

Electron Hybrid model for R/L-waves in cold plasmas

- $\frac{\partial \mathbf{u}_c}{\partial t} + (\mathbf{u}_c \cdot \nabla) \mathbf{u}_c = \frac{q_e}{m_e} (\mathbf{E} + \mathbf{u}_c \times \mathbf{B})$
- $\frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E}$
- $\frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t} = \nabla \times \mathbf{B} - \mu_0 (\mathbf{j}_c + \mathbf{j}_h)$
- $\frac{\partial f_h}{\partial t} + \mathbf{v} \cdot \nabla f_h + \frac{q_e}{m_e} (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \nabla_{\mathbf{v}} f_h = 0$
- $\nabla \cdot \mathbf{B} = 0 \quad \text{at } t = 0$
- $\nabla \cdot \mathbf{E} = \frac{1}{\epsilon_0} [q_i n_i + q_e (n_c + n_h)] \quad \text{at } t = 0$

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- Linearization about an equilibrium state:

- $n_c(\mathbf{x}, t) = n_{c,0}(\mathbf{x}) + \tilde{n}_c(\mathbf{x}, t)$

- $\mathbf{u}_c(\mathbf{x}, t) = \tilde{\mathbf{u}}_c(\mathbf{x}, t)$

- $\mathbf{B}(\mathbf{x}, t) = \mathbf{B}_0(\mathbf{x}) + \tilde{\mathbf{B}}(\mathbf{x}, t)$

- $\mathbf{E}(\mathbf{x}, t) = \tilde{\mathbf{E}}(\mathbf{x}, t)$

→ cold current density $\mathbf{j}_c = q_e (n_{c,0} + \tilde{n}_c) \tilde{\mathbf{u}}_c \approx q_e n_{c,0} \tilde{\mathbf{u}}_c$

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- $\frac{\partial \tilde{\mathbf{j}}_c}{\partial t} = \epsilon_0 \Omega_{pe}^2 \tilde{\mathbf{E}} + \tilde{\mathbf{j}}_c \times \boldsymbol{\Omega}_{ce}$
- $\frac{\partial \tilde{\mathbf{B}}}{\partial t} = -\nabla \times \tilde{\mathbf{E}}$
- $\frac{1}{c^2} \frac{\partial \tilde{\mathbf{E}}}{\partial t} = \nabla \times \tilde{\mathbf{B}} - \mu_0 (\tilde{\mathbf{j}}_c + \mathbf{j}_h)$
- $\frac{\partial f_h}{\partial t} + \mathbf{v} \cdot \nabla f_h + \frac{q_e}{m_e} (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \nabla_{\mathbf{v}} f_h = 0$

$$\Omega_{pe}^2 = \frac{n_{c,0} e^2}{\epsilon_0 m_e}$$

$$\boldsymbol{\Omega}_{ce} = \frac{q \mathbf{B}_0}{m}$$

→ cold plasma model + energetic kinetic electrons

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The dynamical equations conserve exactly the following energy:

$$\begin{aligned}\epsilon &:= \frac{1}{2} \left[\epsilon_0 \int \tilde{\mathbf{E}}^2 d^3\mathbf{x} + \frac{1}{\mu_0} \int \tilde{\mathbf{B}}^2 d^3\mathbf{x} + \frac{1}{\epsilon_0 \Omega_{pe}^2} \int \tilde{\mathbf{j}}_c^2 d^3\mathbf{x} + m_e \int f_h v^2 d^3\mathbf{v} d^3\mathbf{x} \right] \\ &= \epsilon_{em} + \epsilon_c + \epsilon_h\end{aligned}$$

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Dispersion relation for ...

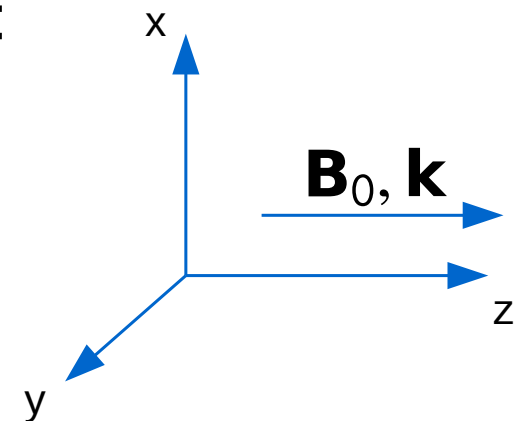
- ... a homogeneous background plasma:
- ... a uniform magnetic background field:
- ... parallel wave propagation:
- ... a uniform equilibrium hot electron distribution:

$$n_{c,0} = \text{const.}$$

$$\mathbf{B}_0 = B_0 \mathbf{e}_z$$

$$\mathbf{k} \parallel \mathbf{B}_0 \parallel \mathbf{e}_z$$

$$f_h(\mathbf{x}, \mathbf{v}, t) = f_h^0(\mathbf{v}) + \tilde{f}_h(\mathbf{x}, \mathbf{v}, t)$$



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- Solution of the fully linearized problem by plane wave ansatz:
 - $\nabla \rightarrow ik\mathbf{e}_z$
 - $\frac{\partial}{\partial t} \rightarrow -i\omega$
- 3 types of solutions of which
 - 2 correspond to electromagnetic waves with $\tilde{\mathbf{E}}, \tilde{\mathbf{B}} \perp \mathbf{k}$
 - 1 corresponds to electrostatic waves with $\tilde{\mathbf{E}} \parallel \mathbf{k}, \quad \tilde{\mathbf{B}} = 0$

Electron Hybrid model for R/L-waves in cold plasmas

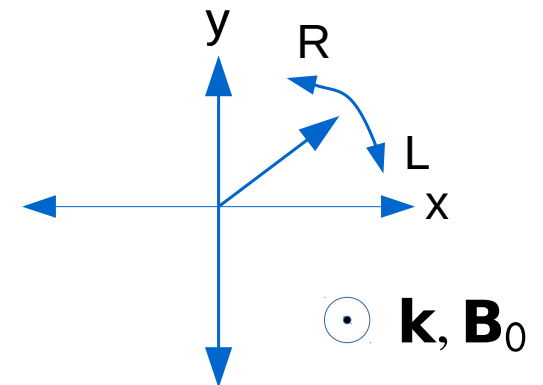
- Dispersion relation for electromagnetic waves: $\tilde{\mathbf{E}}, \tilde{\mathbf{B}} \perp \mathbf{k}$

$$D_{R/L}(k, \omega) = 1 - \frac{c^2 k^2}{\omega^2} - \frac{\Omega_{pe}^2}{\omega(\omega \pm \Omega_{ce})} + \nu_h \frac{\Omega_{pe}^2}{\omega} \int d^3 \mathbf{v} \frac{v_{\perp}}{2} \frac{\hat{G} f_h^0(v_{\parallel}, v_{\perp})}{\omega \pm \Omega_{ce} - k v_{\parallel}} = 0$$

$$\hat{G} = \frac{\partial}{\partial v_{\perp}} + \frac{k}{\omega} \left(v_{\perp} \frac{\partial}{\partial v_{\parallel}} - v_{\parallel} \frac{\partial}{\partial v_{\perp}} \right), \quad \nu_h = n_{h,0}/n_{c,0}, \quad d^3 \mathbf{v} = dv_{\parallel} dv_{\perp} v_{\perp} 2\pi$$

- Wave-particle interaction due to electrons with $k v_{\parallel} = \omega \pm \Omega_{ce}$

- Solutions correspond to
 - Right-hand circularly polarized waves (R)
 - Left-hand circularly polarized waves (L)



Electron Hybrid model for R/L-waves in cold plasmas

Example: anisotropic Maxwellian

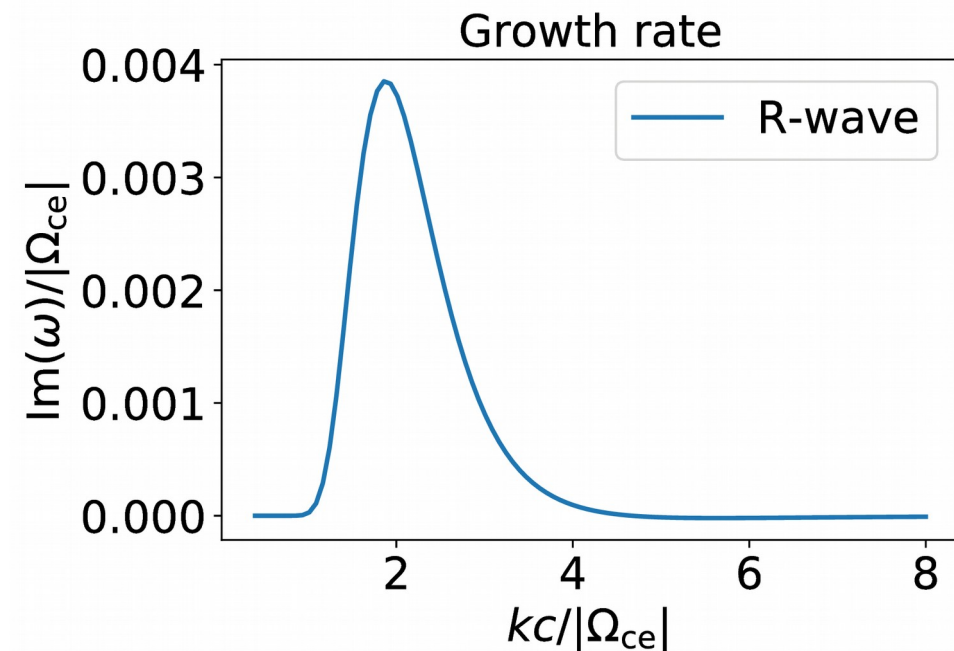
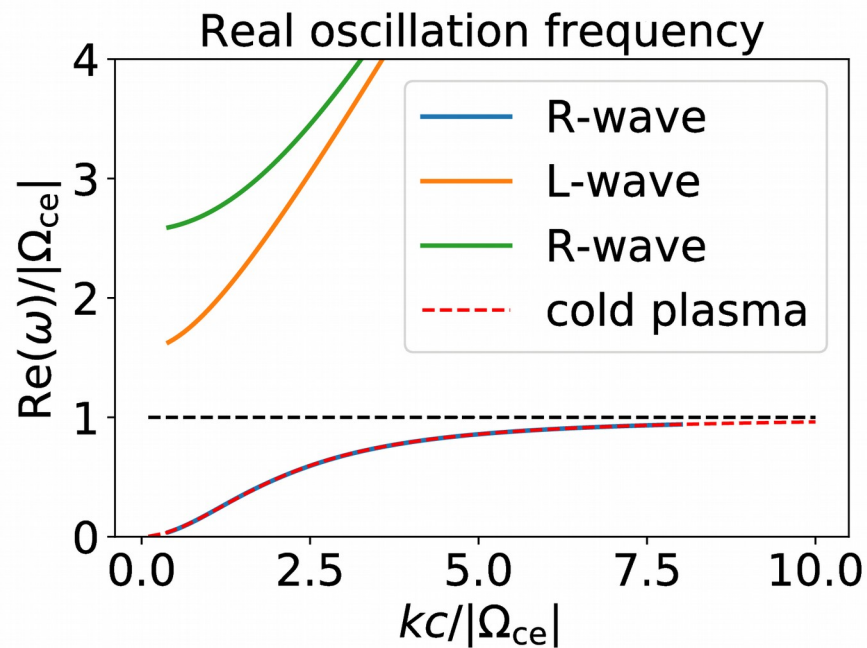
$$f_h^0(v_{\parallel}, v_{\perp}) = \frac{n_{h,0}}{(2\pi)^{3/2} w_{\parallel} w_{\perp}^2} \exp\left(-\frac{v_{\parallel}^2}{2w_{\parallel}^2} - \frac{v_{\perp}^2}{2w_{\perp}^2}\right)$$

$$w_{\parallel} = 0.2 c$$

$$w_{\perp} = 0.5 c$$

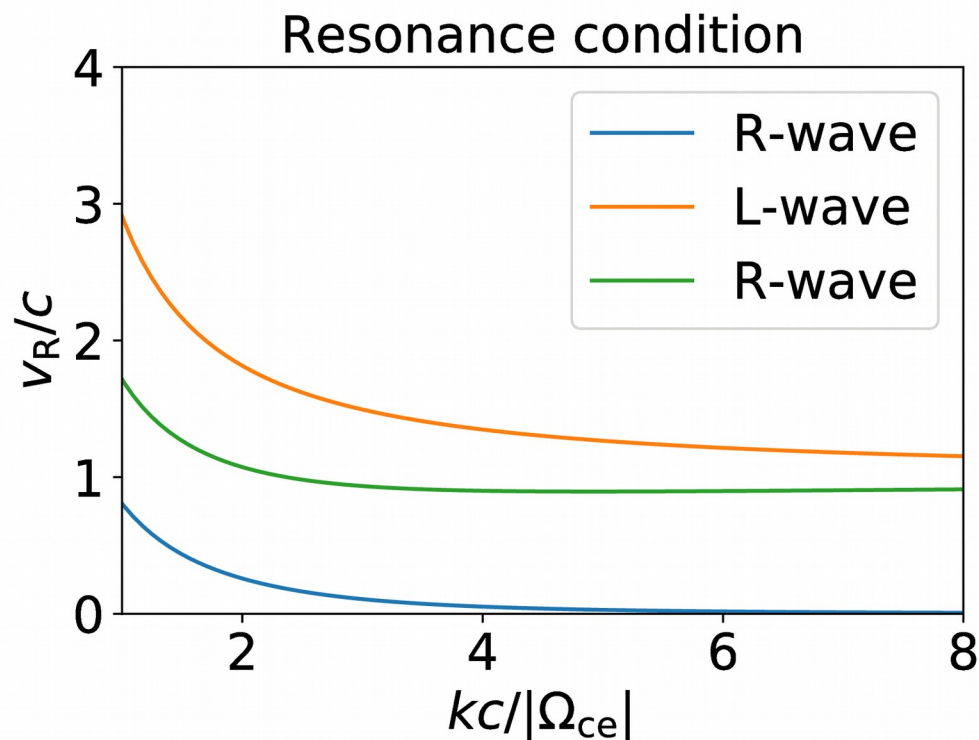
$$\nu_h = 5 \cdot 10^{-3}$$

$$\Omega_{pe} = 2|\Omega_{ce}|$$

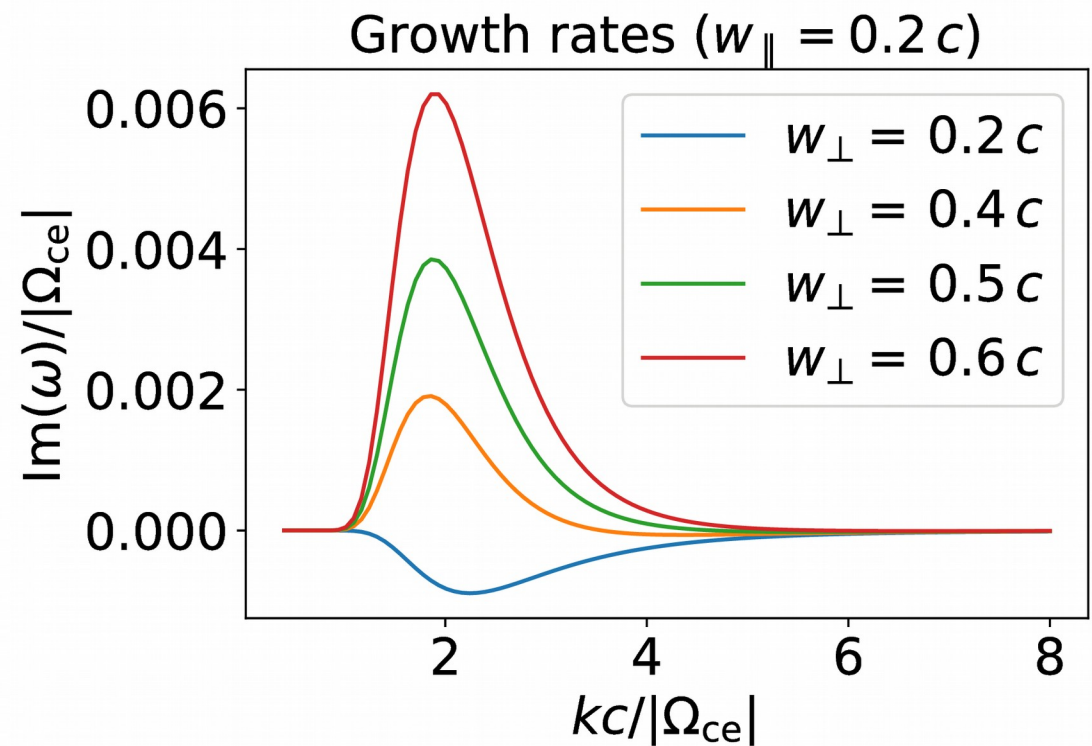


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Example: anisotropic Maxwellian



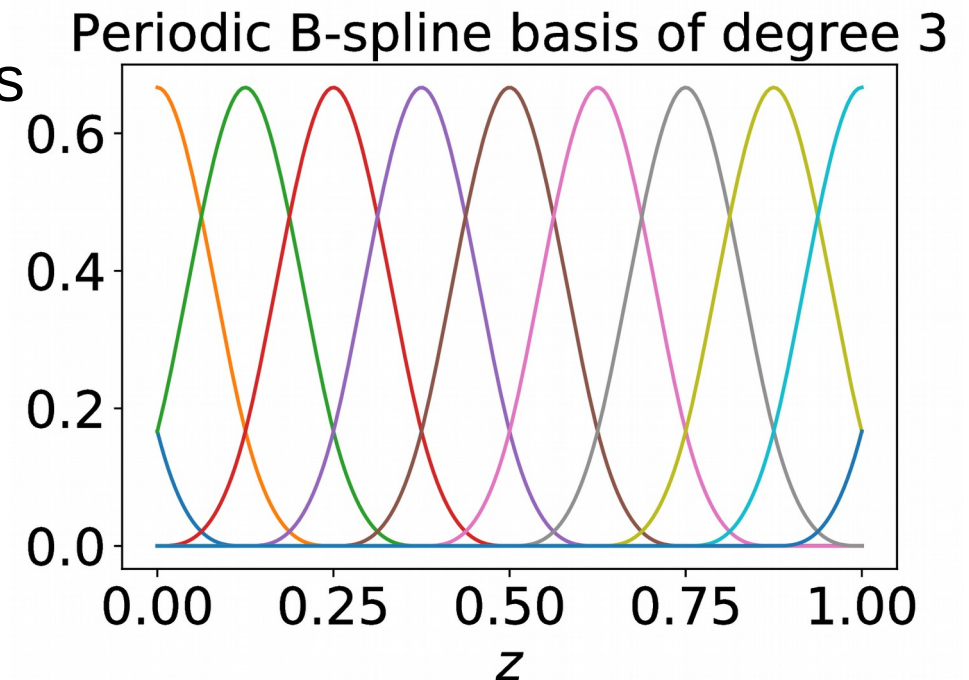
→ no wave-particle
interaction for R/L-waves
with $\omega \gg |\Omega_{ce}|$



Implementation of the model in Python:

- 1d B-spline Finite Elements for fields

$$\mathbf{U} = (E_x, E_y, B_x, B_y, j_{c,x}, j_{c,y})$$

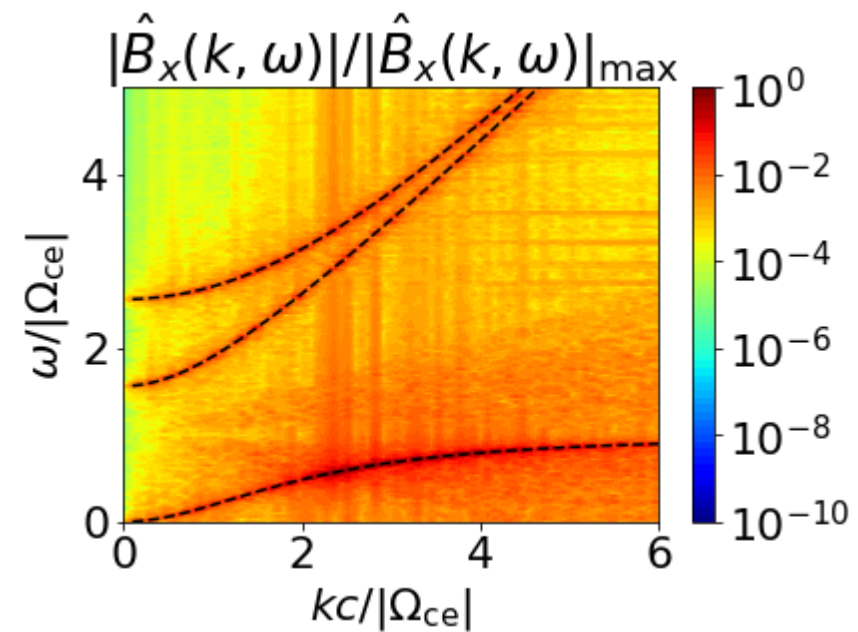
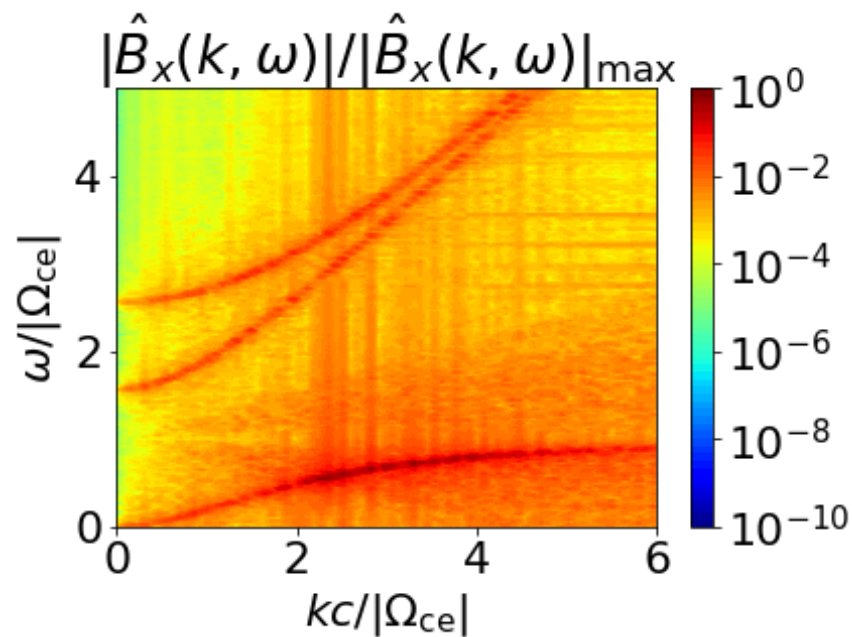


- 1d3v Particle-In-Cell with control variate and Boris particle pusher
- Implicit Crank-Nicolson time stepping scheme

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Simulation Results

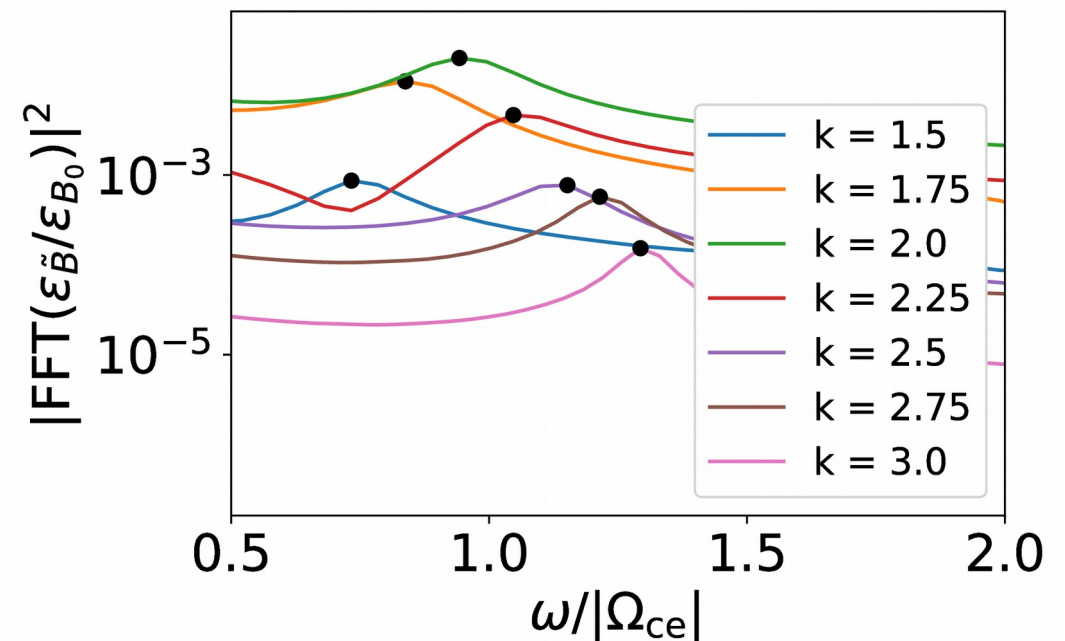
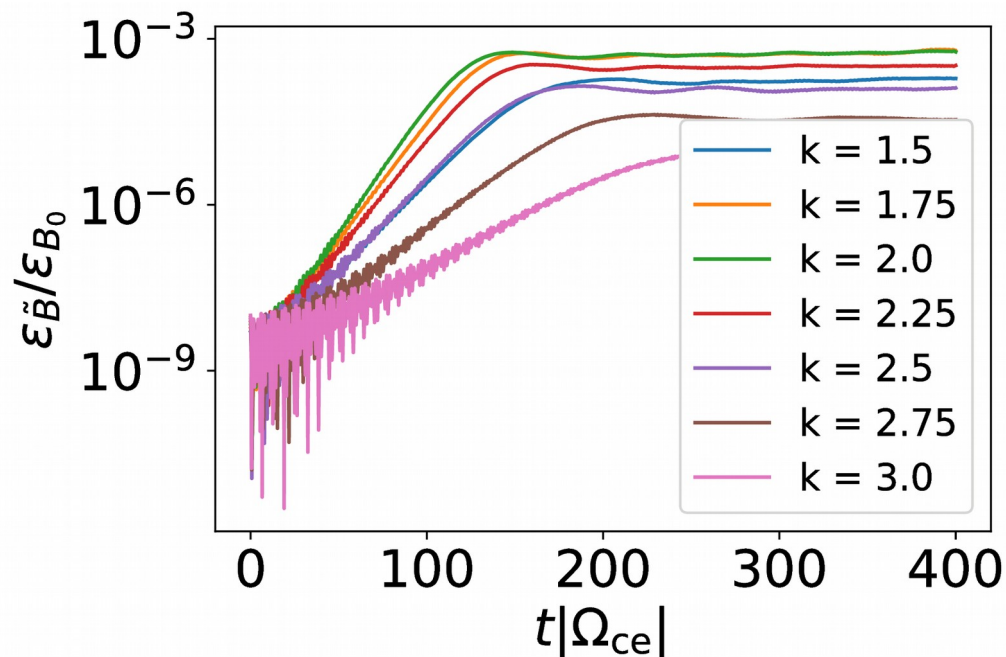
Test run with low density, isotropic Maxwellian to create thermal noise
→ no growth expected



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Simulation Results:

Runs with perturbed magnetic field of the form $B_x^0(z) = a \cdot \sin(kz)$ for different wavenumbers



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