Møller–Plesset Perturbation Theory

Post-HF Methods

- Hartree-Fock often provides a reasonable approximation to energy and electron density in a system
- However, HF is not exact, in particular due to not calculating electron correlation
- Post-HF methods calculate electron correlation using HF energy and wavefunctions as a starting point
- Examples include MP2, CCSD, and CI

Post-HF methods (cont.)

- $E_{Exact} = E_{HF} + E_{corr}$
- Because HF is variational, E_{corr} should always be negative

Perturbation Theory

- Assume total Hamiltonian is equal to:
 - $H = H_o + \lambda V$
 - *H*_o: Unperturbed Hamiltonian
 - V: Perturbation (generally small)
 - λ: Ordering parameter (primarily for bookkeeping)
- Where:
 - $E_i = E_i^{(0)} + \lambda E_i^{(1)} + \lambda^2 E_i^{(2)} + \cdots$
 - $\Psi_i = \Psi_i^{(0)} + \lambda \Psi_i^{(1)} + \lambda^2 \Psi_i^{(2)} + \cdots$

Perturbation Theory (cont.)

- For
 - $H|\Psi_i\rangle = E_i|\Psi_i\rangle$
- Substituting expressions from previous slide:

•
$$(H_o + \lambda V) | \Psi_i^{(0)} + \lambda \Psi_i^{(1)} + \lambda^2 \Psi_i^{(2)} + \cdots \rangle =$$

 $E_i^{(0)} + \lambda E_i^{(1)} + \lambda^2 E_i^{(2)} + \cdots) | \Psi_i^{(0)} + \lambda \Psi_i^{(1)} + \lambda^2 \Psi_i^{(2)} + \cdots \rangle$

Expanding and equating terms with equal order of λ:

•
$$H_{o}|\Psi_{i}^{(0)}\rangle = E^{(0)}|\Psi_{i}^{(0)}\rangle$$

• $H_{o}|\Psi_{i}^{(1)}\rangle + V|\Psi_{i}^{(0)}\rangle = E^{(0)}|\Psi_{i}^{(0)}\rangle + E^{(1)}|\Psi_{i}^{(0)}\rangle$
• $H_{o}|\Psi_{i}^{(2)}\rangle + V|\Psi_{i}^{(1)}\rangle = E^{(0)}|\Psi_{i}^{(2)}\rangle + E^{(1)}|\Psi_{i}^{(1)}\rangle + E^{(2)}|\Psi_{i}^{(0)}\rangle$
• λ^{2}

• ...

Perturbation Theory (cont.)

• Multiplying by $\langle \Psi_i^{(0)} |$ gives the following expressions for energy:

$$\bullet \ E_i^{(0)} = \left\langle \Psi_i^{(0)} \middle| H_o \middle| \Psi_i^{(0)} \right\rangle$$

$$\bullet \ E_i^{(1)} = \left\langle \Psi_i^{(0)} \middle| V \middle| \Psi_i^{(0)} \right\rangle$$

$$\bullet \ E_i^{(2)} = \left\langle \Psi_i^{(0)} \middle| V \middle| \Psi_i^{(1)} \right\rangle$$

• If we take $|\Psi_i^{(1)}\rangle$ to be a linear combination of zeroth order wavefunctions $|\Psi_n^{(0)}\rangle$ ($n\neq i$), $E_i^{(2)}$ can be rewritten:

•
$$E_i^{(2)} = \sum_{n \neq i} \frac{\left| \left\langle \Psi_i^{(0)} \middle| V \middle| \Psi_n^{(0)} \right\rangle \right|^2}{E_i^{(0)} - E_n^{(0)}}$$

Møller–Plesset Perturbation Theory

- In MP2, we use the following Hamiltonian:
 - $H = H_0 + V$
- H_o is the Hartree-Fock Hamiltonian:
 - $H_o = \sum_i H^{core}(i) + v^{HF}(i)$
 - H^{core}: 1 electron integrals
 - *v^{HF}*: Effective electron potential
- *V* is our perturbation:
 - $V = \sum_{i < j} r_{ij}^{-1} \sum_{i} v^{HF}(i)$
 - r_{ij}^{-1} : 2 electron integrals

Zeroth and First Order MP2 Energies

- Zeroth order energy of ground state:
 - $E_0^{(0)} = \left\langle \Psi_0^{(0)} \middle| H_o \middle| \Psi_0^{(0)} \right\rangle = \sum_a \varepsilon_a$
 - ε_a : Energy of occupied orbital a
- First order energy of ground state:

•
$$E_0^{(1)} = \left\langle \Psi_0^{(0)} \middle| V \middle| \Psi_0^{(0)} \right\rangle = \left\langle \Psi_0^{(0)} \middle| \sum_{i < j} r_{ij}^{-1} \middle| \Psi_0^{(0)} \right\rangle - \left\langle \Psi_0^{(0)} \middle| \sum_{i} v^{HF}(i) \middle| \Psi_0^{(0)} \right\rangle = -\frac{1}{2} \sum_{ab} (\langle ab | ab \rangle - \langle ab | ba \rangle)$$

 Sum of zeroth and first order energies gives the HF ground state energy

Second Order MP2 Energy

Second order energy of ground state:

•
$$E_0^{(2)} = \sum_{0 \neq n} \frac{\langle 0|V|n\rangle\langle n|V|0\rangle}{E_0^{(0)} - E_n^{(0)}}$$

• Due to orthogonality rules, $\langle \Psi_0 | r_{ij}^{-1} | \Psi_n \rangle$ equals zero except when Ψ_n is a doubly excited state. Therefore:

•
$$E_0^{(2)} = \sum_{a < b, r < s} \frac{\left| \left\langle \Psi_0 \middle| \sum_{i < j} r_{ij}^{-1} \middle| \Psi_{ab}^{rs} \right\rangle \right|^2}{E_0^{(0)} - E_n^{(0)}} = \sum_{a < b, r < s} \frac{\left| \left\langle ab \middle| rs \right\rangle - \left\langle ab \middle| sr \right\rangle \right|^2}{\varepsilon_a + \varepsilon_b - \varepsilon_r - \varepsilon_s}$$

MP2 Algorithm

- Perform SCF calculation to obtain MO energies and coefficients
- 2) Convert 2-electron integrals from AO basis to MO basis
- Use 2-electron integrals to compute MP2 energy over all occupied and unoccupied orbitals

Using MP2

- MP2 is widely available in almost all quantum chemistry programs
- MP2 scales at N⁵ with system size, making it generally not practical for very large systems
- MP2 requires a reasonably good starting point. If HF does not give good results, neither will MP2 (this can sometimes be addressed by using a starting wavefunction other than HF)
- Will almost always give better results than pure HF, but is generally less accurate (although faster) than other post-HF methods like CCSD(T)

Higher Order Terms

- While we've only considered perturbation theory through second order, higher order terms are sometimes used (MP3, MP4, etc.)
- Computational time increases dramatically with higher order perturbations
- Note that very high order perturbations do not necessarily converge, oscillation or divergence are possible