable) hierarchical triples, quadruples, and higher multiples. They have to be treated by using special stability criteria. Some of them are actually already implemented, but there is ongoing research and development in the field.

A typical way to treat all such special higher subsystems is to define their c.m. to be a pseudo-particle, i.e. a particle with a known sub-structure (very much like nodes in a TREE code). The members of the pseudo-particles will be deactivated by setting their mass to zero (ghost particles). At present there can only exist one chain at a time in the code, while merged KS binaries, and hierarchical subsystems can be more frequent. Details of these procedures are beyond the scope of this introductory manual.

Every subsystem — KS pair, chain or hierarchical subsystem — is perturbed. Perturbers are typically those objects that get closer to the object than $R_{\text{sep}} = R/\gamma_{\text{min}}^{1/3}$, where R is the typical size of the subsystem; for perturbers, the components of the subsystem are resolved in their own force computation as well (routines cmfreg.f, cmfirr.f).

38 12 Nbody-units

12 Nbody-units

The NBODY-code uses Dimensionless units, so-called "Nbody units". They are obtained when setting the gravitational constant G and the initial total cluster mass M equal to 1, and the initial total energy E to -1/4 (see [12], [7]).

Since the total energy E of the system is E = K + W with $K = \frac{1}{2}M\langle v^2 \rangle$ being the total kinetic energy and $W = -(3\pi/32)GM^2/R$ the potential energy of the Plummer sphere, we find from the virial theorem that

$$E = \frac{1}{2}W = -\frac{3\pi}{64} \frac{GM^2}{R}.$$
 (15)

R is a quantity which determines the length scale of a Plummer sphere. Using the specific definitions for G, M, and E above, this scaling radius becomes $R = 3\pi/16$ in dimensionless units. The half mass radius r_h can easily be evaluated by the formula (e.g. [26]):

$$M(r) = M \frac{r^3/R^3}{(1+r^2/R^2)^{3/2}}$$
 (16)

when setting $M(r_h) = \frac{1}{2}M$. It yields $r_h = (2^{2/3} - 1)^{-1/2}R = 1.30R$. The half-mass radius is located at R = 0.766, or about 3/4 "Nbody-radii".

The virial radius of a system is defined by $R_{\rm vir} = GM^2/2|W|$, while the r.m.s. velocity is $\langle v^2 \rangle^{1/2} = 2K/M$. In virial equilibrium |W| = 2K, so it follows for the crossing time

$$t_{\rm cr} := \frac{2R_{\rm vir}}{\langle v^2 \rangle^{1/2}} = \frac{GM^{5/2}}{(2|E|)^{3/2}}.$$
 (17)

The setting of G = M = 1 and E = -0.25 also determines the unit of time; so it follows that $t_{\rm cr} = 2\sqrt{2}$ in *N*-body units. By inversion we have

$$\tau_{\rm NB} = \frac{GM^{5/2}}{(4|E|)^{3/2}},\tag{18}$$

for the unit of time τ_{NB} . The virial radius of Plummer's model is $R_{vir} = 1$ in N-body units.

13 Output

Table 17: Definition of parameters

	Global properties				
TIME	time of simulation				
RSCALE	Half mass radius				
RTIDE	Tidal radius				
RC	Core radius				
NC	Number of stars inside core radius				
MC	Core mass				
VC	r.m.s velocity inside core radius				
CMAX	Maximum number density / half mass mean value				
RDENS(1:3)	Density center position				
RHOD	Density weighted average density $\Sigma RHO^2/\Sigma RHO$				
RHOM	Maximum mass density / half mass mean value				
$\langle M \rangle$	Average mass of star				
M1	Mass of most massive star				
ZMASS	Total mass of cluster				
MODEL	Snapshot counter in output				
NRUN	Run identification index				
TIDAL4	Twice angular velocity for linearised tidel force				
	Energy				
DE	relative energy error				
DELTA	absolute energy error				
BE(1)	Intial total energy				
BE(2)	Last adjust total energy				
BE(3)	Current total energy				
ZKIN	Kinetic energy				
POT	Potential energy				
ETIDE	Tidal energy				
ETOT	Total energy				
Е	Mechanical energy: ZKIN - POT + ETIDE (E(3))				
ESUB	Binding energy of unperturbed triples and quadruples				
EMERGE	Binding energy of mergers (E(9))				
EBIN	Binding energy of KS binaries				
ECOLL	The difference of binding energy of inner binary at the end and begin				
	of hierarchical systems (E(10))				
EMDOT	Mechanical energy of mass loss due to stellar evolution (E(12))				
ECDOT	Energy of velocity kick due to stellar evolution				
ECH	Binding energy of chain				
EBINP	Primordial KS binary energy (E(1))				
EBINN	Energy of new KS binary formed by dynamics (E(2))				
EESCS	Single escaper mechanical energy (E(4))				
EESCPB	Binding energy of primordial KS binary escapers (E(5))				

EESCPC	Mechanical energy of center mass of primordial KS binary escapers
	(E(6))
EESCNB	Binding energy of new formed KS binary escapers (E(7))
EESCNC	Mechanical energy of center mass of new KS binary escapers (E(8))
Scal	ing factors (Astronomical units = N -body units \times scaling factor)
RBAR	PC
TSCALE	Myr
TSTAR	Myr
VSTAR	km/s
RAU	AU
ZMBAR	Solar mass
SU	Solar radius
	Astronomical units
R*	Solar radius
L*	Solar luminosity
M*	Solar mass
	Status number
NTOT	Total number of particles (include all binary components, single stars
	and center of mass)
N	Total number of stars (binary counts as two stars)
NS	Single star number
NPAIRS	Number of KS regularization pairs
NMERGE	Number of mergers (stable triples)
MULT	Number of \geq 4 bodies merger
NZERO	Initial particle number (2* binaries + singles, initial N)
NB0	Primordial binary number
NUPKS	Unperturbed KS
NPKS	Total perturber number of KS
NTYPE(1:10	6) Number of stars with type 1 to 16
	For stars
I	Index of star (position in particle data array)
NAME	Identification of individual star, it's constant and unique value for each
	star (exclude un-physical particles like center mass and ghosts) during
	the whole simulation
K*	KSTAR type:
	-1: Pre-main sequence.
	0: Low main sequence $(M < 0.7)$.
	1: Main sequence.
	2: Hertzsprung gap (HG).
	3: Red giant (RG) branch.
	4: Core Helium burning.
	5: First Asymptotic giant branch (AGB).
	6: Second AGB.
	7: Helium main sequence.
	8: Helium HG.
	9: Helium giant branch.
	7. Honom giant oranon.

10: Helium white dwarf.				
11: Carbon-Oxygen white dwarf.				
12: Oxygen-Neon white dwarf.				
13: Neutron star.				
14: Black hole.				
15: Massless supernova remnant.				
Mass of star				
Three dimension position				
Three dimension velocity				
Current mass loss due to stellar evolution (<i>N</i> –body units)				
Accumulated mass loss due to stellar evolution				
Irregular time step of star				
Regular time step of star				
Kinetic energy				
Potential				
Neighbor number				
Neighbor radius				
Mass density of individual star calculated by nearest 5 neighbors, (only				
avaiable for particles inside core radius				
Stellar evolution of star				
star radius				
luminosity				
effective temperature				
angular velocity of star				
For binaries				
semi-major axis				
eccentricity				
Pericenter distance				
distance between two members of binary				
distance to density center				
velocity of center of mass				
Orbit period				
Index for binary component 1/2 (Not always equal name)				
Index for center mass particle (Not always equal name)				
perturber number				
Binary energy per unit mass				
Binary energy: M(I1)*M(I2)/M(ICM)*H				
Perturbation on KZ binary				
Pair index for binary				
KS time step of binary				
circularization timescale for current pericenter				
Primordial binary indicator1: Primordial bianry; 0: New binary				
Index of new star generated by binary collision or coalescence				
For hierarchical systems				
merger index				
Number count of current merger				

INPAIR I	Index of inner binary
	Inner binary center mass index
NAME(IM) N	Merger center mass name
I1/I2 I	Inner binary two component indexs
I3 (Outer particle index
ECC0 I	Inner binary orbit eccentricity
	Outer orbit eccentricity
EB0 I	Inner binary energy
EB1 C	Outer orbit energy
PO I	Inner binary orbit period
P1 (Outer orbit period
R12 I	Inner binary components seperation
	Separation between inner center mass and outer component
TG I	Inner orbit eccentricity growth timescale
ECCMIN N	Minimum eccentricity of inner binary orbit
ECCMAX N	Maximum eccentricity of inner binary orbit
PERIM s	smallest pericenter distance of outer particle orbit
PCR S	Stability triple system criterion for PERIM (assess.f), the real stability
C	criterion is more complicated and depend on the ECC1
	Inner binary orbit semi-major axis
SEMI1 C	Outer orbit semi-major axis
	Inclination angle in unit degree
FLAG-H I	Hierarchical system indicator1: merger, triple, chaotic binary, tidal
	circularization binary; 0: normal binary
	For quadruple systems
	Outer binary center of mass index
	Index of outer binary
	Outer binary two component index
	Outer binary eccentricity
	Outer binary energy
	Outer binary orbit semi-major axis
	Outer binary components seperation
DP34 I	Difference potential correction for the outer binary
	For chain
	Chain index
	Number of chain members
	Total energy of perturbed system (N-body interface)
	Perturbers of chain
	Total energy of chain
	Sum of all chain distances
	Gravitational radius ((sum M(I)*M(J))/ABS(ENERGY))
	Local crossing time ((sum M(I))**2.5/ABS(2*ENERGY)**1.5)
RMAXS N	Maximum size of unperturbed configuration
	Distance between member i and j
ICM1 I	First binary index after termination
	Second binary index after termination

	For kick				
M0	mass before kick				
MN	mass after kick				
VK	kick velocity after limit check				
VI	Initial velocity of kick star in cluster				
VF	Final velocity of kicked star in cluster				
VK0	Kick velocity generated from Henon's method (Douglas Heggie				
	22/5/97)				
FB	Fallback ratio, $VK = VK*(1-FB)$				
VESC	cluster escape velocity				
VDIS	escape velocity from binary system: SQRT(2*M(ICM)/R12)				

Table 18: Notice for Table 19

File format				
Header-*	The Header of file with line number *, the description of it is shown in			
	the right cell			
H-Label-*	Content labels are shown at the line number *, data begin from the next			
	line			
F-Label	Content labels are shown at the beginning of each line			
I-Label	Content labels are shown before each data			
N-Label	No labels in file			
	Frequency (freq.)			
Tevent	Output when event is triggered			
T_0	Output during initialization			
ΔT_{out}	Output time interval (input parameter DELTAT)			
ΔT_{adj}	Adjust time interval (input parameter DTADJ)			
ΔT_{HR}	Stellar evolution output time interval (input parameter DTPLOT)			
NFIX	Frequency of output (input parameter)			
	option			
#[<i>num</i>]	KZ option [num]			
	logical or			
&&	logical and			
CHAIN	Use chain: $#15 > 0&&(#30 > 0) #30 = -1)$			
USE_GPU	switch on GPU during compiling code			
USE_HDF5	switch on HDF5 during compiling code			

Table 19 show all output files of *NBODY6++*. The filename will be named as "[name].[unit]". The first column [name] with suffix "*" means this file will output as seperated snaphshots split by TIME[NB] (shown as suffix of file name).

Table 19: Output file information

_						
n	ame	unit	code	option	freq.	content

conf*	3	output.F	#3 > 0	ΔT_{out}	Basic data snapshots
				×	
				NFIX	
Heade	r-1	NTOT, MODEL, NRUN, NK			
Heade	er-2	TIME[NB], NPAIRS, RBAR, ZMBAR, RTIDE, TIDAL4, RDENS(1:3),		MBAR, RTIDE, TIDAL4, RDENS(1:3),	
		TIME/TCR, TSCALE, VSTAR, RC, NC, VC, RHOM, CMAX,			
		RSCALE, RSMIN, DMIN1			
N-Lal	bel	M, RHO, X	NS, X(1:3),	V(1:3), P	OT, NAME (All in NB unit)

Notice the file is unformatted (binary file). Each item output continually from 1 to NTOT. All items output in one line after two header lines.

NK: The number of parameters in Header-2, right now is always 20

TCR: Crossing time

RSMIN: Smallest neighbor radius obtained in last output (output.F) time

DMIN1: Smallest two body distance

XNS: The fifth nearest neighbor distance, (only avaiable for particles inside core radius

degen	4	degen.f	#9 ≥ 3	Tevent	Binary with degenerate stars	
Heade	r-1	RBAR, $\langle M$	RBAR, $\langle M \rangle$ [M*], M1[M*], TSCALE, NB0, NZERO			
H-Lab	el-2		ICASE, TIME[Myr], SEMI[AU], ECC, PERI/RS, P[days], RI[PC],			
		M(I1)[M*]	, M(I2)[M*],	K*(I1),K	X*(I2),K*(ICM), NAME(I1),NAME(I2)	

ICASE: 3: normal binary; 4: CE binary; 5: physical collision binary PERI/RS: Pericenter / maximum stellar radius of two members

lagr	7	lagr.f	$#7 \ge 3$	ΔT_{out}	Lagrangian radii, average mass, aver-
					age velocity, velocity dispersion out-
					put (calculation of Lagrangian radii
					use initial total mass of cluster)
Heade	er-1	Labels and column number for each output			
H-Lab	el-2	R_{lagr} , $R_{lagr,s}$, $R_{lagr,b}$, $\langle M \rangle$, N_{Shell} , $\langle V_x \rangle$, $\langle V_y \rangle$, $\langle V_z \rangle$, $\langle V \rangle$, $\langle V_r \rangle$, $\langle V_t \rangle$, σ^2 ,			
		$\sigma_r^2, \sigma_t^2, \langle V_{rot} \rangle$.) (All in NB	units)	

For each items above, there are 18 columns with different mass fraction(%): 0.1, 0.3, 0.5, 1, 3, 5, 10, 20, 30, 40, 50, 60, 70, 80, 90, 95, 99, 100 and inside core radius (exclude $R_{lagr.s}$, $R_{lagr.b}$)

 R_{lagr} : Lagrangian radius

 $R_{lagr,s}$: Single star Lagrangian radius

 $R_{lagr,b}$: KZ binaries Lagrangian radius

 $\langle M \rangle$: Average mass of a spherical shell defined by R_{lagr}

 $\langle V_{x/y/z} \rangle$: Mass weighted average velocity in x/y/z direction

 $\langle V_t \rangle$: Mass weighted average tangential velocity

 $\langle V_r \rangle$: Mass weighted average radial velocity

 σ^2 : Mass weighted velocity dispersion square

 σ_r^2 : Mass weighted radius velocity dispersion square

 σ_t^2 : Mass weighted tangential velocity dispersion square

 $\langle V_{rot.} \rangle$: Mass weighted average rotational velocity projected in x-y plane

bdat	8	ksin2.f	#8 > 0	T_{event}	New hierarchical (B-S)-S binary in-
					formation
		ksinit.F	#8 > 0		New binary information
		ksterm.F	#8 > 0		End binary information

I-Lat	oel	TIME[NB], NAME(I1) NANE(I2), FLAG-PB, M(I1)[NB], M(I2)[NB],				
		EB[NB], SEMI[NB], R12[NB], GAMMA[NB], RI[NB]				
bdat*	9	bindat.f $\#8 \ge 2$ ΔT_{out} KS binary output				
Heade	er-1	NPAIRS, MODEL, NRUN, N, NC, NMERGE, TIME[NB],				
		RSCALE[NB], RTIDE[NB], RC[NB], TIME[Myr], ETIDC[NB],				
		0				
Heade	er-2	EBINP, EBINN, E, EESCS, EESCPB, EESCPC, EESCNB, EESCNC,				
		EMERGE, ECOLL (All in NB unit)				
Heade	er-3	SBCOLL, BBCOLL, ZKIN, POT, EBIN0, EBIN, ESUB, EMERGE,				
		BE(3), ZMASS, ZMBIN, CHCOLL, ECOLL (All in NB unit)				
H-Lab	el-4	NAME(I1), NAME(I2), M1[M*], M2[M*], E[NB], ECC, P[days],				
		SEMI[AU], RI[PC], VI[km/s], K*(I1), K*(I2), ZN[NB], RP[NB],				
		STEP(I1)[NB], NAME(ICM), ECM[NB], K*(ICM)				

ETIDC[NB]: escape energy due to tidal force

SBCOLL: The difference of binding energy of inner binary at the end and begin of unperturbed triples

BBCOLL: The difference of binding energy of inner binary at the end and begin of unperturbed B-B quadruples

ZMBIN: Total KS binary masses

start.F

dat

CHCOLL: The difference of binding energy of inner binary at the end and begin of chain

Basic data after initialization

 T_0

#22 = 1

N-label		M[NB], X(1:3)[NB], V(1:3)[NB]				
esc	11	escape.f	#23 = 2,	ΔT_{adj}	escaping star output	
			4			
H-Lab	el-1	TIME[Myr], M[M*], E	ESC, VI[l	km/s], K*, NAME	
EESC: d	limensi	onless escape	e energy			
hiarch	12	hiarch.f	#18 = 1,	Tevent	New/End stable hierarchical system	
			3		(mergers) information	
Heade	er-1	RBAR, $\langle M \rangle$	$\langle [M^*], M1[N] \rangle$	Л*], TSC.	ALE, NB0, NZERO	
F-Lal	bel	TIME, SEMIO, SMEI1, ECC1, PERIO, PERIOM, P1/P0,				
		M(INCM)/M(I3), PCR/SEMIO, M(INCM)/ <m>, MR, INA, NAME(I1),</m>				
		NAME(I2), NAME(I3), K*(INCM), ECC0, ECCMIN, ECCMAX,				
K*(I1), K*(I2), RSM (All in NB unit)				unit)		
F-Label		TIME RIA	RC, SEMI	0, ECC	O, PERIO, POF/POI, RC/RSCALE,	
		GAMMA(I	NCM), NKI	, NKF, NI	PAIRS, NAME(I2) (All in NB unit)	

PERIO: Inner binary pericenter distance

PERIOM: Inner binary minimum pericenter distance

MR: Mass ratio of inner binary components (> 1)

PSM: Maximum stellar radius of two members of inner binary

P0F/P0I: Period of inner binary at the end of merger over at the beginning of merger

NKI: Orbit number the inner binary during the life of merger over the period of inner binary at the beginning of merger

NKF: Orbit number the inner binary during the life of merger over the period of inner binary at the end of merger

coll	13	mix.f	#19 ≥ 3	Tevent	Mixed star (physical collision of binary without evolved stars) information		
Heade	r-1	RBAR. ⟨ <i>M</i>	<u> </u> }[M*], M1[N		ALE, NB0, NZERO		
H-Lab		TIME[NB], NAME(I1), NAME(I2), K*(I1), K*(I2), K*(INEW), M(I1)[M*], M(I2)[M*], M(INEW)[M*], DM[M*], RS(I1)[R*], RS(I2)[R*], RI/RC, R12[R*], ECC, SEMI[R*], P[days]					
shrink	14	shrink.f		Tevent	Diagnostics for shrink regular time step for incoming high velocity star coming		
F-Lal	oel	I, RN, FI/F	J, DT, STEP	R (All in	NB unit)		
		nce from high	•		Γ		
		ninimum dis					
DT: eval	uated t		num approac	h truncate	ed to next time		
mix	15	mix.f	$#19 \ge 3$	Tevent	Mixed star information for the case NS/BH form		
F-Lal	201	V*(I1) V*	(12) V*(INE	(W) M(I	1)[M*], M(I2)[M*], M(INEW)[M*]		
				1			
hirect	16	hirect.f	#27 = 2	T _{event}	Diagnostics for rectification of hierarchical binary due to the internal en-		
			(hi-		ergy change of system		
			grow.f) #34 > 0		ergy change of system		
			(brake2.f)				
			#28 > 0				
			(brake3.f)				
F-Lal	bel	TIME[Nb].	, ,	ECC,R12/	SEMI,H,DB,DH/H		
H: inner	binary						
DM: cha	inge of	binding ener	gy				
		ksrect.f		Tevent	Diagnostics for rectification of KS orbit.		
F-Lal	nel	TIME(NB)	. IPAIR. R12	L Z/SEMI. F	I, GAMMA, DB, DH/H		
binev	17	binev.f	$#9 \ge 2$	T_{event}	Binary evolution stage, output when		
oniev	17	omev.i	11 / <u>/ </u>	1 event	binary change type		
H-Lab		TIME[Myr], NAME(I1), NAME(I2), K*(I1), K*(I2), K*(ICM), M(I1)[M*], M(I2)[M*], RS(I1)[R*], RS(I2)[R*], RI[PC], ECC, SEMI[R*], P[days], IQCOLL					
IQCOLI	ر: Type	of stage, nee	ed table in the	e future			
pbin	18	binout.f	#8 > 0	ΔT_{out}	Diagnostics for the primordial binary which exchanges members		
I-Lat	nel .	TIMEINBI	NAME(I1) NAM	E(I2), Flag-PB, Flag-H, M(I1)[NB],		
1 Lac	1			-	, ECC, GX, RI[NB], VR[NB]		
GX: max	kimum	perturbation			, = = =, ===, ====, , ===== ,		
VR: radi		_	· T				
bwdat*	19	bindat.f	#8 ≥ 2		Wide Non-KS bianry output		
Heade			TIME[Myr]	, N			
		_[2]	, -L-:-J*J				

H-Labe	H-Label-2 NAME(I1), NAME(I2), M1[M*], M2[M*], E[NB], ECC, P[da SEMI[AU], RI[PC], VI[km/s], K*(I1), K*(I2)								
symb	20	mdot.F	#19 ≥ 3	T_{event}	Symbiotic stars information				
F-Lab	F-Label NAME, K*, TIME[Myr], M[M*], SEMI[R*], DM, DMA??								
		star index							
		s loss from s		_	ion star				
DMA: A	ccrete	d mass from o	companion st	tar					
rocdeg	22	roche.f	#34 > 0	Tevent	Roche overflow binary involving degenerate objects				
F-Lab	bel	NAME(I1).	NAME(I2), K*(I1), K*(I2), M(I1)[M*], M(I2)[M*],				
		TIME[Myr], SEMI[R*]	, P[days],	MD(I1)[M*/Myr], MD(I2)[M*/Myr]				
MD: Ma	ss loss	rate							
ibeigen	23	binpop.F	(#8 = 1	T_0	Initial binary data by using eigen-				
			#8 ≥ 3)		evolution				
			&&						
			#42 = 6						
F-Lab	oel	ITER, I1, M	I(ICM)[M*]	, ECCI, E	CCC, SEMII, SEMIC, P[days]				
		_			satisfy the input criterions				
		centricity fro		istributio	n				
		rized eccentri	•						
		emi-major az	•	•	•				
SEMIC:	Circul	arized semi-r	najor axis ge	nerated b	y ECCC and period				
coal	24	coal.f	#19 ≥ 3	Tevent	Binary coalescence (Stellar type with				
					cores and circular orbit)				
Heade		, ,	,M1,TSCAL						
H-Labe	el-2				$E(I2), K^*(I1), K^*(I2), K^*1, IQ$				
		COLL, M(I1)[M*], M(I2)[M*], M(INEW)[M*], DM[M*], RS(I1)[R*],							
		RS(I2)[R*], RI/RC, R12[R*], ECC, SEMI[R*], P[days], RCOLL[R*],							
			P[NB], VIN						
		••	•		o binary exchanged to single star				
		r type of the		•					
1	•	y distance be							
VINF: V	elocity	at infinity fo		coalescei	VINF: Velocity at infinity for hyperbolic coalescence				
sediag	25	unpert.f	#27 > 0	T_{event}	Diagnostics for the stellar evolution				
		•			next look-up time of unpert KS				
sediag F-Lab		IPAIR, K*	(I1), K*(I2),	K*(ICM	next look-up time of unpert KS I), TEVNXT[NB], STEP(I1)[NB] (No				
F-Lab	pel	IPAIR, K*((I1), K*(I2), t when NWA	K*(ICM ARN≥ 10	next look-up time of unpert KS I), TEVNXT[NB], STEP(I1)[NB] (No				
F-Lab	oel T: Nex	IPAIR, K*0 more output time to chec	(I1), K*(I2), t when NWA ck stellar evo	K*(ICM ARN≥ 10	next look-up time of unpert KS I), TEVNXT[NB], STEP(I1)[NB] (No 00)				
F-Lab	pel	IPAIR, K*((I1), K*(I2), t when NWA	K*(ICM ARN≥ 10	next look-up time of unpert KS I), TEVNXT[NB], STEP(I1)[NB] (No 00) Binary coalescence (Stellar type with				
F-Lab TEVNX	Del T: Nex	IPAIR, K*(more output t time to chec	(I1), K*(I2), t when NWA ck stellar evo $\frac{1}{9}$	$K*(ICM)$ $ARN \ge 100$ lution T_{event}	next look-up time of unpert KS (), TEVNXT[NB], STEP(I1)[NB] (No ()00) Binary coalescence (Stellar type with no cores but circular orbit)				
F-Lab TEVNX coal Heade	T: Nex	IPAIR, K*(more output time to checombody.f	(I1), K*(I2), t when NWA ek stellar evo $\#19 \ge 3$,M1,TSCAL	$K*(ICM)$ $ARN \ge 100$ lution T_{event} $E,NB0,N$	next look-up time of unpert KS (i), TEVNXT[NB], STEP(I1)[NB] (No (iii) (iii) Binary coalescence (Stellar type with no cores but circular orbit) ZERO				
F-Lab TEVNX	T: Nex	IPAIR, K*(more output time to check cmbody.f RBAR,⟨M⟩ TIME[NB]	(I1), K*(I2), t when NWA ck stellar evo $\#19 \ge 3$,M1,TSCAL , NAME(I1	$K*(ICM \\ ARN \ge 100 \\ lution$ T_{event} $E,NB0,N$ NAN	next look-up time of unpert KS (i), TEVNXT[NB], STEP(I1)[NB] (No (00)) Binary coalescence (Stellar type with no cores but circular orbit) ZERO (E(I2), K*(I1), K*(I2), IQCOLL,				
F-Lab TEVNX coal Heade	T: Nex	IPAIR, K*(more output time to chec cmbody.f RBAR,⟨M⟩ TIME[NB] M(I1)[M*]	(I1), K*(I2), t when NWA ck stellar evo $\#19 \ge 3$,M1,TSCAL , NAME(I1	$K*(ICM)$ $ARN \ge 100$ lution T_{event} $E,NB0,N$), NAN], $DM[M]$	next look-up time of unpert KS (), TEVNXT[NB], STEP(I1)[NB] (No ()00) Binary coalescence (Stellar type with no cores but circular orbit) ZERO (ME(I2), K*(I1), K*(I2), IQCOLL, [I*], RS(I1)[R*], RS(I2)[R*], RI/RC,				

48 13 Output

highv	29	hivel.f		Tevent	Diagnostics for high-velocity particle		
					added or removed from LISTV		
F-Label (REMOVE) TIME[NB], I, NAME, RI(NB), VI(NB)							
F-Lal	F-Label (ADD NS, terminated KS/chain) TIME[NB], NHI, I, NAME, K						
VI[NB], RI[NB], STEP[NB]							
F-Label (ADD fast single) TIME[NB], NHI, NAME, IPHASE, VI[NB], RI[HI, NAME, IPHASE, VI[NB], RI[NB],		
		STEP[NB]					
F-Lal	bel	(ADD hyp	erbolic two	-body m	otion) TIME[NB], NHI, NAME(I1),		
	NAME(I2), IPHASE, RIJ[NB]						
NHI: hig	NHI: high-velocity particle number						
IPHASE	· Inter	nal status of c	ode (check r	nhody6 F	for details)		

IPHASE: Internal status of code (check nbody6.F for details)

global	30	output.F		ΔT_{out}	Global feat	ures of cluste	r and event
					counters		
H-Lab	el-1	TIME[NB]	, TIME[My	r], TCR	[Myr], DE	BE(3), RS	CALE[PC],
], RHOI	
		RHOM[M*	$^{\circ}/PC^{3}$], MC[M*], CM	AX, $\langle Cn \rangle$, Ir/	R, RCM[NB],	VCM[NB],
	AZ, EB/E, EM/E, VRMS[km/s], N, NS, NPAIRS, NUPKS, NP					KS, NPKS,	
		NMERGE,	MULT, $\langle N \rangle$	$B\rangle$, NC,	NESC, NST	EPI, NSTEPE	B, NSTEPR,
		NSTEPU,	NSTEPT,	NSTEPQ	, NSTEPC,	NBLOCK,	NBLCKR,
		NNPRED,	NIRRF, 1	NBCORR	, NBFLUX	, NBFULL,	NBVOID,
		NICONV, I	NLSMIN, N	BSMIN,	NBDIS, NBI	DIS2, NCMDI	ER, NFAST,
		NBFAST,	NKSTRY,	NKSREC	G, NKSHYP	, NKSPER,	NKSMOD,
		NTTRY, N	ΓRIP, NQU <i>P</i>	D, NCH	AIN, NMERO	G, NEWHI	

TCR: Crossing time

RDENS: density center to coordinate center distance

 $\langle Cn \rangle$: frequency 1/STEP weighted averaged neighbor number Ir/R: Irregular cost ($\sum NB/STEP$) over regular cost ($N/\sum STEPR$)

RCM: Center mass distance to coordinate center

VCM: Center mass velocity

AZ: Angular momentum in z axis including tidal effect (Chandrasekhar equation 5.530)

VRMS: root mean square velocity of cluster

NESC: Escapers

NSTEPI: Irregular integration steps

NSTEPB: Irregular integration steps of binary center mass particles

NSTEPR: Regular integration steps

NSTEPU: Regularized integration steps

NSTEPT: Triple regularization integration steps (#15 > 0)

NSTEPQ: Quadruple regularization integration steps (#15 > 0)

NSTEPC: Chain regularization steps (# DIFSY calls)

NBLOCK: Number of irregular blocks (block-step version)

NBLCKR: Number of regular blocks (block-step version)

NNPRED: Coordinate & velocity predictions of all particles

NIRRF: Calculated irregular force

NBCORR: Force polynomial corrections

NBFLUX: Number of changes in neighbor lists (NBLOSS+NBGAIN)

NBFULL: Neighbor number overflows with standard criterion

NBVOID: No neighbours inside 1.26 times the basic sphere radius

NICONV: Irregular step reduction (force convergence test)

NLSMIN: Small step neighbours selected from other neighbour lists

NBSMIN: Retained neighbours inside 2*RS (STEP < SMIN)

NBDIS : Second component of recent KS pair added as neighbour (#18) NBDIS2: Second component of old KS pair added as neighbour (#18 > 1)

NCMDER: C.m. values for force derivatives of KS component

NFAST : Fast particles included in LISTV (#18 > 0)

NBFAST: Fast particles included in neighbour list (#18 > 0)

NKSTRY: Two-body regularization attempts

NKSREG: Total KS regularizations

NKSHYP: Hyperbolic KS regularizations NKSPER: Unperturbed KS binary orbits NKSMOD: Slow KS motion restarts (#26 > 0)

NTTRY: Search for triple, quad & chain regularization or mergers

NTRIP : New three-body regularizations (#15 > 0) NQUAD : New four-body regularizations (#15 > 0)

NCHAIN: New chain regularizations (#15 > 0 & #30 > 0)

NMERG : New mergers of stable triples or quadruples (#15 > 0) NEWHI : New hierarchical systems (counted by routine HIARCH)

lagr1	31	lagr2.f	#7 ≥ 5	ΔT_{out}	Two mass group systems Lagrangian radii (first group)		
N-Lal	bel	TIME[NB]	, $R_{lagr}[NB]$ (mass frac	etion: 0.01, 0.02, 0.05, 0.1, 0.2, 0.3, 0.4,		
		0.5, 0.625,	0.75, 0.9) (H	ere calcul	ation of R_{lagr} use the current total mass)		
lagr2	32	lagr2.f	#7 ≥ 5	ΔT_{out}	Two mass group systems Lagrangian		
					radii (second group)		
N-Lal	bel	see Unit 31					
ns	33	degen.f	#9 ≥ 3	T _{event}	Neutron stars (never used)		
F-Lat	oel	I, NAME, IFIRST, K*, TIME[Myr], VI[km/s]					
bh	34	degen.f	#9 ≥ 3	T _{event}	Black holes (never used)		
F-Lat	oel	I, NAME, IFIRST, K*, TIME[Myr], VI[km/s]					
event	35	events.f	#19 >	ΔT_{out}	Stellar evolution and tidal capture		
			0 #27 >		event counter and energy		
			0				
H-Lab	le-1	TIME[Myr	TIME[Myr], NDISS, NTIDE, NSYNC, NCOLL, NCOAL, NDD,				
		NCIRC, NROCHE, NRO, NCE, NHYP, NHYPC, NKICK, EBIN,					
		EMERGE,	EMERGE, ECOLL, EMDOT, ECDOT, EKICK, ESESC, EBESC,				
		EMESC, I	DEGRAV, E	BIND, M	MMAX, NMDOT, NRG, NHE, NRS,		
		NNH, NW	D, NSN, N	NBH, NE	SS, ZMRG, ZMHE, ZMRS, ZMNH,		

NDISS: Tidal dissipations at pericentre (#27 > 0)

NTIDE: Tidal captures from hyperbolic motion (#27 > 0) NSYNC: Number of synchronous binaries (#27 > 0)

ZMWD, ZMSN, ZMDOT, NTYPE(1:16)

NCOLL: Stellar collisions NCOAL: Stellar coalescence NDD: Double WD/NS/BH binaries NCIRC: Circularized bianries (#27 > 0) NROCHE: Roche stage triggered times

NRO: Roche binary events

NCE: Common envelope binaries NHYP: Hyperbolic collision

NHYPC: Hyperbolic common envelope binaries

NKICK: WD/NS/BH kick

NSESC: Escaped single particles (#23 > 0)

NBESC: Escaped binaries (#23 > 0)

NMESC: Escaped mergers (#15 > 0& > 0)

EKICK: KICK energy of WD/NS/BH

ESESC: Single star escaper energy

EBESC: KS Binary star escaper energy

EMESC: Merger escaper energy

DEGRAV: Change of binary energy compared to initial value

EBIND: Binding energy of cluster (E)

MMAX: Maximum stellar mass NMDOT: Stellar mass loss event

NRG: New red giants NHE: New helium stars NRS: New red supergiants NNH: New naked Helium stars

NNH: New naked Helium NWD: New white dwarfs NSN: New neutron stars NBH: New black holes NBS: New blue stragglers

ZMRG: New red giants mass ZMHE: New helium stars mass ZMRS: New red supergiants mass ZMNH: New naked Helium stars mass

ZMWD: New white dwarfs mass

ZMSN: New neutron stars mass

status	36	global	#46 > 0	#47	Global parameters which combine
		_output.F			global output and generalized La-
					grangian radii
N-lal	bel	TIME[NB]	, TIME[My	r], Tcr[N	Iyr], Trh[Myr], TM[M*], TSM[M*],
		TBM[M*],	Q, Rh[pc], Rtid[J	pc], Rden(1:3)[pc], RHOD[M^*/pc^3],
		RHOM[M ²	*/pc ³], Mr	nax[M*],	Etot[NB], Ekin[NB], Epot[NB],
		Ebin[NB],	Etid[NB],	Em[NB]	, Ecol[NB], Ece[NB], Ekick[NB],
		Eesc[NB],	Ebesc[NB],	Emesc[N	NB], N, NS, NB, NM, NP, Lagr*(R,
		N, M, V, V	/x, Vy, Vz,	Vr, Vt, V	7rot, S, Sx, Sy, Sz, Sr, St, Srot, E) &
		(all,single,l	oinary), Mbl	lagr*, Nb	lagr*, Mpblagr*, Npblagr*, Eblagr*,
		Eblagrb*, I	Epblagr*, Ep	blagrb*, A	Alagr*
		IF #19 >=	= 3: MMD	OT, MRC	G, MHE, MRS, MNH, MWD, MSN,
		MKW*(sin	gle, binary w	vith one co	omponent, binary with two components)
		& (K* fro	m -1 to 15)	, NDISS,	NTIDE, NSYNC, NCOLL, NCOAL,
		NCIRC, N	ROCHE, NI	RO, NCE	, NHYP, NHYPC, NKICK, NMDOT,
		NRG, NHE	E, NRS, NNF	I, NWD,	NSN, NBH, NBS, NKW*, LagrK*
	٦ ,				

TIME: Current time

Tcr: Half-mass radius crossing time

Trh: Half-mass radius relaxation time

TM: Total mass

TSM: Total single mass

TBM: Total binary mass including mergers

Q: Virial ratio

Rh: Half-mass radius Rtid: Tidal radius

Rden: 3-dimensional density center position in current coordinate

RHOD: Density weighted average density $\sum \rho^2 / \sum \rho$. RHO: Mass density of individual star calculated by nearest 5 neighbors (only avaiable for particles inside core radius)

RHOM: Maximum mass density / half mass mean value

Mmax: Maximum particle mass

Etot: Cluster total energy without escaping energy

Ekin: Cluster kinetic energy Epot: Cluster potential energy Ebin: Cluster Binary binding energy

Etid: Tidal energy

Em: Cluster mass loss energy

Ecol: Collision energy

Ece: Common envelope energy

Ekick: Neutron star/black hole initial kick energy

Eesc: escapers kinetic and potential energy (binaries/mergers using center-of-mass

Ebesc: binary escapers binding energy Emesc: merger escapers binding energy

N: Total number of particles (binary/merger resolved)

NS: Total number of single particles

NB: Total number of binary/merger particles (unresolved)

NM: Total number of merger particles (unresolved)

NP: Total number of particles (NS+NB; unresolved)

Lagr*: Lagrangian radii and all related parameters. All, single, binary particles are calculated separately. If stellar evolution is switched on, all main type of stars are also calculated individually. Similar as lagr.7, whether the reference total mass is initial cluster mass or current mass is controlled by #7 = 2 or 4. But this only work for all/single/binary, not for different stellar types. The order of Lagr* is organized as hierarchical groups (from high level to low level; lowest level are neighbor data groups in output): [all, single, ***] [R, N, ***] [0.001, 0.01, ***]

The mass fraction list is 0.001, 0.01, 0.1, 0.3, 0.5, 0.7, 0.9, 1.0, Rc. (Rc is inside core radius)

Lagrangian radii related parameters in order: (whether paremeters are calculated in shells or from the center also used the same option as lagr.7 (#7))

R: Lagrangian radii

N: Number of particles (Number counts for Total Lagrangian resolved all binaries and mergers to get correct average mass; Number counts for binaris use center-of-mass)

M: Averaged mass

V: Averaged velocity value

Vx: Averaged x component of velocity

Vy: Averaged y component of velocity

Vz: Averaged z component of velocity

Vr: Averaged radial velocity

Vt: Averaged tangential velocity

Vrot: Averaged rotational velocity along z axis

S*: mass weighted velocity dispersion with similar definitions as V*

E: kinetic and potential energy

Mblagr*: Binary mass within total (global) Lagrangian radii, the ratio follows Lagr* without Rc (thus one data column less)

Nblagr*: Binary number (resolved; including mergers counted as 3) within total (global) Lagrangian radii, similar as Mblagr*

Mpblagr*: Primordial binary (detected by NAME(I1)-NAME(I2)==1) mass within total (global) Lagrangian radii, the ratio follows Lagr* with Rc

Npblagr*: Primordial binary number (resolved within total (global) Lagrangian radii, similar as Mpblagr*

Eblagr*: Binary binding energy within total (global) Lagrangian radii, the ratio follows Lagr*

Eblagrb*: Binary binding energy within binary Lagrangian radii, the ratio follows Lagr* without Rc

Epblagr*: Primordial binary binding energy within total (global) Lagrangian radii, the ratio follows Lagr*

Epblagrb*: Primordial binary binding energy within binary Lagrangian radii, the ratio follows Lagr* without Rc

Alagr*: 3-dimensional angular momentum (binary unresolved) within total (global) Lagrangian radii, the ratio follows Lagr*. Notice different Lagrangian ratios are minimum groups of data (neighbor data) in output, the three component is in high level

MMDOT: Stellar mass loss

MRG: Cumulative new red giants mass

MHE: Cumulative new helium stars mass

MRS: Cumulative new red supergiants mass

MNH: Cumulative new naked Helium stars mass

MWD: Cumulative new white dwarfs mass

MSN: Cumulative new neutron stars mass

MKW*: Total mass of different stellar types (all types from -1 to 15). For each type, a three masses (single particle, binary with one certain stellar type component, binary with both components are the certain stellar types) are grouped together in the output data

NDISS: Cumulative event number of tidal dissipations at pericentre (#27 > 0)

NTIDE: Cumulative event number of tidal captures from hyperbolic motion (#27 > 0)

NSYNC: Cumulative event number of synchronous binaries (#27 > 0)

NCOLL: Cumulative event number of stellar collisions

NCOAL: Cumulative event number of stellar coalescence

NDD: Cumulative event number of ouble WD/NS/BH binaries

NCIRC: Cumulative event number of circularized bianries (#27 > 0)

NROCHE: Cumulative number of Roche stage triggered times

NRO: Cumulative event number of Roche binaries

NCE: Cumulative event number of common envelope binaries

NHYP: Cumulative event number of hyperbolic collision

NHYPC: Cumulative event number of hyperbolic common envelope binaries

NKICK: Cumulative event number of WD/NS/BH kick

NMDOT: Cumulative number of stellar mass loss event

NRG: Cumulative number of new red giants

NHE: Cumulative number of new helium stars

NRS: Cumulative number of new red supergiants

NNH: Cumulative number of new naked Helium stars

NWD: Cumulative number of new white dwarfs

NSN: Cumulative number of new neutron stars

NBH: Cumulative number of new black holes

NBS: Cumulative number of new blue stragglers

NKW*: Total number of different stellar types, similar as MKW*

LagrKW*: Lagrangian radii and all related parameters for main stellar types

The main stellar types in order:

- 1. Low mass main sequence (M < 0.7) (0)
- 2. High mass main sequence (1)
- 3. Hertzsprung gap (HG). (2)
- 4. Red giant. (3)
- 5. Core Helium burning. (HB) (4)
- 6. AGB (5-6)
- 7. Helium types (7-9)
- 8. White dwarf (10-12)
- 9. Neutron star (13)
- 10.Black hole (14)
- 11.Pre main sequence (-1)

sediag	38	expel2.f	#19 ≥ 3	Tevent	Diagnostics for common envelop type	
			&&		change	
			Chain		-	
N-Lal	bel					
		mix.f	#19 ≥ 3	Tevent	Diagnostics for mixed stars	
N-Lal	bel					
		trflow.f	$#19 \ge 3$	Tevent	Diagnostics for iteration convergency	
					check until Roche-lobe overflow	
N-Lal	bel					
stellar ev	olution	health chec	k			
hbin	39	adjust.F	#9 = 1,3	ΔT_{out}	The hardest binary below ECLOSE	
F-Lat	oel	TIME[NB], NAME(I1), NAME(I2), K*, NP, ECC, SEMI[NB], P[days],				
		EB[NB], E	M[NB]			
T		L 3/	L . J			
data*	40	custom_	#46 > 0	#47	HDF5/BINARY or ANSI output of	
data*	40			#47	HDF5/BINARY or ANSI output of basic data. Notice all units here are	
data*	40	custom_		#47	_	
data*	40	custom_		#47	basic data. Notice all units here are	
data*	40	custom_		#47	basic data. Notice all units here are shown as astrophysical units, but it	
data*	40	custom_		#47	basic data. Notice all units here are shown as astrophysical units, but it can be <i>N</i> -body units or input data units	

If HDF5 is complied and #46 = 1 or 3, file names are snap.40_[time].h5part. Snapshot output time interval is controlled by DELTAT (see input parameters chapter 4). If HDF5 is switched off or ANSI output is used, the file names are single/bianry/merger.40_[time], each snapshot file only contain data of one certain time. The time interval is controlled by #47 (see input parameters chapter 4 for details)

	,
H-Label-1	Single particle: NAME, M[M*], X(1:3)[pc], V(1:3)[km/s], POT[NB],
	RS[R*], L[L*], Teff[K], MCORE[M*], RSCORE[R*], K*
	Binary particle: NAME(I1), NAME(I2), NAME(ICM), M(I1)[M*],
	$M(I2)[M^*]$, $XCM(1:3)[pc]$, $VCM(1:3)[km/s]$, $XREL(1:3)[AU]$,
	VREL(1:3)[km/s], POT[NB], SEMI[AU], ECC, P[days], GAMMA,
	$RS(I1)[R^*]$, $RS(I2)[R^*]$, $L(I1)[L^*]$, $L(I2)[L^*]$, $Teff(I1)[K]$,
	Teff(I2)[K], MCORE(I1)[M*], MCORE(I2)[M*], RSCORE(I1)[R*],
	RSCORE(I2)[R*], K*(I1), K*(I2), K*(ICM)
	Merger particle: NAME(I1), NAME(I2), NAME(I3), NAME(ICM),
	M(I1)[M*], M(I2)[M*], M(I3)[M*], XCM(1:3)[pc], VCM(1:3)[km/s],
	XREL0(1:3)[AU], VREL0(1:3)[km/s], XREL1(1:3)[AU],
	VREL1(1:3)[km/s], POT[NB], SEMI0[AU], ECC0, P0[days],
	SEMI1[AU], ECC1, P1[days], RS(I1)[R*], RS(I2)[R*], RS(I3)[R*],
	$L(I1)[L^*]$, $L(I2)[L^*]$, $L(I3)[L^*]$, $Teff(I1)[K]$, $Teff(I2)[K]$,
	Teff(I3)[K], MCORE(I1)[M*], MCORE(I2)[M*], MCORE(I3)[M*],
	RSCORE(I1)[R*], RSCORE(I2)[R*], RSCORE(I3)[R*], K*(I1),
	K*(I2), K*(I3), K*(ICM)
VCM: position	of center of mass

XCM: position of center-of-mass

VCM: velocity of center-of-mass

XREL: relative position of two components in a binary (from 1 to 2: X(I1)-X(I2))

VREL: relative velocity of two components in a binary (from 1 to 2: V(I1)-V(I2))

XREL: relative position of two components in a binary (from 1 to 2: X(I1)-X(I2))

VREL: relative velocity of two components in a binary (from 1 to 2: V(I1)-V(I2))

MCORE: Stellar core mass RSCORE: Stellar core radius

0/1 in mergers: 0 means inner binary, 1 means outer binary (center-of-mass with outer

particles

I1/I2/I3 in mergers: I1/I2 are inner binary components, I3 means outer particle

nbflow	41	fpoly0.F	USE_GPU	T_0	Diagnostics for neighbor list overflow	
					from GPU regular force initialization	
F-Lat	oel	NSTEPR, N	NAME, NB,	RNB[NB], RI[NB]	
		util_gpu.F	USE_GPU	T_{event}	Diagnostics for neighbor list overflow	
					from GPU regular force calculation	
I-Lab	el	I,NAME, N	INPRE, NBN	IEW, RN	B[NB], RI[NB], TIME[NB]	
NBPRE:	NBPRE: previous neighbor number					
NBNEW	NBNEW: new neighbor number that cause overflow					
ibcoll	42	binpop.F	#8 = 1,≥	T_0	Diagnostics for the binary physical	
			3		collision cases when initializing pri-	
					mordial binaries	
I-Lab	el	I1, M(I1)[N	M*], M(I2)[]	M*], EC	C, SEMI[AU], PERI[R*], RS(I1)[R*],	
		RS(I2)[R*]				
sediag	43	mdot.F	#19 ≥ 3	Tevent	Diagnostics of warning when stellar	
					radius expand more than 1.5x	

F-Lat	pel	I, NAME, RS0[R*], F], DT[M	yr], K*0, K*N, M0[M*], MN[M*],
K*0. Pre	evious	stellar type	torille j		
K*N: Ne					
M0: Init		• •			
		ew stellar ma	SS		
		stellar radius			
		ar radius			
hinc	44	induce.f	#27 > 0	T _{event}	Information of high inclinations and
	' '	maaccii		- eveni	$TC2 < 10^7$ yrs of hierarchical binary
F-Lat	nel	ECC0, EC	CMIN. ECC	MAX. K	*(I1), K*(I2), K*(ICM), SEMI0[NB],
			•	-	Myr], TCM[Myr], TIME[Myr]
IN: Indic	cator of			•	(2.5), where AIN is inclination angle
			cale for smal	_	
mbh	45	bhplot.f	#24 = 1	ΔT_{BH}	Mass black hole information
H-Lab		•			X(1:3)[PC], V(1:3)[km/s], NB,
II Luo	OI I		- • -		m/s], DEN[M*/PC ³], RIJMAX[PC],
			1:3)[km^2/s^2]	L (1.5)[III	
Notice th	ne XAV			A is not a	accurate due to the neighbor criterion
					ary system or single
		nassive black		C 15 111 511.	ary system of single
				ole neighb	oors (relative to black hole velocity)
	-			•	hbors (relative to black hole velocity)
	_	•		•	neighbors within RNB (exclude black
hole mas		•		•	· ·
RIJMAX	K: Max	imum distand	ce from neigl	hbor star t	to black hole
VSIGMA	A: 3-di	mensional v	elocity dispe	rsion of b	black hole neighbors (relative to black
hole velo	ocity)				
mbhnb	46	bhplot.f	#24 = 1	ΔT_{BH}	Mass black hole neighbor information
Heade	er-1	TIME[Myr], NB		-
N-Lal	bel	NAME, M	[M*], XREL	(1:3)[AU]], RIJ[AU], VREL(1:3)[km/s], K*
XREL: I	Position	vector relat	ive to black l	nole	
RIJ: Dist	tance to	black hole			
VREL: V	Velocity	vector relat	ive to black l	hole veloc	eity
itid3	52	xtrnl0.F	#14 = 3	<i>T</i> ₀	Initialization of circular velocity in the plane for galaxy tidal force
F-Lat	nel	VC[km/s],	RIIKPC1		plane for garany traditionee
VC: Circ			Iu[III 0]		
hypcep	54	ksint.f	#19 \ge 3	Tevent	Close encounter for hyperbolic mo-
пурсер]]]	KSIIIt.I	1117 2 3	1 event	tion (pericenter $< 5.0 \times$ Maximun
					stellar radius of two stars
F-Lat	nel	TIMFINR	NAMF(I1) NAM	E(I2), K*(I1), K*(I2), VINF[km/s],
· Lat			, TV RVIL(11 , RX[R*], PI		
VINF: V	elocity		r hyperbolic		nce
	-	•	• •		(binary will form)
10111.	captare	Gibtuilee of	ing personic ci		(chiai) will folill)

RX: Max	ximum	stellar radius	s of two stars				
hypcec	55	ksint.f	#19 ≥ 3	Tevent	Close encounter for hyperbolic motion (physical collision case)		
F-Label		TIME[NB], IPAIR, NAME(I1), NAME(I2), K*(I1), K*(I2), K*(ICM), VINF, ECC, H[NB], R12[NB], SEMI[NB], PERI[NB], M(I1)[NB], M(I2)[NB], M(ICM)[M*], RI[RC], VI[km/s], RHOD, RS(I1)[R*], RS(I2)[R*], RCAP[R*], RX/PERI, RCOLL/PERI					
RCOLL:	If #27				, otherelse normal collision criterion		
			s of two stars		,		
*fort	60	ellan.f	#49 > 0	ΔT_{out}	Moments of Inertia ??		
N-Lal	bel	??					
cirdiag	71	spiral.f	#27 > 0	T _{event}	Diagnostics for skip removal of chaos binary if this is member of sin- gle/double (stable quadruple) merger.		
F-Lat	bel	NCHAOS,	NAMEC, NA	AME(IM)	NAME(I3)		
NCHAO	S: Nun	nber of chaos		` /			
NAMEC	: Name	e for chaos b	inaries				
histab	73	impact.f	#15 > 0	Tevent	Diagnostics for checking Zare exchange stability criterion (exchange of outer particle with inner member of binary), but the slingshot still can happen, thus not triple system stablility criterion.		
F-Lat	201	TIMEINDI	M(I2)/M(IN	JCM) E	CC0, ECC1, SEMI0[NB], PERIM[NB],		
r-Lat	961		, M(13)/M(11 TG[Myr], SI	-			
SP: >=	1, no e		, will be exc		51, 11		
cirdiag	75	decide.f	#27 = 2	Tevent	Diagnostics output for large eccentricity (> 0.9) during merger decision (deny stable triple forming if circularization timescale is short).		
F-Lat	oel	NAME, T	IME[NB], E	CCO, E	CC1, EMIN, EMAX, ECCD[1/Myr],		
		EDT[NB],	TG[Myr], To	C[Myr], E	EDAV[1/Myr], PERIM[RSM]		
EDAV: A	Average aximu	icity change eccentricity n stellar radi	rate	rs			
kscrit	77	chmod.f	#16 > 2	Tevent	Diagnostics for increasing or decreas-		
			&& CHAIN		ing regularization parameters in chain		
F-Lat	oel	TIME[NB]		GPERT, R	MIN, RIJ[NB]		
GPERT: RMIN: I	Dimen Distanc	sionless pert e criterion fo	_	hain ion (also i	s input parameter)		
	RIJ: Distance between chain center mass and perturber						

chstab	81	chstab.f	CHAIN	Tevent	New hierarchical system with stability condition for bound close pair (RB > semi) (formed from 4th body escape or perturber make it stable).				
I-Lab	I-Label		TIMEC, RI[NB], NAME(I3), M(I3)/M(INCM), ECC0, ECCMAX,						
		ECC1, SEMI0[NB], SEMI1[NB], PCR/PERIM, INA[deg]							
		cstab2.f	CHAIN	Tevent	Hierarchical stability condition (SEMI1 $> 0 \Rightarrow$ ECC1 < 1).				
N-Label		TIMEC[NB], RI[NB], NAME(I3), ECC0, ECC1, SEMI0[NB], SEMI1[NB], PCR/PERIM, INA[deg]							
		cstab3.f	CHAIN	Tevent	Continued chain integration if outer orbit unstable or large pert.				
N-Lal	bel	TIMEC[NB], RI[NB], NAME(I3), ECC0, ECC1, SEMI0[NB], SEMI1[NB], PCR/PERIM, INA[deg]							
		cstab4.f	CHAIN	Tevent	Check hierarchical stability condition for bound close pair.				
N-Lal	bel	_	TIMEC[NB], RI[NB], NAME(I3), ECC0, ECC1, SEMI0[NB], SEMI1[NB], PCR/PERIM, INA[deg]						
TIMEC:	time v	hen chain fo	ormed						
bev*	82	hrplot.F	#12 > 0	ΔT_{HR}	KS binary stellar evolution data				
Heade	r-1	NPAIRS, T	IME[Myr]						
N-Label		TIME[NB], I1, I2, NAME(I1), NAME(I2), K*(I1), K*(I2), K*(ICM), RI[RC], ECC, log10(P[days]), log10(SEMI[R*]), M(I1)[M*], M(I2)[M*], log10(L(I1)[L*]), Log10(L(I2)[L*]), Log10(RS(I1)[R*]), Log10(RS(I2)[R*]), Log10(Teff(I1)[K]), Log10(Teff(I2)[K])							
sev*	83	hrplot.F	#12>0		Single star stellar evolution data				
Heade	r-1	NS, TIME	[Myr]	I					
N-Lal	bel	TIME[NB], I, NAME, K*, RI[RC], M[M*], log10(L[L*]), log10(RS[R*]), log10(Teff[K])							
merger	84	bindat.f	#8 ≥ 2	ΔT_{out}	Extra mergers information (main merger output is in hidat.87_*)				
F-Label		TIME[NB], NAME[I1], NAME[I3], K*[I1], K*[I3], K*[IM], ECC0, ECC1, PERI(I3)/PCR, PERI(INCM)[RSM], P0[days], P1[days], SEMI1[NB]							
roche	85	roche.f	#34 > 0	Tevent	Roche overflow stage data				
H-Label		NAME(I1), NAME(I2), K*(I1), K*(I2), TIME[Myr], AGE(I1), AGE(I2), M0(I1), M0(I2), M(I1), M(I2), Z, ECC, P[days], JSPIN(I1), JSPIN(I2), STAT							
AGE: Stellar age JSPIN: Angular momentum of star STAT: Type of binary Z: Metallicity *M0: Stellar mass before mass transfer?									
hidat*	87	hidat.f	#18 > 3	ΔT_{out}	Hierarchical data of mergers (stable triples, quadruples)				

Heade	r-1	NPAIRS, NRUN, N, NC, NMERGE, MULT, NEWHI, TIME[NB]				
H-Lab	oel	NAME(I1), NAME(I2), NAME(I3), K*(I1), K*(I2), K*(I3), M(I1)[M*],				
		M(I2)[M*], M(I3)[M*], RI[NB], ECCMAX, ECC0, ECC1, P0[days],				
		P1[days]				
MULT: Number of deeper mergers (4 bodies ((B-S)-S) or 5 bodies ((B-S)-S)-S)						
NEWHI: Counter of new hierarchical systems in chain						
quastab	89	impact.f	#15 ≥ 3	Tevent	Diagnostics for stability criterion of	
					two binaries in quadruple	
F-Label		TIME[NB], NAME(I1), NAME(J1), LQ, RI[NB], ECC1, EB[NB],				
		EB2[NB], EB1[NB], P1[days], PERIM[NB], PCR[NB]				

J1: Index of first member in second binary

LQ: orbit counter for diagnostics output

ECC1: Outer orbit eccentricity in B-B quadruple

EB: Quadruple binding energy EB1: First binary binding energy EB2: Second binary binding energy

P1: Outer orbit period

bs	91	mix.f	#19 ≥ 3	Tevent	Blue stra	aggler inforn	nation	
F-Label		TIME[NB]	, NAME(I	ΛΕ(I2),	M(INEW),	ECC,	P[days],	
		P(I1)[days], P(I2)[days]						
wdcirc	95	spiral.f	#27 > 0	Tevent	T_{event} Diagnostics for recent WD as the		s the sec-	
					ond con	nponent of b	inary s	ystem in-
					volving	tidal circular	rization	
F-Label		TIME[NB], NAME(I2), NAME(I2)			AME(I1),	K*(I1),	ROT	(I1)[NB],
$ROT(I2)[NB]$, $\langle motion \rangle$, $SPIN(I2)[NB]$								
(motion) : sqrt(M(ICM) / (RSM * SEMI)**3) [NB]								
*SPIN: spin of star ?								

cirdiag	96	hut.f	#27 > 0	Tevent	Diagnostics for reducing steps of inte-
					gration equations for eccentricity and
					angular velocites of stars (Equation
					15.22 in Sverre, 2003 book)
F-Label NSTEPS, IT, U, UD, DTAU (All in NB unit)			in NB unit)		

NSTEPS: Step number for integration

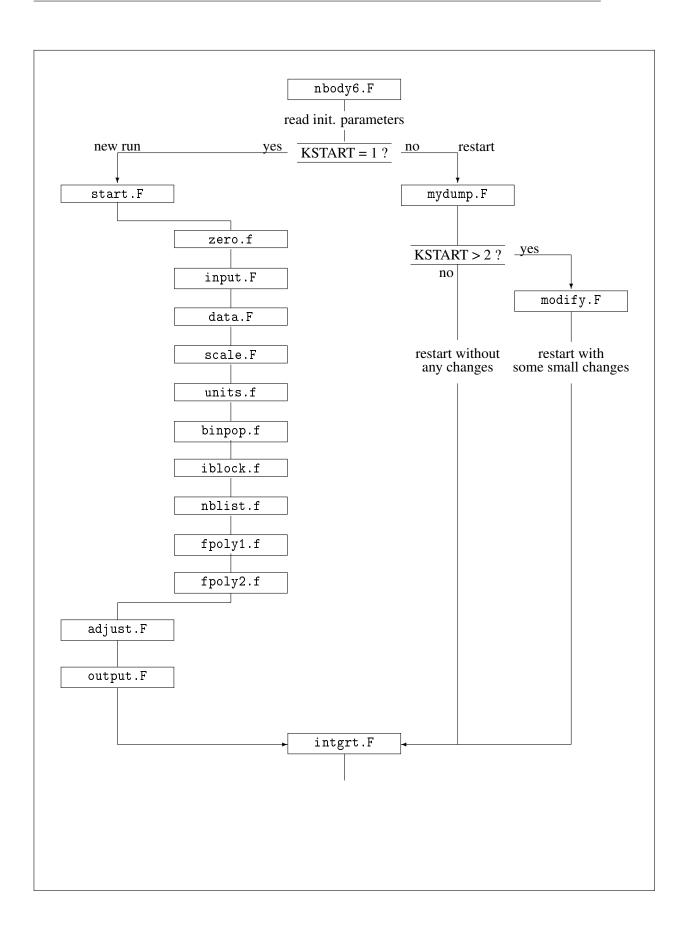
IT: Iteration times for reduction

U: KS vector U

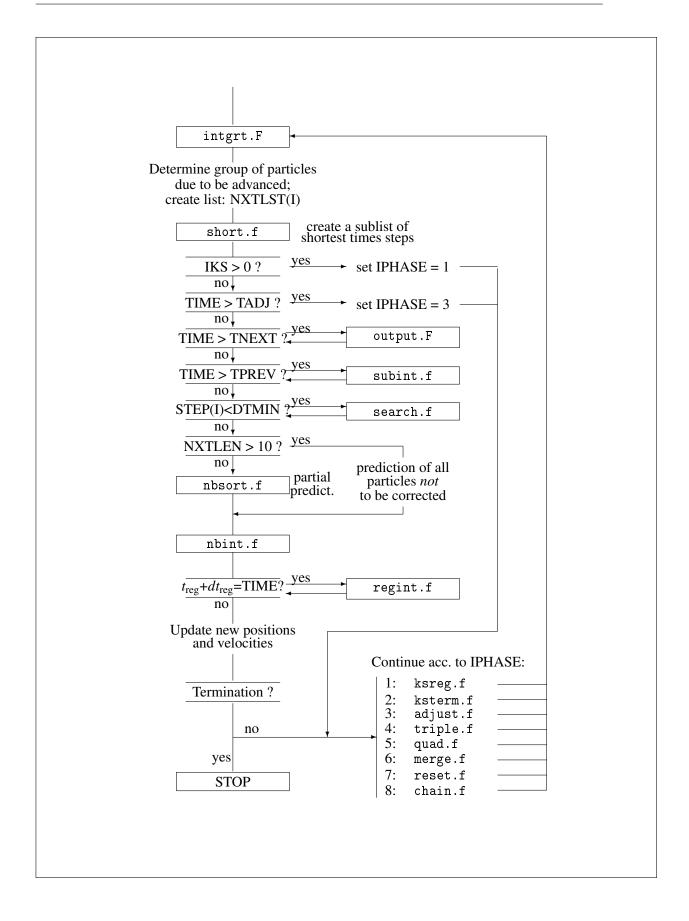
UD: KS vector UDOT

DTAU: KS integration time step

Flow chart 59



60 Flow chart



References 61

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