CONTINUOUS BEAM STRUCTURAL ANALYSIS

Finite element analysis software tools

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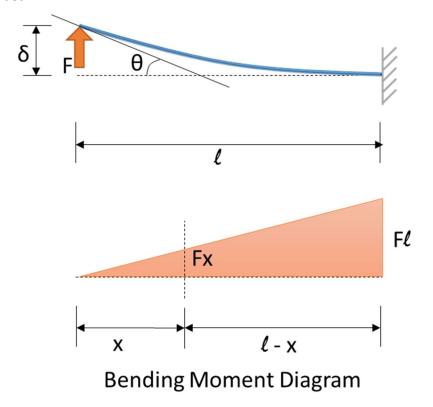
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A. Stiffness matrix of beam element

Stiffness matrix for beam element in this section is derived using Euler – Bernoulli beam theory. The unknowns for a beam element is its end rotation and displacement which is resisted by the beam's bending stiffness. Euler – Bernoulli beam theory relates the deflection of a flexure member to its superimposed load.

$$EI\left(\frac{d^2y}{dx^2}\right) = M$$
 Eq. B1

Let's find the displacement and rotation caused by a Force F acting at the free end of a cantilever beam



$$EI\left(\frac{d^2y}{dx^2}\right) = M$$
 From Eq. B1

From the bending moment diagram,

$$EI\left(\frac{d^2y}{dx^2}\right) = Fx$$

$$EI\left(\frac{dy}{dx}\right) = \frac{Fx^2}{2} + c_1$$
 Eq. B2

$$EIy = \frac{Fx^3}{6} + c_1x + c_2$$
 Eq. B3

Applying the boundary conditions, (i) at fixed end

$$\left(\frac{dy}{dx}\right) = 0 \ at \ x = l$$

$$c_1 = \frac{-Fx^2}{2}$$
 Eq. B4

(ii) at the fixed end

$$y = 0 \text{ at } x = l$$

$$\frac{Fl^{3}}{6} - \frac{Fl^{3}}{2} + c_{2} = 0$$

$$c_{2} = \frac{Fl^{3}}{3}$$
Eq. B5

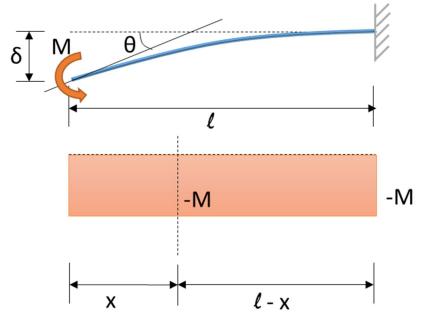
Applying $c_1 \& c_2$ to Eq. D2 & Eq. D3

$$EI\theta_{xF} = \frac{Fx^2}{2} - \frac{Fl^2}{2}$$
 Eq. B6
 $EI\delta_{xF} = \frac{Fx^3}{6} - \frac{Fxl^2}{2} + \frac{Fl^3}{3}$ Eq. B7

The displacement and Slope at x = 0 caused by the force F is

$$\delta_{0F}=rac{Fl^3}{3EI}$$
 Eq. B8
$$\theta_{0F}=-rac{Fl^2}{2EI}$$
 Eq. B9

Similarly we find the displacement and rotation caused by a Moment M acting at the free end of a cantilever beam



Bending Moment Diagram

$$EI\left(\frac{d^2y}{dx^2}\right) = -M$$
 Eq. B10

$$EI\left(\frac{dy}{dx}\right) = -Mx + c_1$$
 Eq. B11

$$EIy = -\frac{Mx^2}{2} + c_1x + c_2$$
 Eq. B12

Applying the boundary conditions, (i) at fixed end

$$\left(\frac{dy}{dx}\right) = 0 \ at \ x = l$$

$$c_1 = Ml$$

(ii) at the fixed end

$$y = 0$$
 at $x = l$

$$-\frac{Ml^2}{2} + Ml^2 + c_2 = 0$$

$$c_2 = -\frac{Ml^2}{2}$$
 Eq. B14

Eq. B13

Applying c₁ & c₂ to Eq. D11 & Eq. D12

$$EI\theta_{xM}=-Mx+Ml$$
 Eq. B15
$$EI\delta_{xM}=-\frac{Mx^2}{2}+Mxl-\frac{Ml^2}{2}$$
 Eq. B16

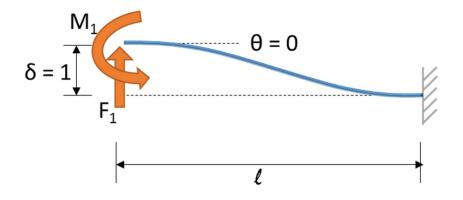
The displacement and Slope at x = 0 caused by the moment M is

$$\delta_{0M} = -\frac{Ml^2}{2EI}$$
 Eq. B17
$$\theta_{0M} = \frac{Ml}{EI}$$
 Eq. B18

So far we have found the end displacement and end rotation caused by Force acting on the free end and moment acting on the free end separately.

We can now find the stiffness by deriving the force and moment required to cause unit displacement and unit rotation.

Phase: 1 => Force and Moment required to cause unit displacement



Form the above beam diagram, the conditions are

$$\delta_{0F1}+\delta_{0M1}=1$$
 Eq. B19

$$\theta_{0F1} + \theta_{0M1} = 0 Eq. B20$$

Using equations D8 & D17 in D19

$$\frac{F_1 l^3}{3EI} - \frac{M_1 l^2}{2EI} = 1$$
 Eq. B21

Using equations D9 & D18 in D20

$$-rac{F_{1}l^{2}}{2EI} + rac{M_{1}l}{EI} = 0$$
 Eq. B22

Solving the Equation D21 & D22 for F_1 and M_1

$$\frac{M_1 l}{EI} = \frac{F_1 l^2}{2EI}$$

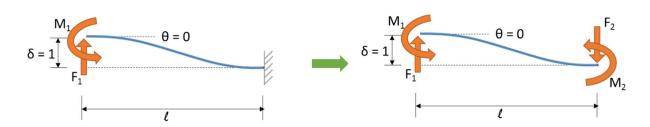
Substituting the above in D21

$$\frac{F_1l^3}{3EI}-\frac{F_1l^2}{2EI}\Big(\frac{l}{2}\Big)=1$$

$$F_1=\frac{12EI}{l^3}$$
 Eq. B23

Substituting the above in D22

$$\frac{M_1 l}{EI} = \frac{l^2}{2EI} \left(\frac{12EI}{l^3}\right)$$
 Eq. B24



Now we find the F_2 and M_2 $F_1 + F_2 = 0$ and $M_2 = F_1 \ell - M_1$

$$F_2 = -\frac{12EI}{l^3}$$
 Eq. B25

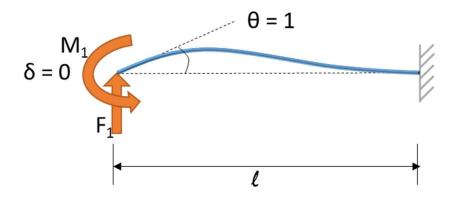
$$M_2 = \left(\frac{12EI}{l^3}\right)l - \frac{6EI}{l^2}$$
 Eq. B26

Now a partial (incomplete) beam stiffness matrix will look like below, using eqns D23, D24, D25, D26

$$\begin{bmatrix} \frac{12EI}{l^3} & \frac{6EI}{l^2} & -\frac{12EI}{l^3} & \frac{6EI}{l^2} \\ \frac{6EI}{l^2} & XX & XX & XX \\ -\frac{12EI}{l^3} & XX & XX & XX \\ \frac{6EI}{l^2} & XX & XX & XX \end{bmatrix} \begin{bmatrix} \delta_1 = 1 \\ \theta_1 = 0 \\ \delta_2 = 0 \end{bmatrix} = \begin{bmatrix} F_1 \\ M_1 \\ F_2 \\ M_2 \end{bmatrix}$$
 Eq. B27

XX - Not derived yet

Phase: 2 => Force and Moment required to cause unit rotation



We do the same exercise as phase 1 for unit rotation Form the above beam diagram, the conditions are

$$\delta_{0F1} + \delta_{0M1} = 0 \hspace{1.5cm} \text{Eq. B28}$$

$$heta_{0F1}+ heta_{0M1}=1$$
 Eq. B29

Using equations D8 & D17 in D28

$$\frac{F_1 l^3}{3EI} - \frac{M_1 l^2}{2EI} = 0$$
 Eq. B30

Using equations D9 & D18 in D29

$$-\frac{F_1 l^2}{2EI} + \frac{M_1 l}{EI} = 1$$
 Eq. B31

Solving the Equation D30 & D31 for F₁ and M₁

$$\frac{F_1 l^3}{3EI} = \frac{M_1 l^2}{2EI} = \frac{3M_1 l}{2EI}$$

$$\frac{3M_1 l}{2EI} = \frac{3M_1 l}{2EI}$$

Substituting the above in Eq. D31

$$-\left(\frac{1}{2}\right)\frac{3M_1l}{2EI} + \frac{M_1l}{EI} = 1$$

$$M_1 = \frac{4EI}{I}$$
 Eq. B32

Substituting the above in Eq. D31

$$F_1 = \frac{6EI}{l^2}$$
 Eq. B33

Now we find the F_2 and M_2 $F_1 + F_2 = 0$ and $M_2 = F_1 \ell - M_1$

$$F_2 = -\frac{6EI}{l^2}$$
 Eq. B34

$$M_2 = \left(\frac{6EI}{l^2}\right)l - \frac{4EI}{l}$$

$$M_2 = \frac{2EI}{I}$$
 Eq. B35

Now a partial (incomplete) beam stiffness matrix will look like below, using eqns D23, D24, D25, D26

$$\begin{bmatrix} \frac{12EI}{l^3} & \frac{6EI}{l^2} & -\frac{12EI}{l^3} & \frac{6EI}{l^2} \\ \frac{6EI}{l^2} & \frac{4EI}{l} & -\frac{6EI}{l^2} & \frac{2EI}{l} \\ -\frac{12EI}{l^3} & -\frac{6EI}{l^2} & XX & XX \\ \frac{6EI}{l^2} & \frac{2EI}{l} & XX & XX \\ \end{bmatrix} \begin{bmatrix} \delta_1 = 0 \\ \theta_1 = 1 \\ \delta_2 = 0 \\ \theta_2 = 0 \end{bmatrix} = \begin{bmatrix} F_1 \\ M_1 \\ F_2 \\ M_2 \end{bmatrix}$$
 Eq. B36

XX - Not derived yet

Either we can do the same derivation for δ_2 and θ_2 unit displacement & rotation or we can arrive at their stiffness by visualizing their free-body diagram.

For δ_2 =1, the force and moment will be the same but opposite direction from the derivation of δ_1 =1 and for θ_2 =1, the moment value is interchanged whereas the force values remains the same.

This leaves us with the stiffness matrix of a beam element,

$$\begin{bmatrix} \frac{12EI}{l^3} & \frac{6EI}{l^2} & -\frac{12EI}{l^3} & \frac{6EI}{l^2} \\ \frac{6EI}{l^2} & \frac{4EI}{l} & -\frac{6EI}{l^2} & \frac{2EI}{l} \\ -\frac{12EI}{l^3} & -\frac{6EI}{l^2} & \frac{12EI}{l^3} & -\frac{6EI}{l^2} \\ \frac{6EI}{l^2} & \frac{2EI}{l} & -\frac{6EI}{l^2} & \frac{4EI}{l} \end{bmatrix} \begin{bmatrix} \delta_i \\ \theta_i \\ \delta_j \\ \theta_j \end{bmatrix} = \begin{bmatrix} F_i \\ M_i \\ F_j \\ M_j \end{bmatrix}$$
 Eq. B37
$$\begin{pmatrix} \frac{EI}{l^3} \end{pmatrix} \begin{bmatrix} 12 & 6l & -12 & 6l \\ 6l & 4l^2 & -6l & 2l^2 \\ -12 & -6l & 12 & -6l \\ 6l & 2l^2 & -6l & 4l^2 \end{bmatrix} \begin{bmatrix} \delta_i \\ \theta_i \\ \delta_j \\ \theta_j \end{bmatrix} = \begin{bmatrix} F_i \\ M_i \\ F_j \\ M_j \end{bmatrix}$$
 Eq. B38

The above equation is the basis for beam analyzer software tool. For the frame analyzer the same element stiffness with the axial stiffness and transformation included as below

Element stiffness matrix in local coordinate for frame element

$$\begin{bmatrix} \frac{AE}{l} & 0 & 0 & -\frac{AE}{l} & 0 & 0\\ 0 & \frac{12EI}{l^3} & \frac{6EI}{l^2} & 0 & -\frac{12EI}{l^3} & \frac{6EI}{l^2}\\ 0 & \frac{6EI}{l^2} & \frac{4EI}{l} & 0 & -\frac{6EI}{l^2} & \frac{2EI}{l}\\ -\frac{AE}{l} & 0 & 0 & \frac{AE}{l} & 0 & 0\\ 0 & -\frac{12EI}{l^3} & -\frac{6EI}{l^2} & 0 & \frac{12EI}{l^3} & -\frac{6EI}{l^2}\\ 0 & \frac{6EI}{l^2} & \frac{2EI}{l} & 0 & -\frac{6EI}{l^2} & \frac{4EI}{l} \end{bmatrix} \begin{bmatrix} u_i \\ \delta_i \\ \theta_i \\ u_j \\ \delta_j \\ \theta_j \end{bmatrix} = \begin{bmatrix} P_i \\ F_i \\ M_i \\ P_j \\ F_j \\ M_j \end{bmatrix}$$

The transformation matrix for the global coordinate system is

$$L = \begin{bmatrix} l & m & 0 & 0 & 0 & 0 \\ -m & l & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & l & m & 0 \\ 0 & 0 & 0 & -m & l & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \end{bmatrix}$$

$$k_g = L^T k_l L$$

All the elements in Frame analyzer is stored in global coordinate system

B. Virtual work principle & stiffness matrix of 1D elements

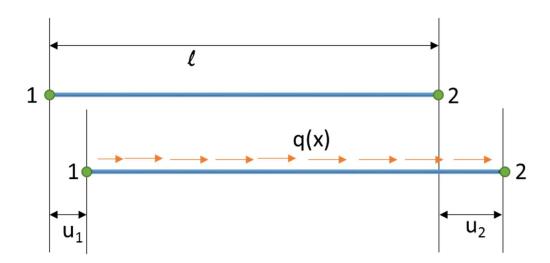
Strain Energy

In mechanics of material it is shown that the strain energy at a point of a linear elastic material subjected to one dimensional state of stress σ & strain ϵ is

$$strain\ energy\ dU = \frac{1}{2}\sigma\epsilon\ dV$$

Total strain energy integrated on the total volume V of the bar

$$W_I = \int\limits_V \frac{1}{2} \sigma \epsilon \ dV$$
 Eq C1



$$\sigma = E\epsilon$$

Uniform area throughout the length of the bar dV = Adx

Eq. E1 becomes

$$W_I = \int\limits_V \frac{1}{2} \epsilon E \epsilon \ A dx$$
 Eq. C2

Strain and displacements are linked by the relation

$$\epsilon(x) = \frac{d}{dx}(u(x))$$

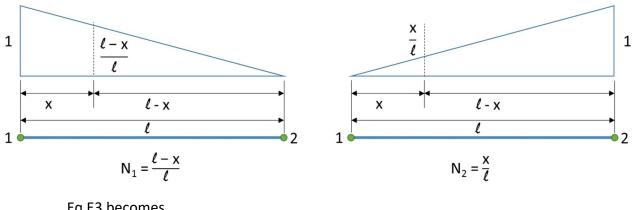
The real variation of displacement along the bar is not known. Instead we will interpolate the value of u(x) from the values of the displacements at the nodes u_1 and u_2 .

The above statement is the fundamental idea behind finite element analysis. A structure's behavior under an external loading is predicted by discretizing it into small elements and summing their individual elements. The accuracy of the prediction is directly proportional to the equations (shape functions) with which the elements are represented. There are lot of pitfalls in using finite element analysis where the picturesque post processing results from commercial software might overshadow the accuracy aspect. I will discuss these topic in the final section.

Now let's continue with selecting a shape function which might rightly represent a bar element. I will discuss only the first order elements in the following sections, the user can extend these idea to utilizing higher order elements. Note that higher order elements doesn't means accurate results. Whichever element closely represents the physical behavior of the structure is the right element which in my opinion none. But with the powerful tool of finite element analysis we can arrive at a solution to make a better design choice.

$$u(x) = N_1(x)u_1 + N_2(x)u_2$$
 Eq. C3

Where $N_1(x)$ and $N_2(x)$ are the shape functions



Eq E3 becomes

$$u(x) = \left(\frac{l-x}{l}\right)u_1 + \left(\frac{x}{l}\right)u_2$$
 Eq. C4
$$u(x) = \left[\frac{l-x}{l} \quad \frac{x}{l}\right] \begin{bmatrix} u_1 \\ u_2 \end{bmatrix}$$

$$u(x) = [N][u]$$

$$\in (x) = \frac{d}{dx}[N][u]$$

$$\in (x) = [B][u]$$
 Eq. C5

where

$$[B] = \frac{d}{dx}[N]$$
 Eq. C6

Is called the strain – displacement matrix

$$[B] = \frac{d}{dx} \left[\frac{l - x}{l} \qquad \frac{x}{l} \right] = \left[\frac{-1}{l} \qquad \frac{1}{l} \right]$$
 Eq. C7

C. Principle of minimum potential energy and stiffness matrix of bar element

For conservative systems, of all kinematically admissible displacement fields, those corresponds to equilibrium extremize the total potential energy. If the extreme condition's total potential energy is minimum, the equilibrium state is stable. In other words, from a super imposed external loading, a structure will lose its potential energy by deforming to a new position to attain equilibrium. The deformed state's potential energy will be minimum.

Total Potential energy

$$\pi = W_I - W_E$$

From principle of minimum potential energy

$$\partial \pi = \partial W_I - \partial W_E = 0$$
 Eq. C8

External work done

$$W_E = \int_L q(x)u(x)dx$$
 Eq. C9

Internal work done

$$W_I = \int_L \frac{1}{2} \epsilon E \epsilon A dx$$
 Eq. C10

Applying Eq.8 & Eq.9 in Eq.7

$$\partial \int_{L} \frac{1}{2} \epsilon E \epsilon A dx = \partial \int_{L} q(x) u(x) dx$$
$$\frac{1}{2} \int_{L} ((\partial \epsilon) E \epsilon + \epsilon E(\partial \epsilon)) A dx = \int_{L} q(x) \partial u(x) dx$$

Note that the loading q(x) is not changing from the initial state to final state, where there is a potential energy change

$$\int_{L} \epsilon E(\partial \epsilon) A dx = \int_{L} q(x) \, \partial u(x) dx$$

$$\int_{L} [B]^{T}[u]E[B][\partial u]Adx = \int_{L} q(x) \, \partial([N][u])dx$$

$$\frac{\partial u}{\partial u}AE \int_{L} [B]^{T}[u][B]dx = \frac{\partial u}{\partial u} \int_{L} q(x) [N]dx$$

$$AE \int_{L} [B]^{T} [B] dx [u] = \int_{L} q(x) [N] dx$$
 Eq. C11

$$[K][u] = [F]$$
 Eq. C12

Now the element stiffness matrix of 1D bar element is given by

$$[k] = AE \int_{L} [B]^{T} [B] dx$$
 Eq. C13

Using Eq.7 in Eq. E13

$$[k] = AE \int_{L} \begin{bmatrix} \frac{-1}{l} \\ \frac{1}{l} \end{bmatrix} \begin{bmatrix} \frac{-1}{l} & \frac{1}{l} \end{bmatrix} dx$$

$$[k] = AE \int_{L} \begin{bmatrix} \frac{1}{l^2} & -\frac{1}{l^2} \\ -\frac{1}{l^2} & \frac{1}{l^2} \end{bmatrix} dx$$
$$[k] = AE \begin{bmatrix} \frac{1}{l} & -\frac{1}{l} \\ -\frac{1}{l} & \frac{1}{l} \end{bmatrix}$$

The element stiffness matrix of 1D bar element is

$$[k] = \frac{AE}{l} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix}$$

Note that the above matrix is the same as Eq. C1

Eq. C14

D. Deriving shape function and stiffness matrix of beam element

We shall do the above exercise to derive the stiffness matrix of beam element. Note that the derivation of shape function is the fundamental behind formulating the finite element.

Hermite Splines

$$f(0) = f_0$$

 $f'(0) = g_0$
 $s = 0$
 $f(1) = f_1$
 $f'(1) = g_1$
 $f'(1) = g_1$

A cubic polynomial s[0,1] is

$$f(s) = a_0 + a_1 s + a_2 s^2 + a_3 s^3$$
 Eq. C15

$$f'(s) = a_1 + 2a_2s + 3a_3s^2$$
 Eq. C16

Let

$$f(0) = a_0 = f_0$$
 Eq. C17

$$f'(0) = a_1 = g_0$$
 Eq. C18

And

$$f(1) = a_0 + a_1 + a_2 + a_3 = f_1$$
 Eq. C19

$$f'(1) = a_1 + 2a_2 + 3a_3 = g_1$$
 Eq. C20

Note that i) $f_0 \& f_1$ are displacement at s = 0 & 1

ii) g_0 and g_1 are slope at s = 0 & 1

Our aim is to express the cubic polynomial coefficients in terms of slope and deflection at the end nodes

Now solving for a₂ and a₃, rewriting Eq. E19 and Eq. E20

$$a_2 + a_3 = f_1 - f_0 - g_0$$
 Eq. C21

$$2a_2 + 3a_3 = g_1 - g_0$$
 Eq. C22

Solving Eq. E21 and E1. E22 for a₂ and a₃

$$a_2 = 3f_1 - 3f_0 - 2g_0 - g_1$$
 Eq. C23

$$a_3 = 2f_0 - 2f_1 + g_0 + g_1$$
 Eq. C24

Using Eq. E23 and Eq. E24 we can write the cubic polynomial equation in terms of

$$f(s) = f_0 + g_0 s + (3f_1 - 3f_0 - 2g_0 - g_1)s^2 + (2f_0 - 2f_1 + g_0)$$
 Eq. C25
+ g_1) s^3

We can re-arrange Eq. E25 to

$$f(s) = (1 - 3s^2 + 2s^3)f_0 + (s - 2s^2 + s^3)g_0 + (3s^2 - 2s^3)f_1$$
 Eq. C26
$$+ (-s^2 + s^3)g_1$$

The above equation is a cubic polynomial in interval [0,1] in terms of end displacement and rotation

Now using Eq. E26, we can write a generic cubic polynomial in interval [0,L], where s=x/l and g_0 , g_1 becomes g_0L , g_1L

The reason why g_0 , g_1 becomes g_0L , g_1L is

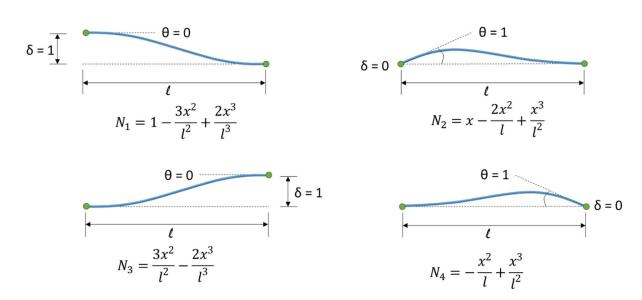
$$f(s) = a_0 + a_1(x/l) + a_2(x/l)^2 + a_3(x/l)^3$$

$$f'(s) = a_1(1/l) + a_2(2/l)(x/l) + a_3(3/l)(x/l)^2$$

$$l * f'(s) = a_1 + 2a_2(x/l) + 3a_3(x/l)^2$$

The generic cubic polynomial in interval [0,L] with end displacement

$$f(x) = \left(1 - \frac{3x^2}{l^2} + \frac{2x^3}{l^3}\right) f_0 + \left(x - \frac{2x^2}{l} + \frac{x^3}{l^2}\right) g_0 + \left(\frac{3x^2}{l^2} - \frac{2x^3}{l^3}\right) f_1$$
 Eq. C27
$$+ \left(-\frac{x^2}{l} + \frac{x^3}{l^2}\right) g_1$$



Strain energy, using Eq. E1

$$dU = \frac{1}{2}\sigma\epsilon \ dV$$

Internal strain energy under pure bending is

$$W_I = \int\limits_V \frac{1}{2} \sigma_{xx} \epsilon_{xx} \, dV$$

Eq. C28

From Euler – Bernoulli the normal stress on a cross- sectional element of area dA at a distance y from neutral axis is

$$\sigma_{xx} = \frac{My}{I}$$

And

$$\epsilon_{xx} = \frac{\sigma_{xx}}{E}$$

Using the above equations in Eq. E.28

$$W_I = \int\limits_V \frac{1}{2} \left(\frac{M^2 y^2}{EI^2}\right) dV$$
 Eq. C29

$$W_{I} = \int_{l} \frac{1}{2} \left(\frac{M^{2}}{EI^{2}}\right) \iint_{A} y^{2} dA dx$$

$$W_{I} = \int_{l} \frac{1}{2} \left(\frac{M^{2}}{EI}\right) dx$$
Eq. C30

Once again using Euler-Bernoulli's beam theory

$$EI\left(\frac{d^2\delta}{dx^2}\right) = M$$

Eq. E30 becomes

$$W_{I} = \int_{I} \frac{1}{2} EI \left(\frac{\partial^{2} \delta}{\partial x^{2}} \right)^{2} dx$$
 Eq. C31

Remember, the displacement $\delta(x)$ is unknown and we use the shape function and end displacement to represent it

Now the external work done is

$$W_E = F_1 \delta_1 + M_1 \theta_1 + F_2 \delta_2 + M_2 \theta_2$$
 Eq. C32

For minimum potential energy at the final state

$$\partial \int_{I} \frac{1}{2} EI \left(\frac{\partial^{2} \delta}{\partial x^{2}} \right)^{2} dx = \partial (F_{1} \delta_{1} + M_{1} \theta_{1} + F_{2} \delta_{2} + M_{2} \theta_{2})$$

The above equation can be written in short form as shown below

$$\frac{1}{2}EI \ \partial \int_{l} [\delta][N'']^{T}[N''][\delta]^{T} dx = \partial [\delta][q]^{T}$$

$$EI \ \frac{\partial [\delta]}{\partial [\delta]} \int_{I} [N'']^{T} [N''] [\delta]^{T} dx = \frac{\partial [\delta]}{\partial [\delta]} [q]^{T}$$

$$EI \int_{l} [N'']^{T} [N''] dx \begin{bmatrix} \delta_{1} \\ \theta_{1} \\ \delta_{2} \\ \theta_{3} \end{bmatrix} = \begin{bmatrix} F_{1} \\ M_{1} \\ F_{2} \\ M_{2} \end{bmatrix}$$
 Eq. C34

The beam element stiffness matrix is

$$k = EI \int_{l} [N'']^{T} [N''] dx$$
 Eq. C35

The below table shows the double derivatives of the beam shape functions Eq. E27

	N	N'	N''
1	$1 - \frac{3x^2}{l^2} + \frac{2x^3}{l^3}$	$-\frac{6x}{l^2} + \frac{6x^2}{l^3}$	$-\frac{6}{l^2} + \frac{12x}{l^3}$
2	$x - \frac{2x^2}{l} + \frac{x^3}{l^2}$	$1 - \frac{4x}{l} + \frac{3x^2}{l^2}$	$-\frac{4}{l} + \frac{6x}{l^2}$
3	$\frac{3x^2}{l^2} - \frac{2x^3}{l^3}$	$\frac{6x}{l^2} - \frac{6x^2}{l^3}$	$-\left(-\frac{6}{l^2} + \frac{12x}{l^3}\right)$
4	$-\frac{x^2}{l} + \frac{x^3}{l^2}$	$-\frac{2x}{l} + \frac{3x^2}{l^2}$	$-\frac{2}{l} + \frac{6x}{l^2}$

Multiplying $[N^{\prime\prime}]^T[N^{\prime\prime}]$ yields

$$\int\limits_{l} \begin{bmatrix} N_{1}'' \\ N_{2}'' \\ N_{3}'' \\ N_{4}'' \end{bmatrix} [N_{1}'' \quad N_{2}'' \quad N_{3}'' \quad N_{4}''] = \int\limits_{l} \begin{bmatrix} N_{1}''^{2} & N_{1}''N_{2}'' & N_{1}''N_{3}'' & N_{1}''N_{4}'' \\ - & N_{2}''^{2} & N_{2}''N_{3}'' & N_{2}''N_{4}'' \\ - & - & N_{3}''^{2} & N_{3}''N_{4}'' \\ - & - & - & N_{4}''^{2} \end{bmatrix}$$

The lower triangle on the matrix is symmetrical to the upper triangle

Further simplification gives

$$\int_{l} \begin{bmatrix} N_{1}^{\prime\prime\prime2} & N_{1}^{\prime\prime\prime}N_{2}^{\prime\prime} & -N_{1}^{\prime\prime\prime2} & N_{1}^{\prime\prime}N_{4}^{\prime\prime} \\ - & N_{2}^{\prime\prime\prime2} & -N_{1}^{\prime\prime\prime}N_{2}^{\prime\prime} & N_{2}^{\prime\prime}N_{4}^{\prime\prime} \\ - & - & N_{1}^{\prime\prime\prime2} & -N_{1}^{\prime\prime\prime}N_{4}^{\prime\prime} \\ - & - & N_{4}^{\prime\prime\prime2} \end{bmatrix}$$

Solving individual elements of the matrix We need to find the elements (1,1) (1,2) (1,4) (2,2) (2,4),(4,4)

$$\int_{0}^{l} N_{1}^{"2} dx = \int_{0}^{l} \left(-\frac{6}{l^{2}} + \frac{12x}{l^{3}} \right) \left(-\frac{6}{l^{2}} + \frac{12x}{l^{3}} \right) dx$$

$$= \int_{0}^{l} \left(\frac{36}{l^{4}} - \frac{144x}{l^{5}} + \frac{144x^{2}}{L^{6}} \right) dx$$

$$= \frac{36x}{l^{4}} - \frac{144x^{2}}{2l^{5}} + \frac{144x^{3}}{3L^{6}} \Big|_{0}^{l}$$

$$= \frac{36}{l^{3}} - \frac{72}{l^{3}} + \frac{48}{L^{3}} = \frac{12}{l^{3}}$$

$$\int_{0}^{l} N_{1}^{"2} dx = \frac{12}{l^{3}}$$

Eq. C36

Eq. C37

Matrix element - (1,2)

$$\int_{0}^{l} N_{1}'' N_{2}'' dx = \int_{0}^{l} \left(-\frac{6}{l^{2}} + \frac{12x}{l^{3}} \right) \left(-\frac{4}{l} + \frac{6x}{l^{2}} \right) dx$$

$$= \int_{0}^{l} \left(\frac{72x^{2}}{l^{5}} - \frac{48x}{l^{4}} - \frac{36x}{L^{4}} + \frac{24}{L^{3}} \right) dx$$

$$= \frac{72x^{3}}{3l^{5}} - \frac{48x^{2}}{2l^{4}} - \frac{36x^{2}}{2L^{4}} + \frac{24x}{L^{3}} \Big|_{0}^{l}$$

$$= \frac{24}{l^{2}} - \frac{24}{l^{2}} - \frac{18}{L^{2}} + \frac{24}{l^{2}} = \frac{6}{l^{2}}$$

$$\int N_{1}'' N_{2}'' dx = \frac{6}{l^{2}}$$

Matrix element - (1,4)

$$\int_{0}^{l} N_{1}^{"}N_{4}^{"}dx = \int_{0}^{l} \left(-\frac{6}{l^{2}} + \frac{12x}{l^{3}}\right) \left(-\frac{2}{l} + \frac{6x}{l^{2}}\right) dx$$

$$= \int_{0}^{l} \left(\frac{72x^{2}}{l^{5}} - \frac{24x}{l^{4}} - \frac{36x}{L^{4}} + \frac{12}{L^{3}}\right) dx$$

$$= \frac{72x^{3}}{3l^{5}} - \frac{24x^{2}}{2l^{4}} - \frac{36x^{2}}{2L^{4}} + \frac{12x}{L^{3}} \Big|_{0}^{l}$$

$$= \frac{24}{l^{2}} - \frac{12}{l^{2}} - \frac{18}{L^{2}} + \frac{12}{l^{2}} = \frac{6}{l^{2}}$$

$$\int_{l}^{l} N_{1}^{"}N_{4}^{"}dx = \frac{6}{l^{2}}$$
Eq. C38

Matrix element - (2,2)

$$\int_{0}^{l} N_{2}^{"2} dx = \int_{0}^{l} \left(-\frac{4}{l} + \frac{6x}{l^{2}} \right) \left(-\frac{4}{l} + \frac{6x}{l^{2}} \right) dx$$

$$= \int_{0}^{l} \left(\frac{36x^{2}}{l^{4}} - \frac{48x}{l^{3}} + \frac{16}{L^{2}} \right) dx$$

$$= \frac{36x^{3}}{3l^{4}} - \frac{48x^{2}}{2l^{3}} + \frac{16x}{L^{2}} \Big|_{0}^{l}$$

$$= \frac{12}{l} - \frac{24}{l} + \frac{16}{L} = \frac{4}{l}$$

$$\int_{0}^{l} N_{2}^{"2} dx = \frac{4}{l}$$
Eq. C39

Matrix element - (2,4)

$$\int_{0}^{l} N_{2}^{"}N_{4}^{"}dx = \int_{0}^{l} \left(-\frac{4}{l} + \frac{6x}{l^{2}} \right) \left(-\frac{2}{l} + \frac{6x}{l^{2}} \right) dx$$

$$= \int_{0}^{l} \left(\frac{36x^{2}}{l^{4}} - \frac{24x}{l^{3}} - \frac{12x}{L^{3}} + \frac{8}{L^{2}} \right) dx$$

$$= \frac{36x^{3}}{3l^{4}} - \frac{24x^{2}}{2l^{3}} - \frac{12x^{2}}{2L^{3}} + \frac{8x}{L^{2}} \Big|_{0}^{l}$$

$$= \frac{12}{l} - \frac{12}{l} - \frac{6}{l} + \frac{8}{l} = \frac{2}{l}$$

$$\int_{0}^{l} N_{2}^{\prime\prime} N_{4}^{\prime\prime} dx = \frac{2}{l}$$
 Eq. C40

Matrix element - (4,4)

$$\int_{0}^{l} N_{4}^{"2} dx = \int_{0}^{l} \left(-\frac{2}{l} + \frac{6x}{l^{2}} \right) \left(-\frac{2}{l} + \frac{6x}{l^{2}} \right) dx$$

$$= \int_{0}^{l} \left(\frac{36x^{2}}{l^{4}} - \frac{24x}{l^{3}} + \frac{4}{L^{2}} \right) dx$$

$$= \frac{36x^{3}}{3l^{4}} - \frac{24x^{2}}{2l^{3}} + \frac{4x}{L^{2}} \Big|_{0}^{l}$$

$$= \frac{12}{l} - \frac{12}{l} + \frac{4}{L} = \frac{4}{l}$$

$$\int_{0}^{l} N_{4}^{"2} dx = \frac{4}{l}$$
Eq. C41

Now the element stiffness matrix by substituting Eq. E36, E37, E38, E39, E40, E41 in Eq. E35 gives us

$$EI\begin{bmatrix} \frac{12}{l^3} & \frac{6}{l^2} & -\frac{12}{l^3} & \frac{6}{l^2} \\ \frac{6}{l^2} & \frac{4}{l} & -\frac{6}{l^2} & \frac{2}{l} \\ -\frac{12}{l^3} & -\frac{6}{l^2} & \frac{12}{l^3} & -\frac{6}{l^2} \\ \frac{6}{l^2} & \frac{2}{l} & -\frac{6}{l^2} & \frac{4}{l} \end{bmatrix} \begin{bmatrix} \delta_i \\ \theta_i \\ \delta_j \\ \theta_j \end{bmatrix} = \begin{bmatrix} F_i \\ M_i \\ F_j \\ M_j \end{bmatrix}$$
 Eq. C42

Note that the above equation is same as that of Eq. B37

$$\begin{pmatrix} \frac{EI}{l^3} \end{pmatrix} \begin{bmatrix} 12 & 6l & -12 & 6l \\ 6l & 4l^2 & -6l & 2l^2 \\ -12 & -6l & 12 & -6l \\ 6l & 2l^2 & -6l & 4l^2 \end{bmatrix} \begin{bmatrix} \delta_i \\ \theta_i \\ \delta_j \\ \theta_i \end{bmatrix} = \begin{bmatrix} F_i \\ M_i \\ F_j \\ M_j \end{bmatrix}$$
 Eq. C43 & Eq. B38

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