

Hybrid fluid/kinetic modeling for plasma

WENO method for plasma simulation

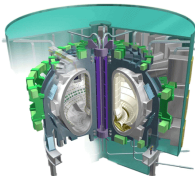
Josselin Massot

Encadré par : Anaïs Crestetto
et Nicolas Crouseilles

2018/12/10

What is plasma or rarefied gas ?

- ▶ System of interacting particles
- ▶ Plasma : hot gas, electrons separated from atoms → electric field
- ▶ Examples of plasmas :
 - ▶ neons, ITER, nebula
- ▶ Examples of rarefied gas :
 - ▶ atmospheric entry (Soyouz, CST-100 Starliner, ...)



Models

Classical models

Hybrid model

Schemes

Time discretization

Space discretization

Numerical test

Validation tests

BoT studying

Conclusion

Outline

Models

- Classical models

- Hybrid model

Schemes

- Time discretization

- Space discretization

Numerical test

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Conclusion

Models

Microscopic model: simulation of all particles

$$(t, x_i(t), v_i(t)), i = 1, \dots, N$$

✓ accuracy ✗ computational time and memory

Macroscopic model: plasma \approx fluid

$$(\rho, u, T)(t, x) \text{ thermodynamic variables}$$

✗ accuracy ✓ computational time and memory

Kinetic model: simulation in phase space

$$f(t, x, v) \text{ distribution of density in phase space}$$

~ accuracy ~ computational time and memory

Study only on macroscopic and kinetic models

Macroscopic model

Euler's equations:

$$\partial_t U + \nabla_x \cdot \mathcal{F}(U) = S_E(U)$$

- ▶ $U = U(t, x)$
- ▶ Vector $U = (\rho, \rho u, e)^T$ thermodynamic variables
- ▶ $S_E(U)$ source term that includes electric field E

Kinetic model

Vlasov-BGK's equation:

$$\partial_t f + v \cdot \nabla_x f + E \cdot \nabla_v f = \frac{1}{\varepsilon} Q(f, f)$$

Transport (x, v) + stiff term $\frac{1}{\varepsilon}$

- ▶ $f = f(t, x, v)$ density distribution in phase space
- ▶ $Q(f, f)$ collision operator (BGK) $Q(f, f) = \mathcal{M}_{[U]} - f$
- ▶ $\mathcal{M}_{[U]}$: velocity distribution at equilibrium
- ▶ $\varepsilon \sim$ mean free path
 - ▶ If $\varepsilon \ll 1 \rightarrow$ Euler equations
 - ▶ If $\varepsilon \gg 1 \rightarrow$ no collision (what I do now)

Relation with macroscopic variables:

$$\int_{\mathbb{R}^d} \begin{pmatrix} 1 \\ v \\ |v|^2 \end{pmatrix} f(t, x, v) \, dv = \int_{\mathbb{R}^d} \begin{pmatrix} 1 \\ v \\ |v|^2 \end{pmatrix} \mathcal{M}_{[U]}(t, x, v) \, dv = U(t, x)$$

Electric field

Poisson's equation

$$\nabla_x \cdot E(x) = \rho(t, x) = \int_{\mathbb{R}^d} f(t, x, v) dv$$

Periodic conditions in space: resolution by FFT

Maxwell's equation (magnetic field) soon! (one day)

Hybrid model

1. Micro-macro model (**mM**)

$$f = \underbrace{\mathcal{M}_{[U]}}_{\text{thermodynamic equilibrium state}} + \underbrace{g}_{\text{gap from equilibrium}}$$

2. Approximation of *micro* part (**mMh**) without interface between models (internship work)

- Transition function $h(t, x)$ between fluid area Ω_F and kinetic area Ω_K

Building hybrid model

1. **mM** consisting of 2 equations:

macro: **mean** in v of kinetic model:

$$\partial_t U + \nabla_x \cdot \mathcal{F}(U) + \nabla_x \cdot \langle vm(v)g \rangle_v = S_E(U)$$

micro: **projection** of kinetic model on image of collision operator $Q(f, f) = \mathcal{M}_{[U]} - f$:

$$\partial_t g + (I - \Pi)[v \cdot \nabla_x (\mathcal{M}_{[U]} + g) + E \cdot \nabla_v (\mathcal{M}_{[U]} + g)] = -\frac{1}{\varepsilon} g$$

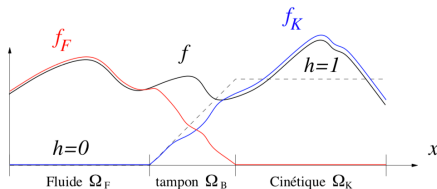
mM model is equivalent to kinetic model

Approximation of *micro* part

2. **Hypothesis:** $f = \mathcal{M}_{[U]}$ on $\Omega_F \Rightarrow g_F = 0$

New model: approximation of **micro**-macro by domain decomposition

$$\Omega = \Omega_F \cup \Omega_K \quad g = (1 - h)g + hg = g_F + g_K$$



Approximated *micro* part

We multiply by h micro part to get:

$$\partial_t g_K + (I - \Pi) \left[v \cdot \nabla_x (\mathcal{M}_{[U]} + g_K) + E \cdot \nabla_v (\mathcal{M}_{[U]} + g_K) \right] = -\frac{1}{\varepsilon} g_K + \frac{g_K}{h} \partial_t h$$

Outside the support of h : $g_K = 0$

Why this model? save computational time (kinetic evaluation **only** on Ω_K)

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Goals

- ▶ Transport in phase space $(x, v) \rightsquigarrow$ scheme easy to use in multi- d (d at least 2)
- ▶ Heavy gradient \rightsquigarrow need high order scheme
- ▶ Stiff term in $\frac{1}{\varepsilon}$, $\varepsilon \in]0, 1]$ \rightsquigarrow need adapted time integrators
- ▶ Long time simulation \rightsquigarrow need stability of space+time scheme

Explicit Euler method

Unstable with 5th-order WENO method:

- ▶ [Wang, R., & Spiteri, R. (2007) SINUM]

Amplification term is small enough to be *controlled* by very small Δt .

If stiff term \rightsquigarrow IMEX:

$$f^{n+1} = f^n - dt(v\partial_x + E\partial_v)(f^n) + \frac{1}{\varepsilon}(\mathcal{M}_{[U^{n+1}]} - f^{n+1})$$

CFL upwind-IMEX: $\frac{\Delta x}{v_{\max}}$

Runge-Kutta 3th order method

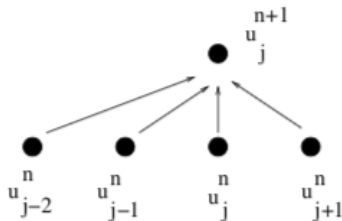
Stable with 5th-order WENO method (see later)

If stiff term \rightsquigarrow exponential formulation:

$$\partial_t(e^{\frac{t}{\varepsilon}}g) + (I - \Pi) \left((v\partial_x + E\partial_v)(e^{\frac{t}{\varepsilon}}(g + \mathcal{M}_{[U]})) \right) = 0$$

and Lawson scheme or IFRK method

Compact scheme



- ▶ High order 6 points scheme ($u_{j-3}^n, \dots, u_{j+2}^n$)
- ▶ Based on **1** polynomial of degree 5

It doesn't work very well (implementation bug ?) and it oscillates with discontinuity (not important for us) or heavy gradient (strong problem for filamentation).

Numerical order

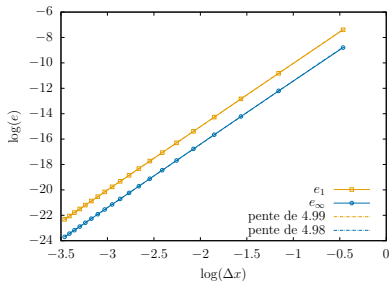
$$\partial_t u + \partial_x u = 0$$

initial: $u(t=0, x) = \cos(x)$

solution: $u(t=t_i, x) = \cos(x - t_i)$

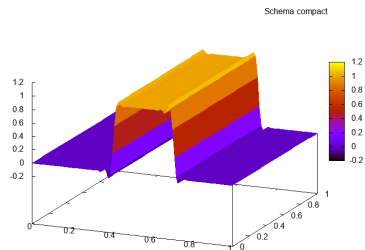
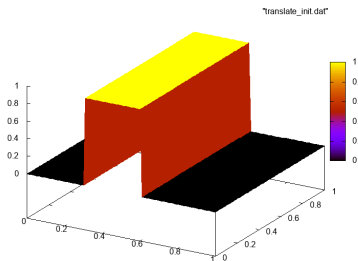
Δx	Δt	T_f
$\frac{2\pi}{N}$	$10^{-5} \Delta x$	1

$N = 10, \dots, 200$



Simplest test

Transport in 1 direction of a discontinuity



WENO method

- ▶ High order 6 points scheme
- ▶ Based on **3** ENO approximations (of lower order) combined with non-linear weights (indicator of smoothness)

BUT: Unstable with explicit Euler method [Wang, R., & Spiteri, R. (2007) SINUM]

Numerical order

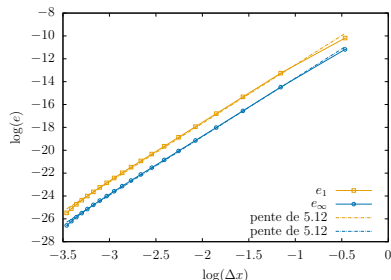
$$\partial_t u + \partial_x u = 0$$

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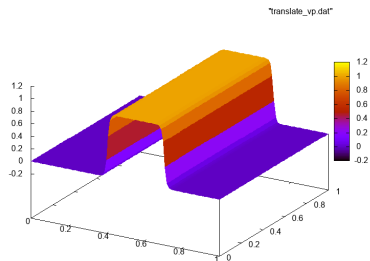
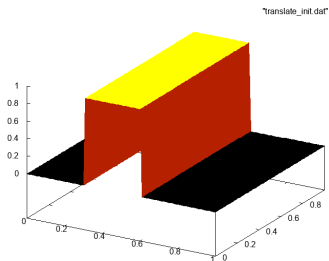
Δx	Δt	T_f
$\frac{2\pi}{N}$	$10^{-5} \Delta x$	1

$N = 10, \dots, 200$



Simplest test

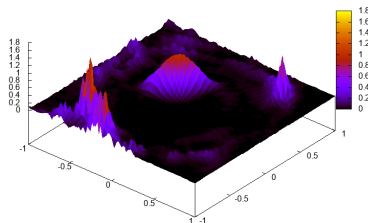
Transport in 1 direction of a discontinuity



Instability illustration

Gaussian rotation

Schéma d'Euler explicite

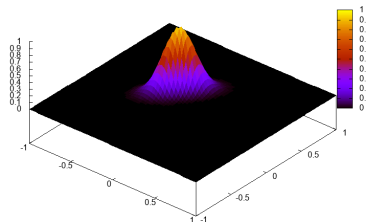


Explicit Euler method

grid: 100×100

[Wang, R., & Spiteri, R. (2007) SINUM]

Schéma RK3



RK3 method

$\Delta t = 0.3 \Delta x$

$T_f = 15$

Steps of proof of instability

1. Von Neumann analysis : $f_{j+k}^n \rightarrow e^{ik\phi}$ ($\phi \equiv \kappa\pi\Delta x$)
2. Linearized WENO scheme (weight $\approx \mathcal{O}(\Delta x^2)$):

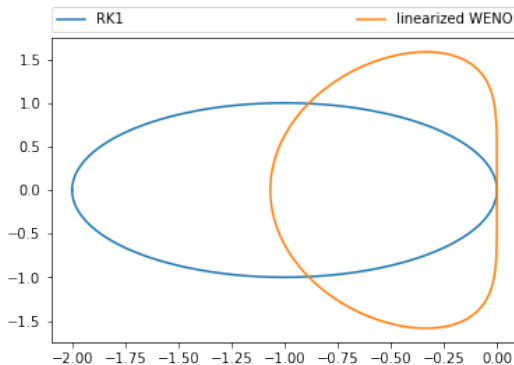
$$z(\phi) = \tilde{z}(\phi) + M(\epsilon_i, \phi)$$

with $M(\epsilon_i, \phi) = \mathcal{O}(\Delta x^2)$

3. Draw $\tilde{z}(\phi)$ with some RK stability curve and conclude.

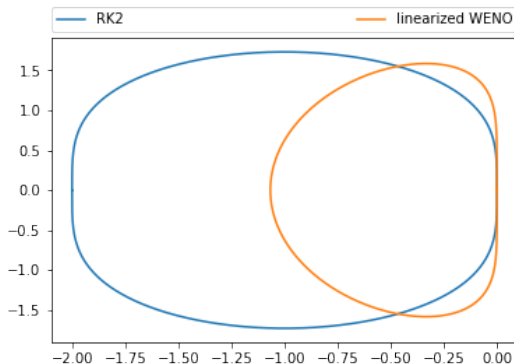
Stability of RKN-WENO

RK1-WENO



Stability of RKN-WENO

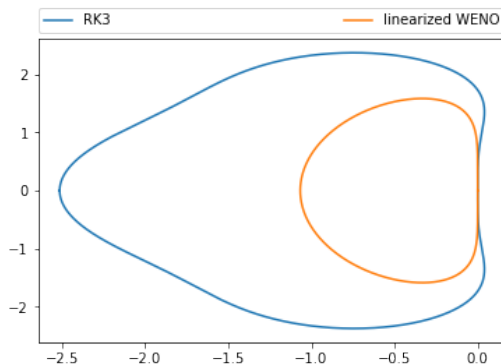
RK2-WENO



possible stability: $\Delta t = 1.73\Delta x^{5/3}$ [Motamed, M., & Macdonald, C. & Ruuth S. (2010) J. Sci. Comput.]

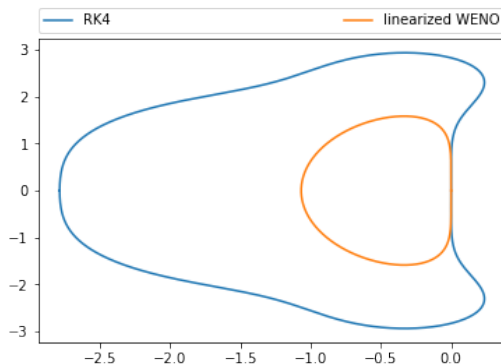
Stability of RKN-WENO

RK3-WENO (CFL: 1.433)



Stability of RKN-WENO

RK4-WENO (CFL: 1.731)



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Sod shock tube

Fluid regime validation

boundary: Neumann in space x , periodic in velocity v

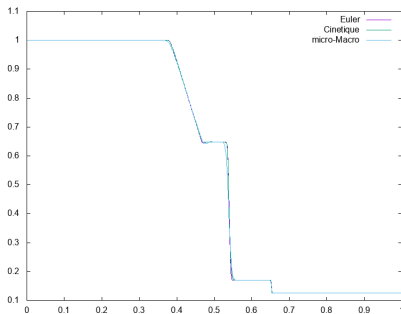
initial condition: discontinuity in ρ and T :

$$U(t=0, x) = \begin{cases} U_L = (\rho_L, u_L, T_L) = (1, 0, 1) & , x \leq \frac{1}{2} \\ U_R = (\rho_R, u_R, T_R) = (0.125, 0, 0.8) & , x > \frac{1}{2} \end{cases}$$

$$f(t=0, x, v) = \mathcal{M}_{[U(t=0, x)]}(x, v) \quad g(t=0, x, v) = 0$$

Models and schemes validation

$$\rho(t = 0.067, x)$$



$$\text{domain: } [0, 1] \times [-18, 18]$$

$$\Delta x = 10^{-3}$$

$$\Delta v = 0.5625$$

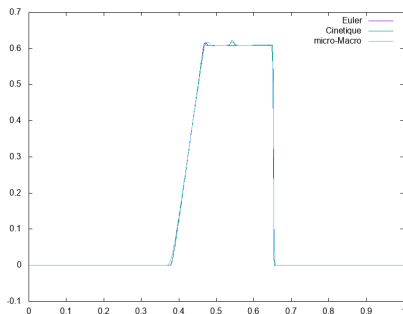
$$\text{grid: } 1000 \times 64$$

$$\Delta t = \frac{1}{2} \frac{\Delta x}{v_{\max}} = 2.77 \cdot 10^{-5}$$

$$\varepsilon = 10^{-4}$$

Models and schemes validation

$$u(t = 0.067, x)$$



$$\text{domain: } [0, 1] \times [-18, 18]$$

$$\Delta x = 10^{-3}$$

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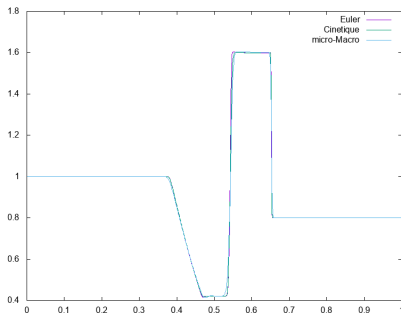
$$\text{grid: } 1000 \times 64$$

$$\Delta t = \frac{1}{2} \frac{\Delta x}{v_{\max}} = 2.77 \cdot 10^{-5}$$

$$\varepsilon = 10^{-4}$$

Models and schemes validation

$$T(t = 0.067, x)$$



$$\text{domain: } [0, 1] \times [-18, 18]$$

$$\Delta x = 10^{-3}$$

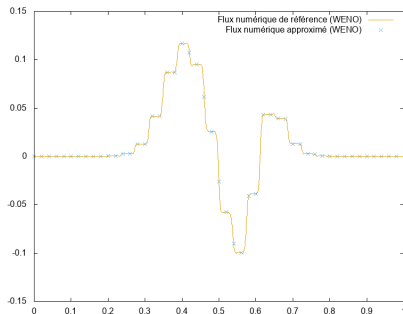
$$\Delta v = 0.5625$$

$$\text{grid: } 1000 \times 64$$

$$\Delta t = \frac{1}{2} \frac{\Delta x}{v_{\max}} = 2.77 \cdot 10^{-5}$$

$$\varepsilon = 10^{-4}$$

Simulation : Sod shock tube, kinetic mode: $\varepsilon = 1$

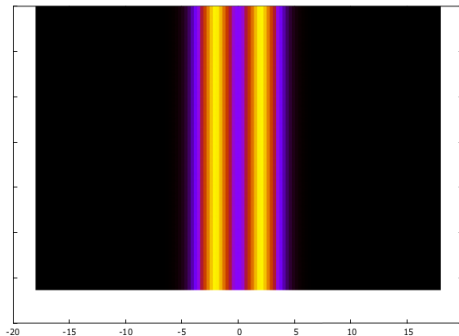


domain: $[0, 1] \times [-18, 18]$
 $\Delta x = 10^{-3}$ $\Delta v = 0.5625$

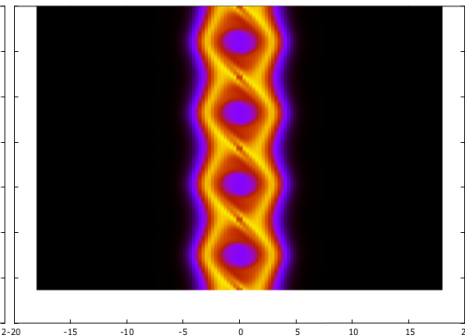
Computing time: divided by 2

grid: 1000×64 $\varepsilon = 1$
 $\Delta t = \frac{1}{2} \frac{\Delta x}{v_{\max}} = 2.77 \cdot 10^{-5}$

Numerical test: two streams



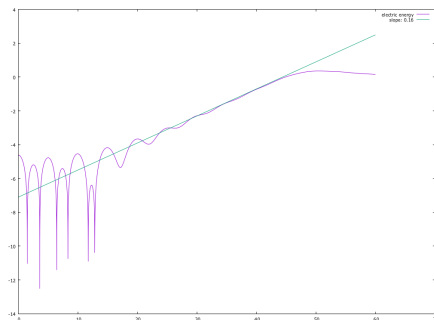
Initial condition



Result

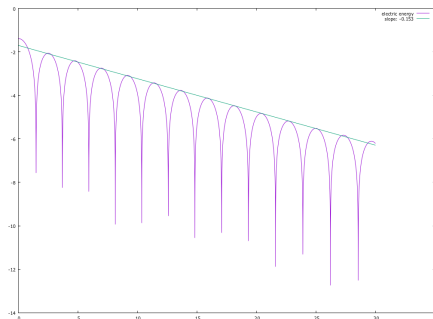
Numerical test: two streams

Electric energy



Numerical test: Landau damping

Electric energy

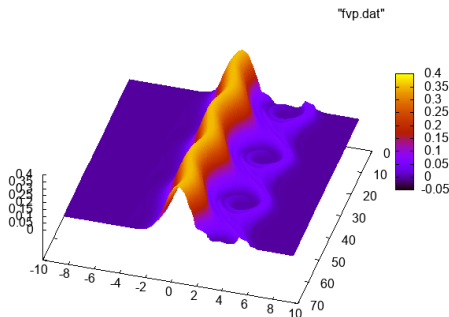


Bump on Tail (BoT)

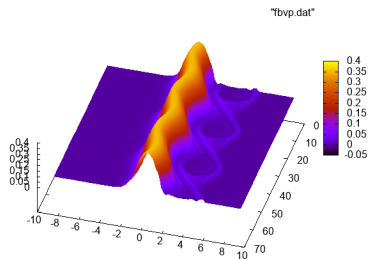
Cold and hot particles splitting (not same as micro-macro splitting)

$$f = f_c + f_h \quad f = \mathcal{M}_{[U]} + g$$

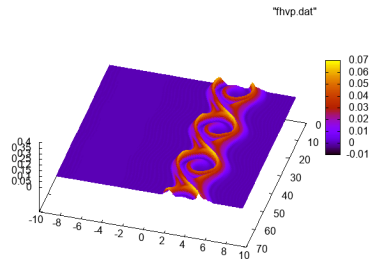
What we expect: $f_c = \mathcal{M}_{[U]}$, $f_h = g \dots$



BoT: f_c , f_h



f_c

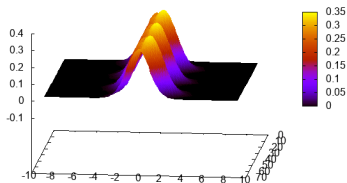


f_h

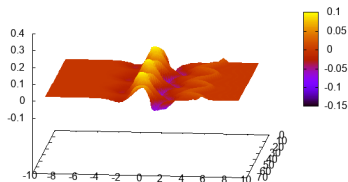
BoT: f_c mM, f_h kinetic

mM on $f_c = \mathcal{M}_{[U_c]} + g_c$, **kinetic** on f_h

"Mbvp.dat"



"gbvp.dat"

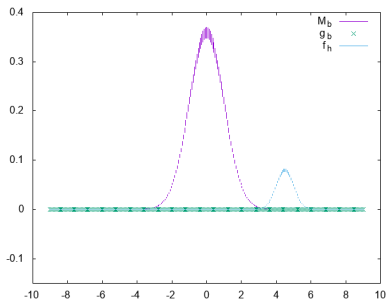


$\mathcal{M}_{[U_c]}$

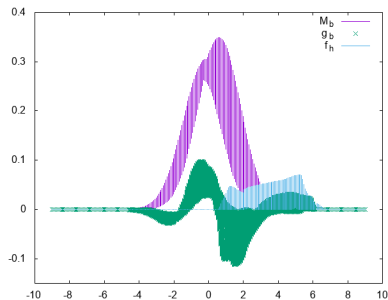
g_c

BoT: f_c mM, f_h kinetic

mM on $f_c = \mathcal{M}_{U_c} + g_c$, **kinetic** on f_h



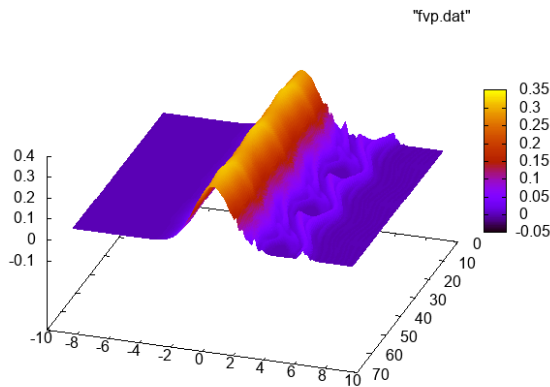
Initial condition



Result

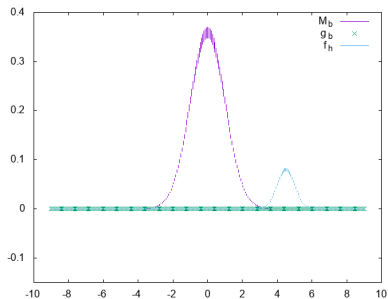
BoT: f_c fluid approximation

approximation : $f_c = \mathcal{M}_{[U_c]}, g_c = 0$

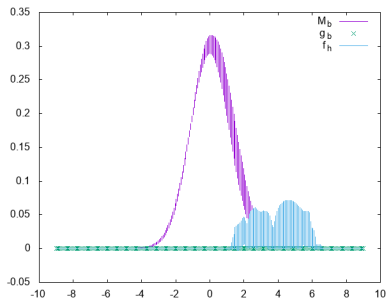


BoT: f_c fluid approximation

approximation : $f_c = \mathcal{M}_{[u_c]}, g_c = 0$



Initial condition



Result

Ideas

- ▶ Physicist idea (IPP Garching): $f_c = \delta_{v-u_c}$
- ▶ Computer scientist idea: reduce grid in v around effective data for f_h

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Conclusion

- ✓ WENO approved!
 - ▶ High order, no oscillation, no problem in multi-D (CFL ?)
 - ▶ Works well with RK3 (even with stiff term)
 - ▶ Could be use for Euler part

Future works

- ▶ Still a grid in phase space for 1Dx 3Dv or 3Dx 3Dv? (MC? PIC?)
- ▶ Code refactoring: some optimization and adaptive grid (no global variables) (Julia? C++?)
- ▶ Better understanding of SSPRK(3,3) diffusion, stability of couple of time-space schemes
- ▶ Automatic study of different schemes (*SymPy* & *NodePy*)
- ▶ Approve approximation of f_c with conservative variables, and integrate Dirac modeling

Tank you for your attention

Outline

Construction de la fonction $h(t, x)$

Euler's equations

WENO scheme

Discrétisation du modèle micro-macro approximé

Détermination *a priori* de la fonction h :

Zone *hors équilibre* au cours du temps :

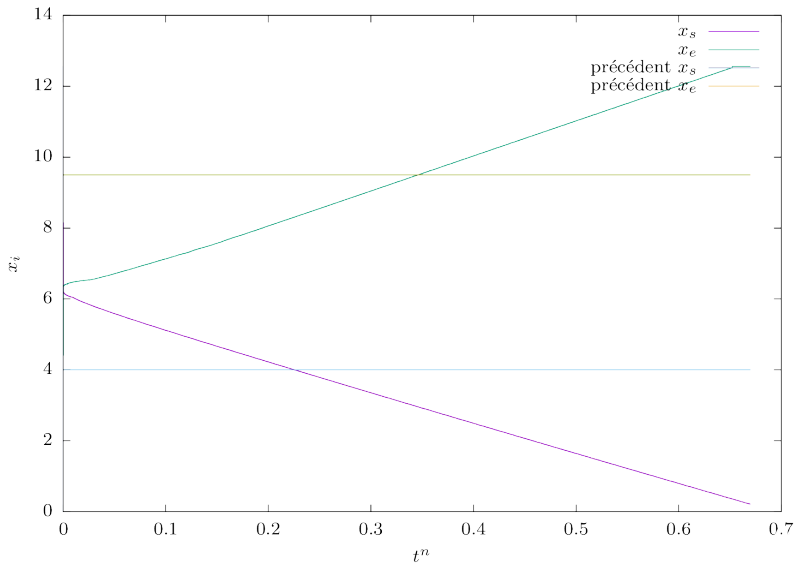
$$\text{supp } h(t, x) = \Omega_K(t)$$

Étude du support numérique (seuil à 10^{-5}) de $(G_i^n)_3$:

$$\mathcal{J}^n = \{i \in \llbracket 0, N_x \rrbracket, |\langle vm(v)g^n \rangle_{x=x_i}| > 10^{-5}\}$$

Étude de g_K limitée à $\llbracket i_s^n, i_e^n \rrbracket = \mathcal{J}^n$

Étude du support de h



Outline

Construction de la fonction $h(t, x)$

Euler's equations

WENO scheme

Euler's equations

$$\partial_t U + \nabla_x \cdot \mathcal{F}(U) = S_E(U)$$

$$\mathcal{F}(U) = \begin{pmatrix} \rho u \\ \rho u \otimes u + p \mathbb{I}_d \\ u(e + p) \end{pmatrix} \quad S(U) = \begin{pmatrix} 0 \\ \rho E \\ 2\rho u E \end{pmatrix}$$

pressure: $p = 2(e - \frac{1}{2}\rho|u|^2)$

Outline

Construction de la fonction $h(t, x)$

Euler's equations

WENO scheme

WENO scheme

Model:

$$\partial_t u + \partial_x (au) = 0$$

We would approximate $\partial_x (au)|_{x=x_i, v=v_k}$:

$$\partial_x (au)|_{x=x_i, v=v_k} \approx \frac{1}{\Delta_X} (\hat{u}_{i+\frac{1}{2},k} - \hat{u}_{i-\frac{1}{2},k})$$

WENO flux

$$\begin{aligned}\hat{u}_{i+\frac{1}{2},k}^+ &= w_0^+ \left(\frac{2}{6} u_{i-2,k}^+ - \frac{7}{6} u_{i-1,k}^+ + \frac{11}{6} u_{i,k}^+ \right) \\ &\quad + w_1^+ \left(-\frac{1}{6} u_{i-1,k}^+ + \frac{5}{6} u_{i,k}^+ + \frac{2}{6} u_{i+1,k}^+ \right) \\ &\quad + w_2^+ \left(\frac{2}{6} u_{i,k}^+ + \frac{5}{6} u_{i+1,k}^+ - \frac{1}{6} u_{i+2,k}^+ \right)\end{aligned}$$

and

$$\begin{aligned}\hat{u}_{i+\frac{1}{2},k}^- &= w_2^- \left(-\frac{1}{6} u_{i-1,k}^- + \frac{5}{6} u_{i,k}^- + \frac{2}{6} u_{i+1,k}^- \right) \\ &\quad + w_1^- \left(\frac{2}{6} u_{i,k}^- + \frac{5}{6} u_{i+1,k}^- - \frac{1}{6} u_{i+2,k}^- \right) \\ &\quad + w_0^- \left(\frac{11}{6} u_{i+1,k}^- - \frac{7}{6} u_{i+2,k}^- + \frac{2}{6} u_{i+3,k}^- \right)\end{aligned}$$

WENO weights

$$w_n^\pm = \frac{\tilde{w}_n^\pm}{\sum_{m=0}^2 \tilde{w}_m^\pm}, \quad \tilde{w}_n^\pm = \frac{\gamma_n}{(\epsilon + \beta_n^\pm)^2}$$

$$\gamma_0 = \frac{1}{10}, \quad \gamma_1 = \frac{3}{5}, \quad \gamma_2 = \frac{3}{10}$$

$\epsilon = 10^{-6}$ numerical value to prevent the denominator from being 0

WENO Indicator of Smoothness

$$\beta_0^+ = \frac{13}{12}(u_{i-2,k}^+ - 2u_{i-1,k}^+ + u_{i,k}^+)^2 + \frac{1}{4}(u_{i-2,k}^+ - 4u_{i-1,k}^+ + 3u_{i,k}^+)^2$$

$$\beta_1^+ = \frac{13}{12}(u_{i-1,k}^+ - 2u_{i,k}^+ + u_{i+1,k}^+)^2 + \frac{1}{4}(u_{i-1,k}^+ - u_{i+1,k}^+)^2$$

$$\beta_2^+ = \frac{13}{12}(u_{i,k}^+ - 2u_{i+1,k}^+ + u_{i+2,k}^+)^2 + \frac{1}{4}(3u_{i,k}^+ - 4u_{i+1,k}^+ + u_{i+2,k}^+)^2$$

and

$$\beta_0^- = \frac{13}{12}(u_{i+1,k}^- - 2u_{i+2,k}^- + u_{i+3,k}^-)^2 + \frac{1}{4}(3u_{i+1,k}^- - 4u_{i+2,k}^- + u_{i+3,k}^-)^2$$

$$\beta_1^- = \frac{13}{12}(u_{i,k}^- - 2u_{i+1,k}^- + u_{i+2,k}^-)^2 + \frac{1}{4}(u_{i,k}^- - u_{i+2,k}^-)^2$$

$$\beta_2^- = \frac{13}{12}(u_{i-1,k}^- - 2u_{i,k}^- + u_{i+1,k}^-)^2 + \frac{1}{4}(u_{i,k}^- - 4u_{i,k}^- + 3u_{i+1,k}^-)^2$$