Simulations with MPS/MPO

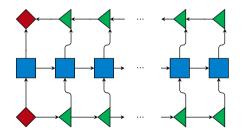
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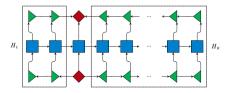
September 20, 2024

DMRG: 1-site update

Initialize MPS with diagonalization center at i = 1.

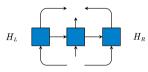


- Sweep at two direction (right left right ...)
 - Calculate the left/right environment H_L/H_R .

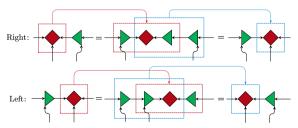


DMRG: 1-site update

- Sweep at two direction (right left right ...)
 - igcap Calculate the effective Hamiltonian $H_{eff}=H_LH_iH_R$

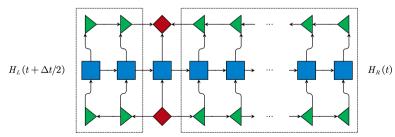


OrientSVD and move the center to next site.



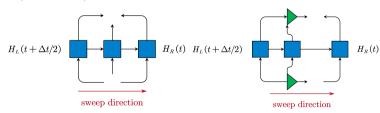
TDVP: 1-site integration

- Initialize MPS with diagonalization center at i = 1.
- Sweep at two direction (right left right ...)
 - Calculate the left/right environment $H_L(t + \Delta t/2)/H_R(t)$.



TDVP: 1-site integration

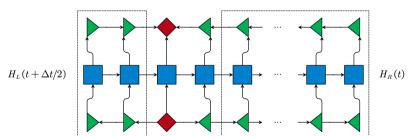
- Sweep at two direction (right left right ...)
 - Calculate the effective Hamiltonian $H_{eff}^{(1)}$
 - Time evolution $A_i(t + \Delta t/2) = \exp(-iH_{eff}^{(1)}\Delta t/2)A_i(t)$
 - OrientSVD and calculate the center with inverse evolution $C_i(t) = \exp\left(iH_{eff}^{(0)}\Delta t/2\right)C_i(t+\Delta t/2)$, then absorb it into nextsite.



TDVP: 2-site integration

Sweep schemes

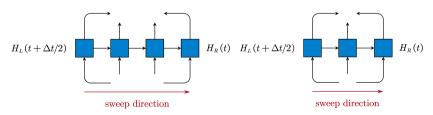
Calculate the left/right environment $H_L(t + \Delta t/2)/H_R(t)$.



TDVP: 2-site integration

Sweep schemes

- lacksquare Calculate the effective Hamiltonian $H_{eff}^{(2)}$
- Time evolution $A_i A_{i+1}(t + \Delta t/2) = \exp\left(-iH_{eff}^{(2)} \Delta t/2\right) A_i A_{i+1}(t)$
- OrientSVD and calculate the center with inverse evolution $A_{i+1}(t) = \exp\left(iH_{eff}^{(1)}\Delta t/2\right)A_{i+1}(t+\Delta t/2)$, then regard it as nextsite.



SETTN

Write ρ in series expansion:

$$\rho = e^{-\beta H} = \sum_{k=0}^{+\infty} \frac{(-\beta H)^k}{k!} \tag{1}$$

calculate the $\{H^k\}$, k = 0, 1, 2...

tanTRG

Suppose $X_{\rho} \in T_{\rho}M$

$$\frac{\mathrm{d}\rho}{\mathrm{d}\beta} = \arg\min\left|X_{\rho} + H\rho\right| \tag{2}$$

i.e., project the H
ho onto the tangent space of M at ho



Model: Transverse Ising Chain

Hamiltonian:

$$H = -J \sum_{i=1}^{N-1} \sigma_i^z \sigma_{i+1}^z - h \sum_{i=1}^{N} \sigma_{i=1}^x + h_z \sum_{i=1}^{N} \sigma_i^z$$
 (3)

we always choose $J = \pm 1.0$.

Exact Solutions: With **Jordan-Wigne** and **Bogoliubov** transformation, we can diagonalize *H* in fermion picture:

$$H = \sum_{k>0} \lambda(k) \gamma_k^{\dagger} \gamma_k, \quad \lambda(k) = \sqrt{\epsilon_k^2 + \Delta_k^2}, \quad \begin{pmatrix} \gamma_k^{\dagger} \\ \gamma_{-k} \end{pmatrix} = \mathbf{P}(k) \begin{pmatrix} c_k^{\dagger} \\ c_{-k} \end{pmatrix}$$
(4)

$$\epsilon_k = -2J\cos ka - 2h, \quad \Delta_k = 2J\sin ka$$
 (5)

Model:MPO

Hamiltonian: In kron form:

$$H = J\sigma_1^z \otimes \sigma_2^z \otimes \mathbb{1} \otimes \mathbb{1} \otimes \cdots + \mathbb{1} \otimes J\sigma_2^z \otimes \sigma_3^z \otimes \mathbb{1} \otimes \cdots + h\sigma_1^x \otimes \mathbb{1} \otimes \mathbb{1} \otimes \cdots + \mathbb{1} \otimes h\sigma_2^x \otimes \mathbb{1} \otimes \cdots + h_z\sigma_1^z \otimes \mathbb{1} \otimes \mathbb{1} \otimes \cdots + \mathbb{1} \otimes h_z\sigma_2^z \otimes \mathbb{1} \otimes \cdots$$

$$(6)$$

▶ MPO

$$H_{1} = \begin{pmatrix} h\sigma_{1}^{x} + h_{z}\sigma_{1}^{z}, & J\sigma_{1}^{z}, & 1 \end{pmatrix},$$

$$H_{i} = \begin{pmatrix} 1, & 0, & 0 \\ \sigma_{i}^{z}, & 0, & 0 \\ h\sigma_{i}^{x} + h_{z}\sigma_{i}^{z}, & J\sigma_{i}^{z}, & 1 \end{pmatrix}, \quad H_{N} = \begin{pmatrix} 1 \\ \sigma_{N}^{z} \\ h\sigma_{N}^{x} + h_{z}\sigma_{N}^{z} \end{pmatrix}$$

$$(7)$$

Model:MPO

Total Magnetic Moment: In kron form:

$$M_z = \sigma_1^z \otimes \mathbb{1} \otimes \mathbb{1} \otimes \cdots + \mathbb{1} \otimes \sigma_2^z \otimes \mathbb{1} \otimes \cdots$$
 (8)

► MPO

$$H_{1} = \begin{pmatrix} \sigma_{1}^{z}, & \mathbb{1} \end{pmatrix}, H_{i} = \begin{pmatrix} \mathbb{1}, & \mathbb{0} \\ \sigma_{i}^{z}, & \mathbb{1} \end{pmatrix}, \quad H_{N} = \begin{pmatrix} \mathbb{1} \\ \sigma_{N}^{z} \end{pmatrix}$$

$$(9)$$

Model:MPS

 $|s_i\rangle$ denote that site i is in eigenstate of σ_i^z with eigenvalue $s_i=\pm 1$. $|s_1s_2s_3\cdots s_N\rangle$ is corresponding many body basis of system.

Random Initialized State: Random initial state

$$\psi_0 = \sum_{\{s_i\}} C_{s_1 s_2 s_3 \cdots s_N} | s_1 s_2 s_3 \cdots s_N \rangle$$
 (10)

equals the random tensor $C_{s_1s_2s_3\cdots s_N}$. Then orientSVD:



Model:MPS

Assigned Initialized State FM corresponds to state $C_{s_1s_2s_3\cdots s_N}$ with

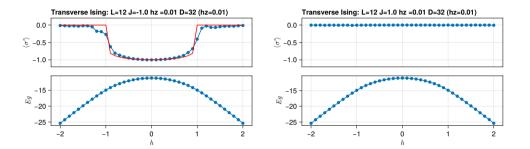
$$C_{111\dots 1} = 1 \quad \text{or} \quad C_{-1-1\dots -1} = 1$$
 (11)

while AFM corresponds to state $C_{s_1s_2s_3\cdots s_N}$ with

$$C_{1-11-1...} = 1$$
 or $C_{-11-11...} = 1$ (12)

Then orientSVD in the same way as that in Random Initialized State.

Simulation results:



Where exact solution can be obtained by calculating $\left\langle \sigma_i^z \sigma_j^z \right\rangle$ and let $|i-j| \to +\infty$, i.e.

$$\langle \sigma^z \rangle = \left[1 - \left(\frac{h}{J} \right)^2 \right]^{1/8} \tag{13}$$

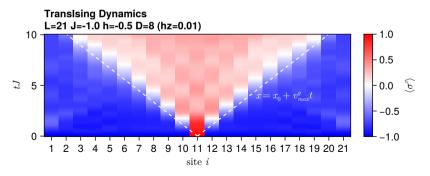


Results: Dynamics

Choose initial state relevant is

$$\Psi_0 = |-1, -1, \dots, -1, 1, -1, \dots, -1, -1\rangle \tag{14}$$

Simulation results: $v_g^{max} = \max \partial \lambda(k)/\partial k$



Model: Heisenberg Chain

Hamiltonian:

$$H = -J \sum_{i=1}^{N-1} \sigma_i \cdot \sigma_{i+1}$$
 (15)

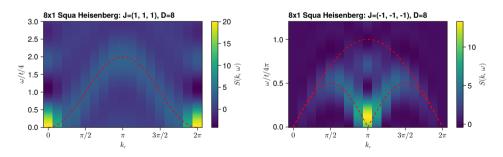
we always choose $J = \pm 1.0$.

Results: Dynamical Spin Structure Factor

Set

$$J = \pm 1.0, \quad S(\mathbf{q}, \omega) = \int_{-\infty}^{+\infty} \left\langle \sigma_{\mathbf{q}}(t) \cdot \sigma_{-\mathbf{q}}(0) \right\rangle e^{i\omega t} dt$$
 (16)

► Simulation results:



Model: Free Fermion on Square Lattice

Hamiltonian:

$$H = t \sum_{\langle i,j \rangle} c_i^{\dagger} c_j + c_j^{\dagger} c_i \tag{17}$$

Parameters: Physics parameters:

$$t = 1 \tag{18}$$

▶ J-W transformation:

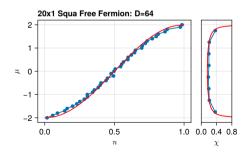
$$c_i = \prod_{j=1}^{i-1} F_j b_i, \quad c_i^{\dagger} = \prod_{j=1}^{i-1} F_j b_i^{\dagger}$$
 (19)

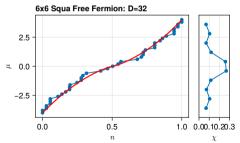
$$F_j = -\sigma_j^z, \quad b_i = \sigma_i^-, \quad b_i = \sigma_i^+ \tag{20}$$

Charge compressibility.

$$\chi = \frac{1}{N} \frac{\partial N}{\partial \mu} \tag{21}$$

Simulation results:

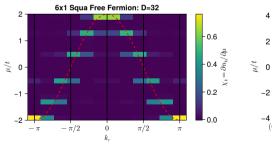


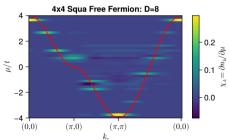


 $\triangleright k$ dependent Charge compressibility.

$$\chi_k = \frac{1}{N} \frac{\partial N_k}{\partial \mu} \tag{22}$$

Simulation results:





Obserables

► Spectrum function

$$S(\mathbf{k}, E) = -\frac{1}{\pi} \text{Im} G^{ret}(\mathbf{k}, E), \quad G^{ret}(\mathbf{k}, E) = \int_{-\infty}^{+\infty} G^{ret}(\mathbf{k}, t) e^{iEt} dt \quad (23)$$

$$G^{ret}(\mathbf{k}, t) = -i\Theta(t) \langle \Psi_0 | \left[c_k(t) c_k^{\dagger}(0) \right] | \Psi_0 \rangle$$

$$= -i \left[e^{iE_0 t} \left(\langle \Psi_0 | c_k \rangle e^{-iHt} \left(c_k^{\dagger} | \Psi_0 \rangle \right) \right) \quad (24)$$

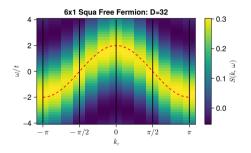
$$+ e^{-iE_0 t} \left(\langle \Psi_0 | c_k^{\dagger} \rangle e^{iHt} \left(c_k | \Psi_0 \rangle \right) \right]$$

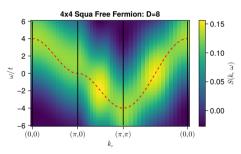
with

$$c_k = \sum_j e^{-i\mathbf{k}\cdot\mathbf{R}_j} c_j \tag{25}$$

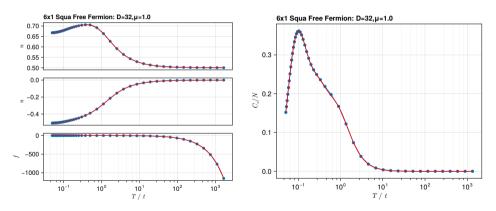


► Simulation Results:





Finite Temperature Results:



Model: Hubbard Model on Square Lattice

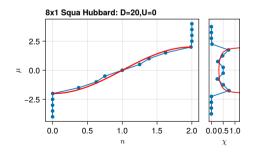
► Hamiltonian:

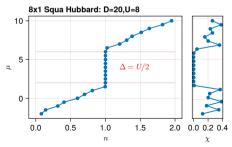
$$H = t \sum_{\langle i,j \rangle} \left(c_i^{\dagger} c_j + c_j^{\dagger} c_i \right) + \frac{U}{2} \sum_{\langle i \rangle} n_{i,\sigma} n_{i,\hat{\sigma}}$$
 (26)

Charge compressibility.

$$\chi = \frac{1}{N} \frac{\partial N}{\partial \mu} \tag{27}$$

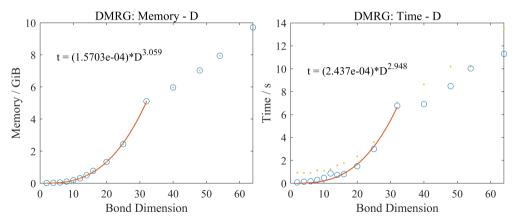
Simulation results:





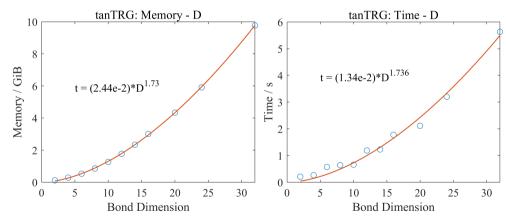
Benchmark

DMRG complexity:



Benchmark

▶ tanTRG + SETTN complexity:



Source code

Can be found at https://github.com/KunyangDU/iMPS.git, where dev branch is the latest version, but is still under developement.