

Whaley group meeting talk

Liwen Ko
09/17/2024

Outline

- Tutorial on input-output relation (board, 30 min)
- Analyzing optical signals (board, 3 min)
- Normal-ordered expansion of the reduced system state (slides, 5 min)
- Electron transfer assisted by off-resonant underdamped vibrations (slides, 7 min)
- Some open questions. My disorganized and speculative thoughts (board, 7 min)
- Conclusion (slides, 3 min)

Normal-ordered expansion for the reduced system state

Conventional time-ordered

$$\begin{aligned} \rho_{\text{sys}}(t) = \text{Tr}_{\text{field}} \bigg(& e^{\mathcal{K}'t} \rho_{\text{tot}}(0) \\ & + \int_0^t dt_1 e^{\mathcal{K}'(t-t_1)} \mathcal{L}(t_1) e^{\mathcal{K}'t_1} \rho_{\text{tot}}(0) \\ & + \int_0^t dt_2 \int_0^{t_2} dt_1 e^{\mathcal{K}'(t-t_2)} \mathcal{L}(t_2) e^{\mathcal{K}'(t_2-t_1)} \mathcal{L}(t_1) e^{\mathcal{K}'t_1} \rho_{\text{tot}}(0) \\ & + \dots \bigg). \end{aligned}$$

$$\mathcal{L}(t) = [-a(t)L^\dagger + a^\dagger(t)L, \bullet]$$

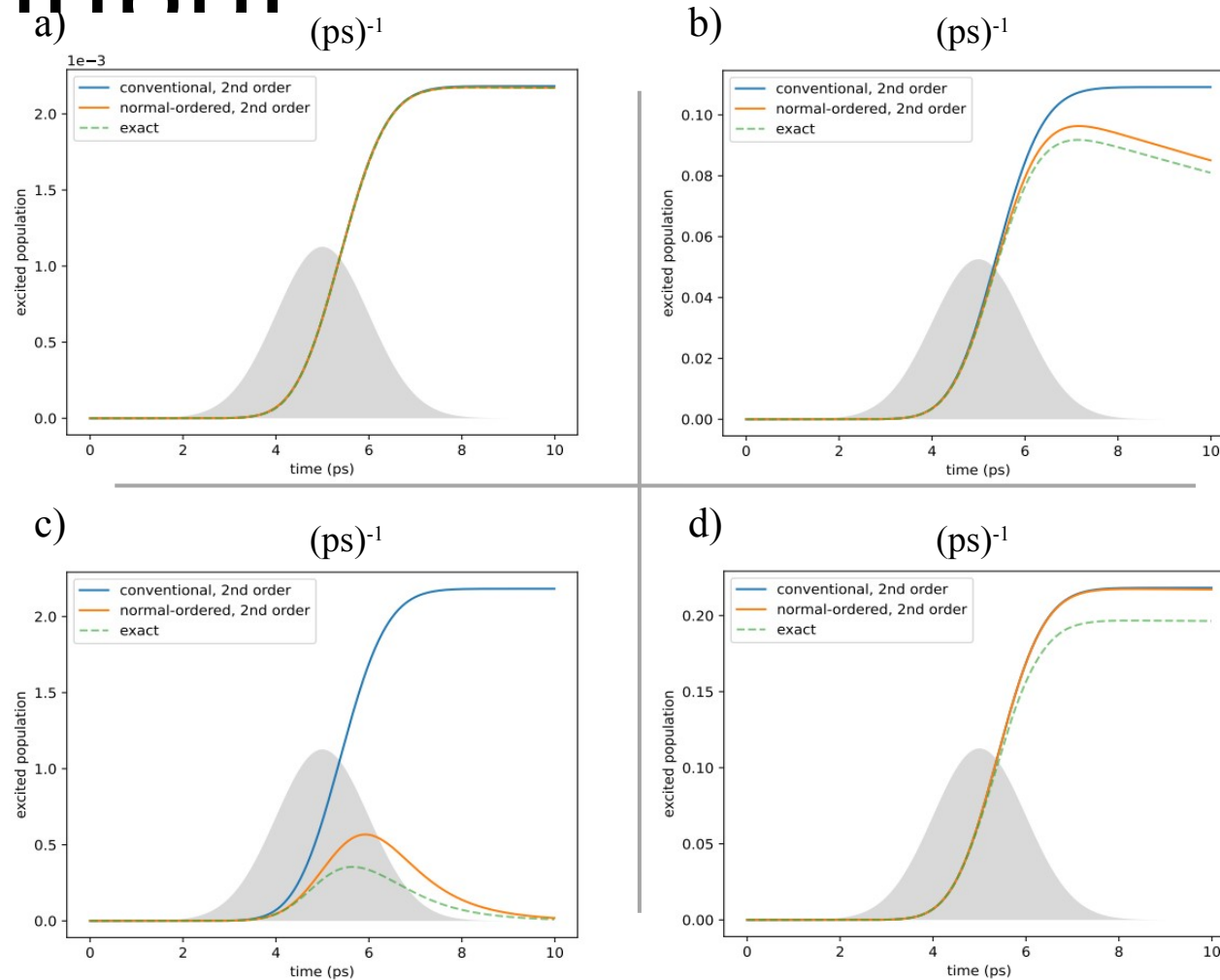
$$\mathcal{K}' = -i[H_{\text{sys}}, \bullet]$$

Normal-ordered

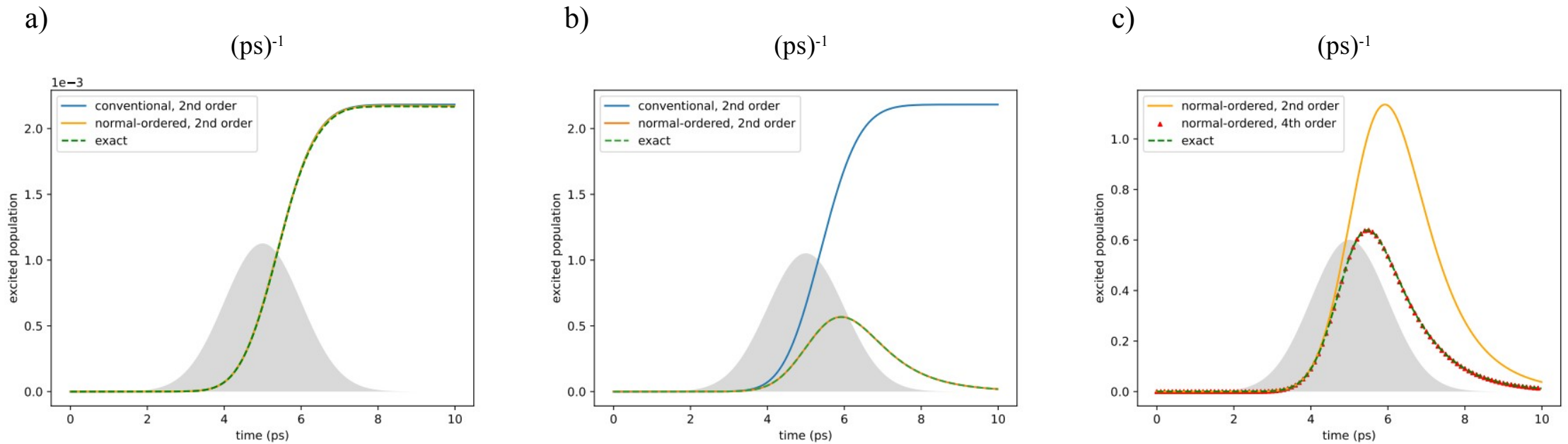
$$\begin{aligned} \rho_{\text{sys}}(t) = \text{Tr}_{\text{field}} \hat{\mathcal{N}} \bigg(& e^{\mathcal{K}t} \rho_{\text{tot}}(0) \\ & + \int_0^t dt_1 e^{\mathcal{K}(t-t_1)} \mathcal{L}(t_1) e^{\mathcal{K}t_1} \rho_{\text{tot}}(0) \\ & + \int_0^t dt_2 \int_0^{t_2} dt_1 e^{\mathcal{K}(t-t_2)} \mathcal{L}(t_2) e^{\mathcal{K}(t_2-t_1)} \mathcal{L}(t_1) e^{\mathcal{K}t_1} \rho_{\text{tot}}(0) \\ & + \dots \bigg). \end{aligned}$$

$$\mathcal{K} = -i[H_{\text{sys}}, \bullet] - \frac{1}{2}\{L^\dagger L \bullet\} + L \bullet L^\dagger$$

Normal-ordered expansion for the reduced system state – coherent state init



Normal-ordered expansion for the reduced system state – m-photon Fock state input

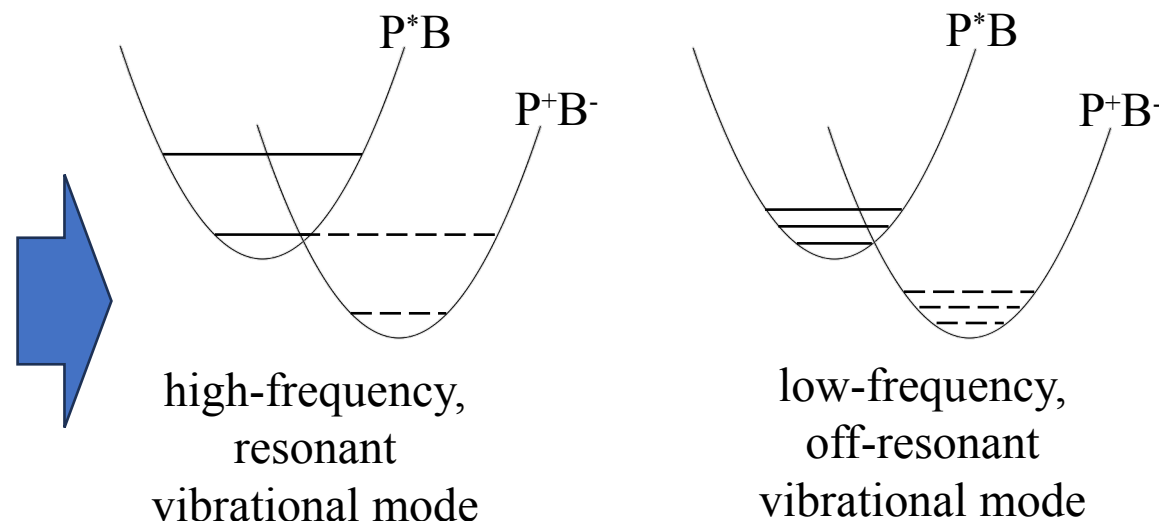
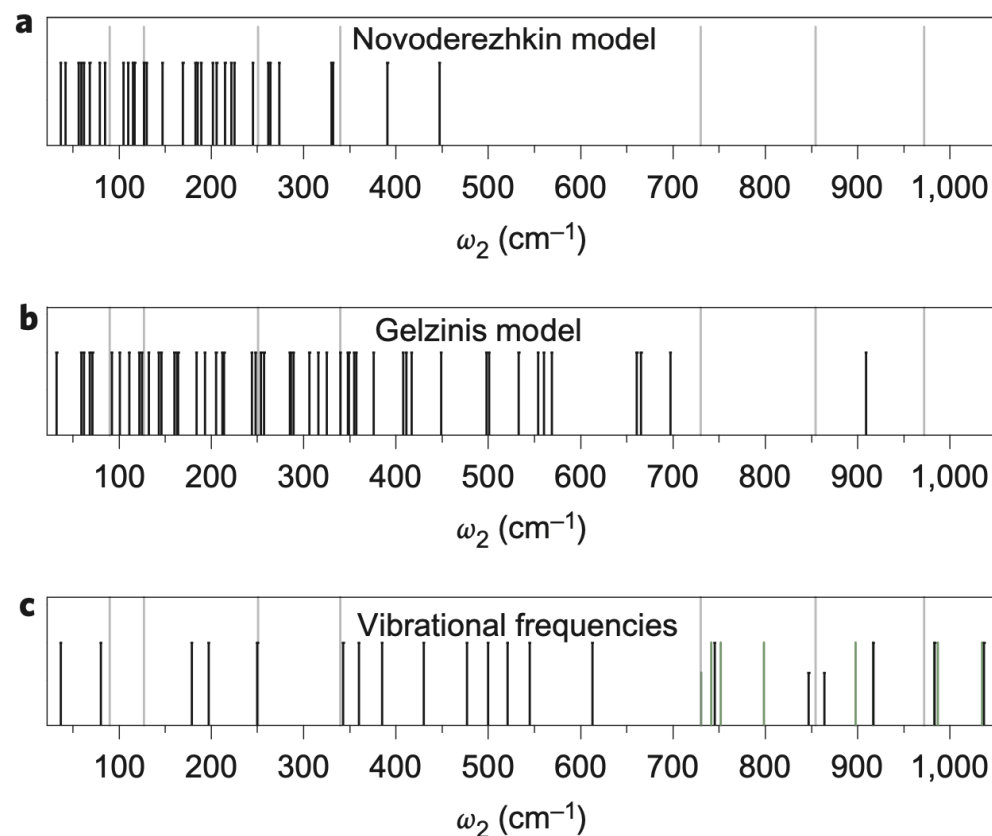


Truncates at $2m$ -th order \neq exact expression

Outline

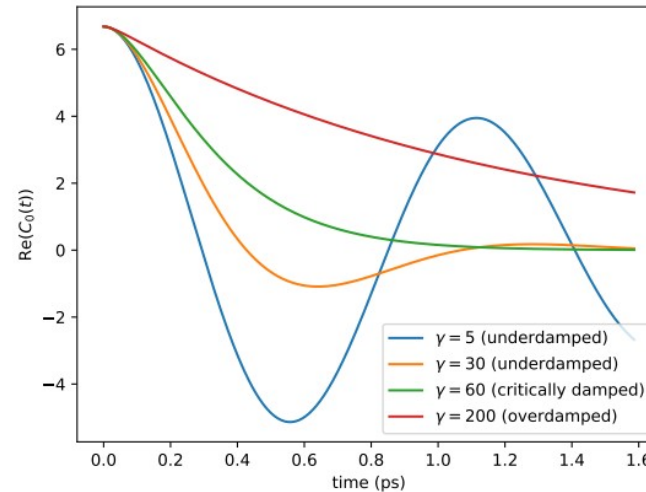
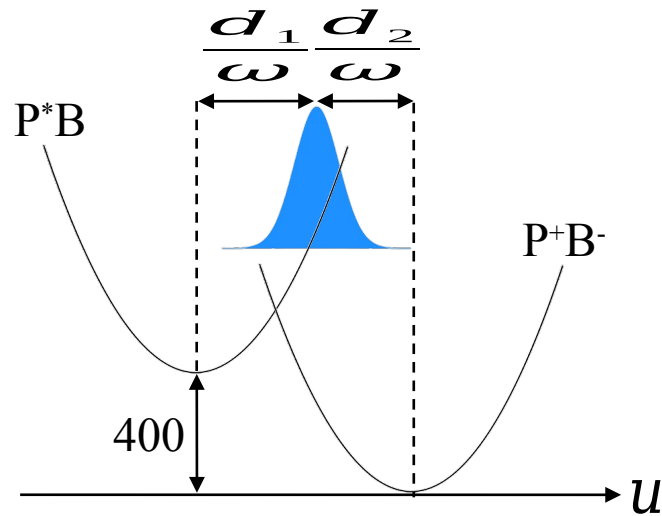
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Electron transfer assisted by underdamped nuclear vibrations



Low-frequency, off-resonant, underdamped vibrational modes can assist electron transfer

Modeling the effect of an underdamped mode



Long vibrational coherence time
 → Memory effect
 → Non-Markovian equations of motion

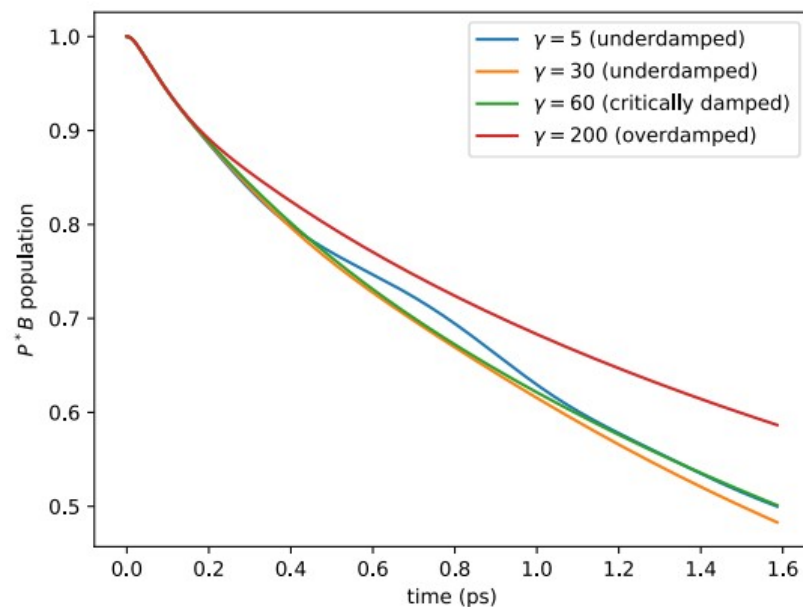
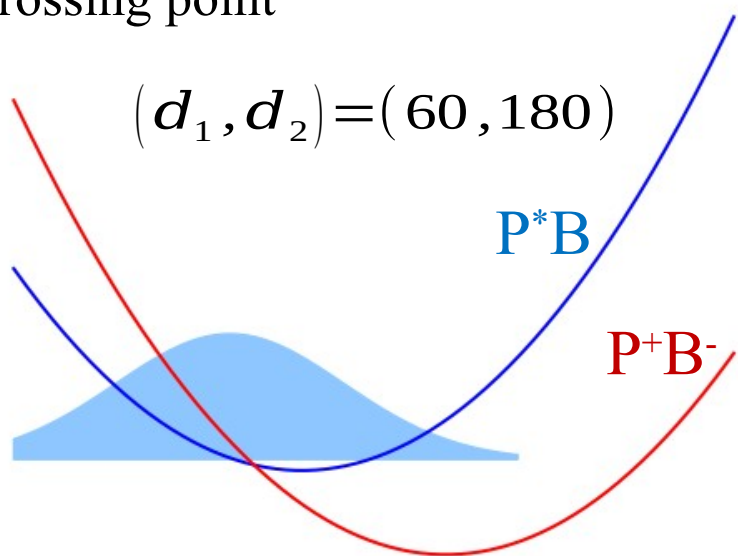
Hierarchical equations of motion (HEOM):

$$\begin{aligned} \frac{d}{dt} \rho^{n,m} = & -i H^\times \rho^{n,m} - \frac{\gamma}{2} (n+m) \rho^{n,m} - S^\times \rho^{n+1,m} \\ & + n c_0(0) S^\times \rho^{n-1,m} + m \left(\text{Re}(c_1(0)) S^\times + i \text{Im}(c_1(0)) S^o \right) \rho^{n,m-1} \\ & + n k_{01} \rho^{n-1,m+1} + m k_{10} \rho^{n+1,m-1} \end{aligned}$$

Damping effects

Set cm^{-1} (i.e., off-resonant, low-frequency) and vary γ . **Underdamped** modes assist the electron transfer the best. Damping effects depend strongly on the position of the initial nuclear position.

Configuration 1: initial nuclear position on the same side as the crossing point

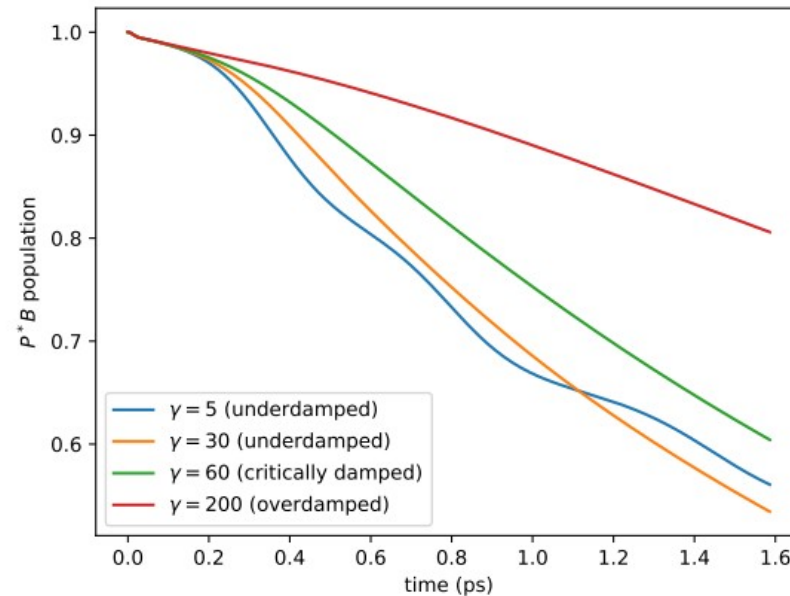
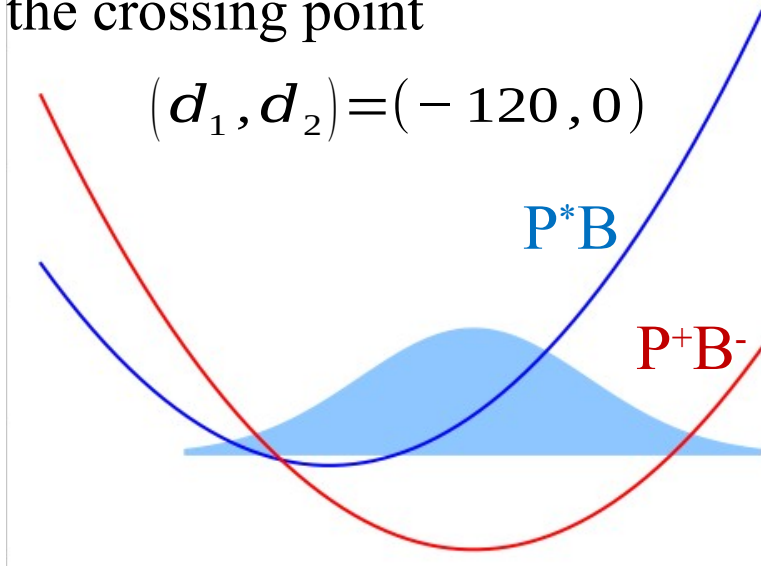


Damping effects

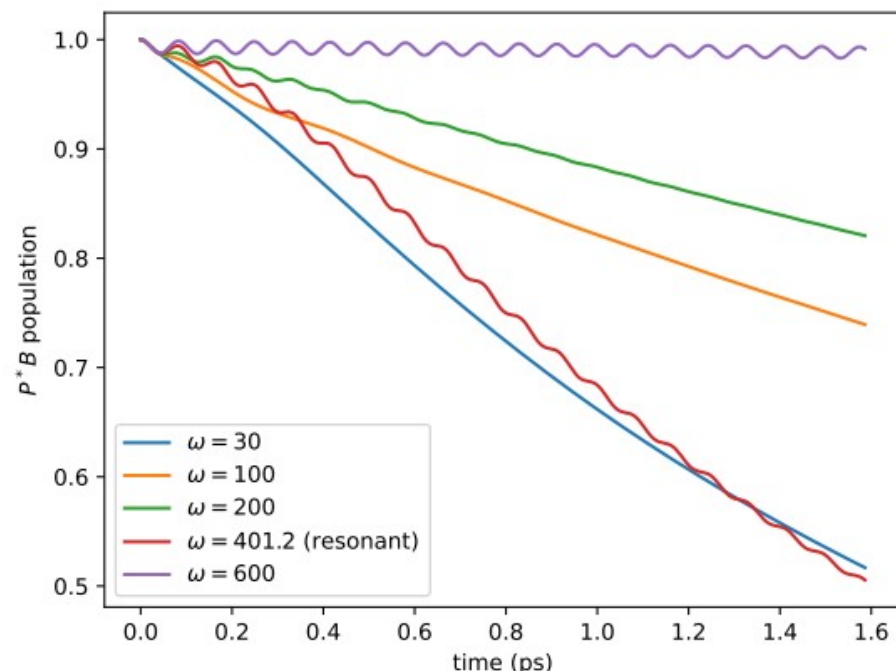
Set cm^{-1} (i.e., off-resonant, low-frequency) and vary γ . **Underdamped** modes assist the electron transfer the best. Damping effects depend strongly on the position of the initial nuclear position.

Configuration 2: initial nuclear position on the opposite side of the crossing point

$$(d_1, d_2) = (-120, 0)$$



Frequency effects



Set cm^{-1} and vary . Electron transfer is enhanced when the mode frequency is **resonant** or **off-resonant at low frequency**.

Low-frequency, off-resonant, underdamped vibrational modes can assist electron transfer

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