(Local) In[3]:=

AppendTo [\$Path , NotebookDirectory []];

<< Matte

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Tensor Spherical Harmonics are made up of *irreducible* rank-n tensors

These are described nicely in Leal Advanced Transport Phenomena Pg. 527 (within Chapter 8 Creeping Flows)

They can be developed from the solutions to the Laplace equation in spherical coordinates. Where you are taking repeated gradients of (1/r).

I.E. $\nabla(\nabla(1/r))$ not Laplacian but a gradient of a gradient.

That should be correct, but constants are really important and magnitude of r is equal to 1, and so it is often written in a different method.

$$\phi_{-(n+1)} = \frac{(-1)^n}{(2n-1)!!} \nabla_1 \dots \nabla_n \left(\frac{1}{r}\right), \ n = 0, 1, 2,...$$

That technically gets you the -(n+1) th decaying harmonic, but then the nth order irreducible tensor just has you substitute |r| = 1

First use a recurrence relation to get the repeated gradient of $\frac{1}{r}$ where invr[1] is $\frac{1}{r}$

(Local) In[5]:=

```
RepeatedPartialR [0] := invr[1];
RepeatedPartialR [n_] := partialr[RepeatedPartialR [n-1], idx[n]] // contract;
```

Now use that to get the decaying harmonics and replace inverse r to any power with 1

So for any vector r, the irreducible rank-n tensor that describes it

IrreducibleRanknTensor [0] // prettyprint

IrreducibleRanknTensor [1] # prettyprint

IrreducibleRanknTensor [2] // prettyprint

IrreducibleRanknTensor [3] // prettyprint

IrreducibleRanknTensor [4] // prettyprint

IrreducibleRanknTensor [5] // prettyprint

(Local) Out[21]=

(Local) Out[22]=

 r_i

(Local) Out[23]=

$$r_i r_j - \frac{\delta_{i,j}}{3}$$

(Local) Out[24]=

$$r_{i} r_{j} r_{k} - \frac{1}{5} r_{k} \delta_{i,j} - \frac{1}{5} r_{j} \delta_{i,k} - \frac{1}{5} r_{i} \delta_{j,k}$$

(Local) Out[25]=

$$r_{i} r_{j} r_{k} r_{l} - \frac{1}{7} r_{k} r_{l} \delta_{i,j} - \frac{1}{7} r_{j} r_{l} \delta_{i,k} - \frac{1}{7} r_{j} r_{k} \delta_{i,l} - \frac{1}{7} r_{i} r_{l} \delta_{j,k} + \frac{1}{35} \delta_{i,l} \delta_{j,k} - \frac{1}{7} r_{i} r_{k} \delta_{j,l} + \frac{1}{35} \delta_{i,k} \delta_{j,l} - \frac{1}{7} r_{i} r_{j} \delta_{k,l} + \frac{1}{35} \delta_{i,j} \delta_{k,l}$$

$$\begin{split} &r_{i}\,r_{j}\,r_{k}\,r_{l}\,r_{p} - \frac{1}{9}\,r_{k}\,r_{l}\,r_{p}\,\delta_{i,j} - \frac{1}{9}\,r_{j}\,r_{l}\,r_{p}\,\delta_{i,k} - \frac{1}{9}\,r_{j}\,r_{k}\,r_{p}\,\delta_{i,l} - \frac{1}{9}\,r_{j}\,r_{k}\,r_{l}\,\delta_{i,p} - \frac{1}{9}\,r_{i}\,r_{l}\,r_{p}\,\delta_{j,k} + \frac{1}{63}\,r_{l}\,\delta_{i,p}\,\delta_{j,k} - \frac{1}{9}\,r_{i}\,r_{k}\,r_{p}\,\delta_{j,l} + \frac{1}{63}\,r_{p}\,\delta_{i,k}\,\delta_{j,l} + \frac{1}{63}\,r_{k}\,\delta_{i,p}\,\delta_{j,l} - \frac{1}{9}\,r_{i}\,r_{k}\,r_{l}\,\delta_{j,p} + \frac{1}{63}\,r_{l}\,\delta_{i,k}\,\delta_{j,p} + \frac{1}{63}\,r_{k}\,\delta_{i,l}\,\delta_{j,p} - \frac{1}{9}\,r_{i}\,r_{j}\,r_{p}\,\delta_{k,l} + \frac{1}{63}\,r_{p}\,\delta_{i,j}\,\delta_{k,l} + \frac{1}{63}\,r_{j}\,\delta_{i,p}\,\delta_{k,l} + \frac{1}{63}\,r_{j}\,\delta_{i,p}\,\delta_{k,l} - \frac{1}{9}\,r_{i}\,r_{j}\,r_{l}\,\delta_{k,p} + \frac{1}{63}\,r_{l}\,\delta_{i,j}\,\delta_{k,p} + \frac{1}{63}\,r_{j}\,\delta_{i,l}\,\delta_{k,p} + \frac{1}{63}\,r_{j}\,\delta_{i,l}\,\delta_{k,p} + \frac{1}{63}\,r_{j}\,\delta_{i,l}\,\delta_{k,p} + \frac{1}{63}\,r_{j}\,\delta_{i,l}\,\delta_{k,p} - \frac{1}{9}\,r_{i}\,r_{j}\,r_{k}\,\delta_{l,p} + \frac{1}{63}\,r_{j}\,\delta_{i,p}\,\delta_{l,p} + \frac{1}{63}\,r_{j}\,\delta_{i,k}\,\delta_{l,p} + \frac{1}{63}\,r_{i}\,\delta_{j,k}\,\delta_{l,p} +$$

Here are the repeated gradients of 1/r to show you that they are equivalent, just divided by leading constants.

(Local) In[14]:=

x1 = partialr[invr[1], idx[1]] // contract; x2 = partialr[x1, idx[2]] // contract; x3 = partialr[x2, idx[3]] // contract; x1 // prettyprint x2 // prettyprint x3 // prettyprint

(Local) Out[17]=

$$-\frac{r_i}{r^3}$$

$$\frac{3 r_i r_j}{r^5} - \frac{\delta_{i,j}}{r^3}$$

$$-\frac{15\;r_{i}\;r_{j}\;r_{k}}{r^{7}}+\frac{3\;r_{k}\;\delta_{i,j}}{r^{5}}+\frac{3\;r_{j}\;\delta_{i,k}}{r^{5}}+\frac{3\;r_{i}\;\delta_{j,k}}{r^{5}}$$