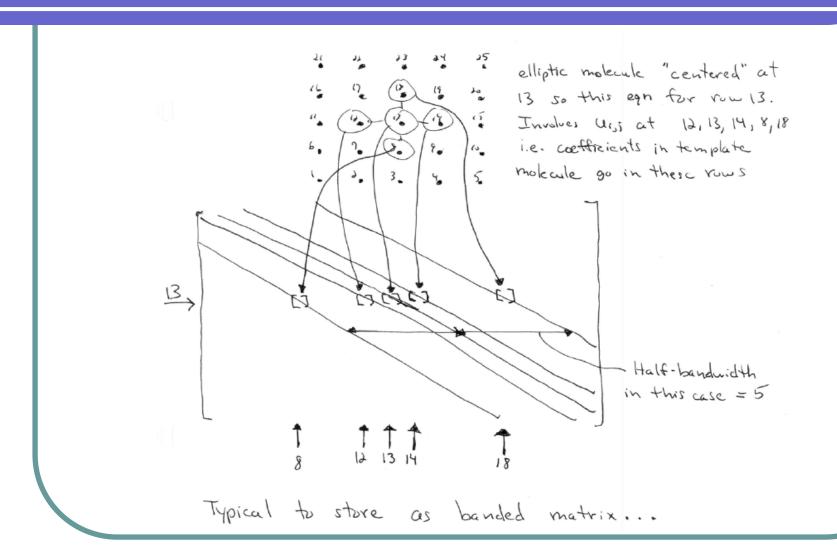
FDM Matrix Structure

- Need a mapping between (i,j) template location and (i,j) matrix entry in A ... assign a unique number to each mesh point ... generates pentadiagonal structure provided some "natural" ordering is used

- each (i,i) template "center" maps to unique

FDM Matrix Structure



LT: Introduction to the Method of Weighted Residuals

- Consider $\nabla^2 \mathbf{u} + \mathbf{f} \mathbf{u} = \mathbf{g} \rightarrow (\nabla^2 + \mathbf{f}) \mathbf{u} = \mathbf{g}$
- Finite Difference Approach
 Approximate L w/ L_{ij} → (δ² + f)u_{ij} = g_{ij}
 - i.e. replace "differential" operator w/ "difference" operator → get "exact solution" to "approximate operator"
 - Limitations:
 - cumbersome on irregular meshes
 - curved boundaries difficult to handle
 - u only found (i,j) points need interpolation strategy

- Weighted Residuals • Approximate ${\bf u}$ as $\hat{u}=\sum_{j=1}^N c_j\Phi_j(x,y,z)$ known function - "basis" func. coefficients - "trial" func.
 - Define "Residual" $\rightarrow R(\hat{u}) = L\hat{u} g$
 - For exact solution: $R(\hat{u}) = 0$ everywhere then

- "expansion" func.

- Want $R(\hat{u}) = 0$ to vanish in average way, one way is in weighted integral sense $\iiint R(\hat{u})W(x,y)dxdydz = 0$ any function of position
- " $R(\hat{u})$ orthogonal to all W(x,y,z)"

• So for $\hat{\mathbf{u}} : \mathbf{R}(\hat{\mathbf{u}}) \neq 0$... choose 'N' \mathbf{c}_{j} 's such that $\langle \mathbf{R}(\hat{\mathbf{u}}), \mathbf{W}_{i} \rangle = 0$ for i = 1, 2, 3, ..., N "Inner Product" $\Rightarrow \langle \mathbf{a}, \mathbf{b} \rangle \equiv \iint \mathbf{a} \cdot \mathbf{b} \, d\mathbf{x} \, d\mathbf{y} \, d\mathbf{z}$

- W_i's set of "weighting" functions → finite!
 local "testing" functions
- Use 'N' independent W_i's → generate 'N' equations in 'N' unknown c_i's

 $\left\langle R(\hat{u}), W_i \right\rangle = \sum_{i=1}^{N} c_j \left\langle L\left(\varphi_j\right) W_i \right\rangle = \left\langle g, W_i \right\rangle \ \, \text{for each $W_i(x,y)$ } \ \, i=1,2,...N$

• Necessary, but not sufficient for $\hat{\mathbf{u}} = \mathbf{u}$

- Continuous function must be zero if it is orthogonal to every member of a complete set ∴WRM can be thought of as a technique which enforces orthogonality between basis and weighting function sets
 - General idea is that the basis is a subset of a complete set that can represent any function, i.e. the true solution
 - As number of coefficients goes up, the approximation approaches the complete set

- Take home interpretation:
 - PDE is determined from first principles
 - Introduce the idea of a weighting function which is a known function of space – on its introduction, the physics behavior is prescribed to a region
 - Integration formulation
 - Introduction of basis as an approximation to the solution – on its introduction, it provides a description of how the solution will be treated behaviorally within in a local region

- Weighted Residual Methods Summary
 - L is typically a differential operator
 - W_i not complete in practice (N finite), but make "R(u) orthogonal to 1st N members of a complete set"
 - "Approximate solution" which exactly satisfies "differential relations" in PDE
 - Is an "integral" formulation

- What's involved numerically?
 - Expand unknown solution as $\hat{\mathbf{u}} = \sum_{j=1}^{n} \mathbf{c}_{j} \phi_{j}(\mathbf{x}, \mathbf{y}, \mathbf{z})$
 - Finite sum
 - $\phi_j(x,y,z)$ known function of (x,y,z)
 - c_j's unknown coefficients to determine
 - Generate system of equations in unknowns $R(\hat{u}) = \langle (L(\hat{u}) g), W_i(x,y,z) \rangle = 0$ for i = 1,2,3,...,N

$$\sum_{j=1}^{N} c_{j} \left\langle L(\phi_{j}), W_{i}(x, y, z) \right\rangle = \left\langle g, W_{i}(x, y, z) \right\rangle \text{ for } i = 1, 2, 3, \dots, N$$

- W_i(x,y,z) known functions of (x,y,z)
- For each i, generate algebraic equation in c_i's
- Creates 'N' equations in 'N' unknowns

- But analytic methods...
 - Require special knowledge of how to choose basis and weighting functions
 - Different choices are needed for different problems
 - Usually need an infinite # of them
 - Can't find them for many practical problems
- Numerically want ...
 - Basis and Weighting functions to be simple ... easy to integrate
 - Single choice suitable for many problems
 - Can only use finite #, but want convergence as number used increases

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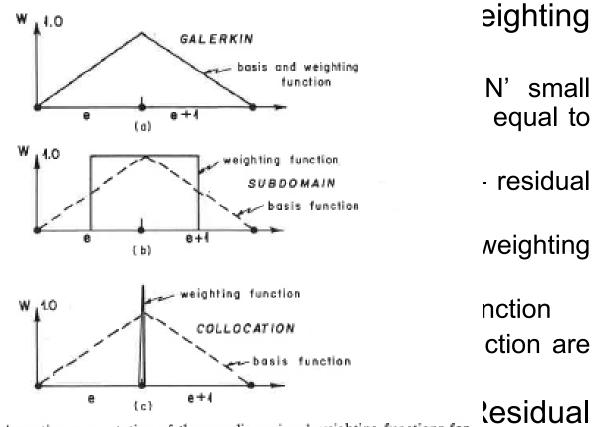


Figure 2.4. Schematic representation of the one-dimensional weighting functions for the Galerkin, subdomain and collocation methods. (It is assumed here that the chapeau function is used as a basis for all methods.)

Galerkin Weighted Residual Method

- Key Feature:
 - Basis function and weighting function are selected to be the same function form
- So what do we choose?
 - Polynomials may be a nice choice
 - Other's exist

- Why?
 - Easy to differentiate
 - Represent a complete set for continuous functions ... e.g. Taylor Polynomial
- Lagrange Polynomials
 - Nth order polynomial for N+1 pts
 - Easily automated
 - Handy on uneven grids
 - $\phi_i(x) = 0$ at $x = x_i$ for $i \neq j$ $\phi_i(x) = 1$ at $x = x_i$ for i = j

Lagrange Polynomials

•
$$\phi_j(\mathbf{x}_j) = 1 \Rightarrow \mathbf{u}(\mathbf{x}) = \sum_{j=1}^{N} \mathbf{c}_j \phi_j(\mathbf{x}, \mathbf{y}) \Rightarrow \mathbf{u}(\mathbf{x}_i) = \mathbf{c}_i$$

$$\therefore \quad \mathbf{u}(\mathbf{x}) = \sum_{j=1}^{N} \mathbf{u}_{j} \phi_{j}(\mathbf{x}, \mathbf{y})$$

Coefficients are solution at nodes ... similar to FD but have functional form specified in between

• everywhere $\sum_{i=1}^{N} \phi_i(\mathbf{x}, \mathbf{y}) = 1$

$$\sum_{j=1}^{N} \frac{\partial \phi_{j}(\mathbf{x}, \mathbf{y})}{\partial \mathbf{x}} = \frac{\partial}{\partial \mathbf{x}} \sum_{j=1}^{N} \phi_{j}(\mathbf{x}, \mathbf{y}) = 0$$

- Global Polynomials
 - Potential for disaster ... "polynomial wiggle"
 - Can have N zeros
 - Sensitive to all u_i "Global Support"
 - Can have large variation between points
- Local Interpolation
 - Use subset of nodes to represent solution over localized areas – so-called finite elements
 - Possibilities : Overlapping Domains, Nonoverlapping Domains

- Overlapping Domains
 - Nonunique
 - Somewhat difficult to implement
 - Further study of this in meshless methods
- Non-overlapping Domains
 - Idea of finite "element" as collection of nodes, i.e. unit of local interpolation
 - Unequal spacing of nodes, i.e. conformity of shape
 - Basic building block of FEM
 - Typically use same type of element throughout for programming ease (unless reason not to do so)

- Continuity of û In 1D, if N+1 nodes/element
 - ûis locally Nth order polynomial
 - On element interior ... 1st N derivatives continuous
 - ... but, at element boundaries only $\hat{\mathbf{u}}$ is continuous; $\partial \hat{\mathbf{u}}_{\partial \mathbf{x}}$ changes abruptly
 - "Co" continuity → continuous in 0th derivative
- Higher order continuity possible
 e.g. "C¹" continuity ... need Hermite polynomial
 - Simplest local unit: Hermite cubic
 - ∂û/_⊘ becomes "nodal parameter"

FEM Steps – 1D Example Problem

FEM Steps

- Determine PDE expression for physics
- On Paper
 - Weight each expression with "trial" function
 - Integrate the expression
 - Approximate solution as expansion containing coefficients and "basis" functions
 - Select the "weighting" function
- In Code
 - Assemble matrices by a process of multi-dimensional numerical integration
 - Apply boundary conditions
 - Solve for unknown expansion coefficients
 - Construct approximate solution
 - Perhaps construct derived quantities

Example:
$$\frac{d^2U}{dx^2} + \int U = g$$
 $\frac{dU}{dx} = \int \frac{dU}{dx} = \int \frac{dU}{d$

Step 1: Generate weighted residual equation

$$\left\langle \frac{d}{dx} \left(\frac{du}{dx} \right), \phi \right\rangle$$

$$\int u dv = uv - \int v du$$

$$\mathbf{u} = \mathbf{\phi}$$

$$u = \phi$$
 $dv = \frac{d}{dx} \left(\frac{du}{dx} \right) dx$

$$\mathbf{du} = \frac{\mathbf{d\phi}}{\mathbf{dx}} \mathbf{dx} \qquad \mathbf{v} = \frac{\mathbf{du}}{\mathbf{dx}} \Big|_{\mathbf{c}}^{\mathbf{L}}$$

$$v = \frac{du}{dx} \Big|_{0}^{L}$$

SO

$$\left\langle \frac{\mathbf{d}^2 \mathbf{u}}{\mathbf{dx}^2}, \mathbf{\phi} \right\rangle = \mathbf{\phi} \frac{\mathbf{du}}{\mathbf{dx}} \bigg|_{0}^{L} - \left\langle \frac{\mathbf{du}}{\mathbf{dx}} \frac{\mathbf{d\phi}}{\mathbf{dx}} \right\rangle$$

First form of Green's theorem

$$\left\langle \mathbf{f} \nabla^2 \mathbf{g} + \nabla \mathbf{f} \bullet \nabla \mathbf{g} \right\rangle = \oint \mathbf{f} \nabla \mathbf{g} \bullet \hat{\mathbf{n}} dS$$

let

$$\mathbf{f} = \mathbf{\phi}, \quad \mathbf{g} = \mathbf{u}$$

$$\left\langle \Phi \frac{d^2 u}{dx^2} \right\rangle = \Phi \frac{du}{dx} \bigg|_{0}^{L} - \left\langle \frac{du}{dx} \frac{d\Phi}{dx} \right\rangle$$

Why... reduces continuity requirements needed on Osi, wi ... integrand may be piecewise discontinuous with finite discontinuities.

Abother sense in which the approach is "weak"

Step 3: Chase weighting function; W:= Dij Galerkin

$$0_{iS} = \left\langle -\frac{dd_{i}}{dx} \frac{dd_{i}}{dx} + f O_{i} O_{i} \right\rangle$$

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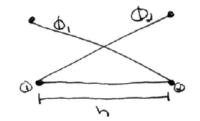
$$0_{iS} = \left\langle -\frac{dd_{i}}{dx} \frac{dd_{i}}{dx} + f O_{i} O_{i} \right\rangle$$

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$$0_{iS} = \left\langle -\frac{dd_{i}}{dx} \frac{dd_{i}}{dx} + f O_{i} O_{i} \right\rangle$$

$$0_{iS} = \left\langle -\frac{dd_{i}}{dx} \frac{dd_{i}}{dx} - \frac{dd_{i}}{dx} \frac{dd_{i}}$$



$$= -\frac{1}{1} \left(\frac{3}{x_{3}} - \frac{3}{x_{5}} (x_{1} + x_{1-1}) + x + x_{1} x_{1-1} \right) dx$$

$$= -\frac{1}{1} \left(\frac{3}{x_{1}} - \frac{3}{x_{1}} (x_{2} + x_{1}) + x_{1} x_{1-1} \right) dx$$

$$= -\frac{1}{1} \left(\frac{3}{x_{1}} - \frac{3}{x_{1}} (x_{2} + x_{1}) + x_{1} x_{1-1} \right) dx$$

$$= -\frac{1}{1} \left(\frac{3}{x_{1}} - \frac{3}{x_{1}} (x_{2} + x_{1}) + x_{1} x_{1-1} \right) dx$$

who loss of generality ... pick
$$x_{i-1} = \frac{1}{3}$$

$$= -\frac{1}{3} \left[\frac{x_i^3}{3} - \frac{x_i^3}{3} \right] = \frac{1}{3} \left[\frac{h_i^3}{3} - \frac{h_i^3}{3} \right] = \frac{h_i}{6}$$

$$\Phi' = \frac{x^2 - x}{y} \qquad \frac{qx}{qx} = \frac{-1}{y}$$

$$\Phi_3 = \frac{x - x}{h}$$
 $\frac{dx}{d\phi} = \frac{1}{h}$

$$= -\frac{1}{12} \left[\frac{3}{x_{13}} - \frac{3}{x_{13}} \right] = \frac{1}{12} \left[\frac{3}{y_{13}} - \frac{3}{y_{13}} \right] = \frac{9}{12}$$

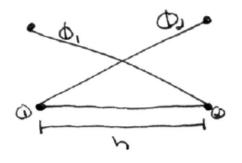
... conclude we have exact integration/differentiation in our WR method... approximation is in the assumption of linear variation of solution between nodes

and our RHS terms become:

$$\frac{f_{0}}{6} \cdot \frac{h_{1} + h_{2}}{6} \cdot \frac{h_{2} + h_{2}}{6} \cdot \frac{h_{2}}{6} \cdot$$

Very similar to FO method!

Integration Formulas for 10 Elements



$$\langle () \rangle = \int_{x_{i}}^{x_{i}} () dx$$

$$\Phi' = \frac{x^2 - x}{y} \qquad \frac{qq'}{qx} = \frac{-1}{y}$$

$$\Phi_{a} = \frac{h}{x - x}$$
 $\frac{dx}{dx} = \frac{1}{h}$

$$\langle a \rangle = h$$

 $\langle a \rangle = h$
 $\langle a \rangle = h$

Integration Fermulas for 1-0 Quadratic Elements

$$\begin{array}{c}
\omega / & \text{Equally spoced nodes} \\
0 & \omega / & \text{Equally spoced nodes}
\end{array}$$

$$\begin{array}{c}
\phi_3 \\
\phi_1 = \frac{\lambda}{\lambda} \\
\phi_1 = \frac{\lambda}{\lambda} \\
\phi_2 = \frac{\lambda}{\lambda} \\
\phi_3 = \frac{\lambda}{\lambda} \\
\phi_$$

$$\langle a, \gamma = h |_3$$

 $\langle a, \gamma = 4h |_3$
 $\langle a, \gamma = h |_3$
 $\langle a, \gamma = h |_15$
 $\langle a, \alpha = 2h |_{15}$
 $\langle a, \alpha = 2h |_{15}$

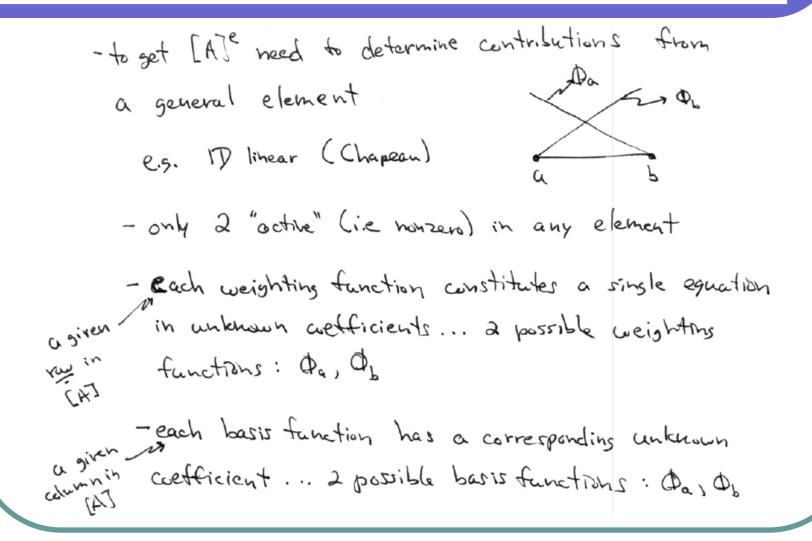
$$\langle \dot{a}_{1} \rangle = 1$$
 $\langle \dot{a}_{1} \rangle = -1/3$
 $\langle \dot{a}_{1} \rangle = 0$
 $\langle \dot{a}_{1} \rangle = 1$
 $\langle \dot{a}_{1} \rangle = 1$
 $\langle \dot{a}_{1} \dot{a}_{2} \rangle = 1/3$
 $\langle \dot{a}_{1} \dot{a}_{3} \rangle = 1/3$
 $\langle \dot{a}_{2} \dot{a}_{3} \rangle = 1/3$
 $\langle \dot{a}_{3} \dot{a}_{3} \rangle = 1/3$
 $\langle \dot{a}_{3} \dot{a}_{3} \rangle = 1/3$
 $\langle \dot{a}_{3} \dot{a}_{3} \rangle = 1/3$

MORE INTECRATION FORMULAS FOR QUADRATEC

$$\langle \phi_{1}^{3} \rangle - \langle \phi_{3}^{3} \rangle = \frac{39L}{420}$$
 $\langle \phi_{1}^{2} \phi_{1} \rangle = \langle \phi_{2} \phi_{3}^{2} \rangle = \frac{20L}{420}$
 $\langle \phi_{1}^{2} \phi_{3} \rangle = \langle \phi_{1} \phi_{3}^{2} \rangle = -3L/420$
 $\langle \phi_{1} \phi_{3} \rangle = \langle \phi_{1} \phi_{3}^{2} \rangle = 16L/420$
 $\langle \phi_{1} \phi_{3} \rangle = -8L/420$
 $\langle \phi_{3}^{3} \rangle = \frac{192L}{420}$

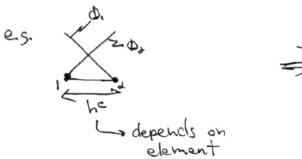
- · In FO common to proceed "molecule-by-molecule" with each representing single difference equation
- · In FE more natural to proceed "element-by-element" need to integrate over problem domain => union of non overlapping elements ...

where
$$a_{ij} = \left(-\frac{d\Phi_{ij}}{dx} \frac{dd_{i}}{\partial x}, f\Phi_{i}, \Phi_{ij}\right)$$
; $b_{ij} = \langle g\Phi_{ij} \rangle - \frac{du}{\partial x}\Phi_{i}|_{i}^{k}$



All possible combinations ... Weighting function (Rou) Basis function (column) Da Da Op6 Da Oh, So [A]e has only 4 nonzero coefficients: ae aa aa ab ab

so [A]e may be stored as a 2xd submotrix where we use a "local" node numbering scheme



$$\alpha'' = \langle \frac{d\alpha'}{d\alpha'}, \frac{d\alpha'}{d\alpha'} + t\alpha'\alpha' \rangle = \frac{1}{1}e + \frac{3}{t\gamma_6}$$

structure the same regardless of PDE for linear ID element... only details of coefficients differ

be
$$= \langle g \phi, \gamma^e - \frac{gh^e}{d} \rangle$$

assume g constant... and

hestectins boundary term

for the moment...

So

$$[A]^e = \begin{bmatrix} -\frac{1}{he} + \frac{fh^e}{3} & \frac{1}{he} + \frac{fh^e}{6} \\ \frac{1}{he} + \frac{fh^e}{6} & -\frac{1}{he} + \frac{fh^e}{3} \end{bmatrix}; \{b\}^e = \begin{bmatrix} gh^e \\ gh^e \\ \frac{1}{he} + \frac{fh^e}{6} & -\frac{1}{he} + \frac{fh^e}{3} \end{bmatrix}$$